

Title: Quantum Error Correction 1B

Date: Jan 09, 2007 05:00 PM

URL: <http://pirsa.org/07010015>

Abstract: 9-qubit Shor code, definition of a quantum error-correcting code, correcting linear combinations of errors, quantum error correction conditions, definition of distance

9-qubit code (Shor code) independently on n qubits

$$|0\rangle = (|1000\rangle + |1111\rangle)$$

Then: Any independent channel

$S(p)$ has fidelity $1 - 3p$
where fidelity $f = \langle S(p) | I \rangle$

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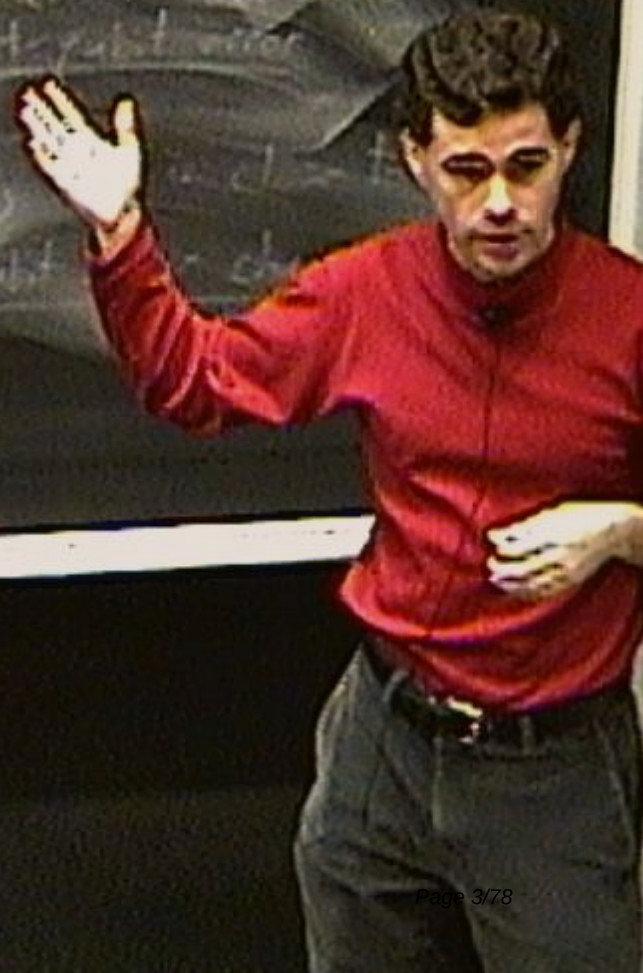
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Correct bit flip (X) errors on each set of 3 qubits

Thm: Any independent channel $S = \sum_i p_i S_i$, S_i is a channel, I is identity, p_i is probability, $\sum_i p_i = 1$, S_i is a qubit error channel, where fidelity $f_p, S_i(p) \geq 1 - f$

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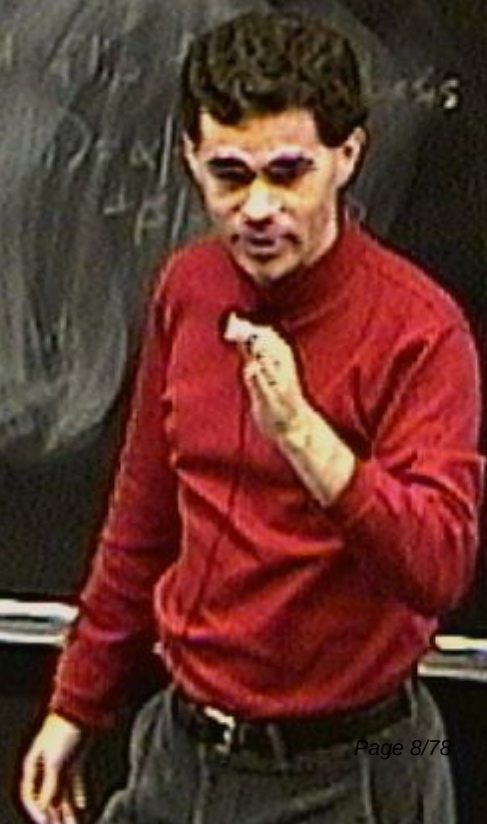
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✓
 $|\Psi\rangle |no\ error\rangle_{enc}$

↓ syndrome measurement
 $Z^{(1)} |\Psi\rangle |Z^{(1)}\rangle_{enc}$

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$$\rightarrow \cos \theta |\Psi\rangle |0 \text{ error}\rangle_{\text{anc}} - i \sin \theta Z^{(1)} |\Psi\rangle |Z^{(1)}\rangle_{\text{anc}}$$

Measure ancilla: Prob. $\cos^2 \theta : |\Psi\rangle |0 \text{ error}\rangle_{\text{anc}}$

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syndrome measurement

Measurement in Z basis

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$\rightarrow$ correct $|\Psi\rangle |Z^{(1)}\rangle_{\text{anc}}$

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syndrome measurement

$$\rightarrow \cos \theta |\bar{\Psi}\rangle |0 \text{ error}\rangle_{\text{anc}} - i \sin \theta Z^{(1)} |\Psi\rangle |Z^{(1)}\rangle_{\text{anc}}$$

Measure ancilla: Prob. $\cos^2 \theta$: $|\bar{\Psi}\rangle |0 \text{ error}\rangle_{\text{anc}}$

Prob. $\sin^2 \theta$: $Z^{(1)} |\Psi\rangle |Z^{(1)}\rangle_{\text{anc}}$

$\rightarrow |\bar{\Psi}\rangle |Z^{(1)}\rangle_{\text{anc}}$
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syndrome measurement

$$\rightarrow \cos \theta |\Psi\rangle |no error\rangle_{anc} - \sin \theta \hat{Z} |\Psi\rangle |yes\rangle_{anc}$$

Measure ancilla: Prob $(\cos^2 \theta) = |\langle \bar{\Psi} | \psi \rangle|^2$ error

Prob. sn't $\theta : \mathbb{Z}^{(n)} / \Psi / \mathbb{Z}^{(n)}$

→ correct $|\psi\rangle_{12}$ HNC

Thm. If a RECC corrects A, B , then it corrects $\alpha A + \beta B$
 " " " C , then it corrects $\text{span } C$

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Note: $\text{span}\{I, X, Y, Z\} = \text{all } 2 \times 2 \text{ matrices}$

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Correct bit flip errors on each set of 3 qubits

Correct phase flip error by comparing phases to see

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X & Z correction are independent correct 1 of each, such as Y

By this

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Correct phase flip (Z) error by comparing phases to see which is different

X & Z correction are independent $\Rightarrow$ can correct 1 of each, such as Y

By thm. 9-qubit code can correct any single-qubit error

Repetition code

$0 \mapsto 000 \xleftarrow{\text{majority}} 010$

Problems w/ Quantum error correction

- ✓ ① No cloning says cannot copy quantum states
- ✓ ② Measuring state collapses it — How can we correct the error without measuring the data?
- ✓ ③ Must correct X, Y, Z
- ✓ ④ What about continuous rotations, decoherence, etc.?

Pauli group: tensor Correction

Quantum state $|\psi\rangle$

Global quantum

$p \rightarrow 5/p$

Denit, not a positive

matrix

positive

the way

the way

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Pauli group: n -qubit tensor products of I, X, Y, Z
 overall phase $\pm 1, \pm i$

Quantum state $|\psi\rangle$, density matrix ρ
 General quantum operation $\mathcal{E}$
 $\rho \mapsto \mathcal{E}(\rho)$ linear map, completely positive
 trace preserving (CPTP)
 Kraus representation: $\mathcal{E}(\rho) = \sum_k K_k \rho K_k^\dagger$
 where $\sum_k K_k^\dagger K_k = I$
 (Kraus operators)

Pauli group: n -qubit tensor products of I, X, Y, Z
 w/ overall phase $\pm 1, \pm i$

and $X \otimes X, iY \otimes Z$

Quantum state $|\psi\rangle$, density matrix

Global quantum operation

$\rho \rightarrow S(\rho)$

completely positive

trace preserving

(Kraus operators)

Pauli group $\mathcal{P}_n$ in n -qubit tensor products of I, X, Y, Z

overall phase $\pm 1, \pm i$
Any $P \in \mathcal{P}_n$, $P^2 = \pm I$

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eg $X \otimes X, iY \otimes Z$

$P, Q \in \mathcal{P}_n$, either $PQ = QP$ or $PQ = -QP$

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$\text{span } \mathcal{P}_n = \text{all } n\text{-qubit matrices}$

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- Weight $\pm$ Pauli errors form a basis for all n -qubit errors

Sufficient condition to correct errors:
Encoded basis states $|\bar{\psi}_i\rangle$



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Basis errors $\{E_a\}$



Sufficient condition to correct errors:

Encoded basis states $|\bar{\Psi}_i\rangle$

Basis errors $\{E_a\}$

$\{E_a|\bar{\Psi}_i\rangle\}$ all orthogonal

Measure Subspaces $\{E_a|\bar{\Psi}_1\rangle, E_a|\bar{\Psi}_2\rangle, \dots\}$

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Measure outcomes $s_a = \{E_a|\bar{\psi}_1\rangle, E_a|\bar{\psi}_2\rangle, \dots\}$

Given a , map $s_a \rightarrow H_L$ $E_a|\bar{\psi}_i\rangle \mapsto |\bar{\psi}_i\rangle$

$$\langle \bar{\psi}_i | E_a^\dagger E_b | \bar{\psi}_j \rangle = \delta_{ab} \delta_{ij}$$

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True for some basis E_a of $\mathcal{E}$

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a) $\{|\bar{\Psi}_i\rangle\}$ are a basis for a QECC correcting $\mathcal{C}$ (or $\text{span}(\mathcal{C})$)

$$b) \langle \bar{\Psi}_i | E_a^\dagger E_b | \bar{\Psi}_j \rangle = C_{ab} \delta_{ij} \quad \forall E_a, E_b \in \mathcal{C} \quad \forall i, j$$

(C_{ab} does not depend on i, j)



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a) $\{|\bar{\Psi}_i\rangle\}$ are a basis for a QECC correcting $\mathcal{C}$ (or $\text{span } \mathcal{C}$)

$$b) \langle \bar{\Psi}_i | E_a^\dagger E_b | \bar{\Psi}_j \rangle = C_{ab} \delta_{ij} \quad \forall E_a, E_b \in \mathcal{C} \quad \forall i, j$$

(C_{ab} does not depend on i, j)

$$c) \langle \bar{\Psi} | E^\dagger E | \bar{\Psi} \rangle = C(E) \quad \forall |\bar{\Psi}\rangle \text{ in code} \quad \forall E \in \text{span } \mathcal{C}$$

Sufficient condition to correct errors:

Encoded basis states $|\bar{\psi}_i\rangle$

Basis errors $\{E_a\}$

$\{E_a|\bar{\psi}_i\rangle\}$ all orthogonal

Measure subspaces $S_a = \{E_a|\bar{\psi}_1\rangle, E_a|\bar{\psi}_2\rangle, \dots\}$

Given a , map $S_a \rightarrow H_L$ $E_a|\bar{\psi}_i\rangle \mapsto |\bar{\psi}_i\rangle$

$$\langle \bar{\psi}_i | E_a^\dagger E_b | \bar{\psi}_j \rangle = \delta_{ab} \delta_{ij}$$

True for some basis E_a of $\mathcal{E}$

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Proof: b) $\Rightarrow$ a) Diagonalize C_{ab} by choosing basis E_i for $\mathcal{E}$, get sufficient condition

$$a) \Rightarrow c) \quad U \text{ recovery: } UE|\bar{\Psi}\rangle = a|\bar{\Psi}\rangle|anc_1\rangle, \quad UE|\bar{\Phi}\rangle = b|\bar{\Phi}\rangle|anc_2\rangle$$

$$UE(|\bar{\Psi}\rangle + |\bar{\Phi}\rangle) = C(|\bar{\Psi}\rangle + |\bar{\Phi}\rangle|anc_1\rangle) = a|\bar{\Psi}\rangle|anc_1\rangle + b|\bar{\Phi}\rangle|anc_2\rangle \Rightarrow a=b$$

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a) $\Rightarrow$ c) U recovery: $UE|\bar{\psi}\rangle = a|\bar{\psi}\rangle|a\rangle_{\mathcal{C}_1}$, $UE|\bar{\phi}\rangle = b|\bar{\phi}\rangle|b\rangle_{\mathcal{C}_1}$

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c) $\Rightarrow$ b): Consider $E = E_a \pm E_b, E_a \pm E_b$
 $\langle \bar{\Psi} | E_a^\dagger E_b | \bar{\Psi} \rangle = c_{ab} \quad \forall |\bar{\Psi}\rangle$

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Cor.: A QECC of distance $2t+1$ will correct t -qubit errors

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 " " " " " will correct $(d-1)$ -qubit erasure errors

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Prop.: A QECC of distance d will detect $d-1$ errors

Surfactant concentration

$$\langle \Psi | \epsilon | \Psi \rangle = \text{tr}(\rho \epsilon) \quad \rho = |\Psi\rangle\langle\Psi|$$

all orthogonal

Many eigenstates $\epsilon_1, \epsilon_2, \dots, \epsilon_n$
 G H L Ψ
 K, H, ϵ

True for some basis ϵ of $\mathcal{E}$

Self-consistent

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \rho = |\Psi\rangle\langle\Psi|$$

For distance d code, interested in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

$$G = \{K_1, \dots, K_n\} \quad E = \sum_{i=1}^n \alpha_i K_i$$

True for some basis E of $\mathcal{E}$

Surf. code and its context

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \rho = |\Psi\rangle\langle\Psi|$$

For distance of code, t in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E)$$

$$K \Psi | E | \Psi \rangle = \text{tr}(\rho E)$$

True for some cases of E

Superficially, one can see that

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \rho = |\Psi\rangle\langle\Psi|$$

For distance d code, interested in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}_{d-1}(\rho_{d-1} E) \quad (\rho_{d-1} = \text{tr}_{d-1} \rho)$$

$$|\Psi\rangle \in \mathcal{C} \quad \text{and} \quad \rho_{d-1}$$

True for some ρ_{d-1} of $\mathcal{C}$

Superficially, one can think of the code as

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \bar{\Psi} \quad \rho = |\Psi\rangle\langle\Psi|$$

For distance d code, interested in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

independent of $|\Psi\rangle$ $\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}_{d-1}(\rho_{d-1} E)$ $(\rho_{d-1} = \text{tr}_{d-1} \rho)$

True for some basis C of E

Superficially, one can think of the code as

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \bar{\Psi} \quad \rho = |\Psi\rangle\langle\Psi|$$

For distance d code, interested in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

independent of $|\Psi\rangle$ — $\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}_{d-1}(\rho_{d-1} E) \quad (\rho_{d-1} = \text{tr}_{\text{other}} \rho)$

Let $E = |i\rangle\langle i| \Rightarrow \rho_{d-1}$ independent of $|\Psi\rangle$

True for some basis E_i of $\mathcal{E}$



Suppose that ρ is a state on n qubits

$$\langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}(\rho E) \quad \bar{\Psi} = |\Psi\rangle \langle \Psi|$$

For distance d code, interested in E acting on $d-1$ qubits

Focus on some particular set of $d-1$ qubits

independent of $|\Psi\rangle$ $\Rightarrow \langle \bar{\Psi} | E | \bar{\Psi} \rangle = \text{tr}_{d-1}(\rho E) \quad (\rho_{d-1} = \text{tr}_{\text{other}} \rho)$

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True for some QECC has distance d