Title: The power of forgetting

Date: Mar 26, 2007 11:00 AM

URL: http://pirsa.org/07030039

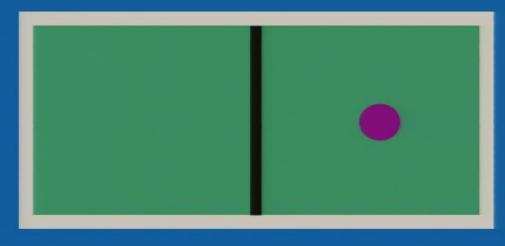
Abstract: Thermodynamics places surprisingly few fundamental constraints on information processing. In fact, most people would argue that it imposes only one, known as Landauer's Principle: a process erasing one bit of information must release an amount kT ln 2 of heat. It is this simple observation that finally led to the exorcism of Maxwell's Demon from statistical mechanics, more than a century after he first appeared. Ignoring the lesson implicit in this early advance, however, quantum information theorists have been surprisingly slow to embrace erasure as a fundamental primitive. Over the past couple of years, however, it has become clear that a detailed understanding of how difficult it is to erase correlations leads to a nearly complete synthesis and simplification of the known results of asymptotic quantum information theory. As it turns out, surprisingly many of the tasks of interest, from distilling high-quality entanglement to sending quantum data through a noisy medium to many receivers, can be understood as variants of erasure. I'll sketch the main ideas behind these discoveries and end with some speculations on what lessons the new picture might have for understanding information loss in real physical systems.

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The power of forgetting

Patrick Hayden (McGill University)

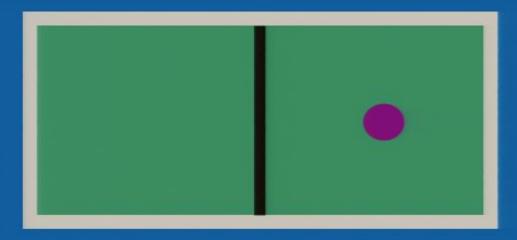




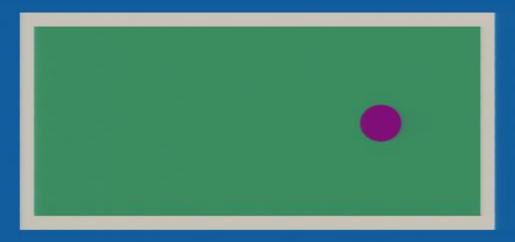
State 1

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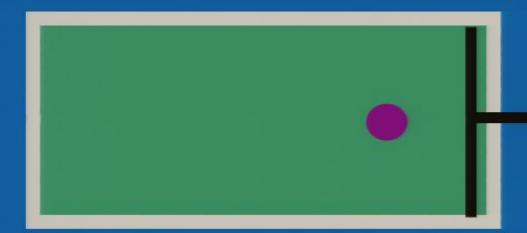




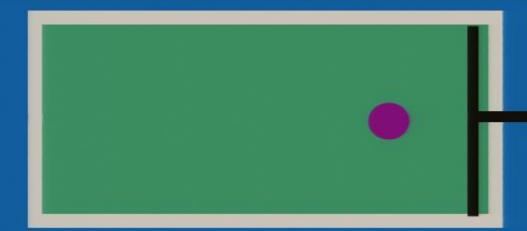
State 1



State 1

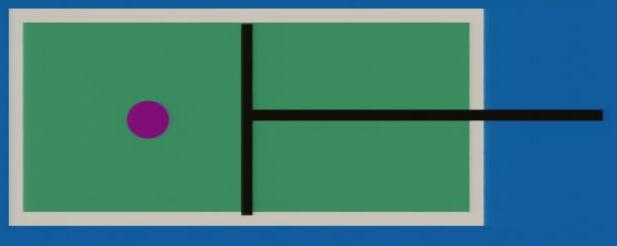


State 1

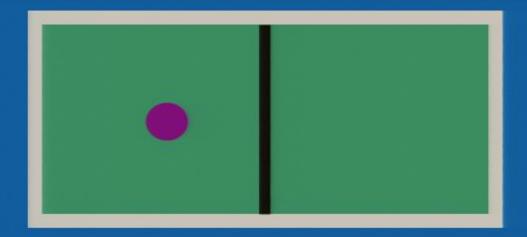


State 1

Erasure: A process that, regardless of the input state, results in output state 0

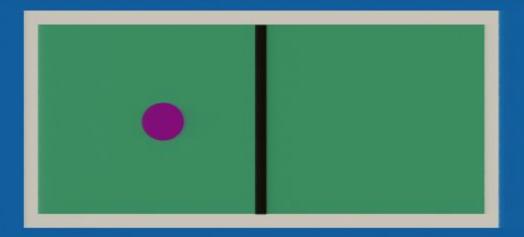


State 1



State 0

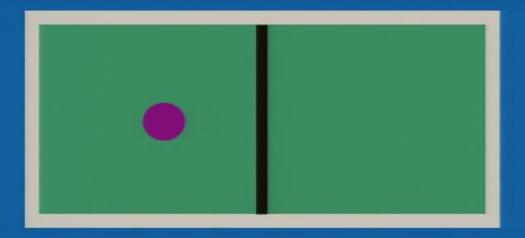
Erasure: A process that, regardless of the input state, results in output state 0



State 0

Erasing information requires a process that can reduce in uncertainty

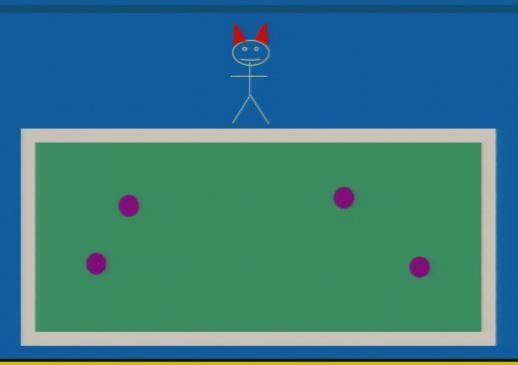
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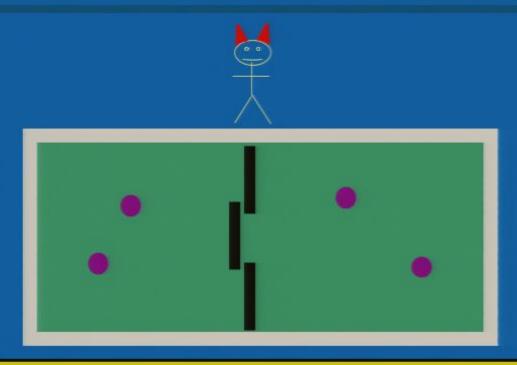
State 0

Erasing information requires a process that can reduce in uncertainty

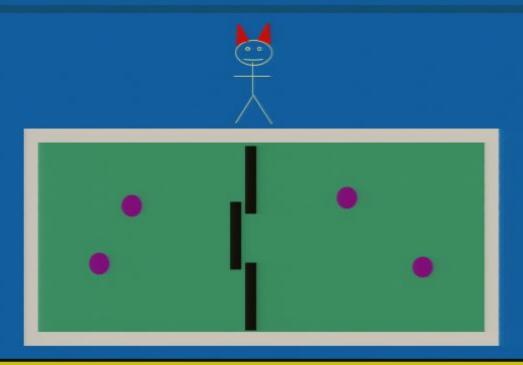
Thermodynamic entropy of the gas is reduced by *k ln* 2 so amount *kT ln* 2 of heat released to the environment.



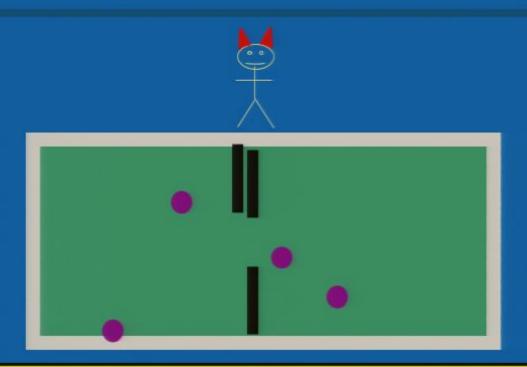
Gas originally at equilibrium. Demon inserts partition.



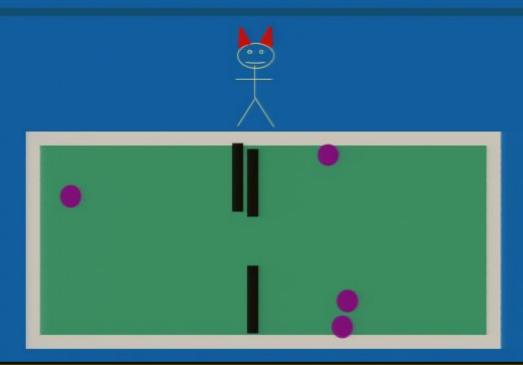
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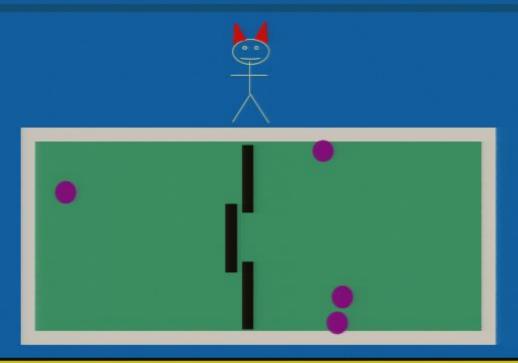
Gas originally at equilibrium. Demon inserts partition.



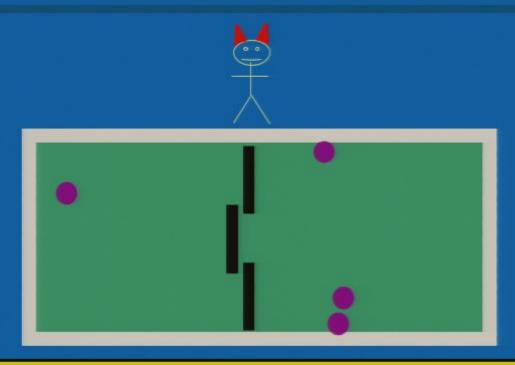
Gas originally at equilibrium. Demon inserts partition.



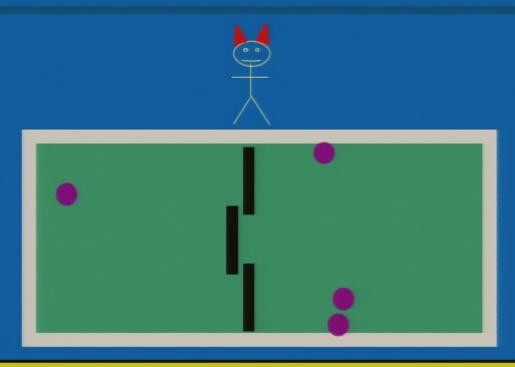
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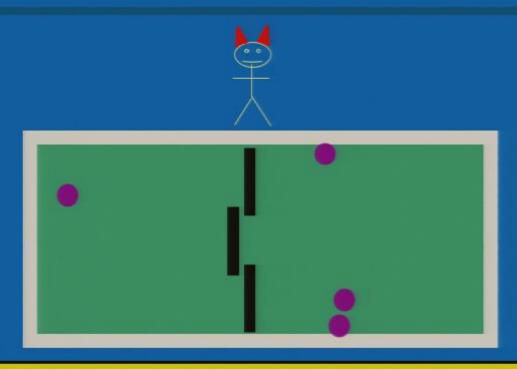
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Gas originally at equilibrium. Demon inserts partition.

Demon uses door to allow particles to move left to right, but not right to left.

Entropy decreases.



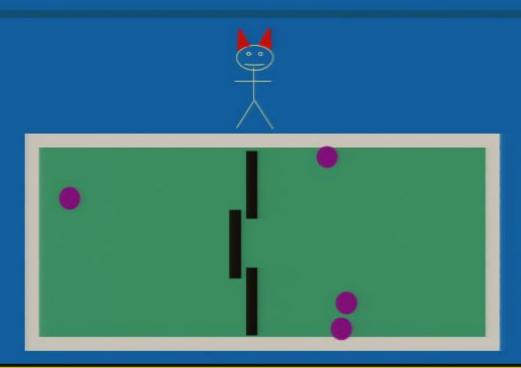
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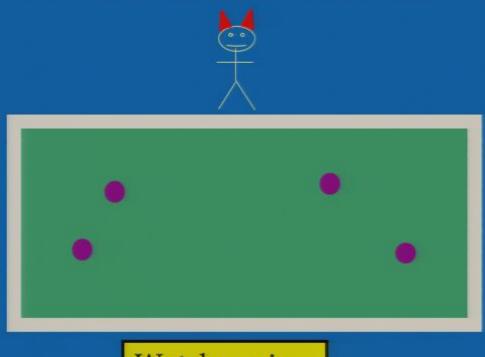
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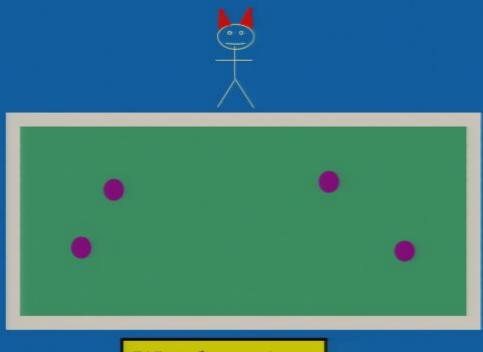
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What gives? Stumped Maxwell, von Neumann, Brillouin, Szilard..



Watch again...

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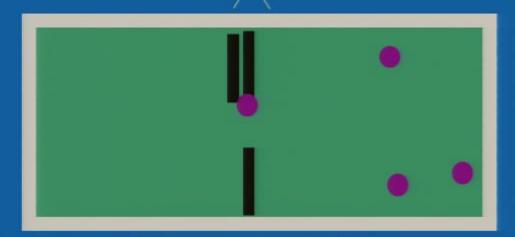
Watch again...

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Particle coming!



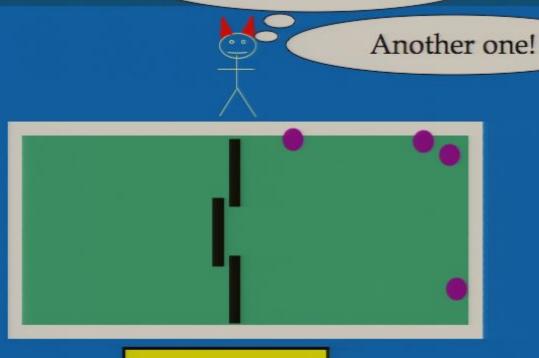
Another one!



Watch again...

R

Particle coming!

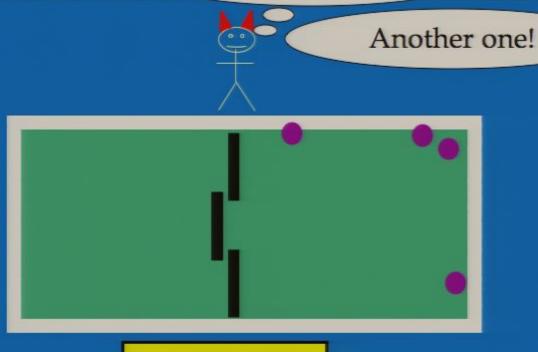


Watch again...

Demon must process information about incoming particles.

P

Particle coming!



Watch again...

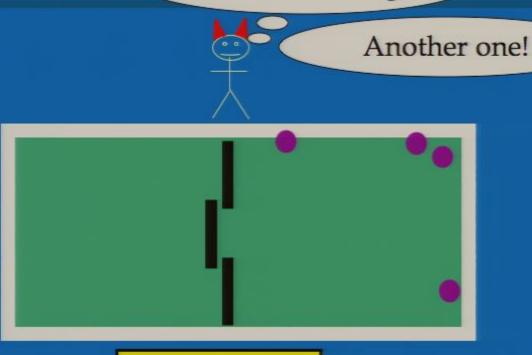
Demon must process information about incoming particles.

Finite memory: he must eventually start *forgetting*.

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A

Particle coming!



Watch again...

Demon must process information about incoming particles.

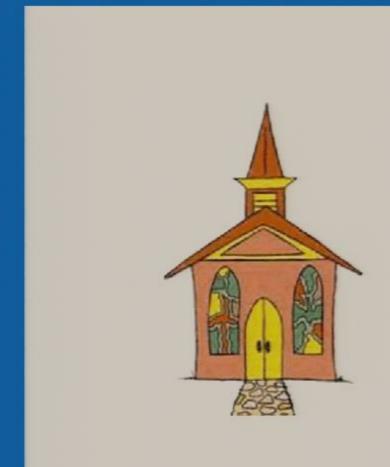
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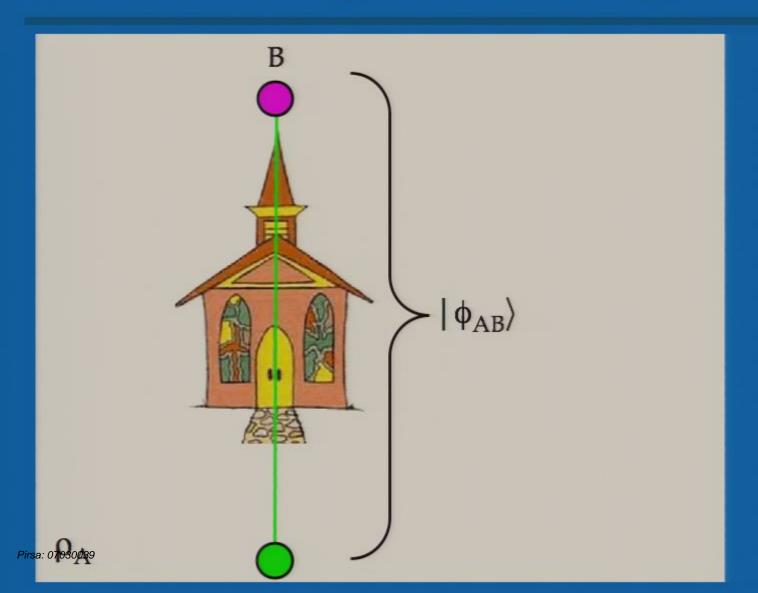
Landauer's principle restores global increase of entropy.

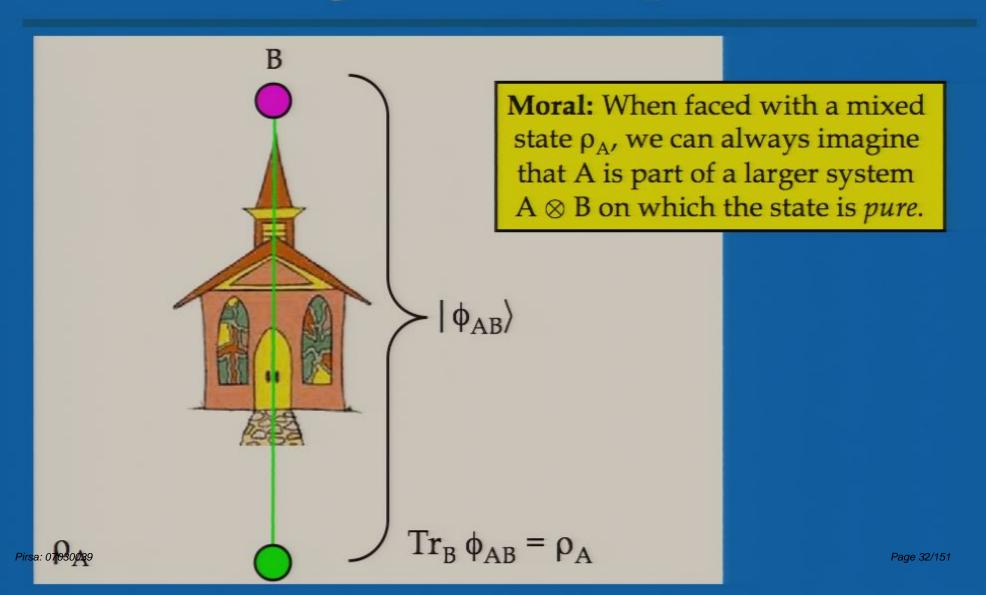
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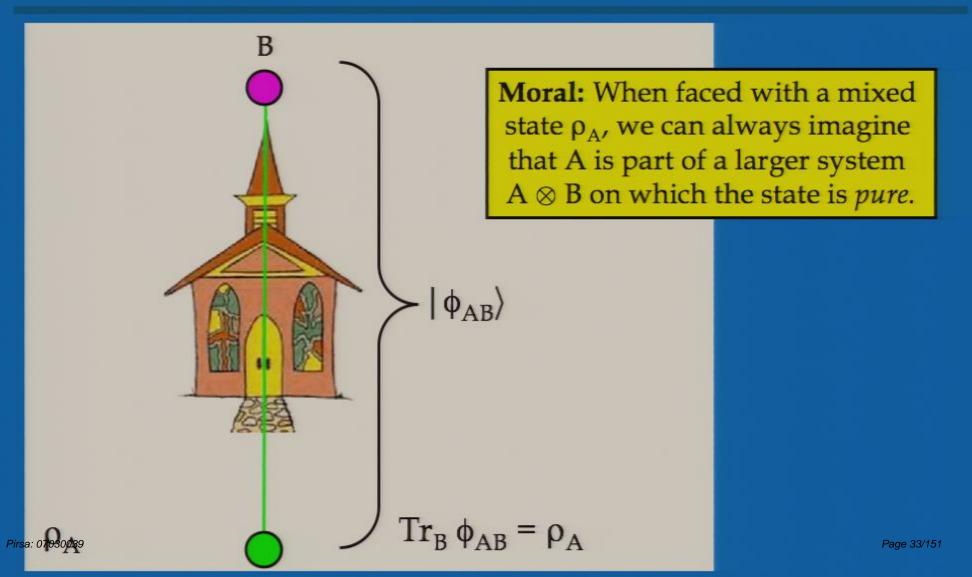




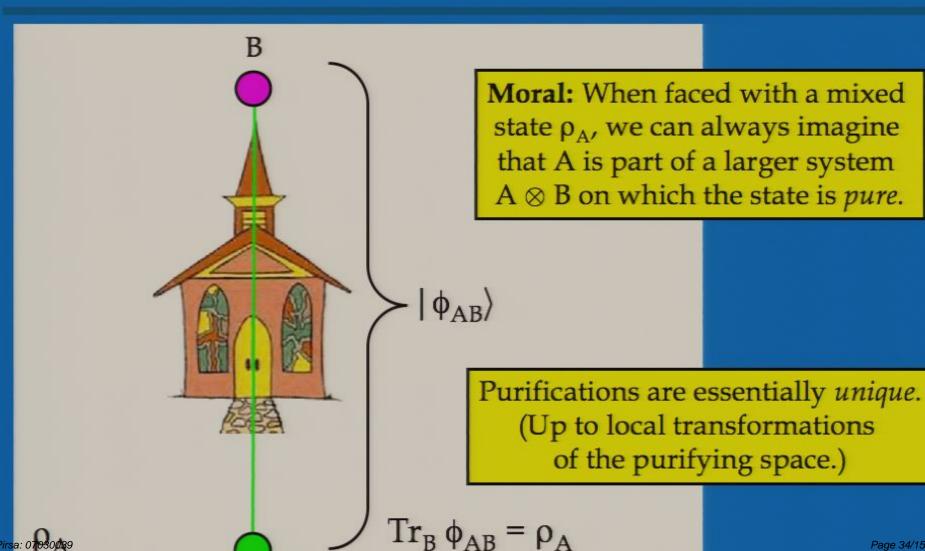












Purification and correlation



Moral: When faced with a mixed state ρ_A , we can always imagine that A is part of a larger system $A \otimes B$ on which the state is *pure*.

Purifications are essentially *unique*. (Up to local transformations of the purifying space.)



Purification and correlation

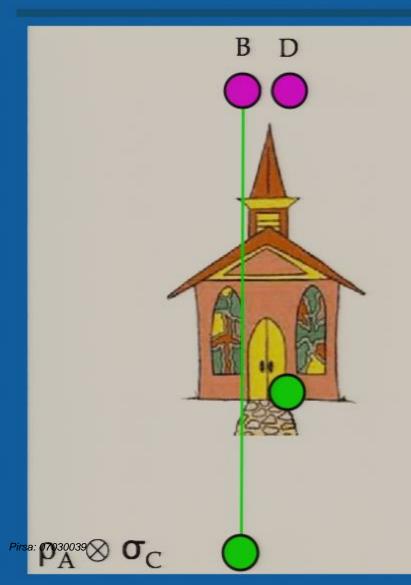


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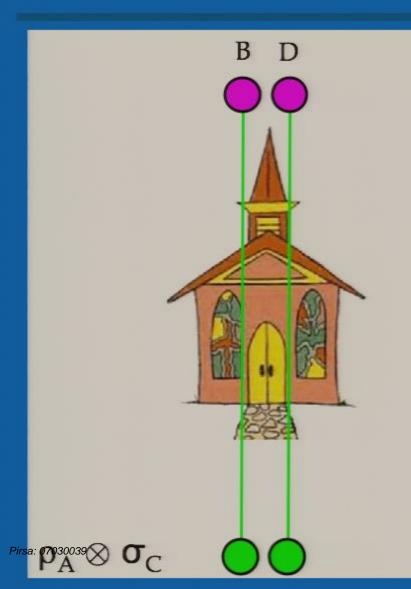
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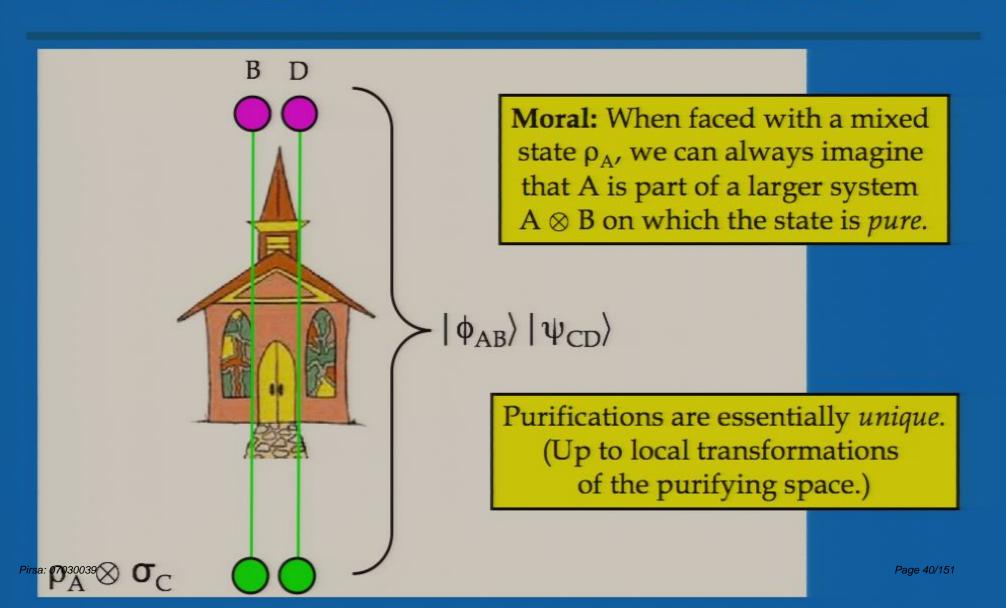
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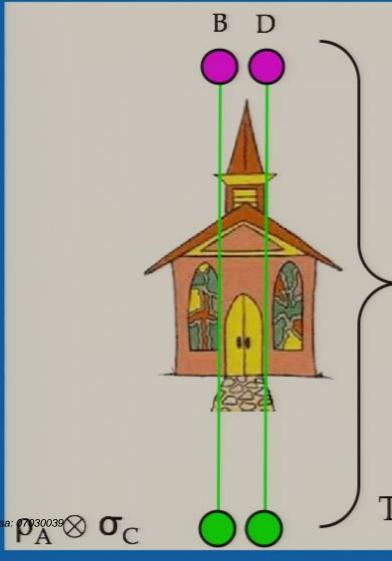


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 $|\phi_{AB}\rangle |\psi_{CD}\rangle$

$$Tr_{BD} \phi_{AB} \otimes \psi_{CD} = \rho_A \otimes \sigma_C$$



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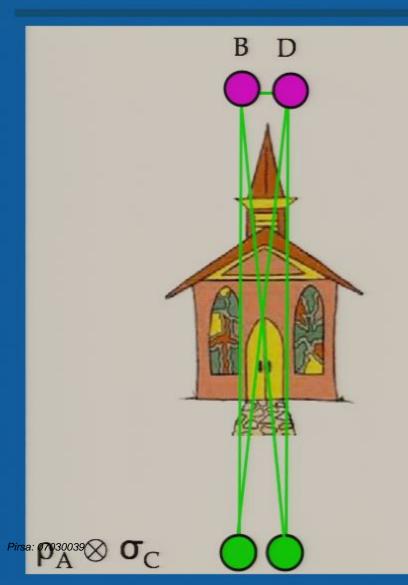




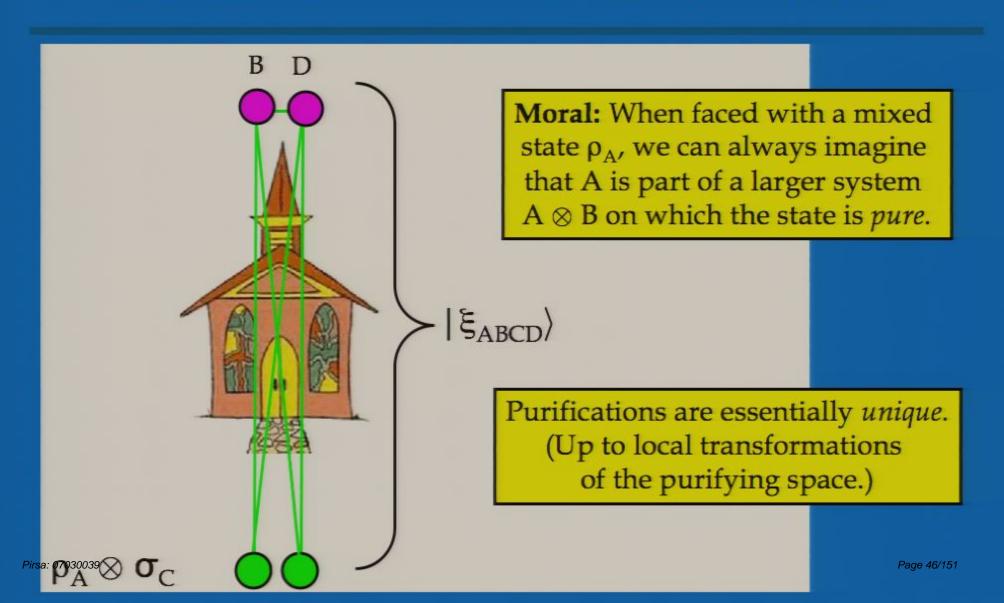
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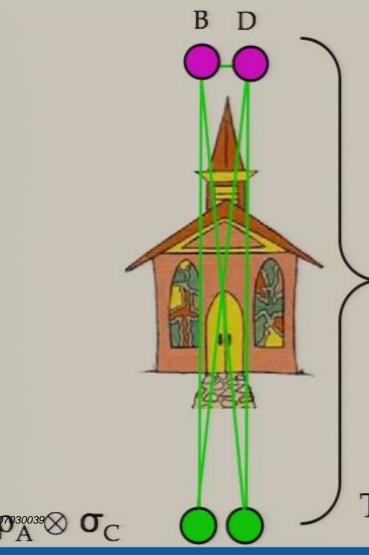


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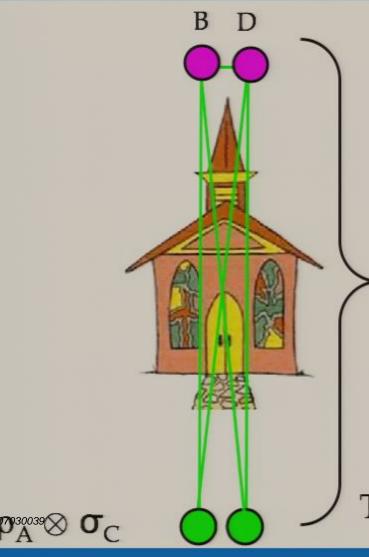




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 $|\xi_{ABCD}\rangle$

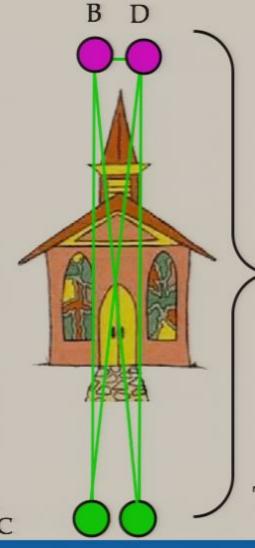
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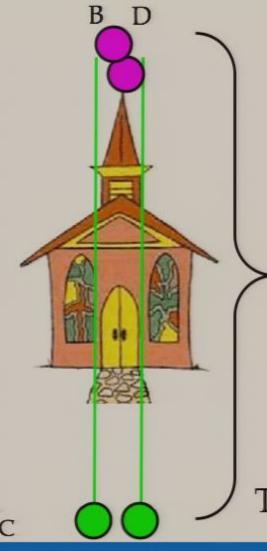
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 $|\xi_{ABCD}\rangle = (id_{AC} \otimes U_{BD}) |\phi_{AB}\rangle |\psi_{CD}\rangle$

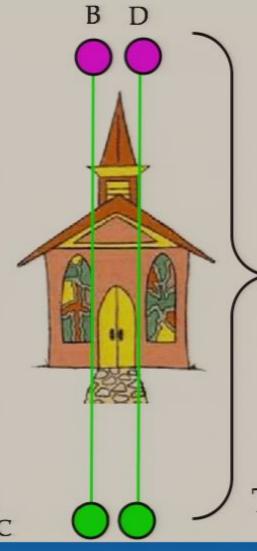
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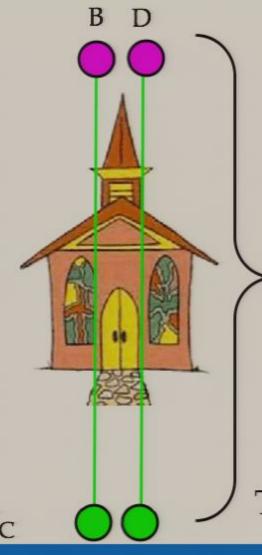
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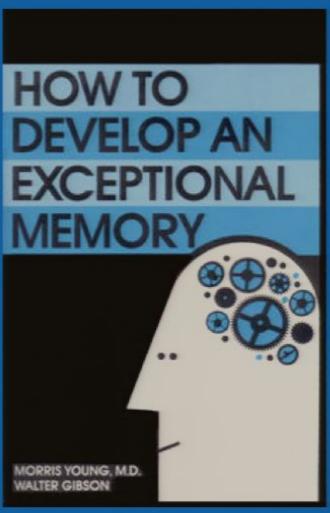
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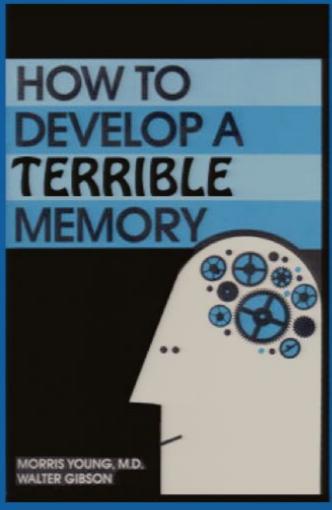
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 $|\phi_{AB}\rangle |\psi_{CD}\rangle = (id_{AC}\otimes U_{BD}^{-1}) |\xi_{ABCD}\rangle$

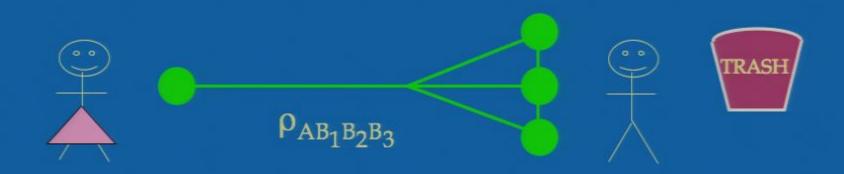
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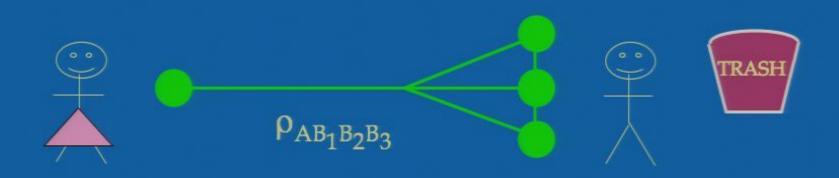
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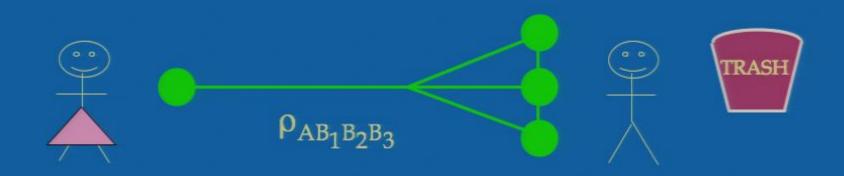


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How can Bob unilaterally destroy his correlation with Alice?

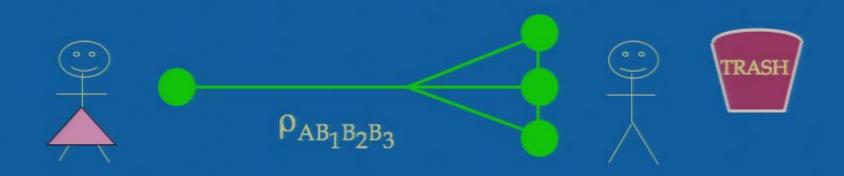
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How can Bob unilaterally destroy his correlation with Alice?

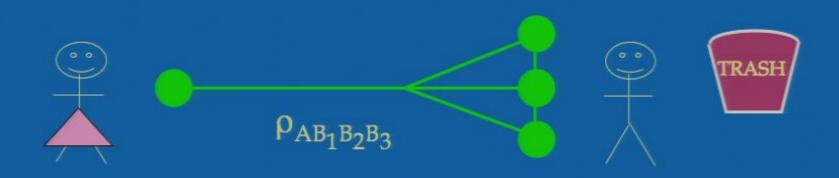
What is the minimal number of particles he must discard before the remaining state is uncorrelated?

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How can Bob unilaterally destroy his correlation with Alice?

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How can Bob unilaterally destroy his correlation with Alice?

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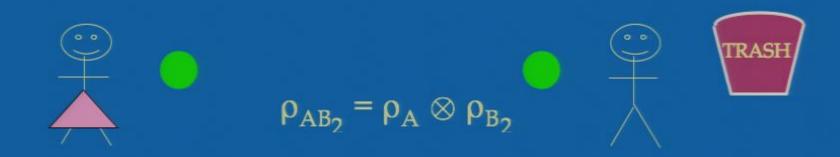
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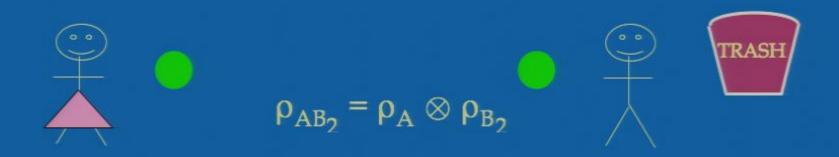
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How can Bob unilaterally destroy his correlation with Alice?

What is the minimal number of particles he must discard before the remaining state is uncorrelated?

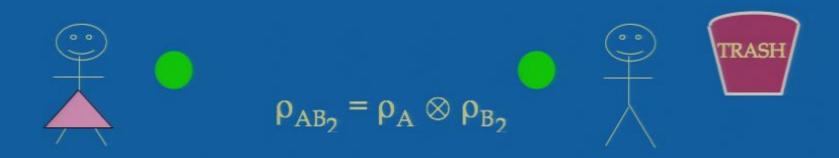
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How can Bob unilaterally destroy his correlation with Alice?

What is the minimal number of particles he must discard before the remaining state is uncorrelated?

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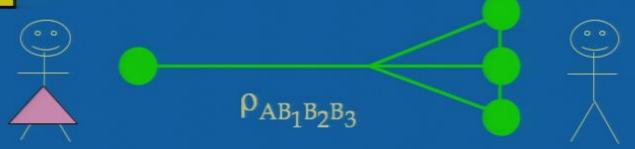
How can Bob unilaterally destroy his correlation with Alice?

What is the minimal number of particles he must discard before the remaining state is uncorrelated?

In this case, by discarding 2 particles, Bob succeeded in eliminating all correlations with Alice's particle

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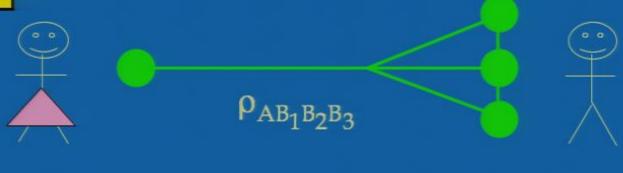
Watch again:

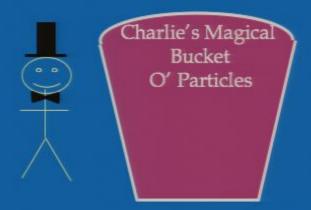




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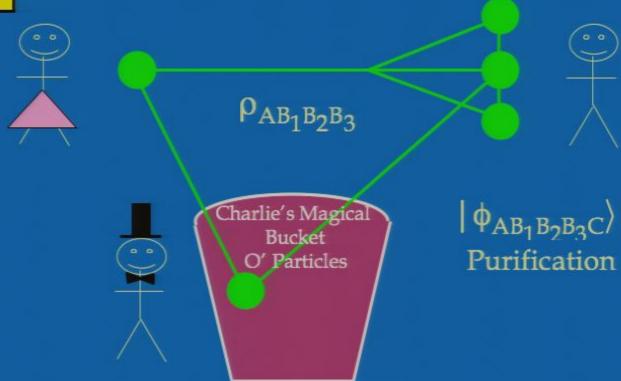
Watch again:





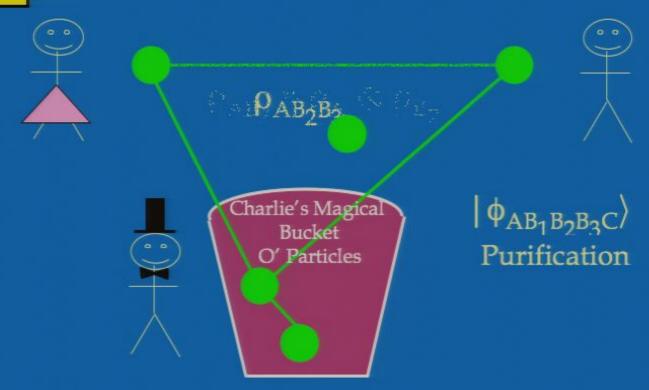
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Watch again:



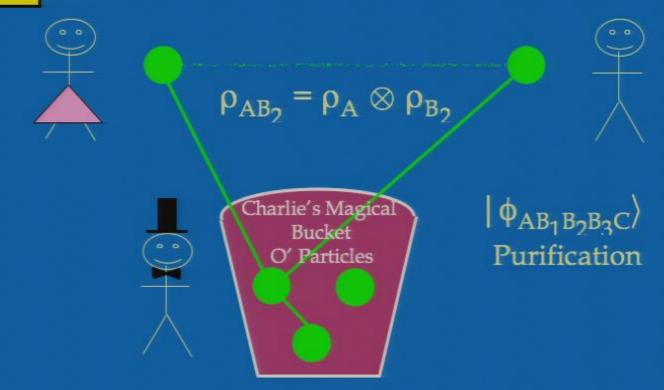
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Watch again:



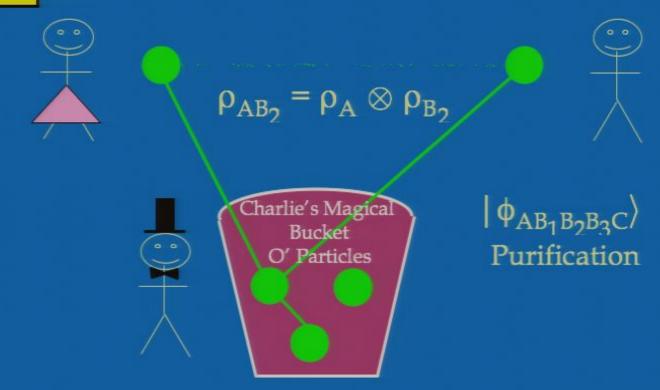
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Watch again:



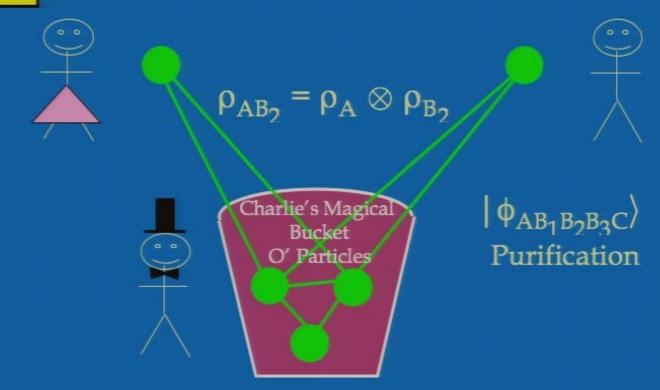
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Watch again:



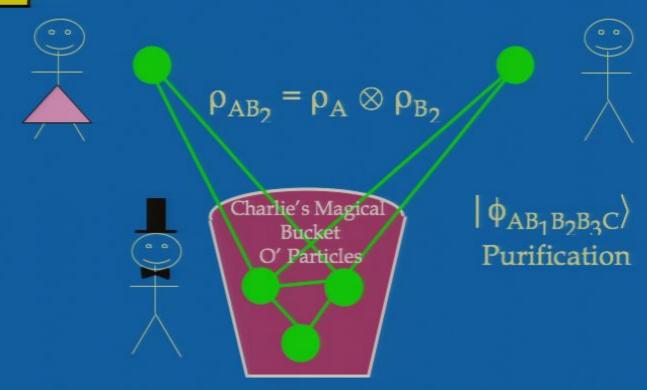
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Watch again:



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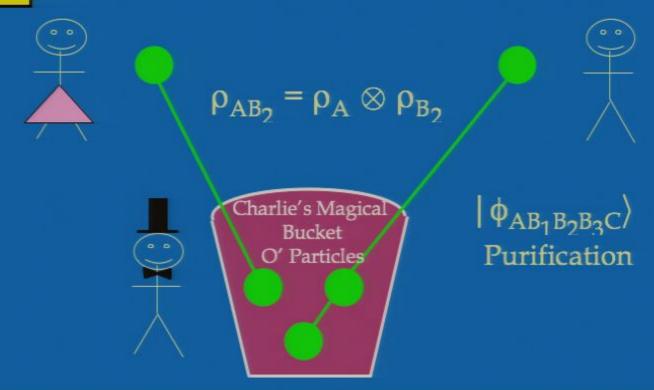
Watch again:



All purifications equivalent up to a local transformation in Charlie's lab.

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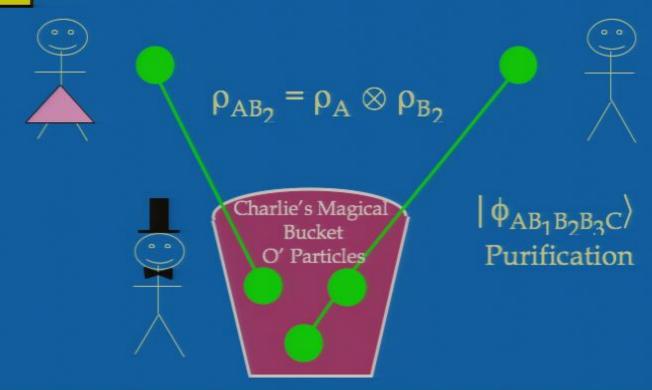
Watch again:



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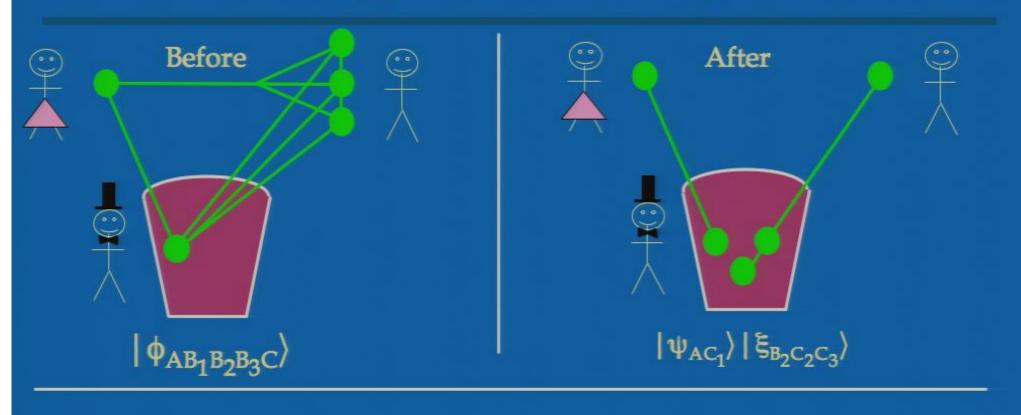
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Watch again:

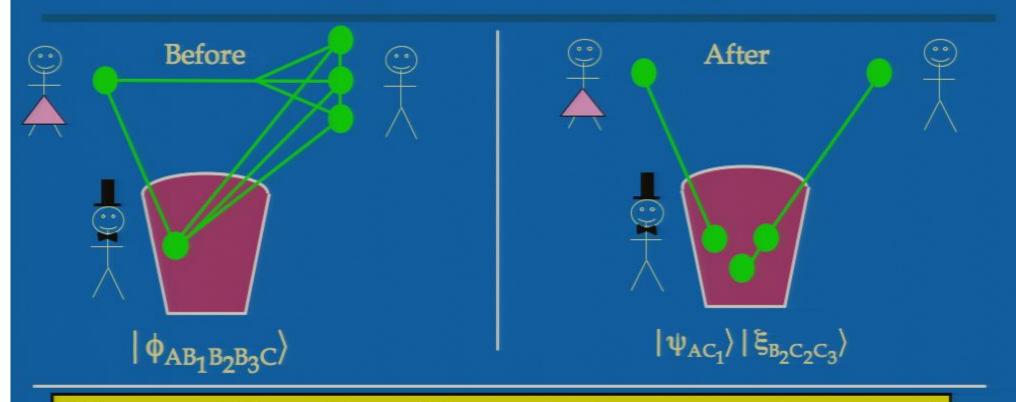


All purifications equivalent up to a local transformation in Charlie's lab.

Charlie holds **uncorrelated** purifications of **both** Alice's particle and Bob's remaining particles.

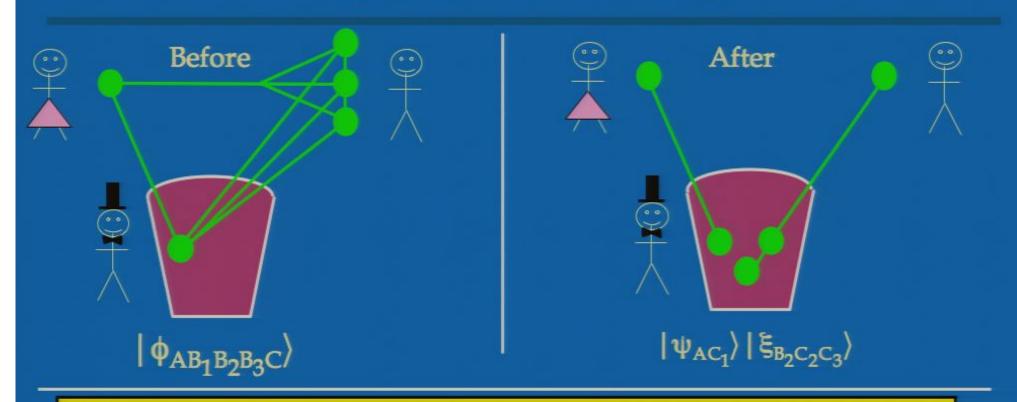


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Alice never did anything \Rightarrow Her marginal state $\phi_A = \psi_A$ is unchanged

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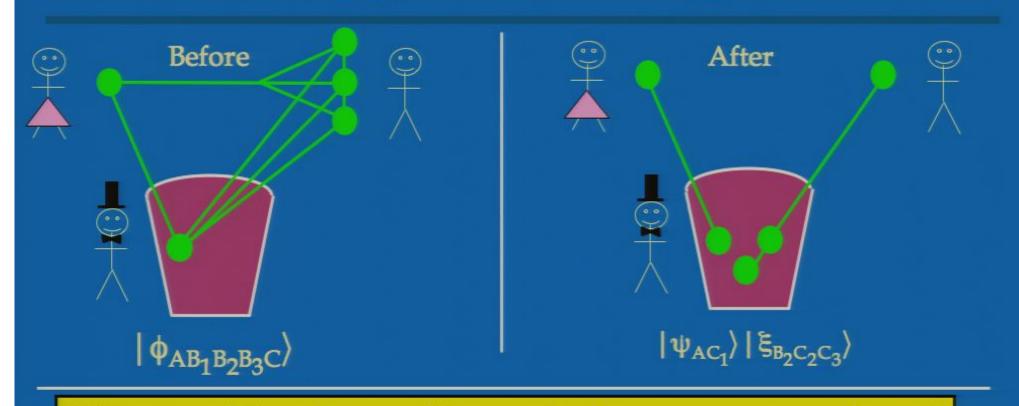


Alice never did anything \Rightarrow Her marginal state $\phi_A = \psi_A$ is unchanged

Originally, her purification is held by both Bob and Charlie.

Afterwards, entirely by Charlie.

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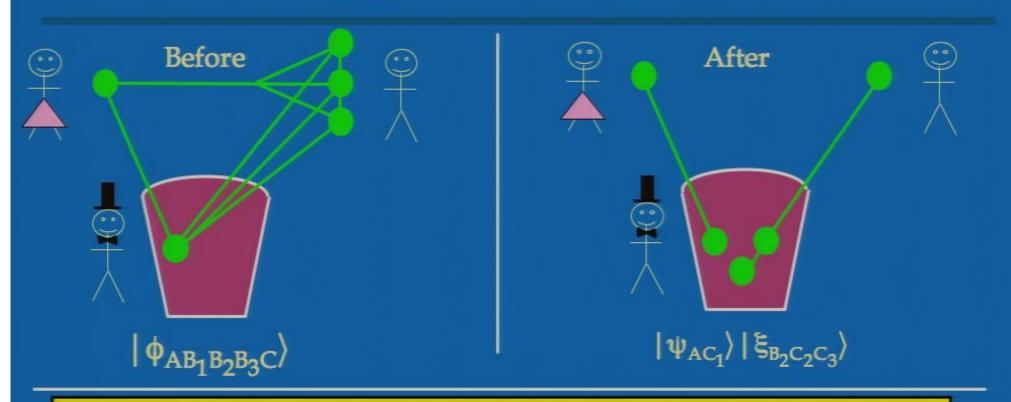


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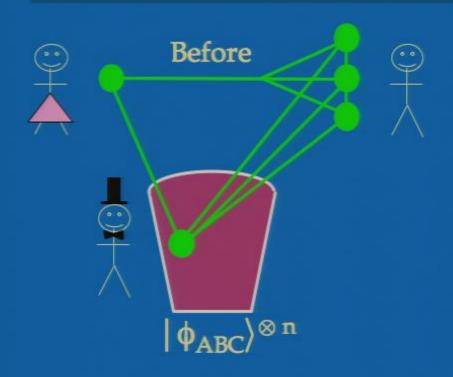


Alice never did anything \Rightarrow Her marginal state $\phi_A = \psi_A$ is unchanged

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Afterwards, entirely by Charlie.

and **distilled** entanglement with Charlie, just by discarding particles!



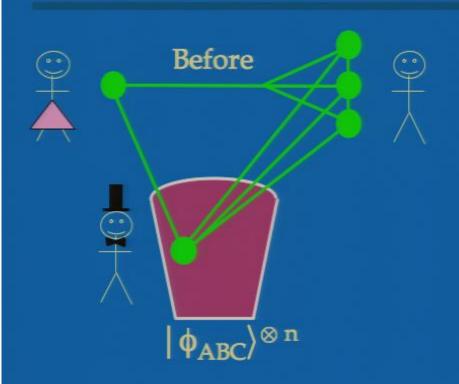
Uncertainty: von Neumann entropy $H(A)_{o} = H(\rho_{A}) = - tr[\rho_{A} \log \rho_{A}]$

Correlation: mutual information $I(A;B)_{\rho} = H(A)_{\rho} + H(B)_{\rho} - H(AB)_{\rho}$

$$I(A;B)_{\rho} = \begin{cases} 0 & \text{if and only if } \rho_{AB} = \rho_{A} \otimes \rho_{B} \\ m & \text{for } m \text{ pairs of correlated bits} \\ 2m & \text{for } m \text{ ebits (maximal)} \end{cases}$$

Pirsa: 07030039 Page 80/151

(10) 103+117 A(1) B)

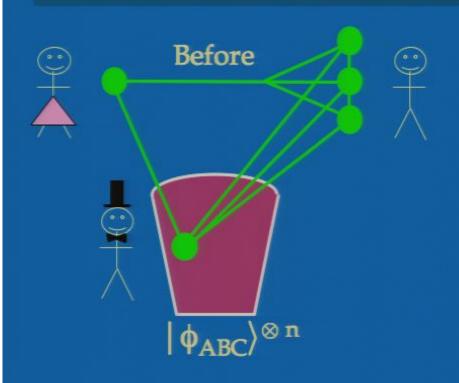


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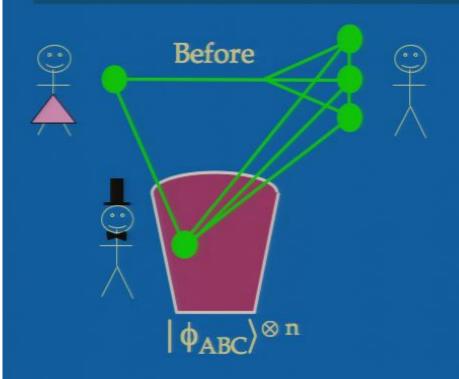


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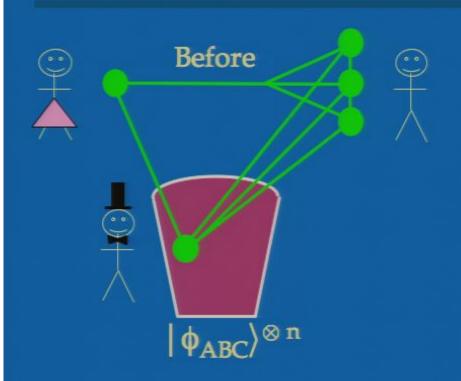
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Final mutual information: ε



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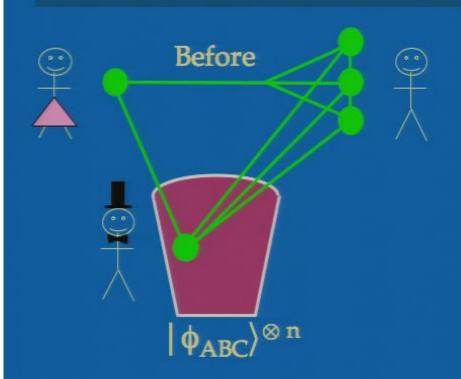
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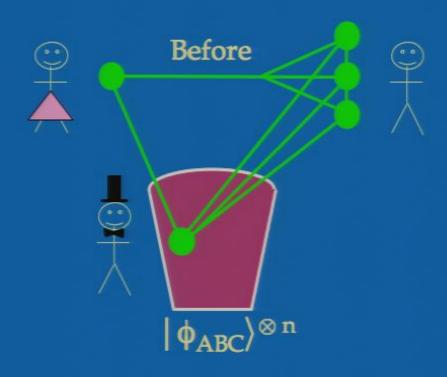
Initial mutual information: n I(A;B)_o

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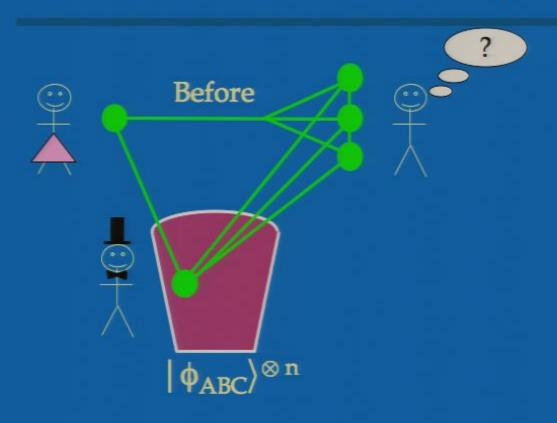
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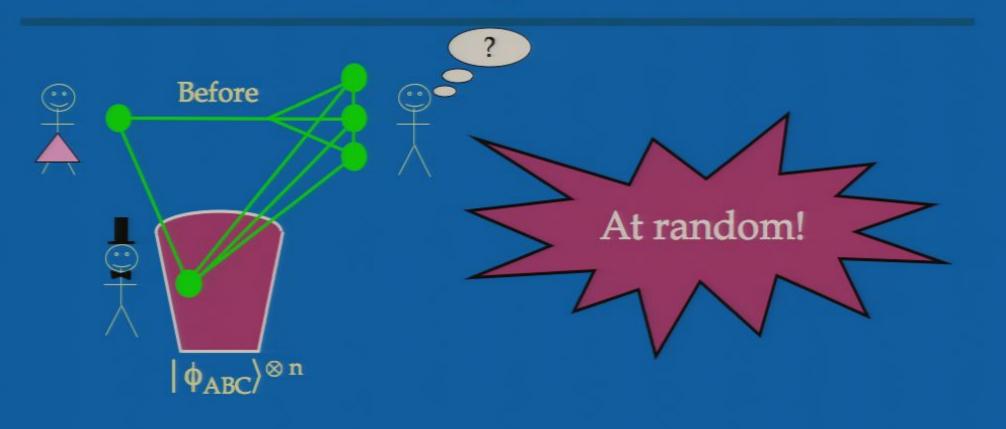
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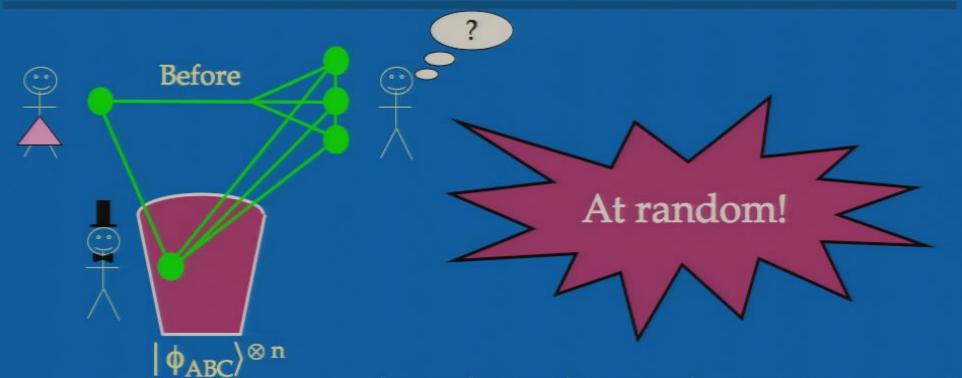
Pirsa: 07030039 Page 87/151



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(According to the unitarily invariant measure on the high-probability subspace of $B^{\otimes n}$.)

Pirsa: 07030039 Page 90/151



(According to the unitarily invariant measure on the high-probability subspace of B^{⊗n}.)

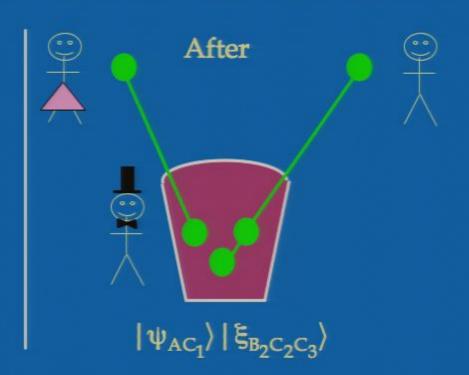
Bob can ignore the correlation structure of his state!

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Final accounting

Investment:

Bob sends Charlie $\sim n[I(A;B)_{\phi}]/2$ qubits



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Final accounting

Investment:

Bob sends Charlie ~n[I(A;B),]/2 qubits

Rewards:

- 1) Charlie holds Alice's purification
- 2) B and C establish $\sim n[I(B;C)_{\phi}]/2$ ebits



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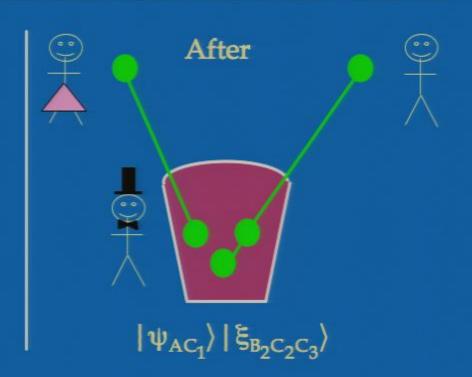
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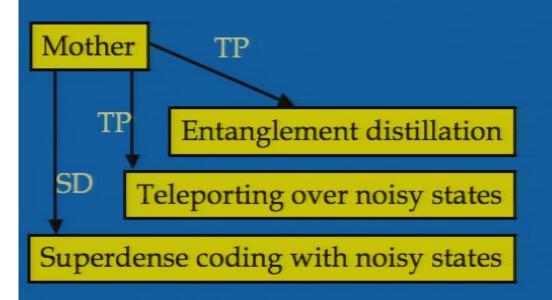
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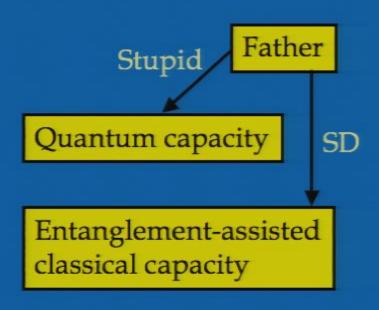
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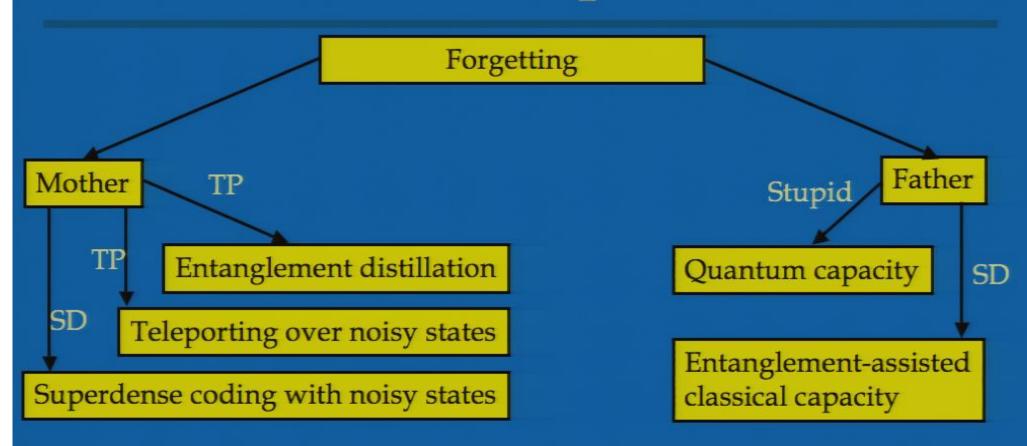
OK - but what good is it?

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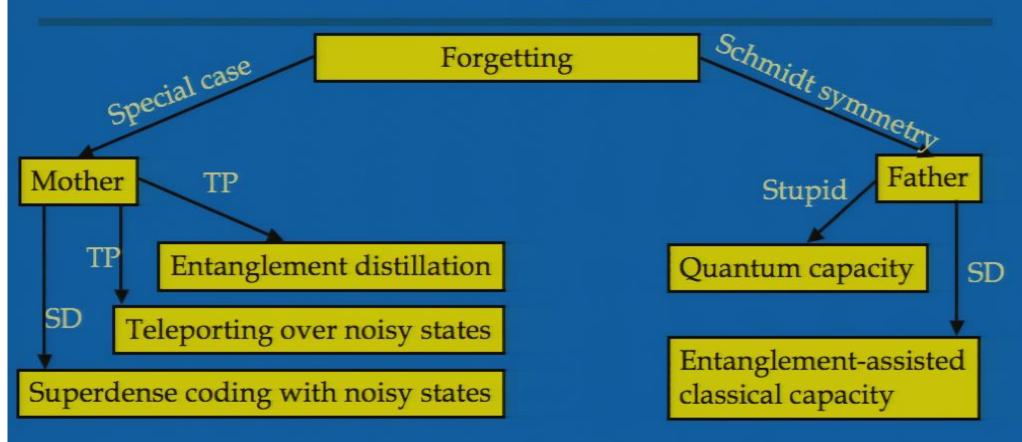




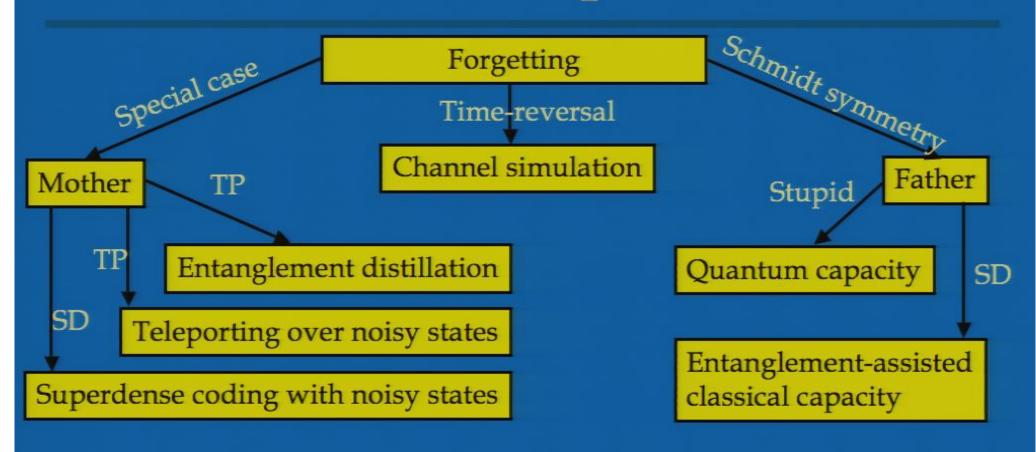
Pirsa: 07030039 Page 95/151



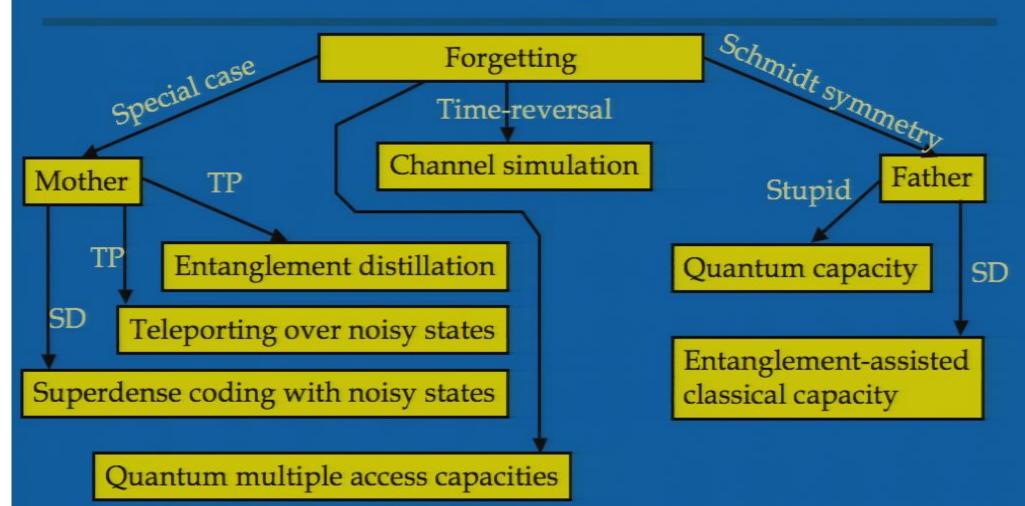
Pirsa: 07030039 Page 96/151



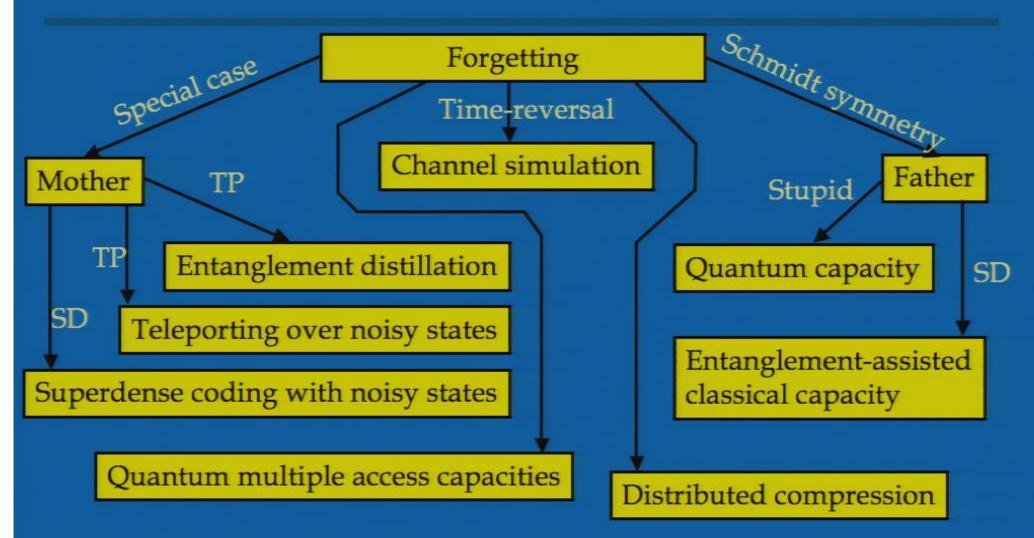
Pirsa: 07030039 Page 97/151



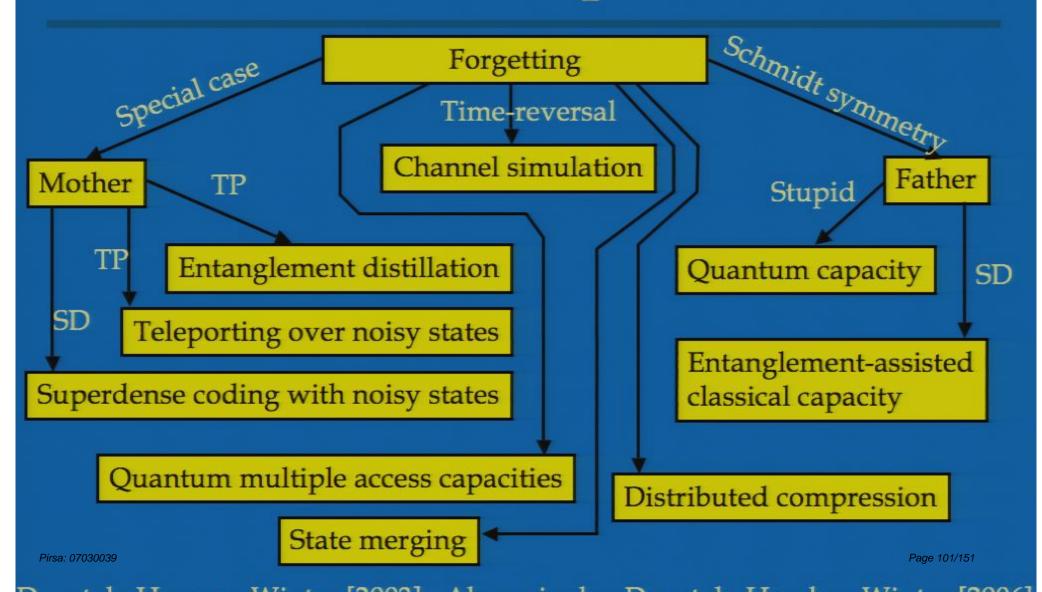
Pirsa: 07030039 Page 98/151

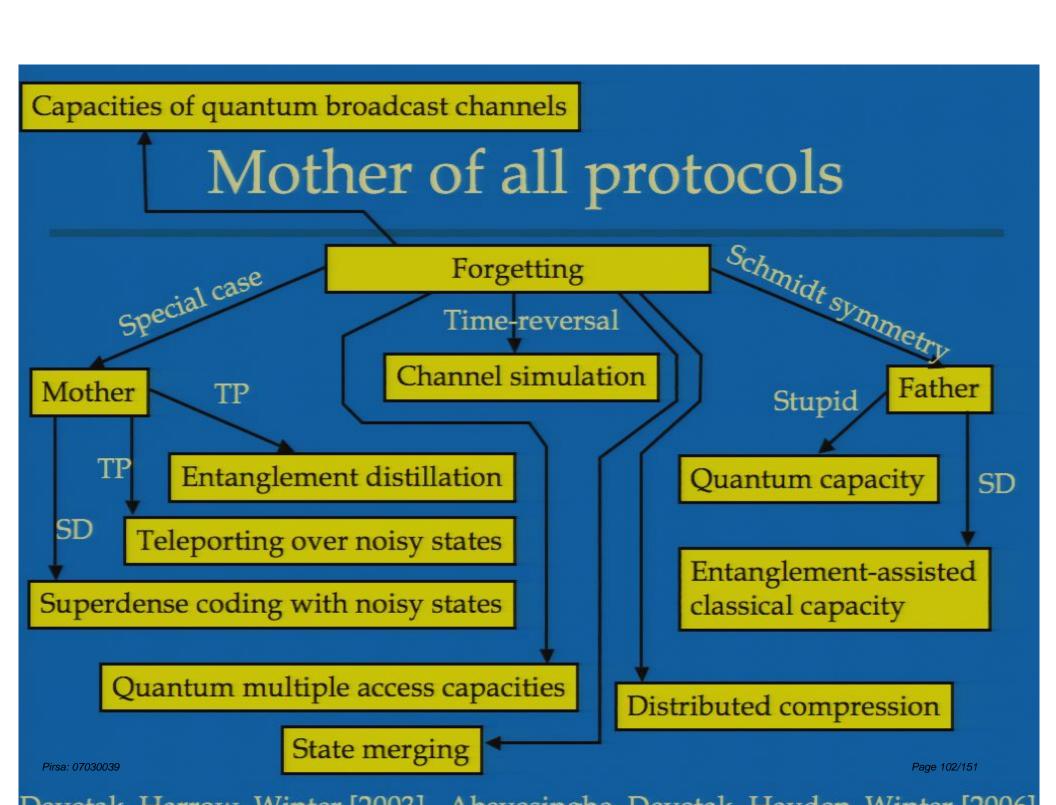


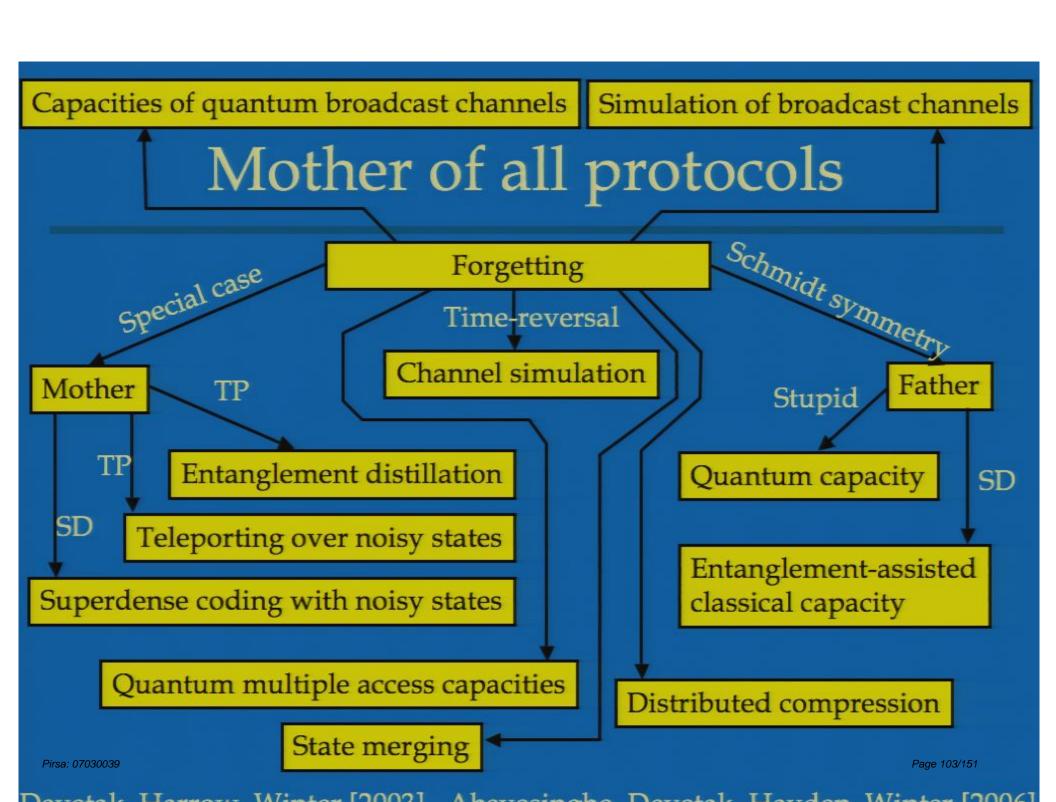
Pirsa: 07030039 Page 99/151

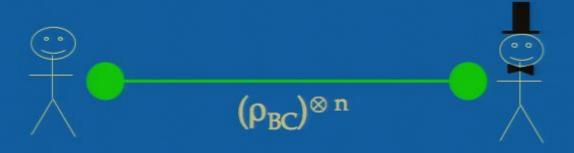


Pirsa: 07030039 Page 100/151

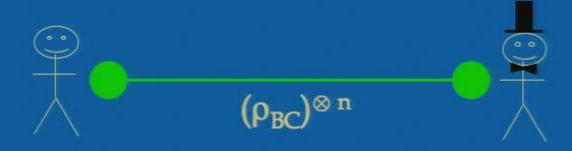








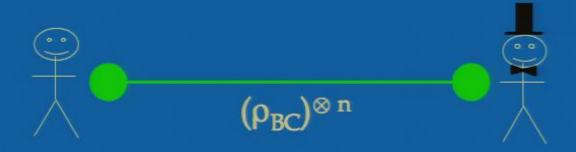
Pirsa: 07030039 Page 104/151



Bob and Charlie share many copies of a noisy entangled state and would like to convert it to ebits.

Only local operations and *classical communication* are allowed. Forgetting protocol good but uses *quantum communication*

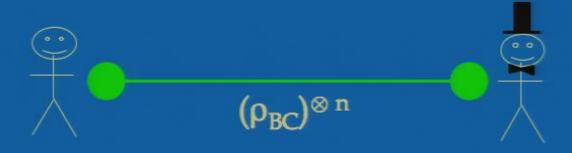
Pirsa: 07030039 Page 105/151



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Pirsa: 07030039 Page 106/151

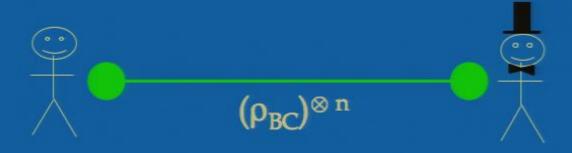


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Pirsa: 07030039 Page 107/151



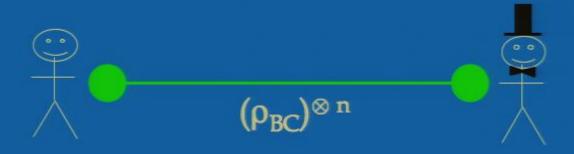
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Scientist knows less than nothing

A MATHEMATICS expert from Bristol has come up with a baffling theory to help explain the complex world of quantum physics.

Or Andreas Winter of Bristol University said it was possible to know less than nothing" and discovered that negative knowledge" does exist in the quantum world of extremely small things

But while the breakthrough may make sense to experts in quantum physics, it will leave the man to the sense completely haffled.

Dr Winter, a lecturer at the university's department of mathematics has been a ratching his head to comup with an explanation of his theory.

And his example will still leave mos seople wondering why he bothered

He said "If I tell you something, you will then know less than before I toll you. Such strange situations can occu because what it means to know some thing is very different in the quantum world."

Dr Winter worked with two other experts, Jonathan Oppenheim from Cambridge University and Michal Horodocki from the Institute for Theorem and Physics, in Calansk, Poland to

by Ian Turner

Lturner@bepp.co.uk

develop the theory. Their work has now been published in the leading science fournal Nature.

Quantum physics tries to explain the behaviour of particles, such as electrons protons, and neutrons, which are smaller than an atom

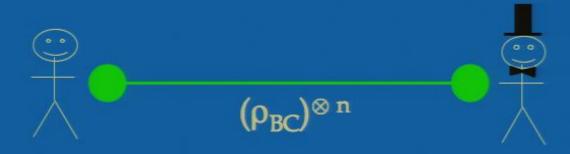
Scientists claim that in the quantum world, there are things we just cannot know no matter how clever we are

But confusingly they also say there are situations where someone can know more than everything. This is known as quantum entanglement, and when two people share entanglement, there can be negative information.

Mr Horadecki said the idea of negarive information can be put on a rig-

He said We can quantify information in terms of how much stuff I need to send you before you get to know something in the case of negative quantum information, you can get to know something without me sending you any quantum particles. Page 110/151 will sain the potential to learn more

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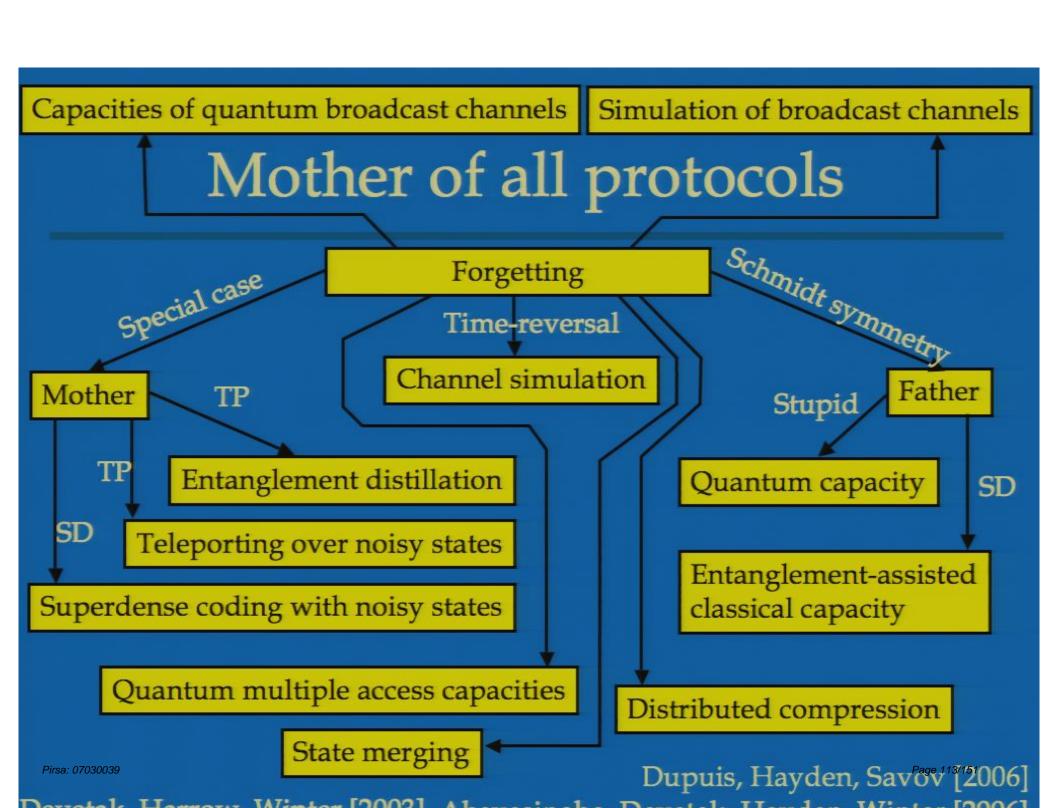
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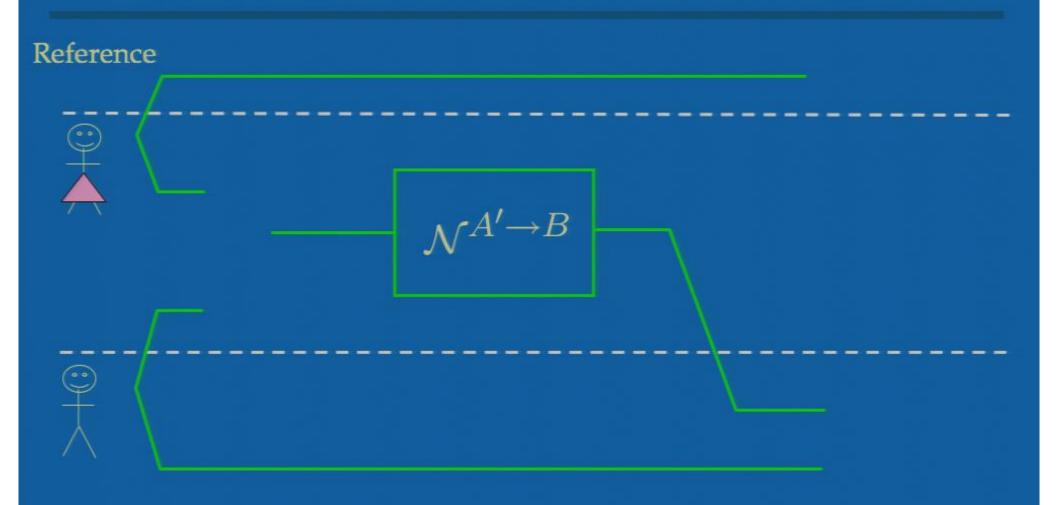


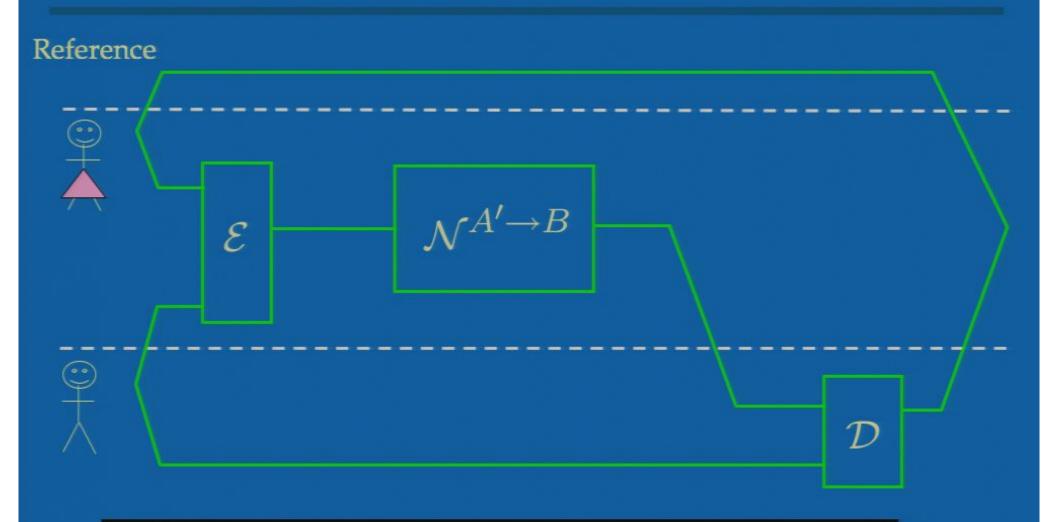
Reference



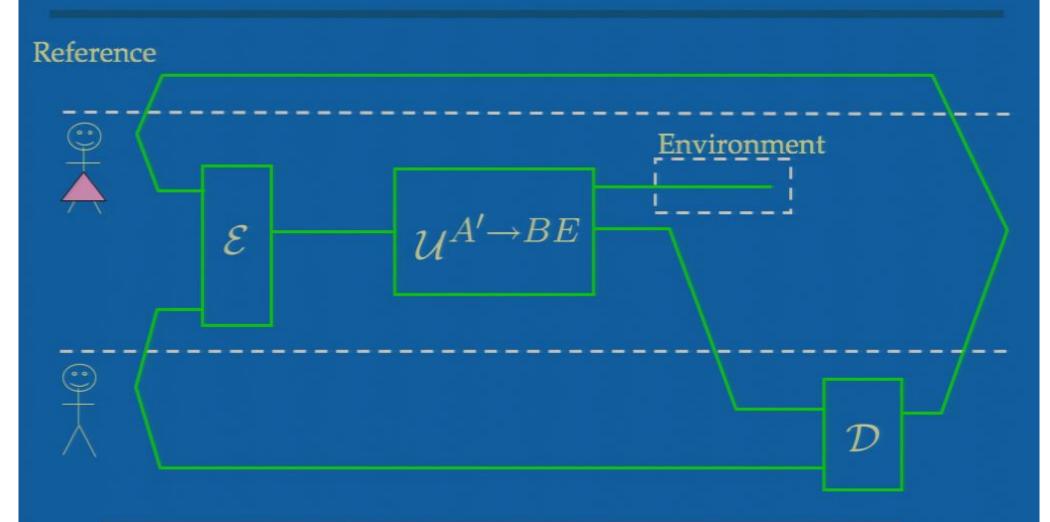




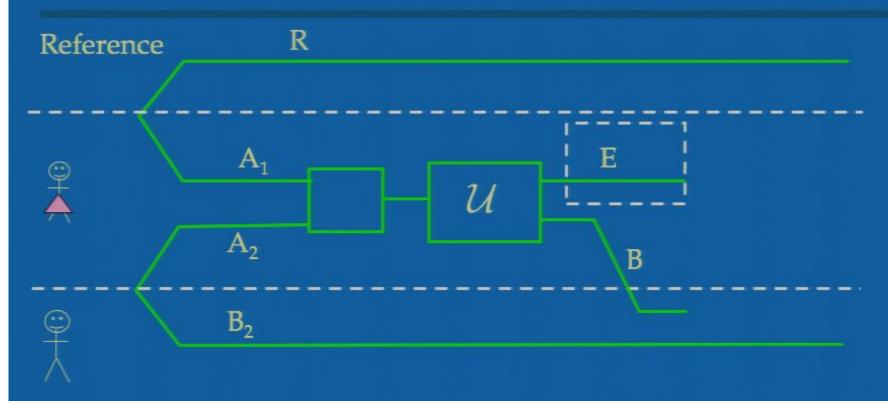


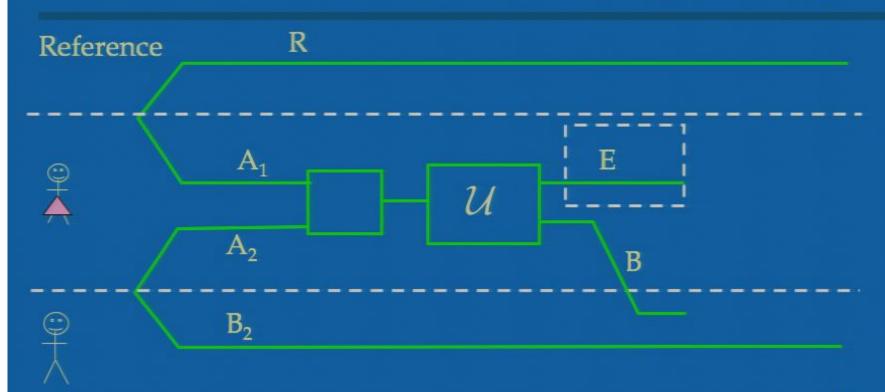


Pirsa: 0703003Objective is to transfer Alice's reference entanglement to Bob

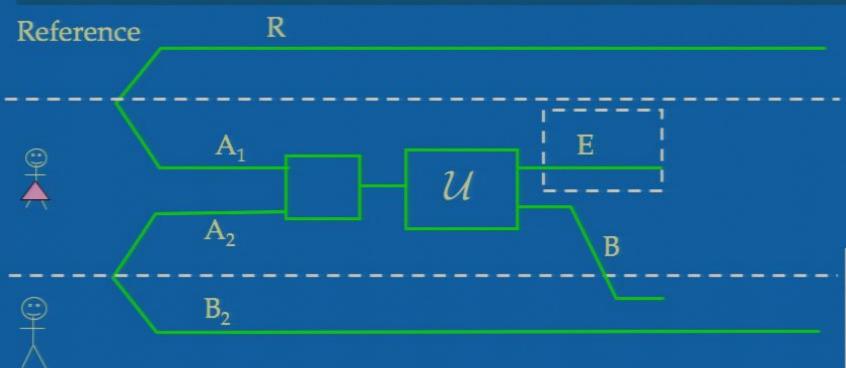


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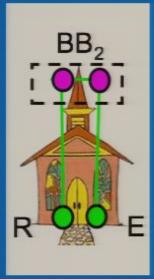


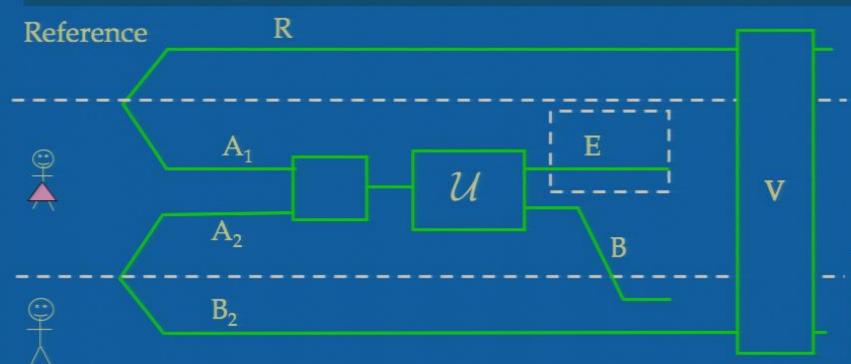


Who needs to do the forgetting? It is sufficient to ensure that there is a product state on $R \otimes E$



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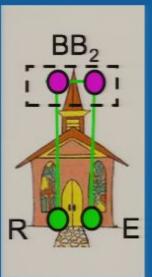




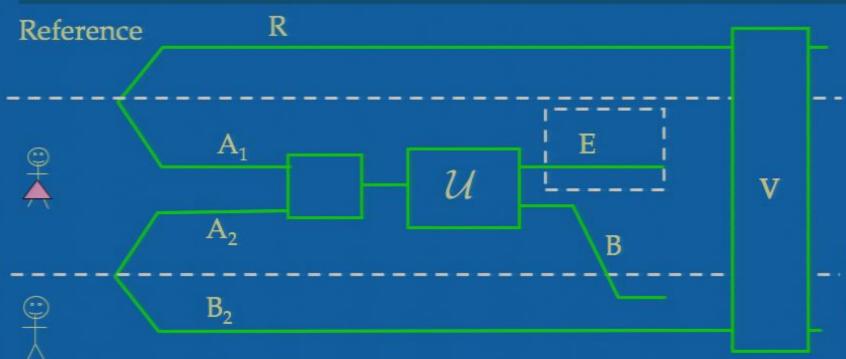
Who needs to do the forgetting? It is sufficient to ensure that there is a product state on $R \otimes E$

Imagine: apply a random unitary V to RB₂.

Pirsa: 07030039 Result: For sufficiently large B_2 , product state on $R \otimes E$!



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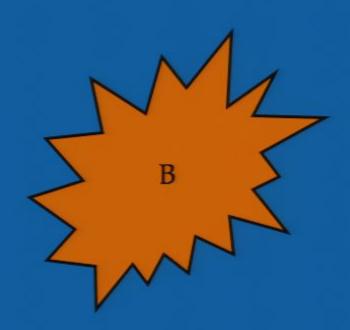
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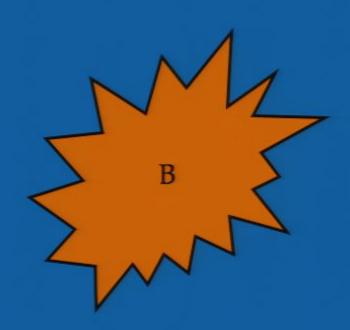
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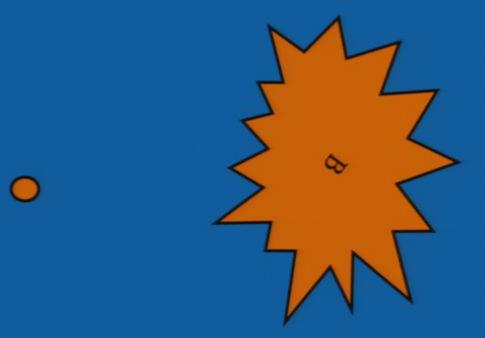
Page 123/151

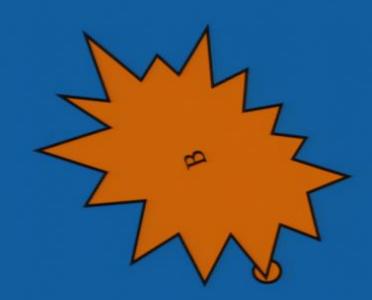


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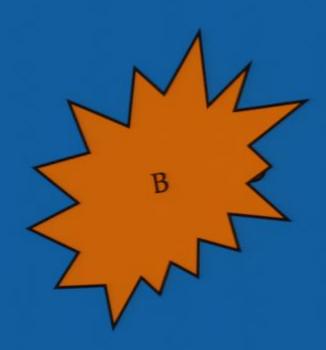


Pirsa: 07030039 Page 125/151

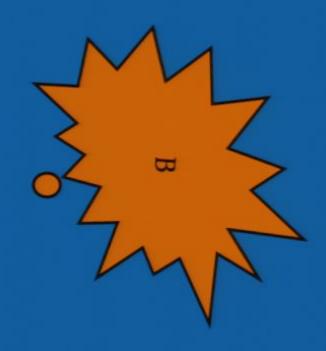




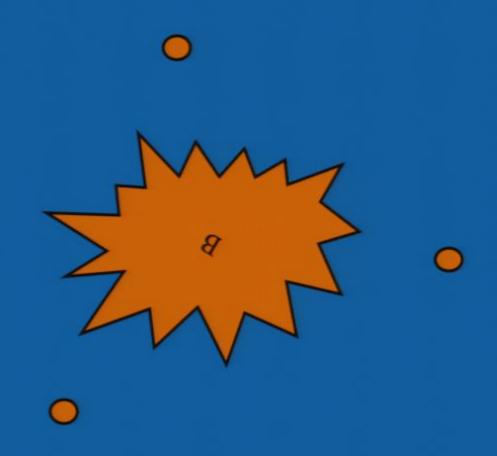
Pirsa: 07030039 Page 127/151



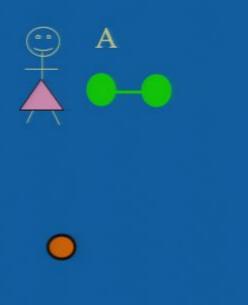
Pirsa: 07030039 Page 128/151

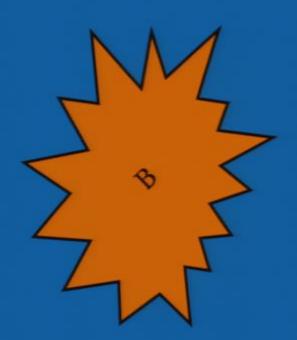


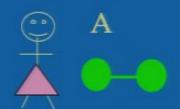
Pirsa: 07030039 Page 129/151

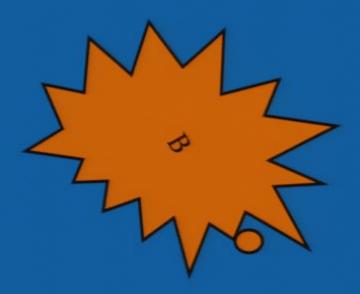


Pirsa: 07030039 Page 130/151

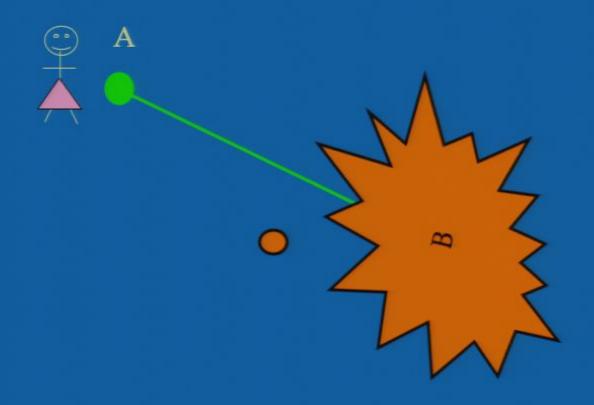


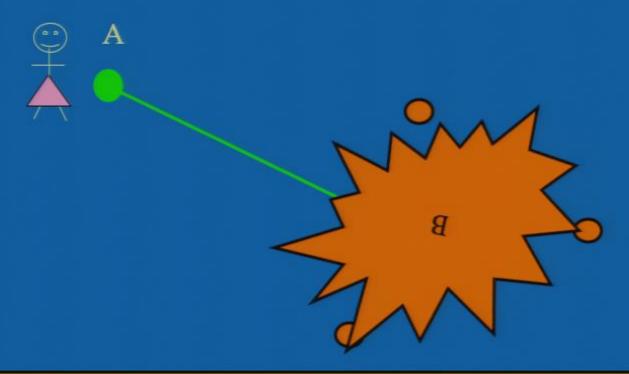






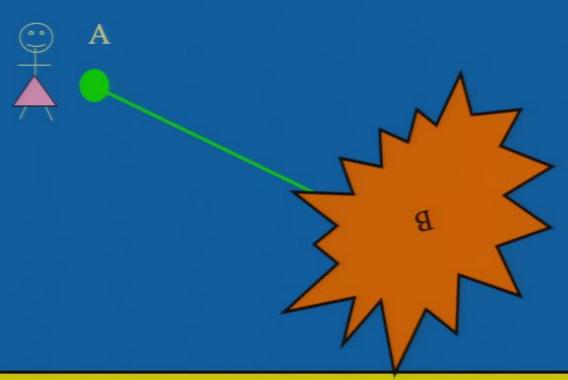
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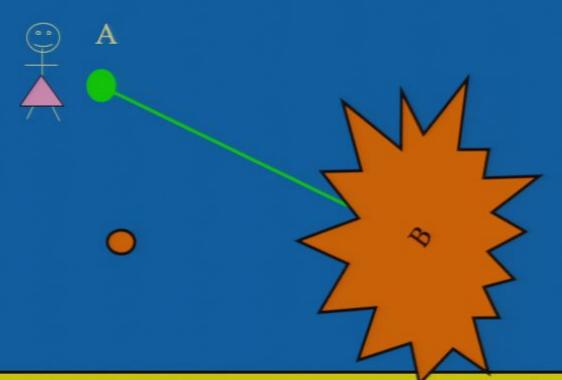
How long must Alice wait until the "information about A" gets radiated?

Pirsa: 07030039 Page 134/151



How long must Alice wait until the "information about A" gets radiated?

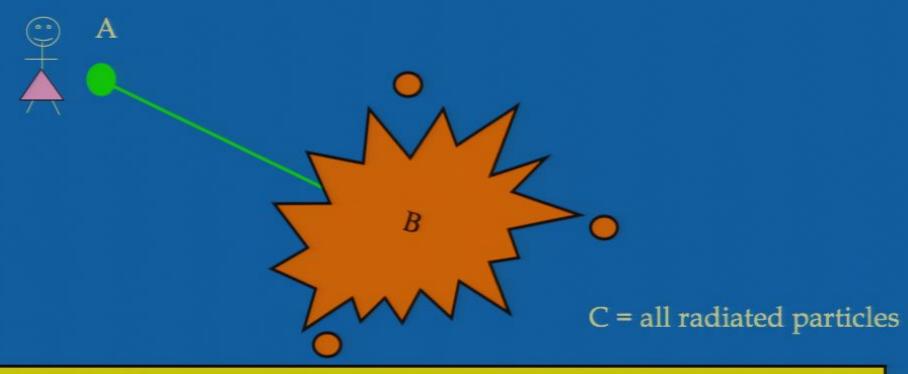
Equivalently, how long until the orange blob has forgotten about A?



C = all radiated particles

How long must Alice wait until the "information about A" gets radiated?

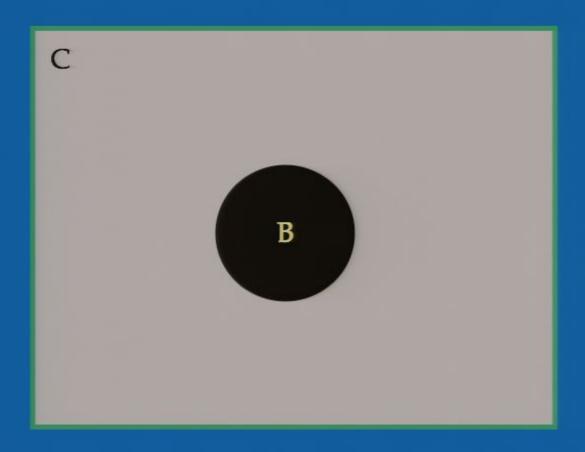
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How long must Alice wait until the "information about A" gets radiated?

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For sufficiently mixing dynamics, information about A is released almost immediately.

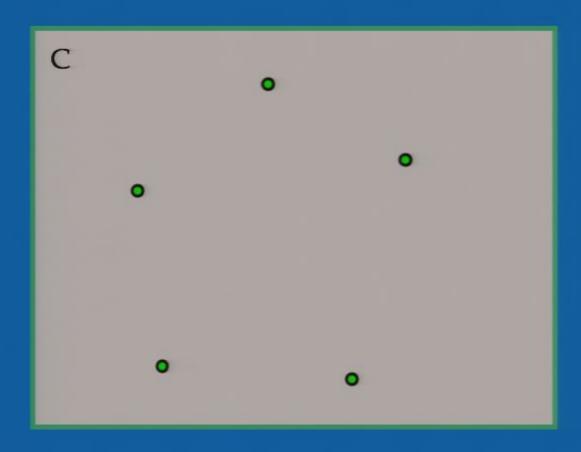


t₀: Pure state:

B: Black hole

C: Rest of universe

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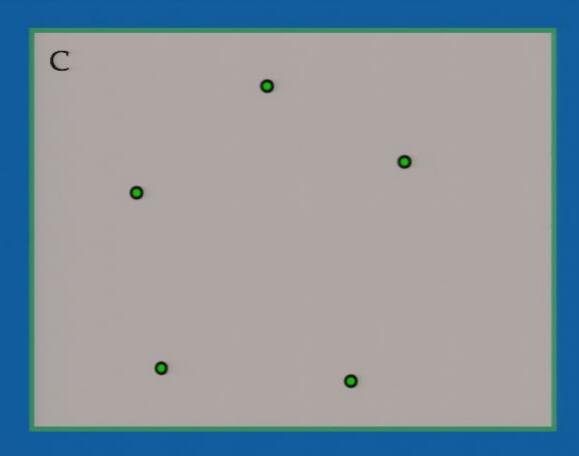
t₀: Pure state:

B: Black hole

C: Rest of universe

t₁: Thermal Hawking radiation

Pirsa: 07030039 Page 139/151



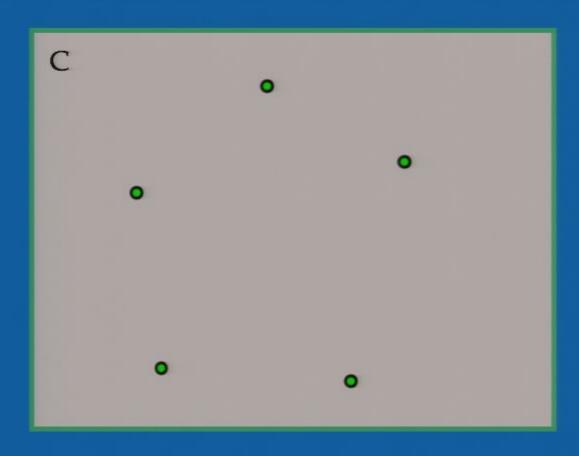
t₀: Pure state:

B: Black hole

C: Rest of universe

t₁: Thermal Hawking radiation

Pirsa: 07030039 Page 140/151



t₀: Pure state:

B: Black hole

C: Rest of universe

t₁: Thermal Hawking radiation

t₂: Radiation but no black hole

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t₀: Pure state:

B: Black hole

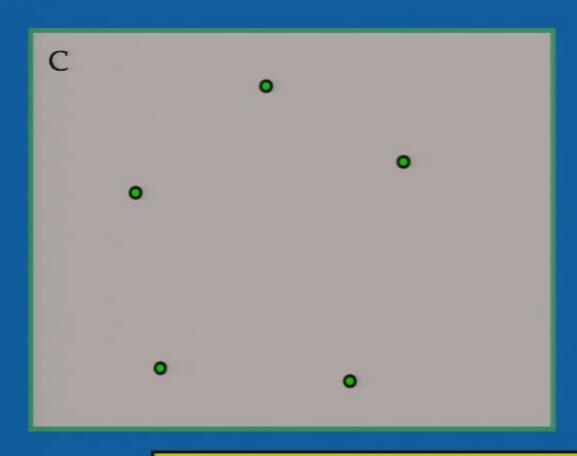
C: Rest of universe

t₁: Thermal Hawking radiation

t₂: Radiation but no black hole

Standard question: Is final state mixed or pure?

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t₀: Pure state:

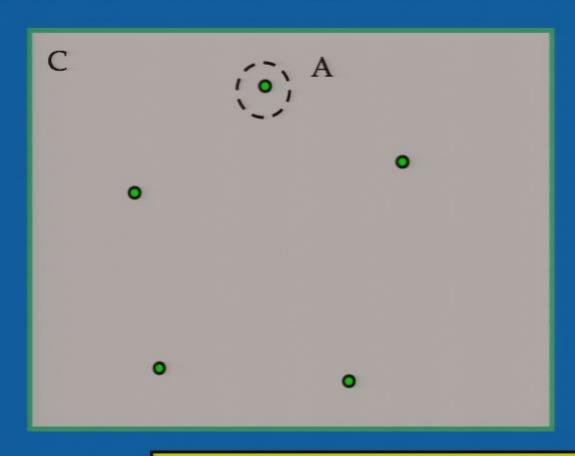
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t₁: Thermal Hawking radiation

t₂: Radiation but no black hole

Standard question: Is final state mixed or pure?



t₀: Pure state:

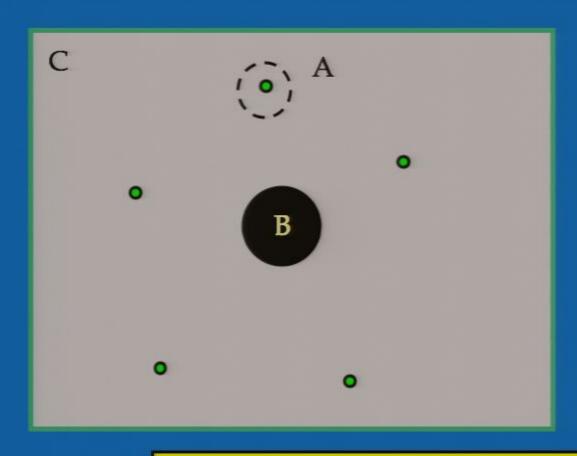
B: Black hole

C: Rest of universe

t₁: Thermal Hawking radiation

t₂: Radiation but no black hole

Standard question: Is final state mixed or pure?



t₀: Pure state:

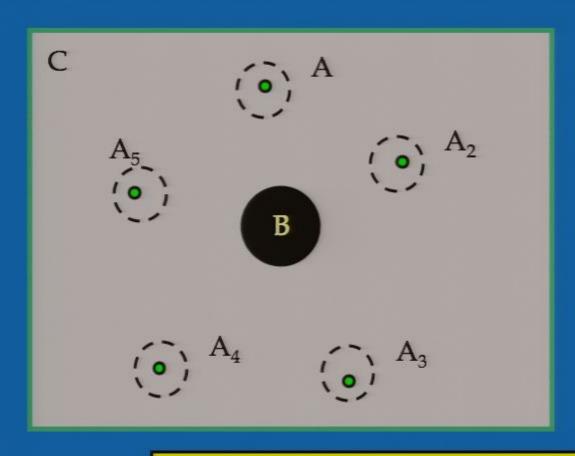
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t₀: Pure state:

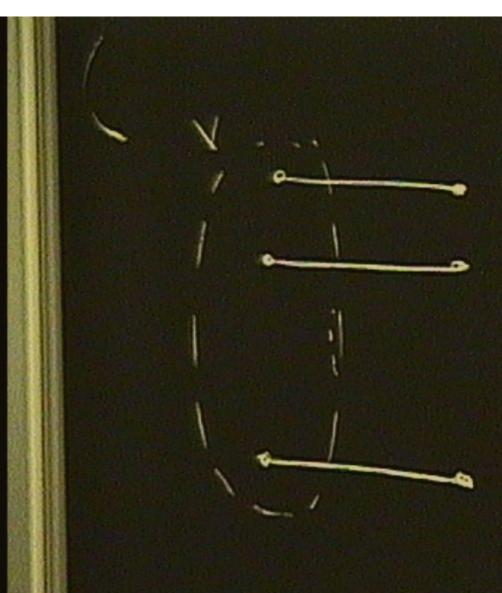
B: Black hole

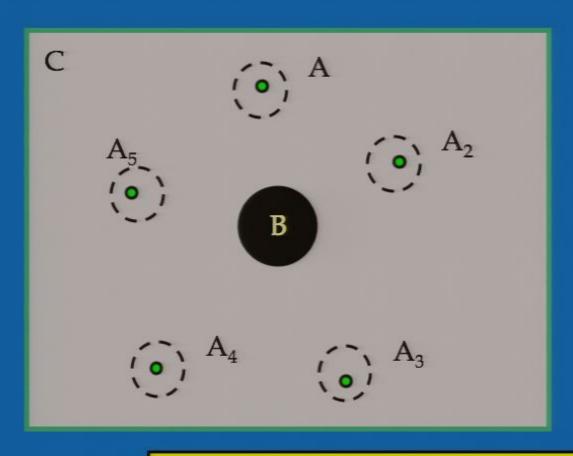
C: Rest of universe

t₁: Thermal Hawking radiation

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t₀: Pure state:

B: Black hole

C: Rest of universe

t₁: Thermal Hawking radiation

t₂: Radiation but no black hole

Standard question: Is final state mixed or pure?

Summary

- Forgetting is a basic primitive for quantum information theory
- Detailed understanding of how to do it most efficiently
- These methods are generated by generic unitary transformations: could be useful for understanding real physics

http://arxiv.org/abs/quant-ph/0606225

