

Title: Cusp Anomaly, Integrability and ADS/CFT

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Abstract: TBA

CUSP ANOMALY, INTEGRABILITY AND ADS/CFT

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Princeton University

ANOMALOUS DIMENSIONS AND SPIN CHAINS

HEISENBERG MODEL FOR THE $SU(2)$ SECTOR

BETHE ANSATZ

$PSU(2,2|4)$ SYMMETRY IN THE FULL $N=4$ SYM

BES/BHL "DRESSING" PHASE

INTEGRAL EQUATION

NUMERICAL STUDY

AdS/CFT, ANOMALOUS DIMENSIONS AND STRINGS

BASIC EXAMPLE:

$$N=4 \text{ SYM} \iff \text{AdS}_5 \times S^5 \text{ TYPE IIB}$$

SCALING DIMENSION
OF
GAUGE INVARIANT
OPERATORS \iff

ENERGIES
OF
STRING STATES

$$\Delta(\lambda)$$

$$E(\lambda)$$

PERTURBATIVE
IF $\lambda \ll 1$

PERTURBATIVE
IF $\lambda \gg 1$

$$\text{tr}(-x y D^2 y x z \dots)$$

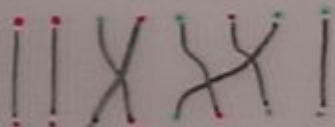


(WE WILL ALWAYS BE IN THE LARGE N LIMIT)

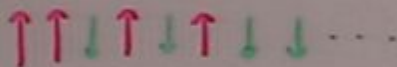
$$\lambda = g_{\text{YM}}^2 N$$

WE WANT TO COMPUTE
THE ANOMALOUS DIMENSION
OF OPERATORS LIKE
 $\text{tr}(x x y x y x y y)$

FEYNMAN
DIAGRAMS



WE CONSIDER THE QUANTUM MECHANICS OF J SPINS



$SU(2)$ SPIN CHAIN OF LENGTH J

HAMILTONIAN:

$$\mathcal{H} = \lambda \sum_{i=1}^J (1 - P_{i,i+1}) + \leftarrow \text{1-LOOP: HEISENBERG MODEL}$$

$$+ \lambda^2 \sum (4\mathbb{1} + 6 P_{i,i+1} - P_{i,i+1} P - P \cdot P)$$

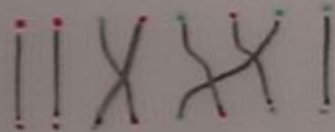
$$+ \lambda^3 \dots$$

FROM CFT'S TO SPIN CHAINS

WE WANT TO COMPUTE
THE ANOMALOUS DIMENSION
OF OPERATORS LIKE

$$\text{tr}(x \times y \times y \times y)$$

FEYNMAN
DIAGRAMS



WE CONSIDER THE QUANTUM MECHANICS OF J SPINS



$SU(2)$ SPIN CHAIN OF LENGTH J

HAMILTONIAN:

$$\begin{aligned} \mathcal{H} = & \lambda \sum_{i=1}^J (1 - P_{i,i+1}) + \leftarrow \text{1-LOOP: HEISENBERG MODEL} \\ & + \lambda^2 \sum (4\mathbb{1} + 6 P_{i,i+1} - P_{i,i+1} P - P \cdot P) \\ & + \lambda^3 \dots \end{aligned}$$



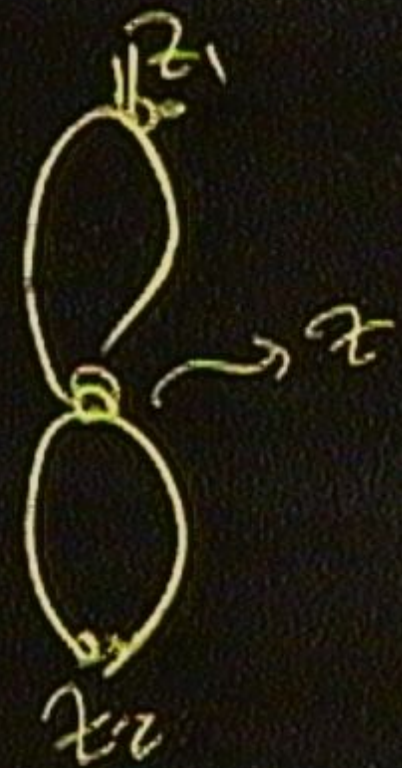
$n \times n$



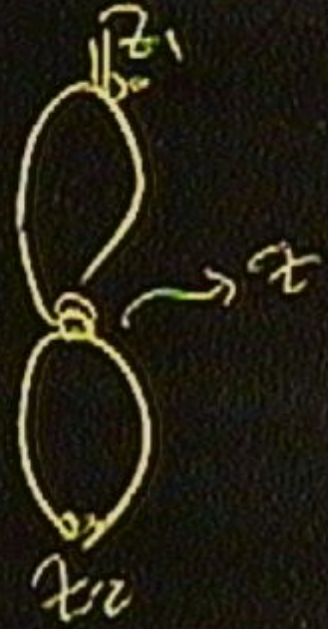
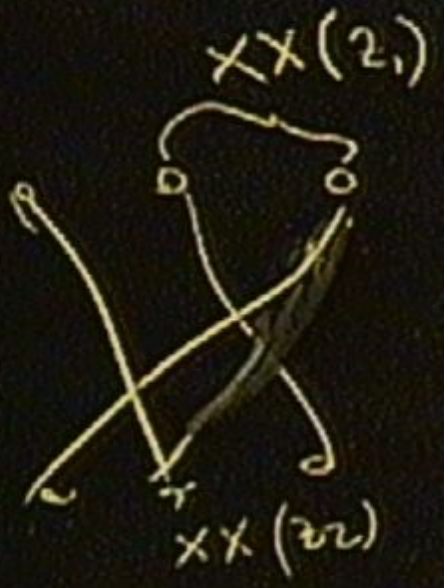
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SPECTRUM OF SPIN CHAINS: BETHE ANSATZ

GROUND STATE: $\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow$

1ST EXCITATIONS: $\uparrow\uparrow\uparrow\downarrow\uparrow\uparrow\uparrow$ IMPURITY TRAVELS
(SINGLE MAGNON) WITH MOMENTUM p

$$E(p) = 4 \sin^2(p/2) \quad 0 \leq p < 2\pi$$

2-MAGNON STATES:

$$|\Psi\rangle = \sum_{x_1, x_2} \psi(x_1, x_2) |\uparrow\uparrow\downarrow\uparrow\uparrow\uparrow\downarrow\downarrow\downarrow\rangle$$

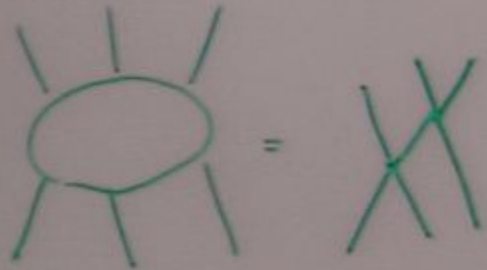
$$\psi(x_1, x_2) = e^{i p_1 x_1 + i p_2 x_2} + \underset{\substack{\uparrow \\ \text{S-MATRIX}}}{S(p_1, p_2)} e^{i p_1 x_2 + i p_2 x_1}$$

$$S(p_1, p_2) = \frac{2e^{i p_1} - e^{i p_1 + i p_2} - 1}{2e^{i p_2} - e^{i p_1 + i p_2} - 1}$$

THIS DESCRIBES $2 \rightarrow 2$ SCATTERING
OF MAGNONS

MULTI-MAGNON STATES

IF THE HAMILTONIAN IS INTEGRABLE,
THE SCATTERING OF MANY MAGNONS
FACTORIZES IN $2 \rightarrow 2$ SCATTERING



THE ALLOWED MOMENTA
ARE OBTAINED BY THE
BETHE EQUATIONS

$$e^{ip_n L} = \prod_{\substack{j=1 \\ j \neq n}}^M S(p_n, p_j) \quad k=1, \dots, M$$

FULL $N=4$ SYM

BEISERT
08

$PSU(2,2|4)$ SYMMETRY

$$16 = 8_B + 8_F \text{ EXCITATIONS}$$

- DISPERSION RELATION:

$$E(p) = -1 + \sqrt{1 + 8\lambda \sin^2(p/2)}$$

VALID AT LARGE J
FOR ANY λ

- S-MATRIX ($16^2 \times 16^2$)

FIXED BY SYMMETRIES
MODULO A SCALAR FACTOR

$$S(p_1, p_2; \lambda) = e^{i\theta}$$

↑
"DRESSING" PHASE

MAIN REQUIREMENT: CROSSING SYMMETRY

JANIK
05

FULL $N=4$ SYM

BEISERT
08

PSU(2,2|4) SYMMETRY

$$16 = 8_B + 8_F \text{ EXCITATIONS}$$

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FOR ANY λ

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FIXED BY SYMMETRIES
MODULO A SCALAR FACTOR

$$\sigma(p_1, p_2; \lambda) = e^{i\theta}$$

↑

"DRESSING" PHASE

MAIN REQUIREMENT: CROSSING SYMMETRY

JANIK
06

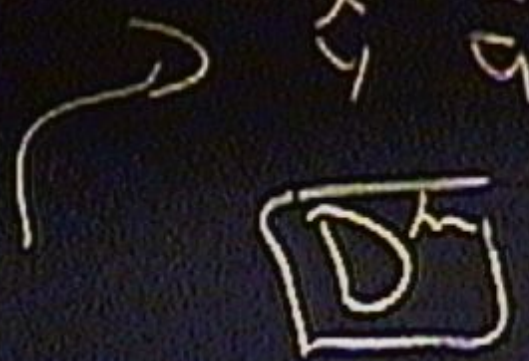
$t_r(z^3)$ \rightarrow x^3 y^3

http://disk.maa


z_1

z_2

Dam

$t_r(z^3)$  x, y z^5

http://disk.maa

z_1




"Damian

FULL $N=4$ SYM

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05

PSU(2,2|4) SYMMETRY

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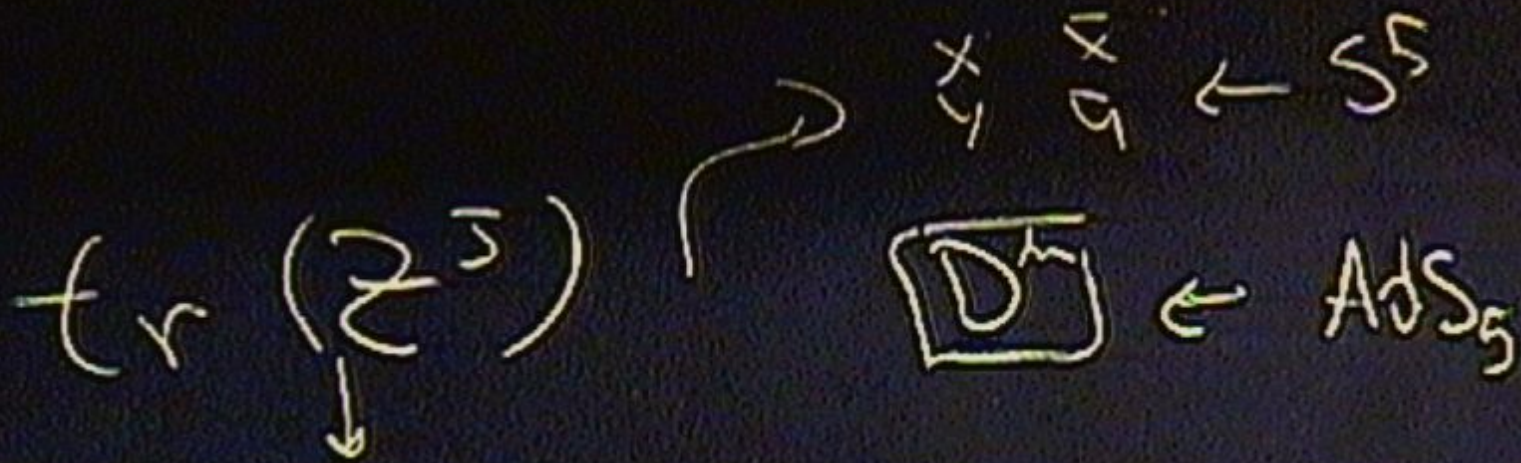
$$S(p_1, p_2; \lambda) = e^{i\theta}$$

↑

"DRESSING" PHASE

MAIN REQUIREMENT: CROSSING SYMMETRY

JANIK
05



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"DamianDKM"

FULL $N=4$ SYM

BEISERT
05

PSU(2,2|4) SYMMETRY

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MODULO A SCALAR FACTOR

$$S(p_1, p_2; \lambda) = e^{i\theta}$$

↑

"DRESSING" PHASE

MAIN REQUIREMENT: CROSSING SYMMETRY

JANIK
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THE DRESSING PHASE

• UP TO 3 LOOPS: $\Theta = 0$

• IN CLASSICAL STRINGS: A.F.S.

$$\Theta(p_1, p_2) = 2 \sum_{r=2}^{\infty} q_r(p_1) q_r(p_2) - q_r(p_2) q_{r+1}(p_1)$$

ARUTYUNOV
FROLON
STAUDACHER
04

$$q_r(p) = \frac{2 \sin((r-1)p/2)}{r-1} \left[\frac{-1 + \sqrt{1 + g^2 \sin^2(p/2)}}{2g \sin(p/2)} \right]^{r-1}$$

• STRING 1 LOOP HERNANDEZ-LOPEZ 06

• CONJECTURE FOR LARGE- λ
ASYMPTOTIC EXPANSION

BEISERT-HERNANDEZ-LOPEZ 06

$$\Theta = \sum_{r,s} c_{r,s}(g) [q_r(p_1) q_s(p_2) - q_r(p_2) q_s(p_1)]$$

• EXACT PROPOSAL

BEISERT-EDEN-STAUDACHER 06

$$c_{r,s}(g) = 2 \cos\left[\frac{\pi}{2}(s-r-1)\right] (r-1)(s-1) \int_0^{\infty} \frac{J_{r-1}(2gt) J_{s-1}(2gt)}{t(e^t-1)} dt$$

THE CUSP ANOMALY

IN GAUGE THEORY WE CONSIDER
THE OPERATORS:

$$\text{tr}(\phi D^S \phi)$$

THEIR ANOMALOUS DIMENSION AT LARGE S IS

$$\Delta - S = f(g) \ln(S) + \dots$$

AT SMALL g

$$f(g) = 8g^2 - \frac{8}{3} \pi^2 g^4 + \frac{88}{45} \pi^4 g^6 \dots$$

IN STRING THEORY WE CONSIDER
A FOLDED STRING ROTATING IN AdS_5

$$ds^2 = dp^2 - \cosh^2 p dt^2 + \sinh^2 p d\theta^2 \quad (AdS_3)$$

$$\begin{cases} t = \kappa \tau \\ \theta = \kappa \tau \\ p = \sigma \end{cases} \quad K = \frac{1}{\pi} \ln\left(\frac{S}{\sqrt{\alpha'}}\right)$$

1-LOOP

2-LOOP

GUBSER
KLEBANOV
POLYAKOV

FROLOV
TSOYTLIS

THE CUSP ANOMALY

IN GAUGE THEORY WE CONSIDER
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$$\text{tr}(\phi D^S \phi)$$

THEIR ANOMALOUS DIMENSION AT LARGE S IS

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IN STRING THEORY WE CONSIDER
A FOLDED STRING ROTATING IN AdS_5

$$ds^2 = dp^2 - \cosh^2 p dt^2 + \sinh^2 p d\theta^2 \quad (AdS_3)$$

$$\begin{cases} t = k\tau \\ \theta = k\tau \\ p = \sigma \end{cases} \quad K = \frac{1}{\pi} \ln\left(\frac{S}{\sqrt{\lambda}}\right)$$

$$f(g) = 4g - \overset{\text{1-LOOP}}{\underbrace{3 \ln(\lambda)}} + \overset{\text{2-LOOP}}{\mathcal{O}(1)}$$

GUBSER
KLEBANUM
POLYAKOV &
FROLON
TSYBULIN &

• IN CLASSICAL STRINGS: A.F.S.

$$\Theta(p_1, p_2) = 2 \sum_{r=2}^{\infty} q_r(p_1) q_r(p_2) - q_r(p_2) q_{r+1}(p_1)$$

ARUTYUNOV
FROLON
STAUDACHER
06

$$q_r(p) = \frac{2 \sin((r-1)p/2)}{r-1} \left[\frac{-1 + \sqrt{1 + g^2 \sin^2(p/2)}}{2g \sin(p/2)} \right]^{r-1}$$

• STRING 1 LOOP HERNANDEZ-LOPEZ 06

• CONJECTURE FOR LARGE- λ
ASYMPTOTIC EXPANSION

BEISERT-HERNANDEZ-LOPEZ 06

$$\Theta = \sum_{r,s} C_{r,s}(g) [q_r(p_1) q_s(p_2) - q_r(p_2) q_s(p_1)]$$

• EXACT PROPOSAL

BEISERT-EDEN-STAUDACHER 06

$$C_{r,s}(g) = 2 \cos\left[\frac{\pi}{2}(s-r-1)\right] (r-1)(s-1) \int_0^{\infty} \frac{J_{r-1}(2gt) J_{s-1}(2gt)}{t(e^t-1)} dt$$

BES EQUATION

BASED ON THE BEHL'S S-MATRIX,
B.E.S. STUDIED THE BETHE ANSATZ
FOR THE CUSP ANOMALY.

INTEGRAL EQUATION FOR $f(\beta)$:

$$f(\beta) = 16g^2 \hat{\sigma}(0)$$

$$\hat{\sigma}(t; \beta) = \frac{t}{e^t - 1} \left(K(2g^2 t, 0) - 4g^2 \int_0^\infty dt' K(2g^2 t, 2g^2 t') \hat{\sigma}(t') \right)$$

$$K(t, t') = K^m(t, t') + 2K^c(t, t')$$

$$K^m(t, t') = \frac{J_1(t) J_0(t') - J_0(t) J_1(t')}{t - t'}$$

$$K^c(t, t') = 4g^2 \int_0^\infty dt'' K_1(t, 2g^2 t'') \frac{t''}{e^{t''} - 1} K_0(2g^2 t', t'')$$

$$K_0(t, t') = \frac{t J_1(t) J_0(t') - t' J_0(t) J_1(t')}{t^2 - t'^2}$$

$$K_1(t, t') = \frac{t' J_1(t) J_0(t') - t J_0(t) J_1(t')}{t^2 - t'^2}$$

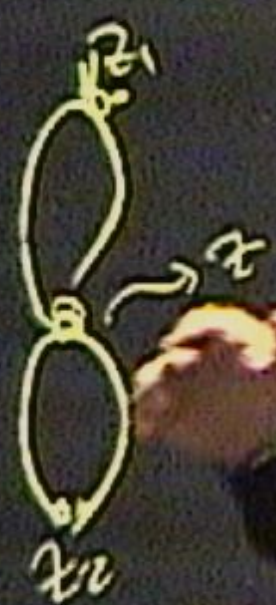
$$SL(2) \quad \text{tr}(\rho^S)$$

$$\text{tr}(\phi D^S \phi)$$

$$x \quad y \quad z \leftarrow S^3$$

$$\square \leftarrow AdS_5$$

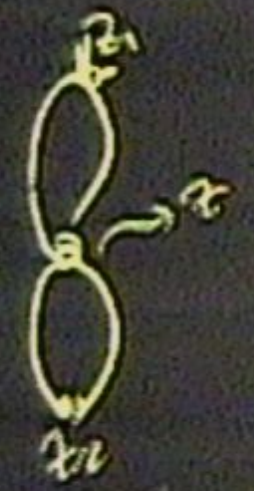
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amianDk

$SL(2)$
 $tr(\phi D^s \phi)$ $tr(z^3)$ \rightarrow $\begin{matrix} x & y \\ y & x \end{matrix} \leftarrow S^1$
 $\square \leftarrow ADS_5$

<http://disk.mac.com/dangerousrobot-Public>



"DamianDKmatter"

$e^{iPL} = \Pi S$

$$SL(2) \quad \text{tr}(\phi D^s \phi) \quad \text{tr}(Z^3) \quad \begin{matrix} x, y \leftarrow S^3 \\ \square \leftarrow ADS_5 \end{matrix}$$

<http://disk.mac.com/dangerousrobot-Public>



"DamianDKmatter"

$$e^{iPL} = \sum \Pi S$$



BES EQUATION

BASED ON THE BEHL'S S-MATRIX,
B.E.S. STUDIED THE BETHE ANSATZ
FOR THE CUSP ANOMALY.

INTEGRAL EQUATION FOR $f(\sigma)$:

$$f(\sigma) = 16g^2 \hat{\sigma}(0)$$

$$\hat{\sigma}(t; \sigma) = \frac{t}{e^t - 1} \left(K(2g^2 t, 0) - 4g^2 \int_0^\infty dt' K(2g^2 t, 2g^2 t') \hat{\sigma}(t') \right)$$

$$K(t, t') = K^m(t, t') + 2K^c(t, t')$$

$$K^m(t, t') = \frac{J_1(t) J_0(t') - J_0(t) J_1(t')}{t - t'}$$

$$K^c(t, t') = 4g^2 \int_0^\infty dt'' K_1(t, 2g^2 t'') \frac{t''}{e^{t''} - 1} K_0(2g^2 t', t'')$$

$$K_0(t, t') = \frac{t J_1(t) J_0(t') - t' J_0(t) J_1(t')}{t^2 - t'^2}$$

$$K_1(t, t') = \frac{t' J_1(t) J_0(t') - t J_0(t) J_1(t')}{t^2 - t'^2}$$

PERTURBATIVE EXPANSION OF $f(\beta)$ FROM BES EQUATION

$$\begin{aligned}
 f(\beta) = & 8g^2 - \frac{8}{3}\pi^2 g^4 + \frac{88}{45}\pi^4 g^6 \\
 & - 16\left(\frac{73}{630}\pi^6 + 4z(\beta)^2\right)g^8 \\
 & + 32\left(\frac{887}{14175}\pi^8 + \frac{4}{3}\pi^2 z(\beta)^2 + 40z(\beta)z(5)\right)g^{10} \\
 & - \dots
 \end{aligned}$$

4-LOOP:
 • BERK, CZARNOY,
 DIXON, KOSOWER,
 SHIRMAN 05
 • CACIARLO
 SPADACIN
 VELUVICH 06

TRANSCENDENTALITY PRINCIPLE

(KOTIKOV
LIPATOV)

THE TERM g^{2k} HAS
TRANSCENDENTALITY $2k-2$

$$t(\pi^n) = n$$

$$t(z(n)) = n$$

CONJECTURE FOR QCD:

HIGHER TRANSCENDENTALITY TERMS IN QCD

PERTURBATIVE EXPANSION OF $\beta(\beta)$ FROM BES EQUATION

$$\begin{aligned}\beta(\beta) = & 8g^2 - \frac{8}{3}\pi^2 g^4 + \frac{88}{45}\pi^4 g^6 \\ & - 16\left(\frac{73}{630}\pi^6 + 4Z(3)^2\right)g^8 \\ & + 32\left(\frac{887}{14175}\pi^8 + \frac{4}{3}\pi^2 Z(3)^2 + 40Z(3)Z(5)\right)g^{10} \\ & - \dots\end{aligned}$$

4-LOOP:
• BERN, CZARNO,
DIXON, ROSSNER,
SHIMAMU 05
• CACHAZ,
SPRADIN
VELINICH 06

TRANSCENDENTALITY PRINCIPLE

(KOTIKOV
LIPATOV)

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$$t(\pi^n) = n$$

$$t(Z(n)) = n$$

CONJECTURE FOR QCD:

HIGHER TRANSCENDENTALITY TERMS IN QCD
COINCIDE WITH THE QCD

QUANTITATIVE STUDY OF $f(g)$ AT FINITE g .

THE BES INTEGRAL EQUATION
ALLOWS, FOR THE FIRST TIME,
TO COMPUTE ANOMALOUS DIMENSIONS
AT FINITE VALUES OF THE COUPLING

EXPANDING

$$\hat{G}(t) = \frac{t}{e^t - 1} \sum_{n=1}^{\infty} S_n(g) \frac{J_n(2gt)}{2g^t}$$

ONE OBTAINS A MATRIX EQUATION FOR $S_n(g)$

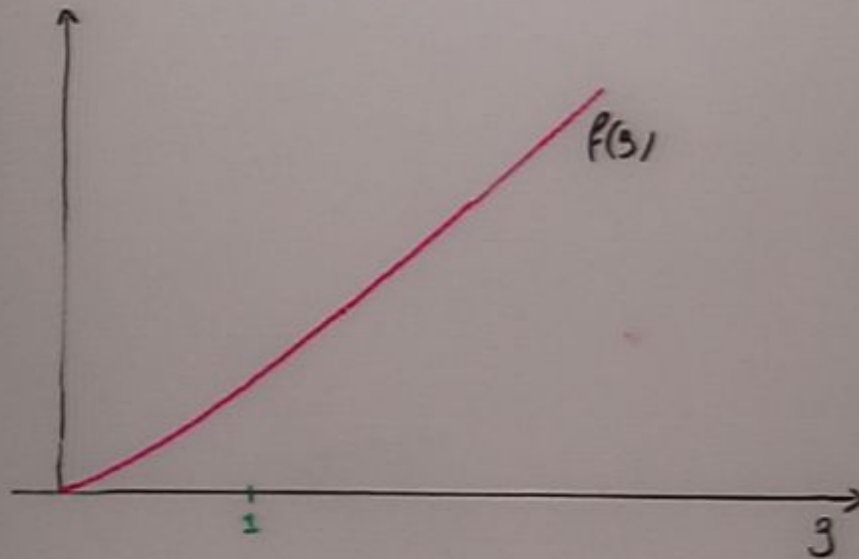
TO COMPUTE $S_n(g)$ AT FINITE g IT IS
ENOUGH TO CONSIDER FINITE MATRICES $K \times K$
WITH $K \sim g$

WITH THIS METHOD IT'S POSSIBLE TO COMPUTE

$$f(g) = 16g^2 S_1(g)$$

NUMERICALLY

BENNA, S.B.,
KLEBANOV
SCARDICCHIO
06



AT LARGE g :

$$f(g) = (4 \pm 10^{-6})g + \leftarrow \text{GKP}$$

$$-(2.661907(2)) + \leftarrow \text{FROLOV-TSEYTLIN}$$

$$- (0.0232(1)) \frac{1}{g} \leftarrow \frac{3 \ln(2)}{16}$$

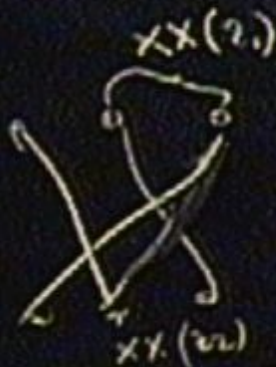
2 LOOP PREDICTION

4th TERM OBTAINED ANALYTICALLY

ALDAY, ARUTYUNOV,
BENNA, EDEN,
KLEBANOV ET

$$SL(2) \quad \text{tr}(\phi D^s \phi) \quad \text{tr}(z^3) \quad \begin{matrix} \nearrow \\ \downarrow \end{matrix} \quad \begin{matrix} x & y & \leftarrow SS \\ \boxed{D^h} & \leftarrow & AdS_5 \end{matrix}$$

<http://disk.mac.com/dangerous>



"DamianDKmatter"

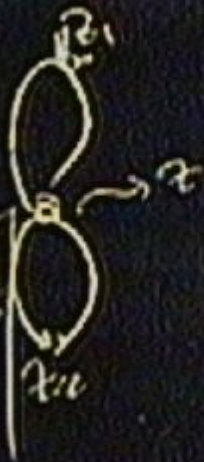
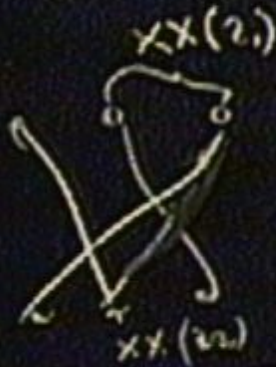
$$S = \sqrt{1 - \delta^2}$$

$$e^{iPL} = \prod S$$

$$\sum$$

$$SL(2) \quad \text{tr}(\phi D^s \phi) \quad \text{tr}(z^3) \quad \begin{matrix} \nearrow x, y \leftarrow SS \\ \searrow \square \leftarrow AdS_3 \end{matrix}$$

<http://disk.mac.com/danger>



"DamianDKmatter"

$$S = \sqrt{\lambda} \int_4 \delta z$$

$$= \pi S$$

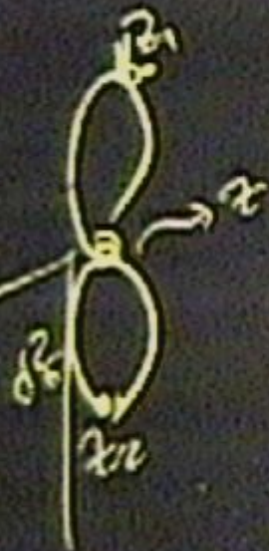
$$\sum S$$

$SL(2)$
 $tr(\phi D^s \phi)$ $tr(z^3)$ $\rightarrow x, y \leftarrow S^5$
 $\square \leftarrow AdS_5$

<http://disk.mac.com/dangerousrob>
 Public

"anDkMatter"

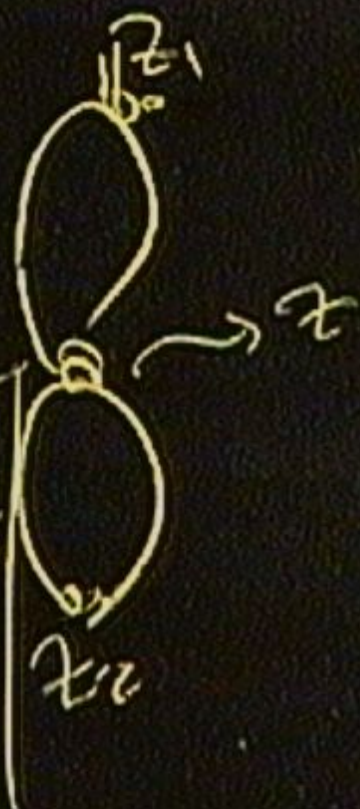
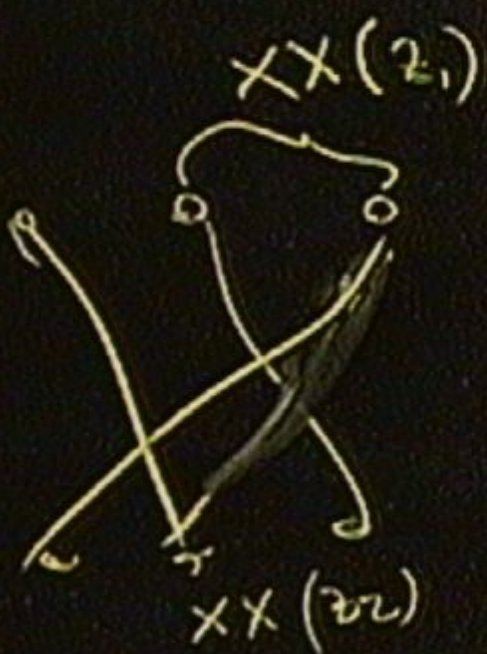
$= \Pi S$



$S = \mathbb{R}P^3$
 4
 δ_2
 z_2

$$\text{tr}(\phi D^S \phi)$$

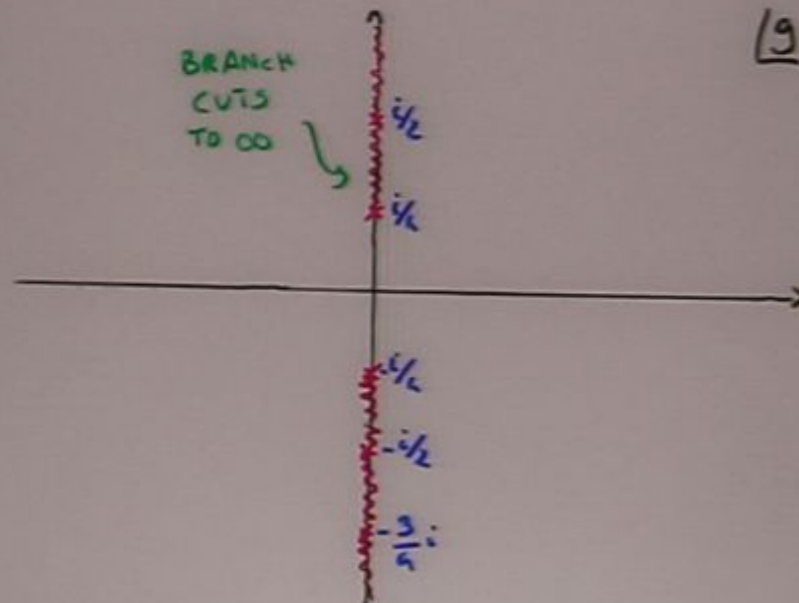
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$$S = \int \mathcal{L}(\dot{x}, x) dt$$

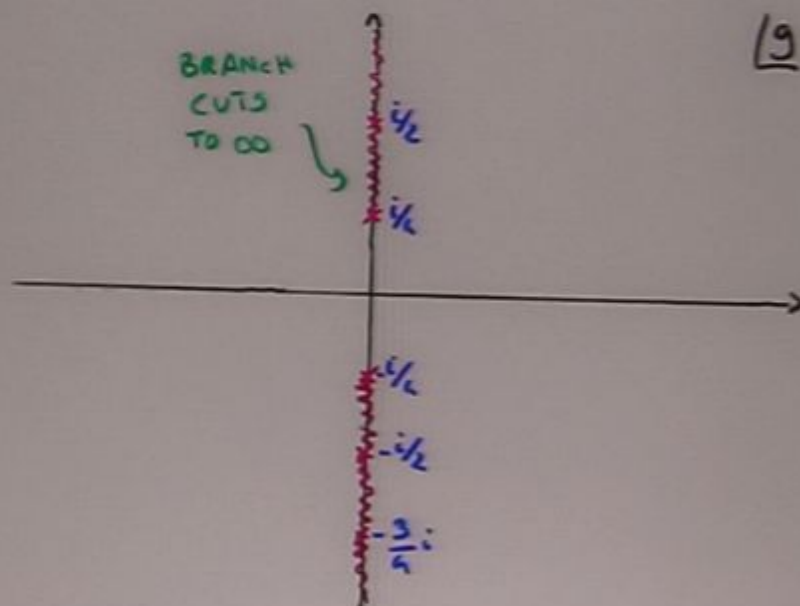
ANALYTIC STRUCTURE

- AT SMALL g , $f(g)$ HAS A
FINITE RADIUS OF CONVERGENCE, $|g| < 1/4$
- AT LARGE g THE EXPANSION
IS ONLY ASYMPTOTIC. ESSENTIAL SINGULARITY.



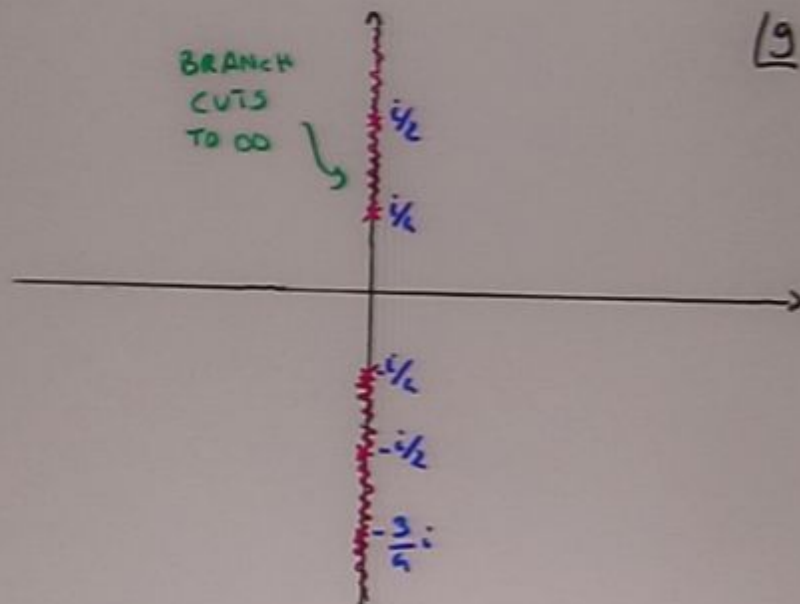
ANALYTIC STRUCTURE

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- AT LARGE g THE EXPANSION IS ONLY ASYMPTOTIC. ESSENTIAL SINGULARITY.

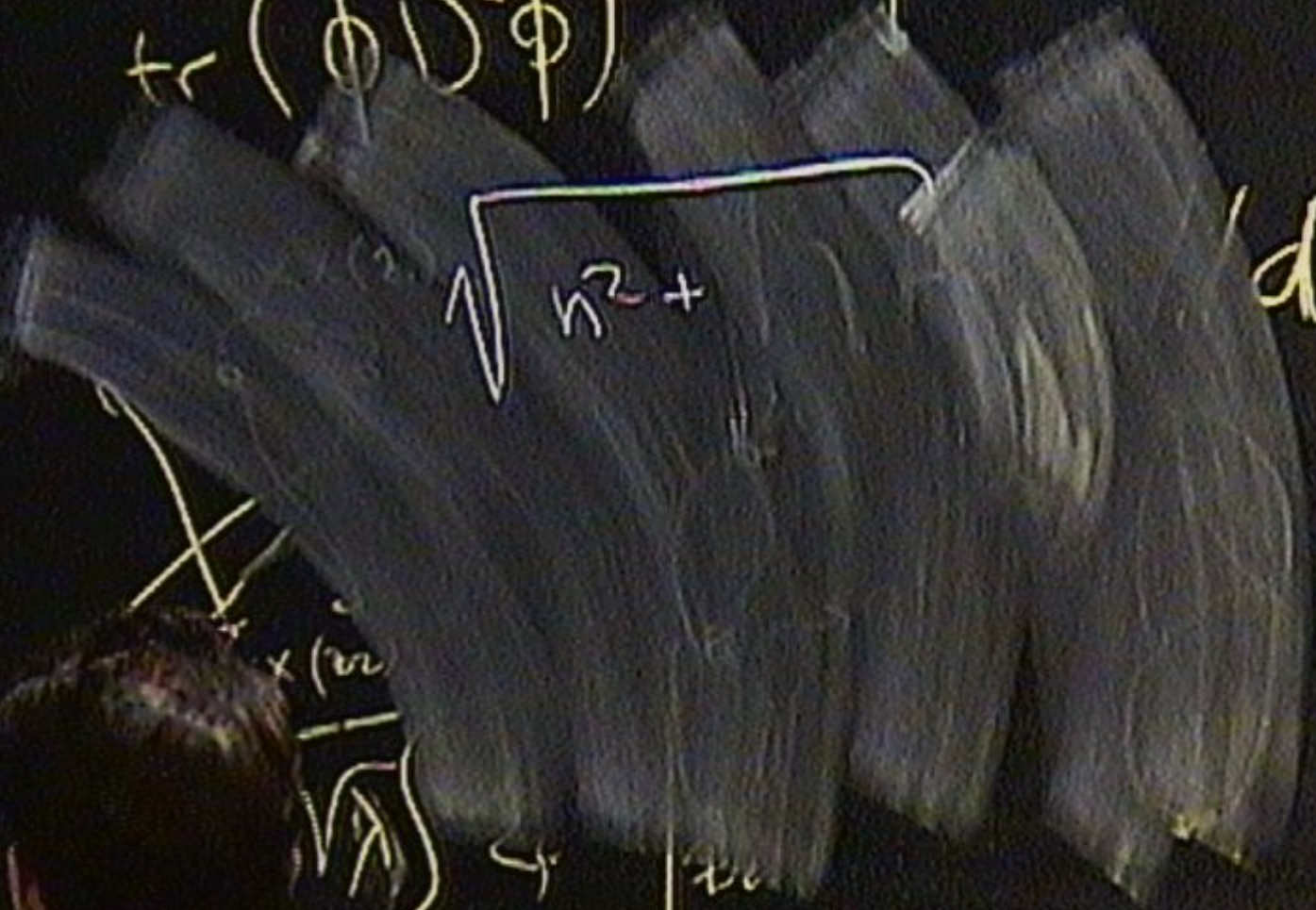


ANALYTIC STRUCTURE

- AT SMALL g , $f(g)$ HAS A FINITE RADIUS OF CONVERGENCE, $|g| < 1/4$
- AT LARGE g THE EXPANSION IS ONLY ASYMPTOTIC. ESSENTIAL SINGULARITY.



$$\text{tr}(\phi D^s \phi) \text{tr}(\epsilon^2) \quad \boxed{D} \leftarrow \text{Ads}$$

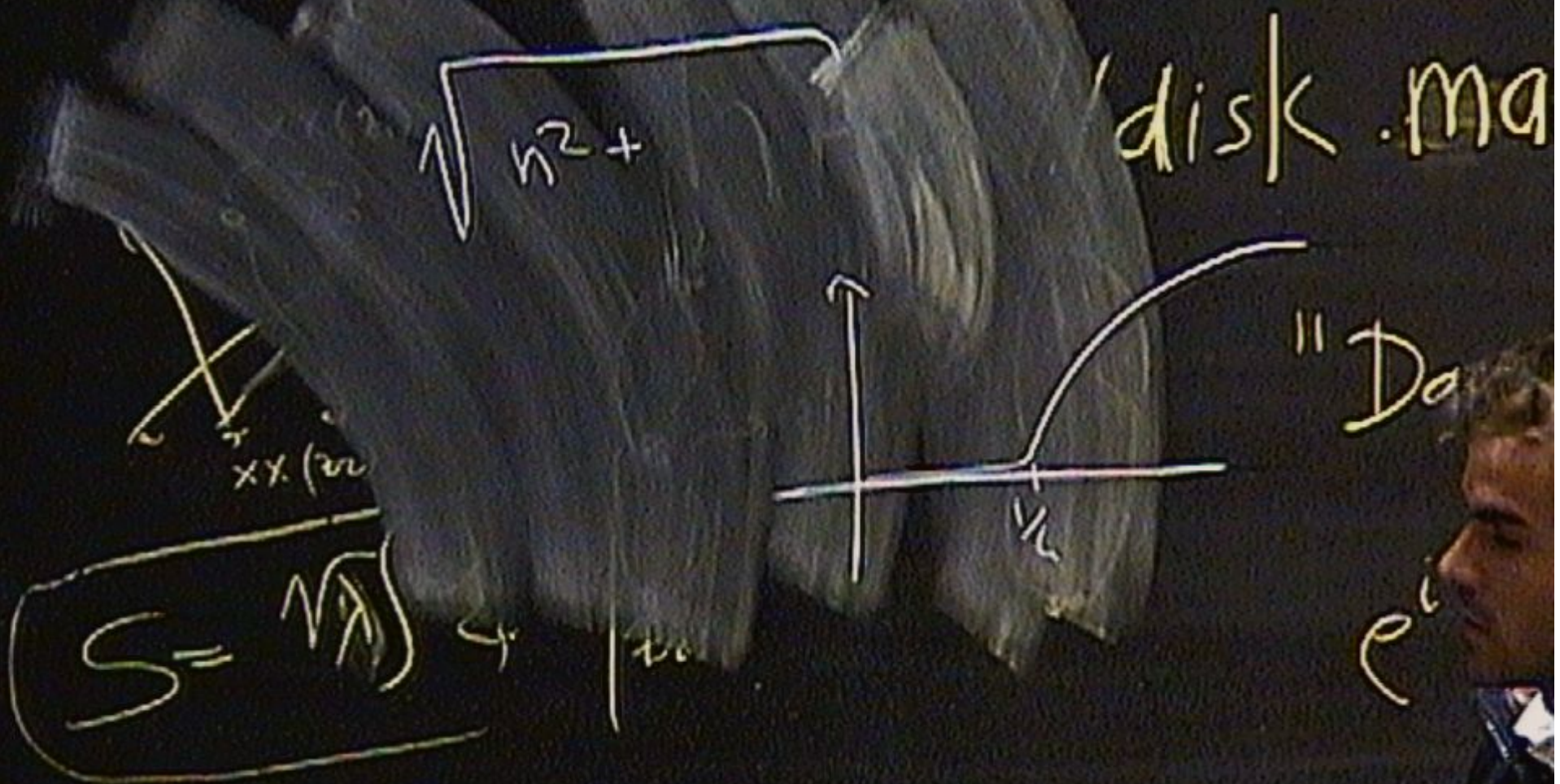


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$$\text{tr}(\phi D^s \phi) \text{tr}(\epsilon^2) \quad [D] \leftarrow \text{Ads}$$

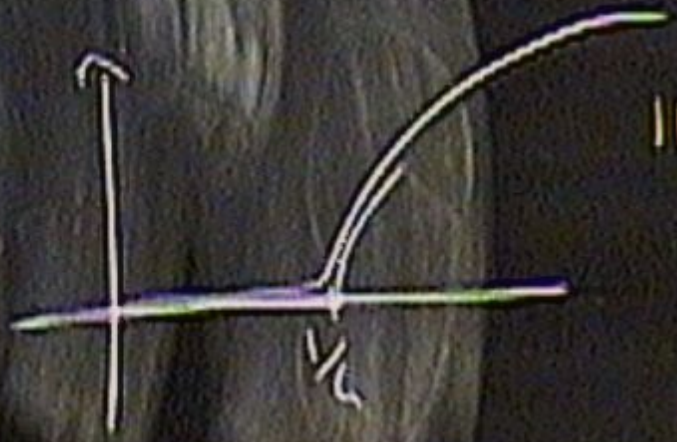


$$\text{tr}(\phi D^s \phi) \text{tr}(\epsilon^2)$$

$\boxed{D} \leftarrow \text{Ads}$

$$\sqrt{k^2 + 16g^2 \sin^2(\phi/2)}$$

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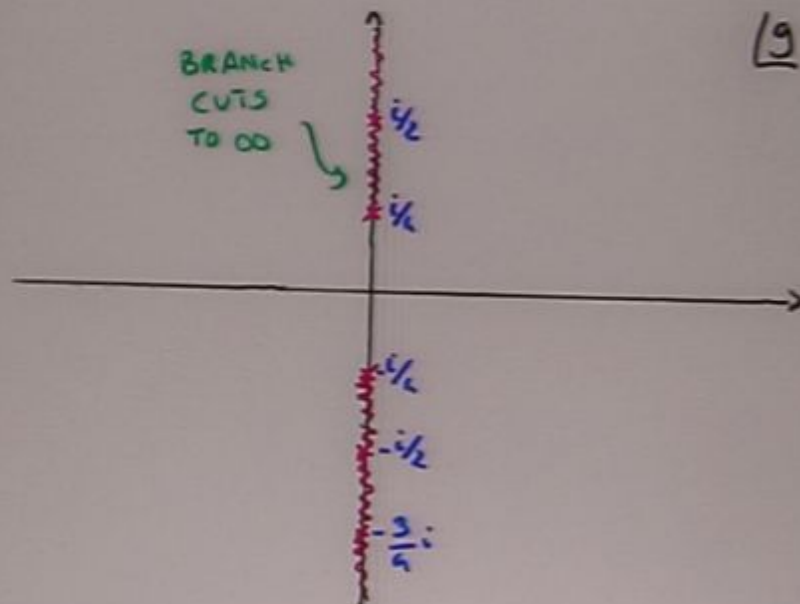


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• IN CLASSICAL STRINGS: A.F.S.

$$\Theta(p_1, p_2) = 2 \sum_{r=2}^{\infty} q_r(p_1) q_{2r}(p_2) - q_r(p_2) q_{2r}(p_1)$$

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$$q_r(p) = \frac{2 \sin((r-1)p/2)}{r-1} \left[\frac{-1 + \sqrt{1 + g^2 \sin^2(p/2)}}{2g \sin(p/2)} \right]^{r-1}$$

• STRING 1 LOOP HERNANDEZ-LOPEZ 06

• CONJECTURE FOR LARGE- λ ASYMPTOTIC EXPANSION BEISERT-HERNANDEZ-LOPEZ 06

$$\Theta = \sum_{r,s} c_{r,s}(g) [q_r(p_1) q_s(p_2) - q_r(p_2) q_s(p_1)]$$

• EXACT PROPOSAL BEISERT-EDEN-STAUDACHER 06

$$c_{r,s}(g) = 2 \cos\left[\frac{\pi}{2}(s-r-1)\right] (r-1)(s-1) \int_0^{\infty} \frac{J_{r-1}(2gt) J_{s-1}(2gt)}{t(e^t-1)} dt$$

BASED ON THE BEHL'S S-MATRIX,
 B.E.S. STUDIED THE BETHE ANSATZ
 FOR THE CUSP ANOMALY.

INTEGRAL EQUATION FOR $f(\beta)$:

$$f(\beta) = 16g^2 \hat{\sigma}(0)$$

$$\hat{\sigma}(t; \beta) = \frac{t}{e^t - 1} \left(K(2g^2 t, \beta) - 4g^2 \int_0^{\infty} dt' K(2g^2 t, 2g^2 t') \hat{\sigma}(t') \right)$$

$$K(t, t') = K^m(t, t') + 2K^c(t, t')$$

$$K^m(t, t') = \frac{J_1(t) J_0(t') - J_0(t) J_1(t')}{t - t'}$$

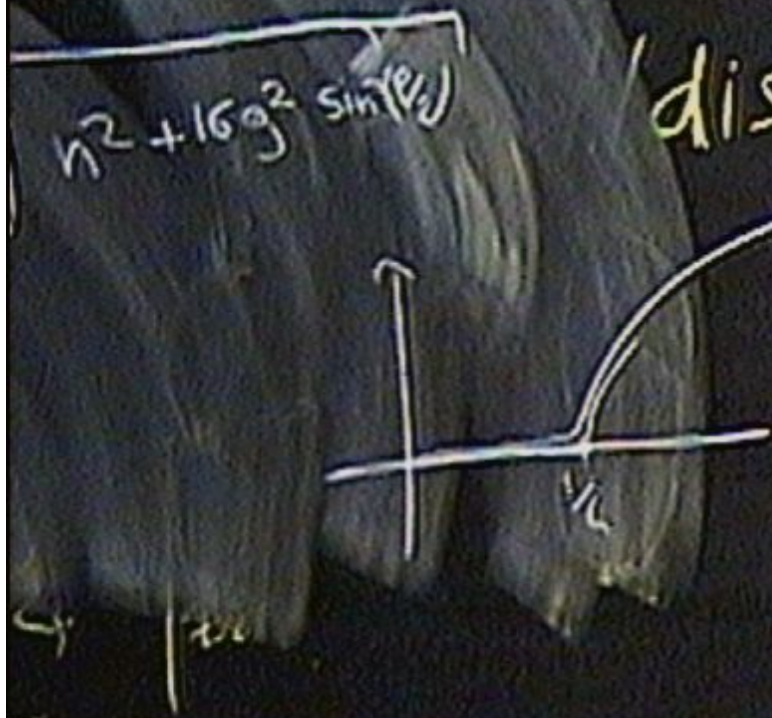
$$K^c(t, t') = 4g^2 \int_0^{\infty} dt'' K_1(t, 2g^2 t'') \frac{t''}{e^{t''} - 1} K_2(2g^2 t'', t')$$

$$K_0(t, t') = \frac{t J_1(t) J_0(t') - t' J_0(t) J_1(t')}{t^2 - t'^2}$$

$$K_1(t, t') = \frac{t' J_1(t) J_0(t') - t J_0(t) J_1(t')}{t^2 - t'^2}$$

$$\phi) \text{tr}(z^3) \rightarrow \begin{matrix} x \\ y \\ z \end{matrix} \leftarrow SS$$

$$\boxed{D^n} \leftarrow \text{Ads}_5 \frac{1}{t} \sum J_n(t') J_m(t)$$



disk.mac.com/dangerousrobot-
Public

"DamianDKMatte"

$$e^{iPL} = \prod \sum S$$



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• UP TO 3 LOOPS: $\Theta = 0$ BASED ON THE BEHL'S S-MATRIX,

• IN CLASSICAL STRINGS: THE BETHE ANSATZ

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$\Theta(p_1, p_2) \equiv \text{THE } \sum_{r=2}^{\infty} \int_{\mathbb{R}} ds q_r(p_1) q_r(p_2) - q_r(p_2) q_r(p_1)$
INTEGRAL EQUATION FOR $f(s)$:

$$q_r(p) = \frac{2 \sin((r-1)\pi/2)}{f(s)^{r-1} (1 + g^2 \sin^2(\pi/2))} \left[-1 + \sqrt{1 + g^2 \sin^2(\pi/2)} \right]^{r-1} f(s)$$

$f(s) = 1 + g^2 \hat{c}(s)$

• STRING $\hat{c}(t) \text{ COOP} = \frac{1}{e^t - 1} \int_{-\infty}^{\infty} dt' K(z, t; z, t') \hat{c}(t')$

• CONJECTURE FOR $K(z, t; z, t') = 2 K_0(z, t; z, t') + 2 K_1(z, t; z, t')$ HERNANDEZ-LOPEZ DE ASYMPTOTIC EXPANSION

$$K_0(t, t') = \frac{J_1(t) J_0(t') - J_0(t) J_1(t')}{t - t'}$$

$$\Theta = \sum_{r,s} C_{r,s}(g) \int_{\mathbb{R}} ds \left[q_r(p_1) q_s(p_2) - q_s(p_1) q_r(p_2) \right] \hat{c}(s)$$

• EXACT PROPOSAL $K_0(t, t') = \frac{t J_1(t) J_0(t') - t' J_0(t) J_1(t')}{t^2 - t'^2}$ BEISERT-EDEN-STAUDACHER DE

$$K_1(t, t') = \frac{t' J_1(t) J_0(t') - t J_0(t) J_1(t')}{t^2 - t'^2}$$

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• IN CLASSICAL STRINGS: A.F.S.

$$\Theta(p_1, p_2) = 2 \sum_{r=2}^{\infty} q_r(p_1) q_{2r}(p_2) - q_r(p_2) q_{r+1}(p_1)$$

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SUMMARY

- SPECTRUM OF GAUGE THEORIES IN TERMS OF SPIN CHAINS
- BETHE ANSATZ FOR $N=4$ SYM AND $AdS_5 \times S^5$
- CUSP ANOMALY: NON BPS OBSERVABLE

OPEN QUESTIONS

- FIND FURTHER EVIDENCE FOR THE BEHL'S S-MATRIX
- STUDY ANALYTICAL STRUCTURE OF ANOMALOUS DIMENSIONS.
- FINITE-LENGTH EFFECTS: HAMILTONIAN?

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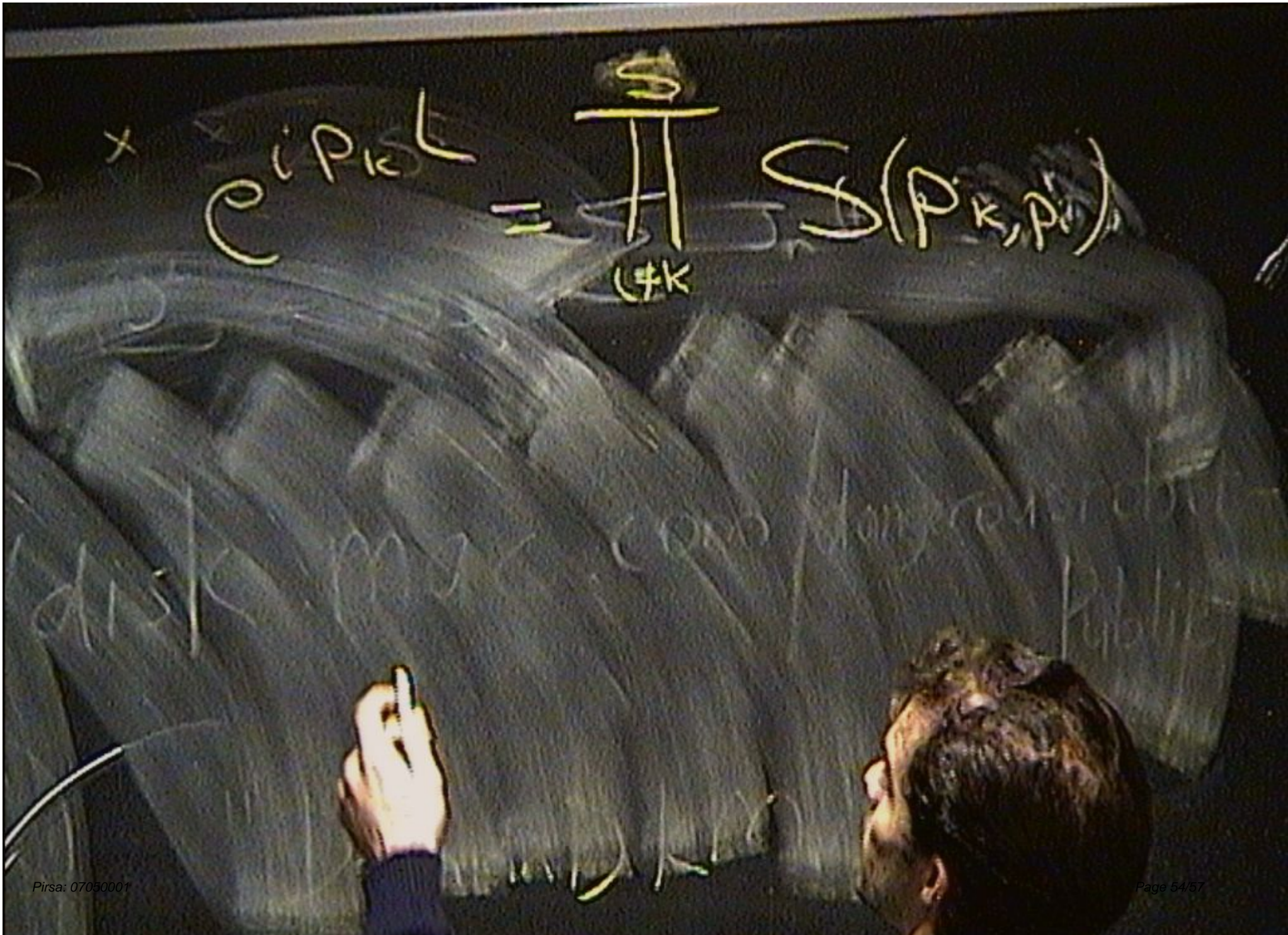
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$$x \quad \times \quad \rho_i p_k L = \prod_{l \neq k} S(p_k, p_l)$$

$$x \cdot \prod_{i \in P_k} L = \prod_{\substack{L \\ L \neq k}}^M S(P_k, P_i)$$

$$x \cdot \prod_{i \neq k} p_i^{\alpha_i} = \prod_{i \neq k} p_i^{\alpha_i} \cdot S(p_k, p_i)$$



$$\sum_i P_{ik} L = \sum_i H_{ik} S(P_k, P_i)$$

$(i \neq k)$

$$x \quad \rho_i p_k L = \prod_{l \neq k} S(p_k, p_l)$$

$$i p_k L = \sum \ln S$$

$$S \rightarrow \infty$$

$$x \quad c_i p_k L = \prod_{l \neq k} S(p_k, p_l)$$

$$c_i p_k L = \sum \ln S$$

$$S \rightarrow \infty$$

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