

Title: The Cyclic Model and the Cosmological Conundrums

Date: Nov 12, 2006 10:30 AM

URL: <http://pirsa.org/06110061>

Abstract: Joel Erickson

THE CYCLIC MODEL AND THE COSMOLOGICAL CONUNDRUMS

JOEL ERICKSON

ISCAP, Columbia University

with S. T. Gratton, P. J. Steinhardt and N. Turok

hep-th/0607164

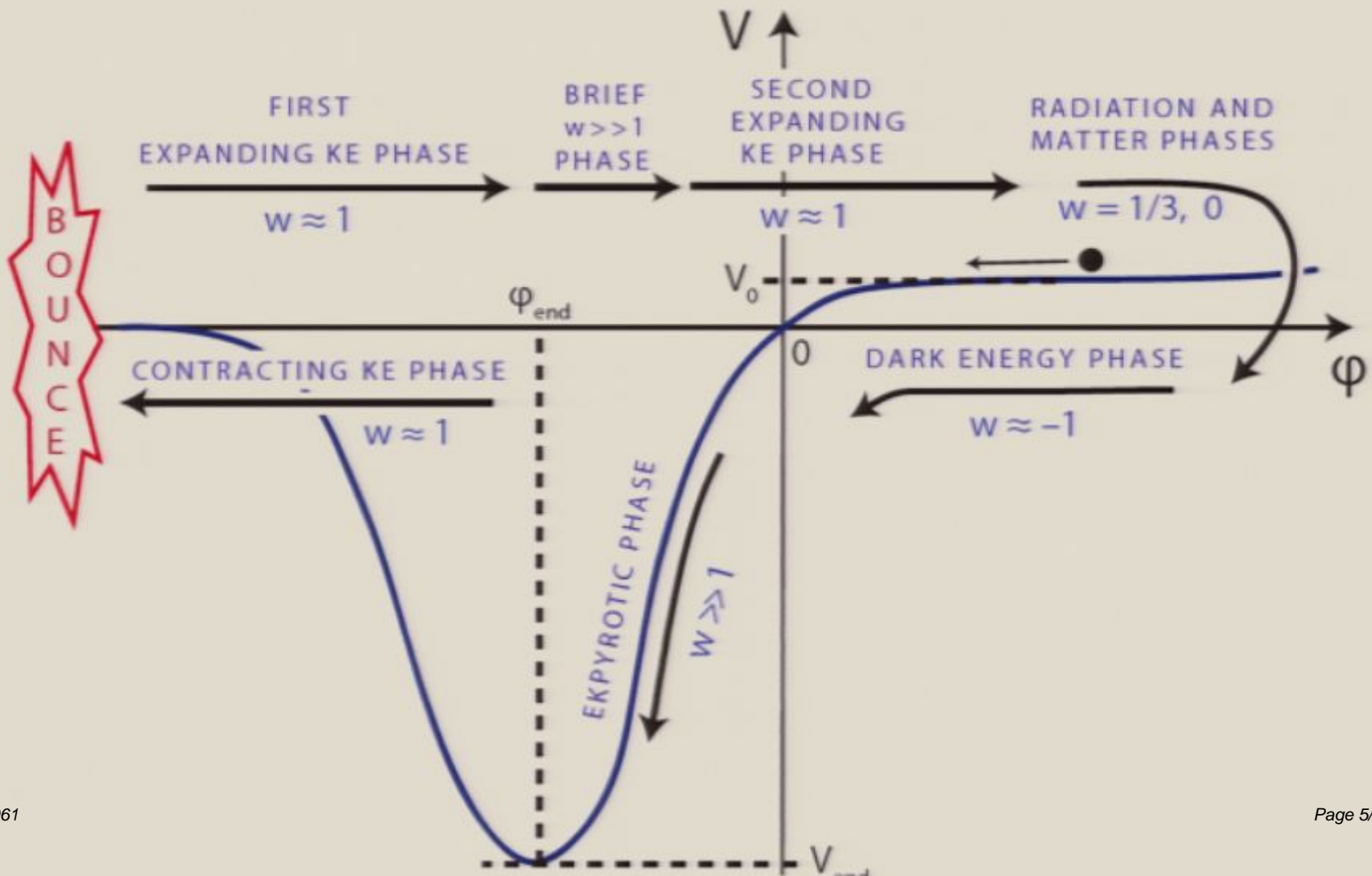
OUTLINE

- I. The cyclic model
- II. Evolution of scales
- III. The cosmological conundrums
- IV. Global structure

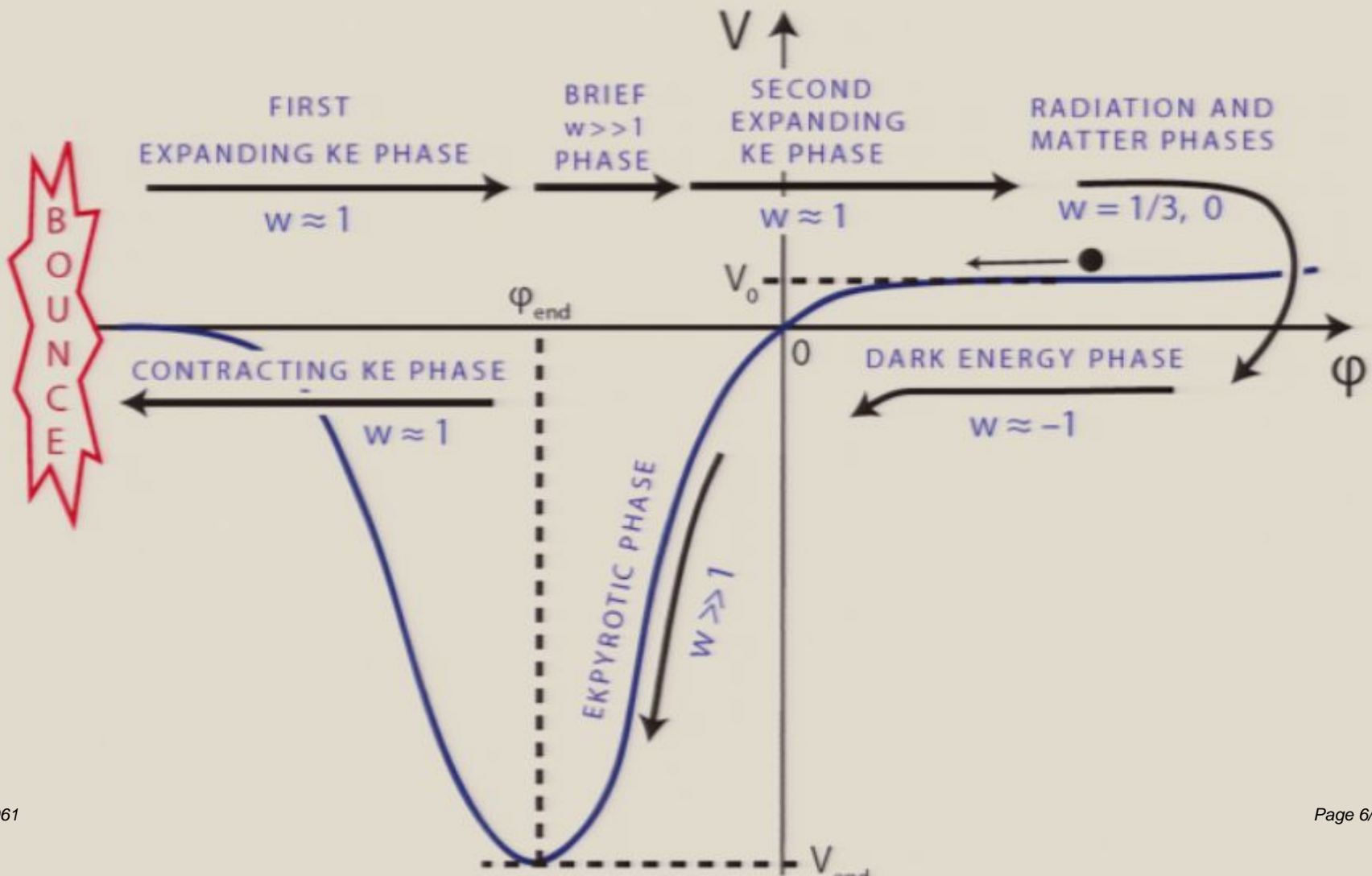
OUTLINE

- I. The cyclic model
- II. Evolution of scales
- III. The cosmological conundrums
- IV. Global structure

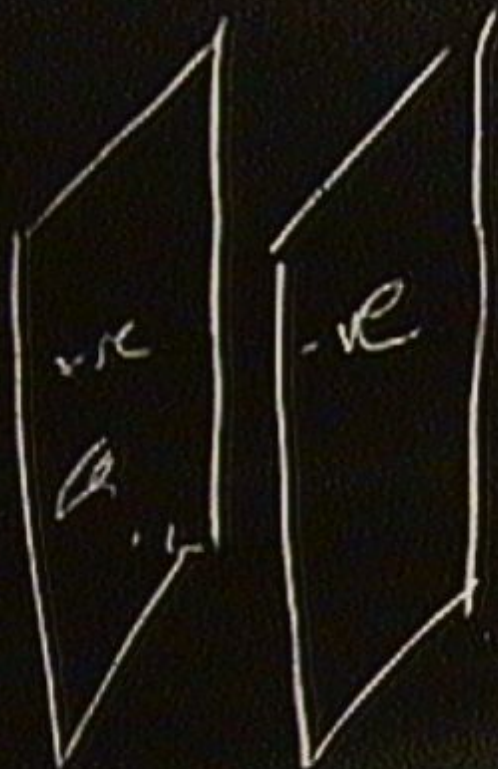
THE POTENTIAL



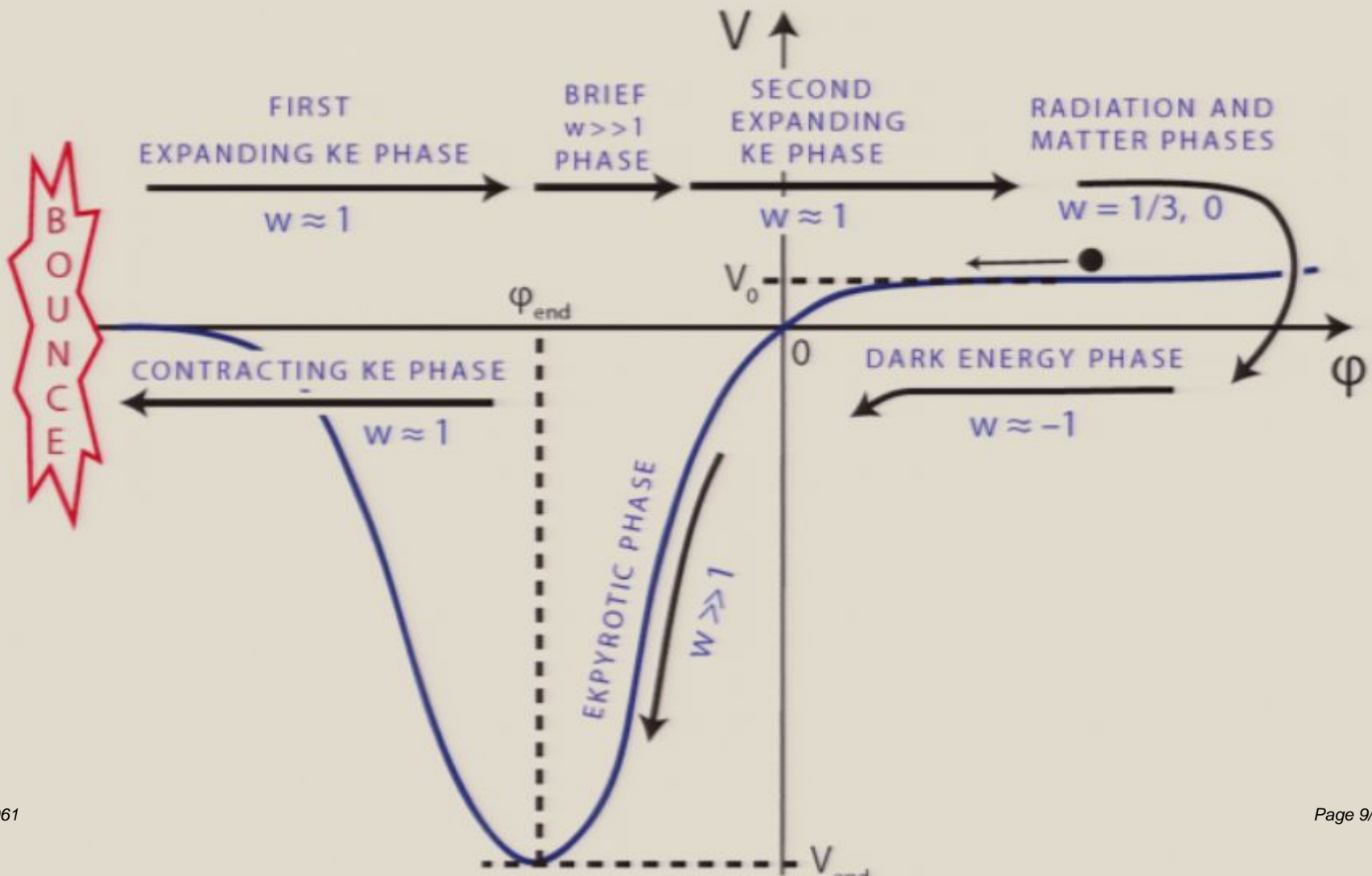
THE POTENTIAL

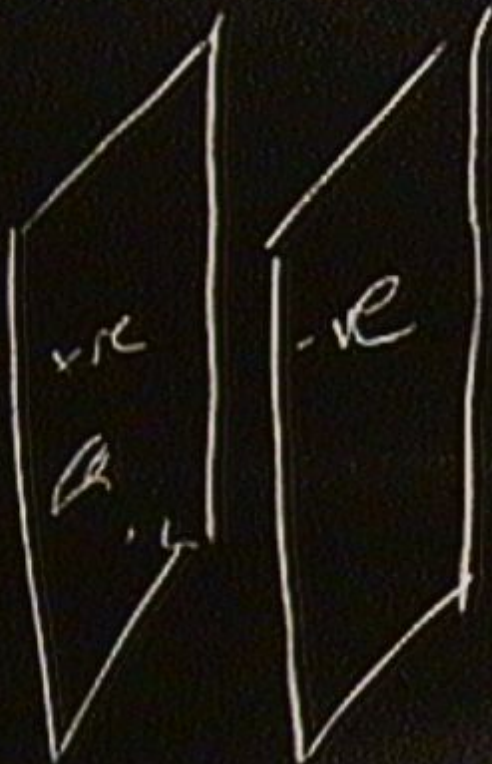






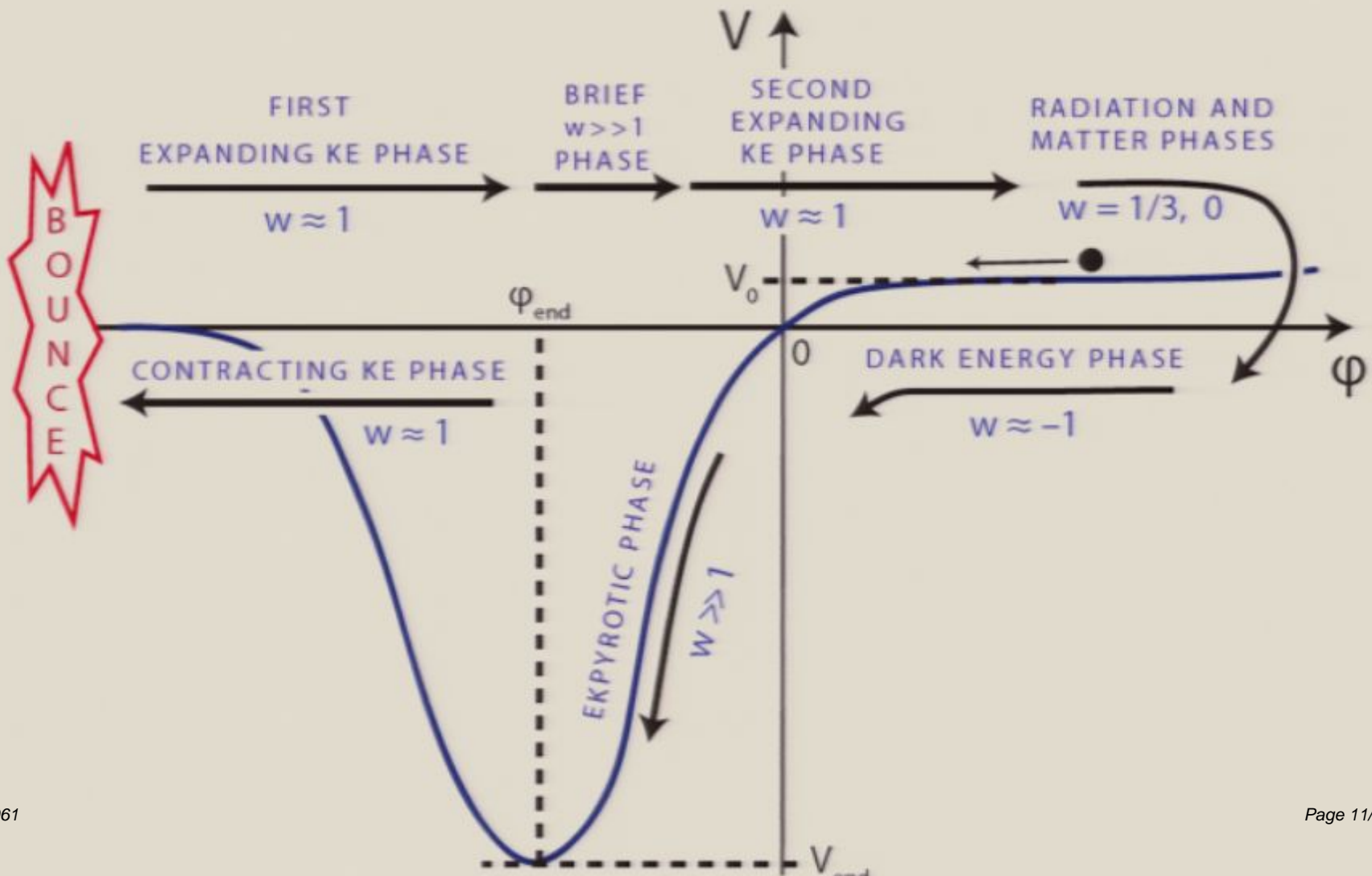
THE POTENTIAL



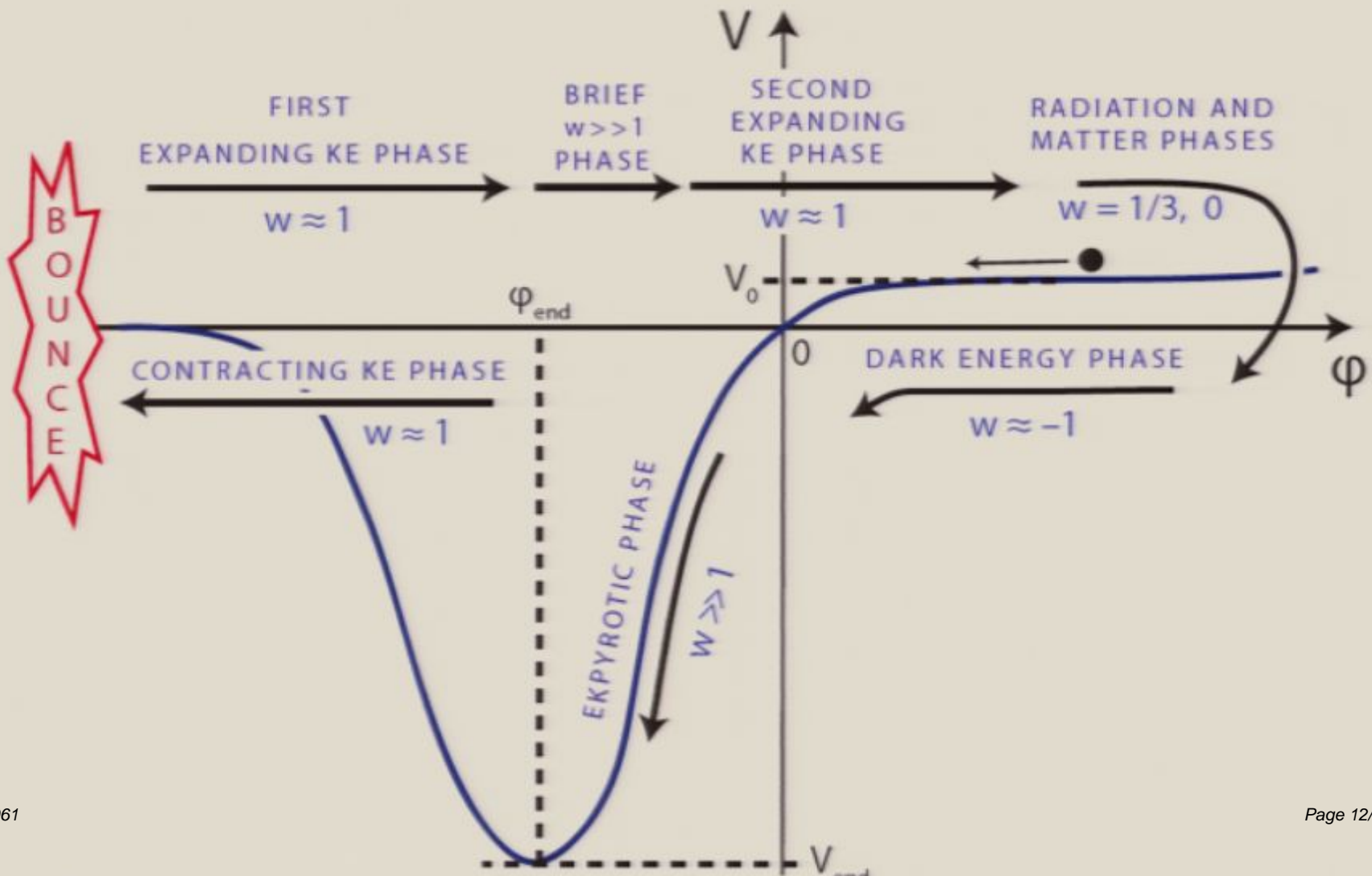


$$\frac{1}{aH}$$

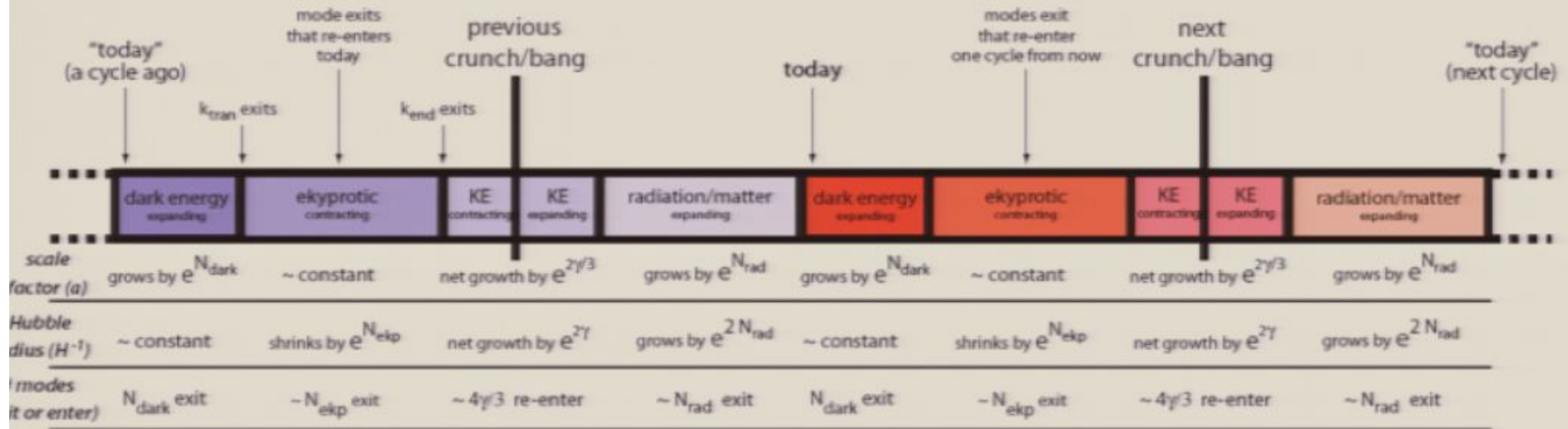
THE POTENTIAL



THE POTENTIAL



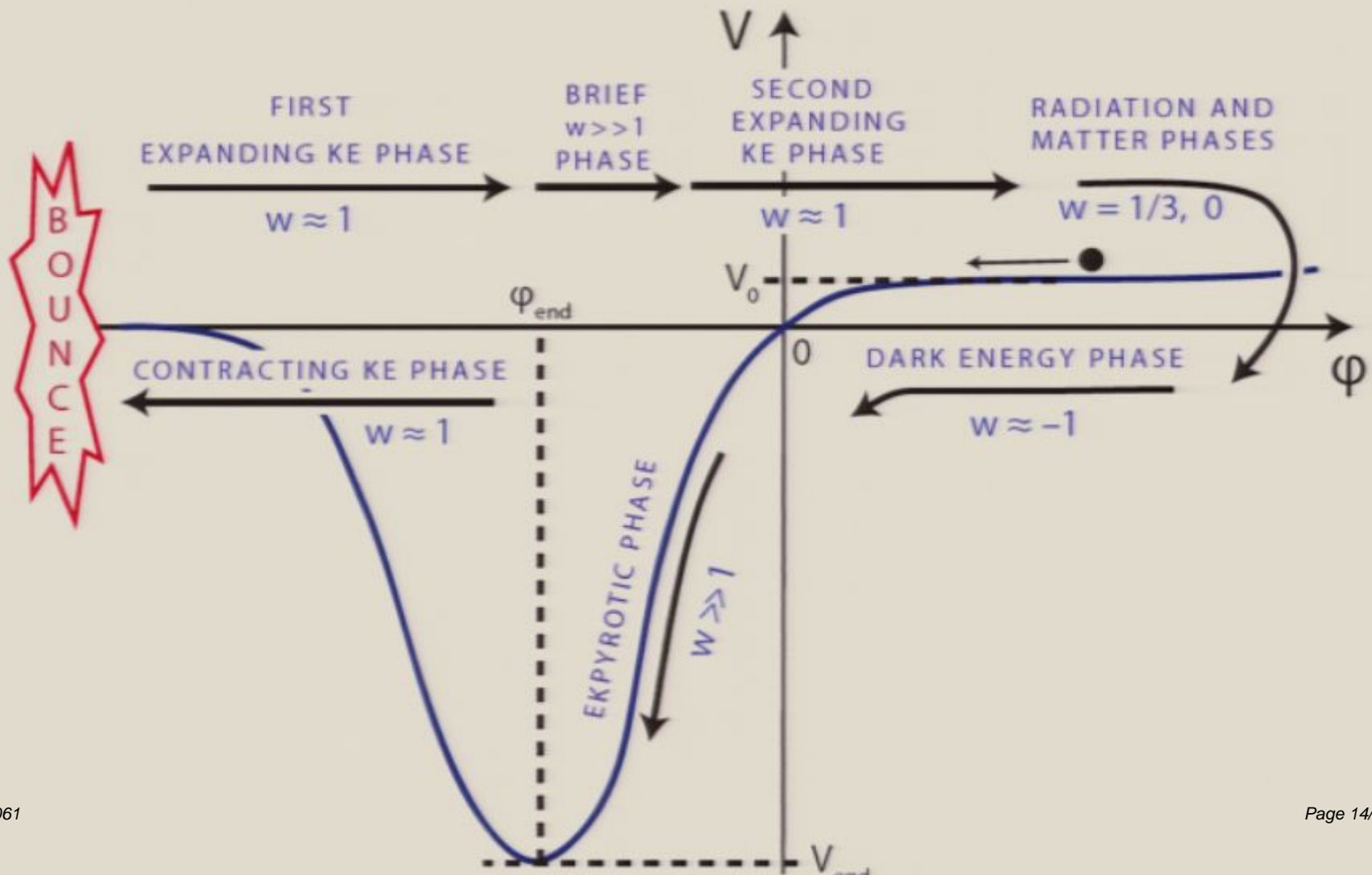
SCALES I



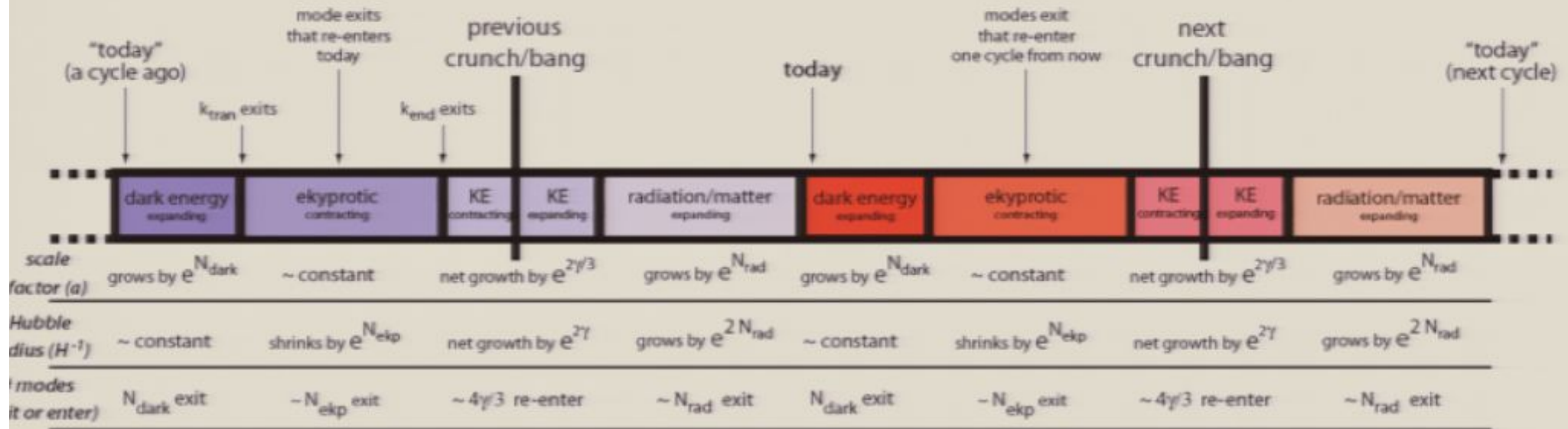
$$\gamma = \log \frac{(-V_{\text{end}})^{1/4}}{T_{\text{reheat}}}, \quad N_{\text{ekp}} = \log \sqrt{\frac{-V_{\text{end}}}{V_0}}, \quad N_{\text{rad}} = \log \frac{T_{\text{rh}}}{T_0}$$

$$N_{\text{ekp}} = 2(\gamma + N_{\text{rad}})$$

THE POTENTIAL



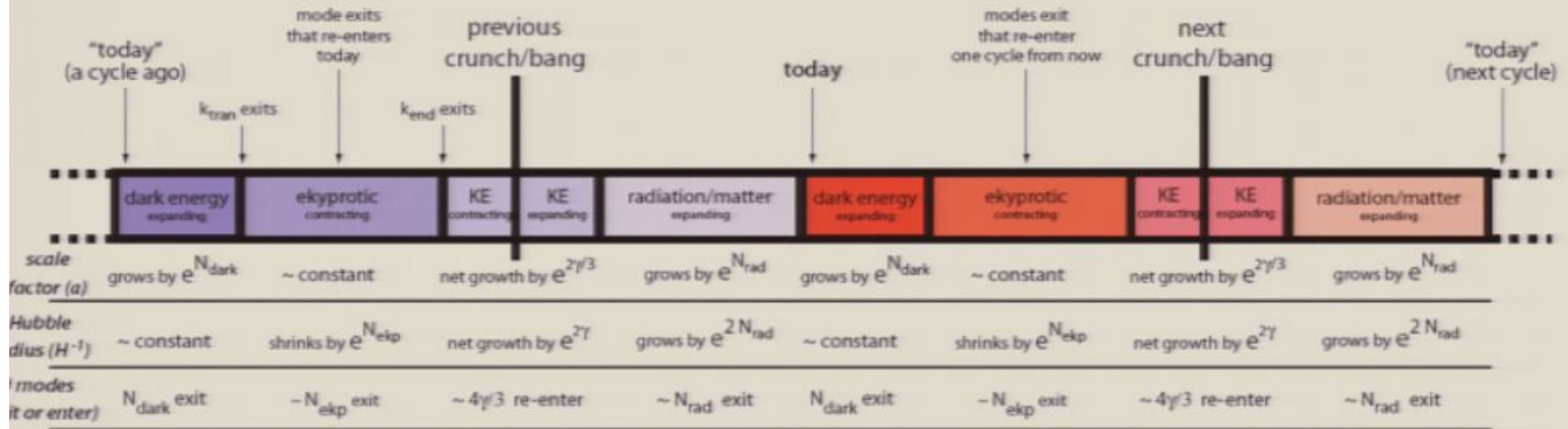
SCALES I



$$\gamma = \log \frac{(-V_{\text{end}})^{1/4}}{T_{\text{reheat}}}, \quad N_{\text{ekp}} = \log \sqrt{\frac{-V_{\text{end}}}{V_0}}, \quad N_{\text{rad}} = \log \frac{T_{\text{rh}}}{T_0}$$

$$N_{\text{ekp}} = 2(\gamma + N_{\text{rad}})$$

SCALES I



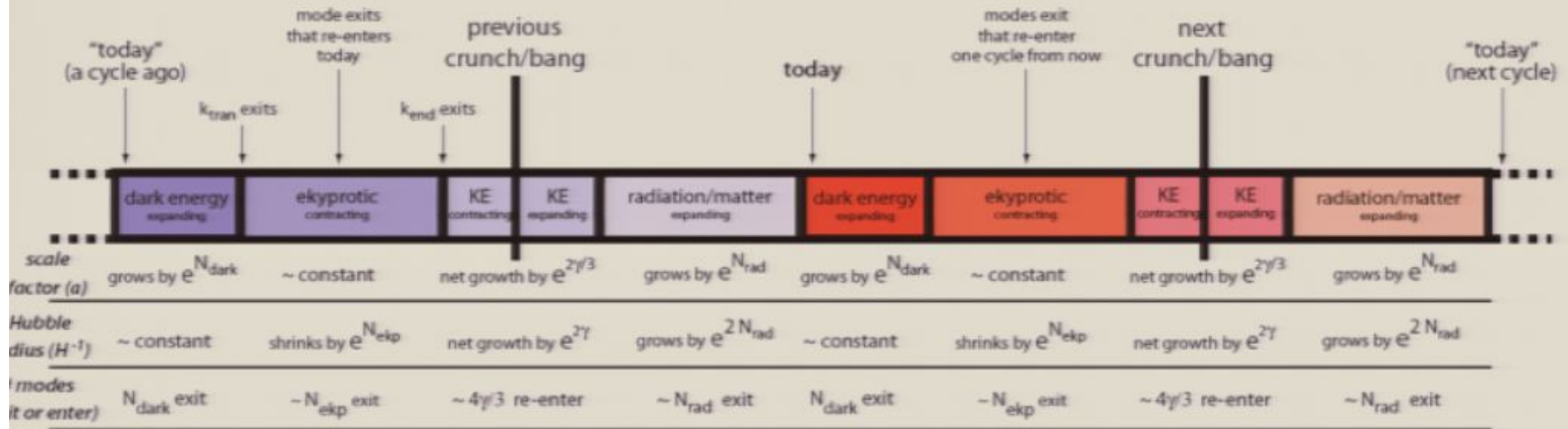
$$\gamma = \log \frac{(-V_{\text{end}})^{1/4}}{T_{\text{reheat}}}, \quad N_{\text{ekp}} = \log \sqrt{\frac{-V_{\text{end}}}{V_0}}, \quad N_{\text{rad}} = \log \frac{T_{\text{rh}}}{T_0}$$

$$N_{\text{ekp}} = 2(\gamma + N_{\text{rad}})$$

$$H \sim \sqrt{Y}$$

dn caist

SCALES I



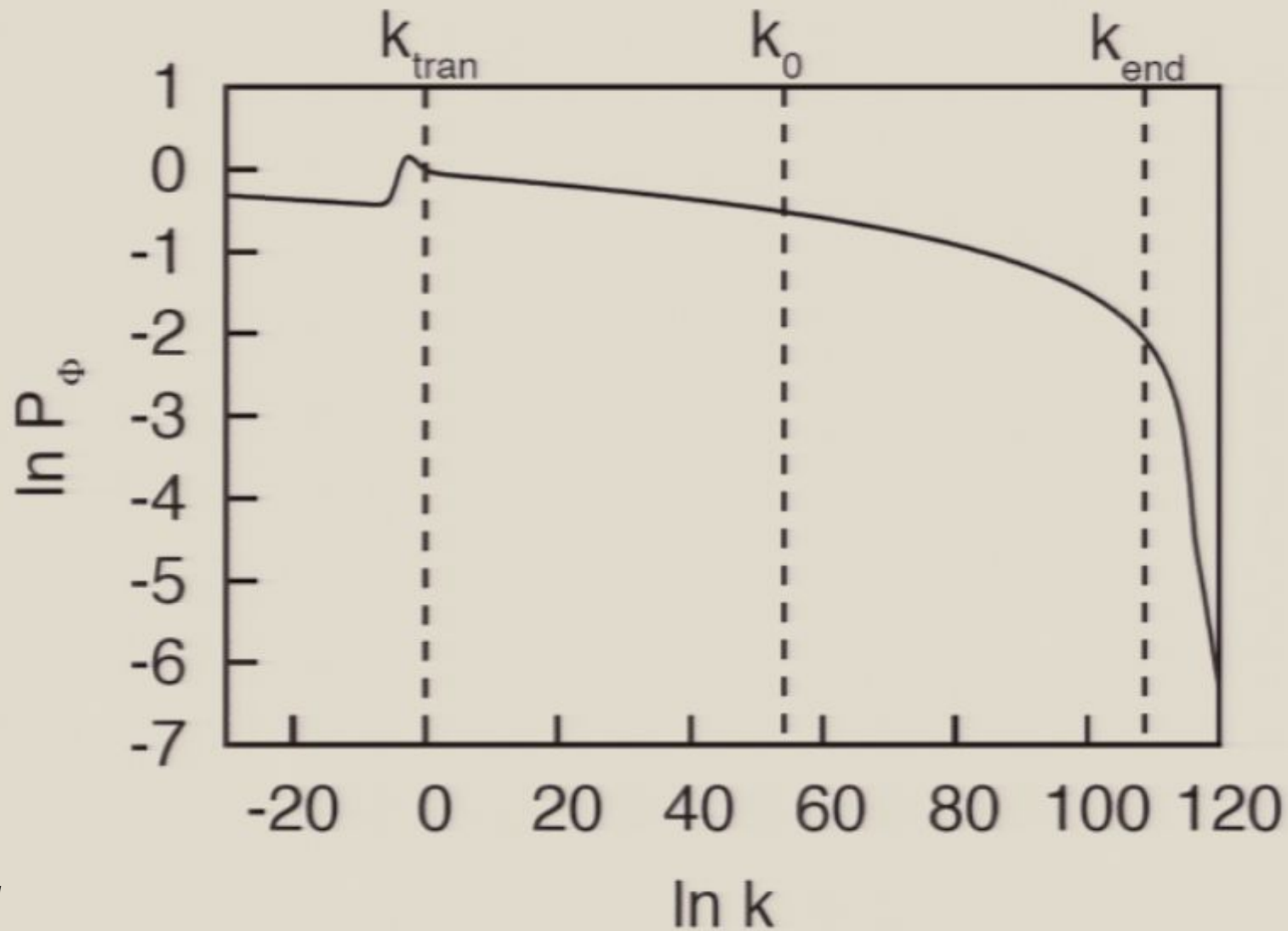
$$\gamma = \log \frac{(-V_{\text{end}})^{1/4}}{T_{\text{reheat}}}, \quad N_{\text{ekp}} = \log \sqrt{\frac{-V_{\text{end}}}{V_0}}, \quad N_{\text{rad}} = \log \frac{T_{\text{rh}}}{T_0}$$

$$N_{\text{ekp}} = 2(\gamma + N_{\text{rad}})$$

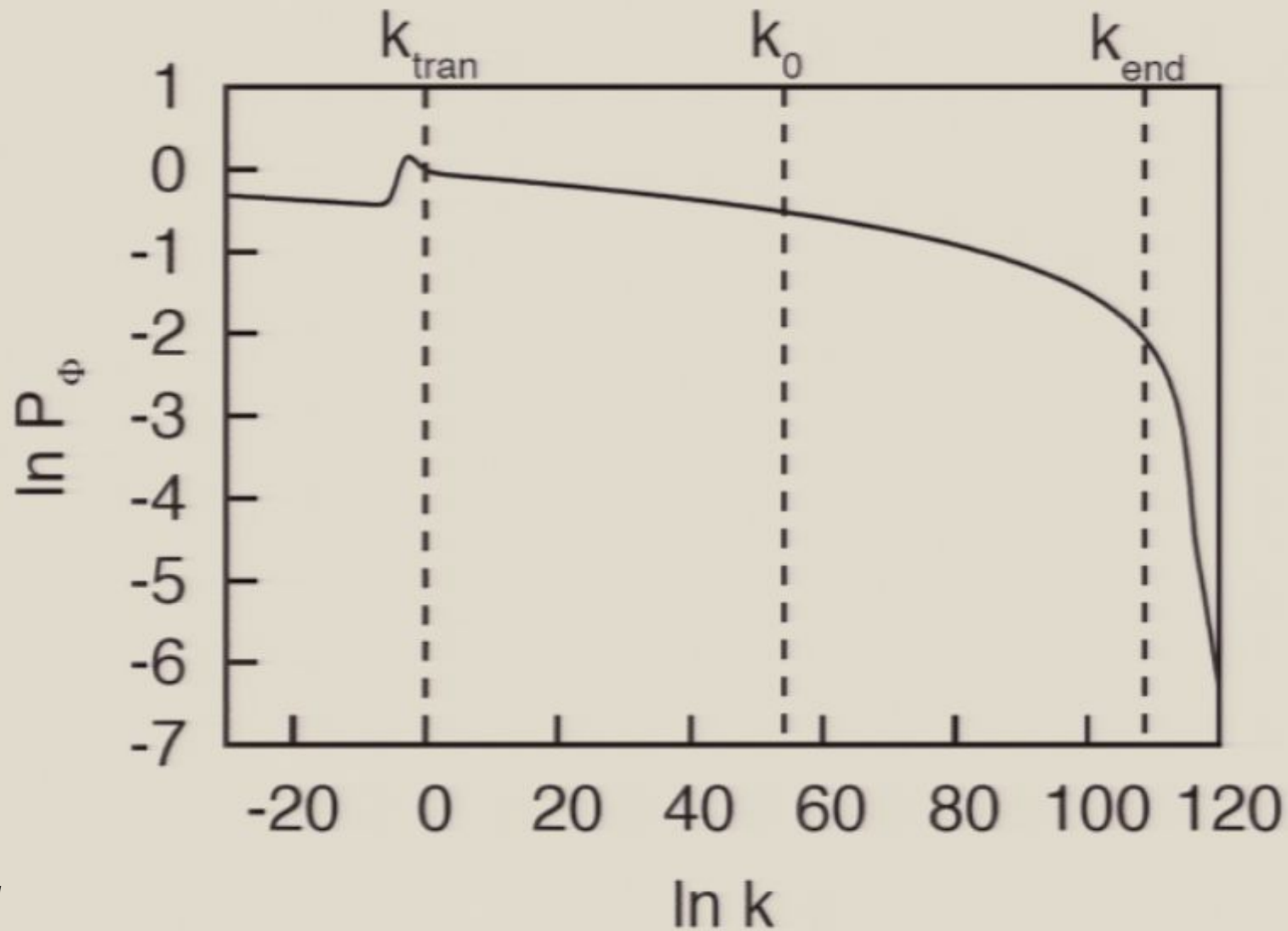
SCALES II

Length scale	Size relative to today's horizon
Today's horizon	1
Current wavelength of the modes that: ...were on the horizon one cycle ago	$e^{N_{\text{dark}} + 2\gamma/3 + N_{\text{rad}}}$
...will be on the horizon one cycle from now	$e^{-N_{\text{dark}} - 2\gamma/3 - N_{\text{rad}}}$
...were the first ones to go ultralocal during the ekpyrotic phase one cycle ago	$e^{2\gamma/3 + N_{\text{rad}}}$
...were the last ones to go ultralocal during the ekpyrotic phase one cycle ago	$e^{-N_{\text{ekp}} + 2\gamma/3 + N_{\text{rad}}}$
...will be the first ones to go ultralocal during the coming ekpyrotic phase of this cycle	$e^{-N_{\text{dark}}}$
...will be the last ones to go ultralocal during the coming ekpyrotic phase of this cycle	$e^{-N_{\text{dark}} - N_{\text{ekp}}}$

POWER SPECTRUM



POWER SPECTRUM

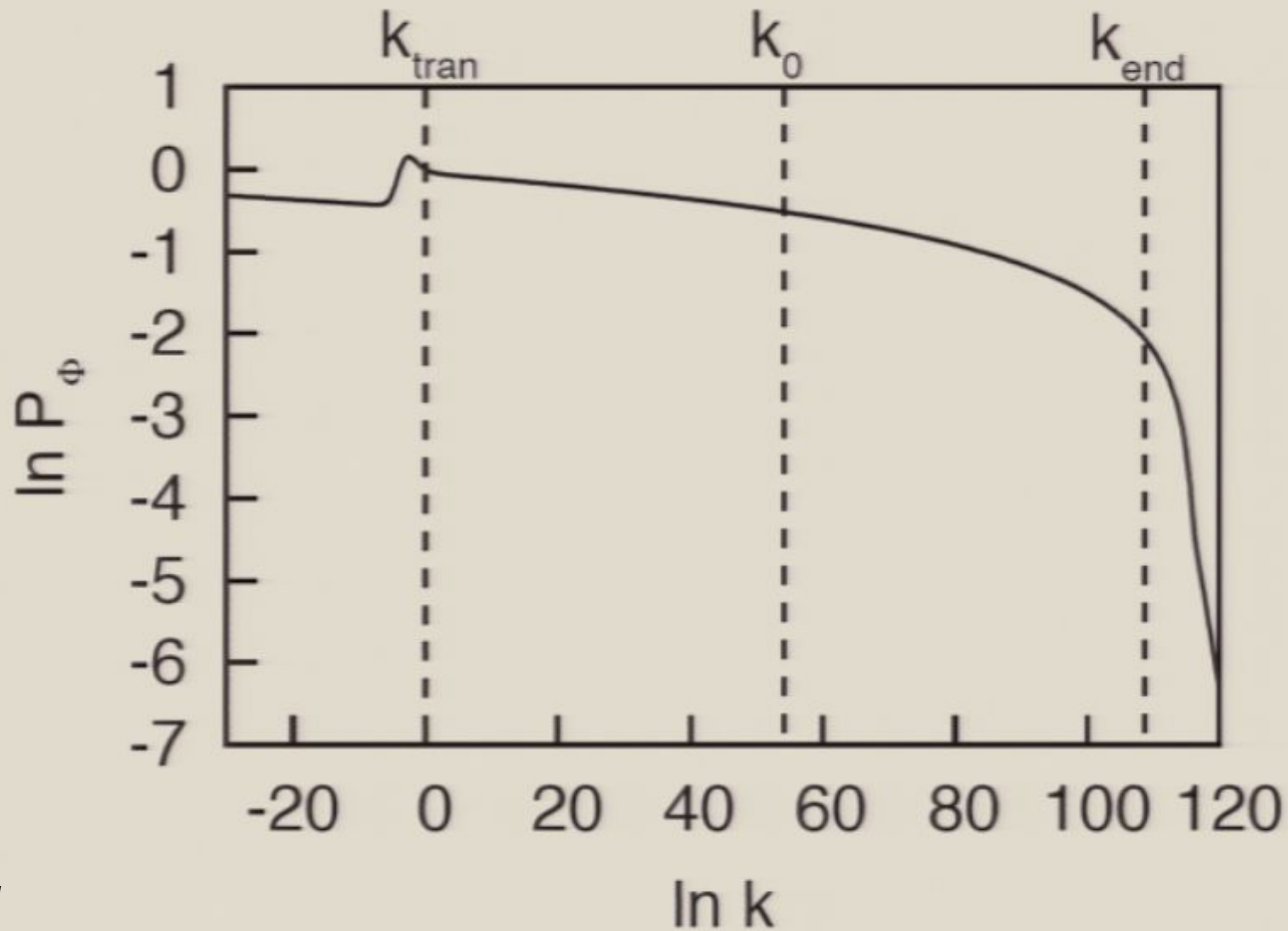


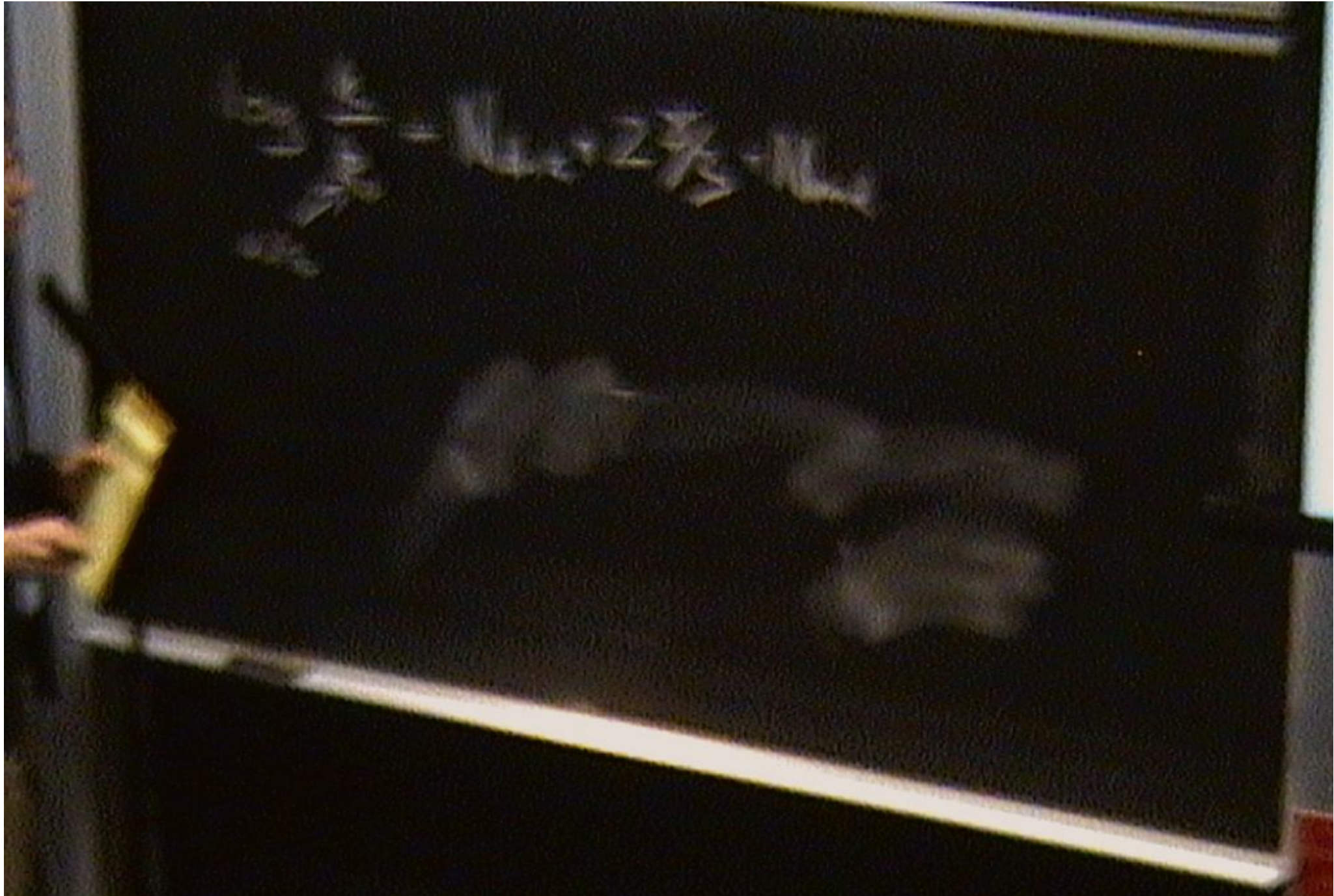
$$H \sim \sqrt{Y}$$

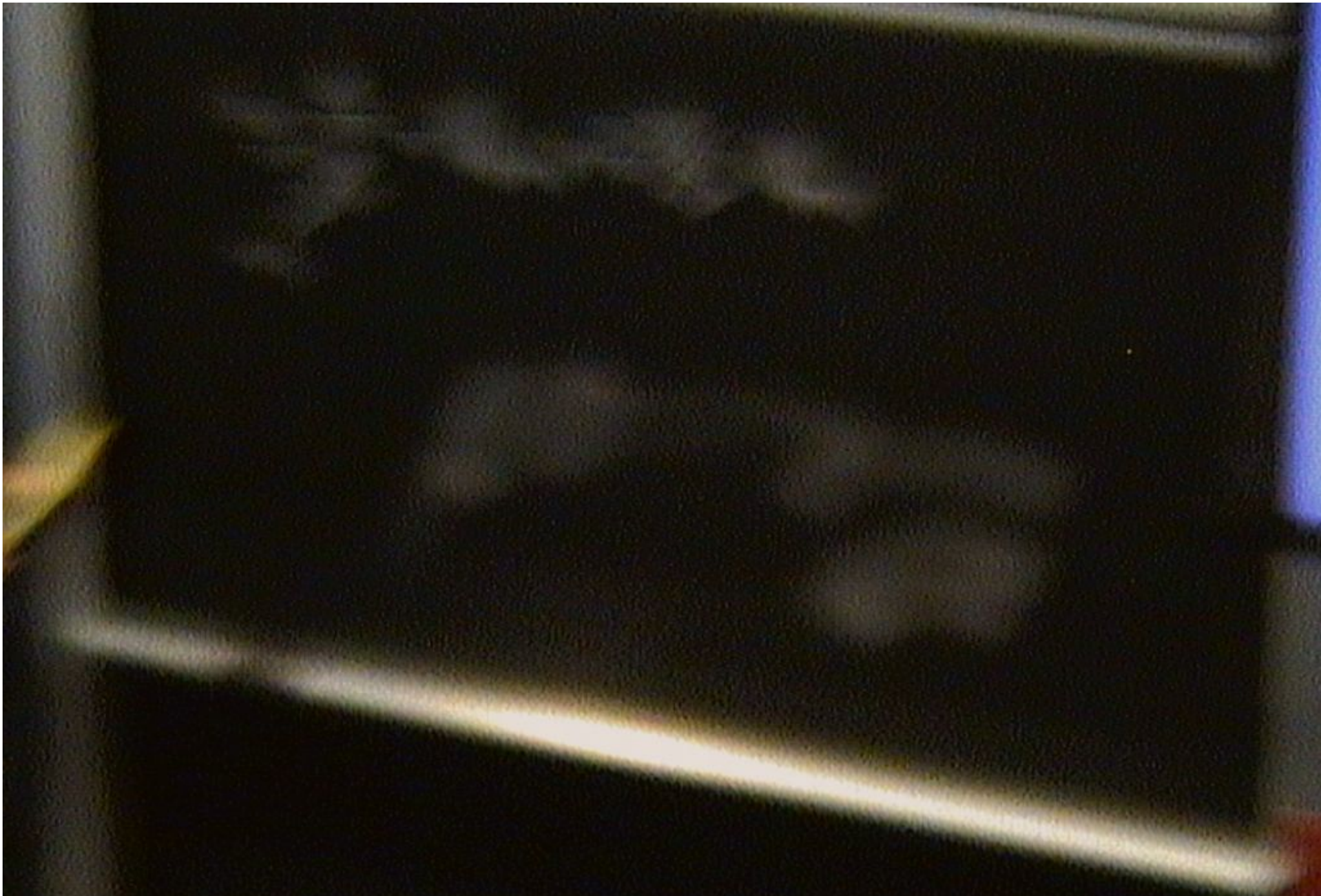
du raist

$$10^{55}$$

POWER SPECTRUM







$$\log \frac{k}{k_0} = N_{\text{db}} + 2\frac{\sigma}{3} + N_{\text{rad}}$$

\rightarrow
 $\sigma \cdot k_0$

$$\log \frac{k}{k_0} = N_{\text{drift}} + 2\frac{\gamma}{3} + N_{\text{rad}}$$

\rightarrow
 $\sigma \cdot t_0$

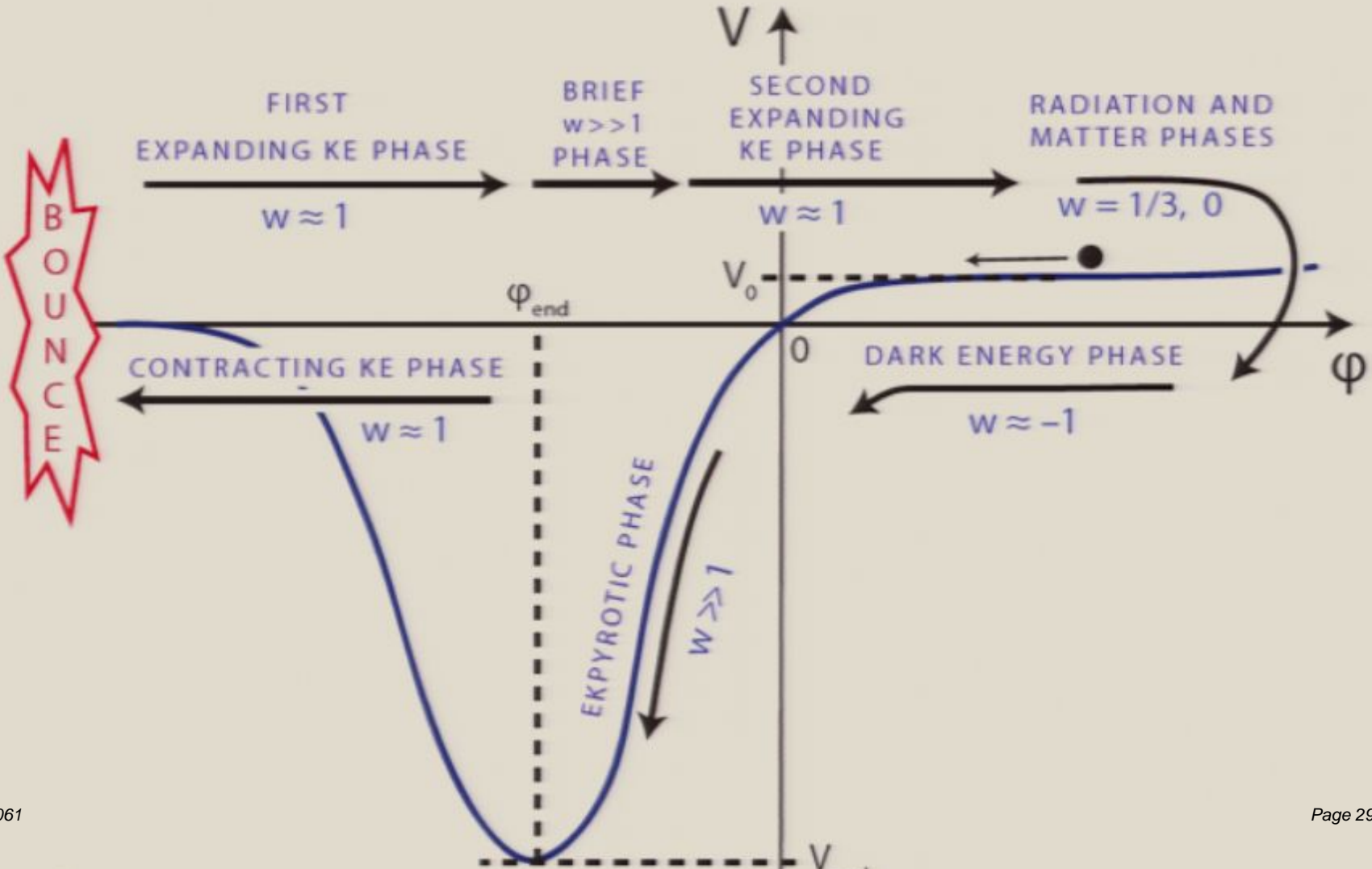
$$\log \frac{k}{k_0} = N_{\text{rad}} + 4\frac{\gamma}{3}$$

$$\log \frac{k}{k_0} = N_{\text{drift}} + 2\frac{\gamma}{3} + N_{\text{rad}}$$

\nearrow
at t_0

$$\log \frac{k}{k_0} = N_{\text{rad}} + 4\frac{\gamma}{3}$$

THE POTENTIAL



$$\log \frac{k}{k_0} = N_{\text{dark}} + 2\gamma/3 + N_{\text{rad}}$$

↖

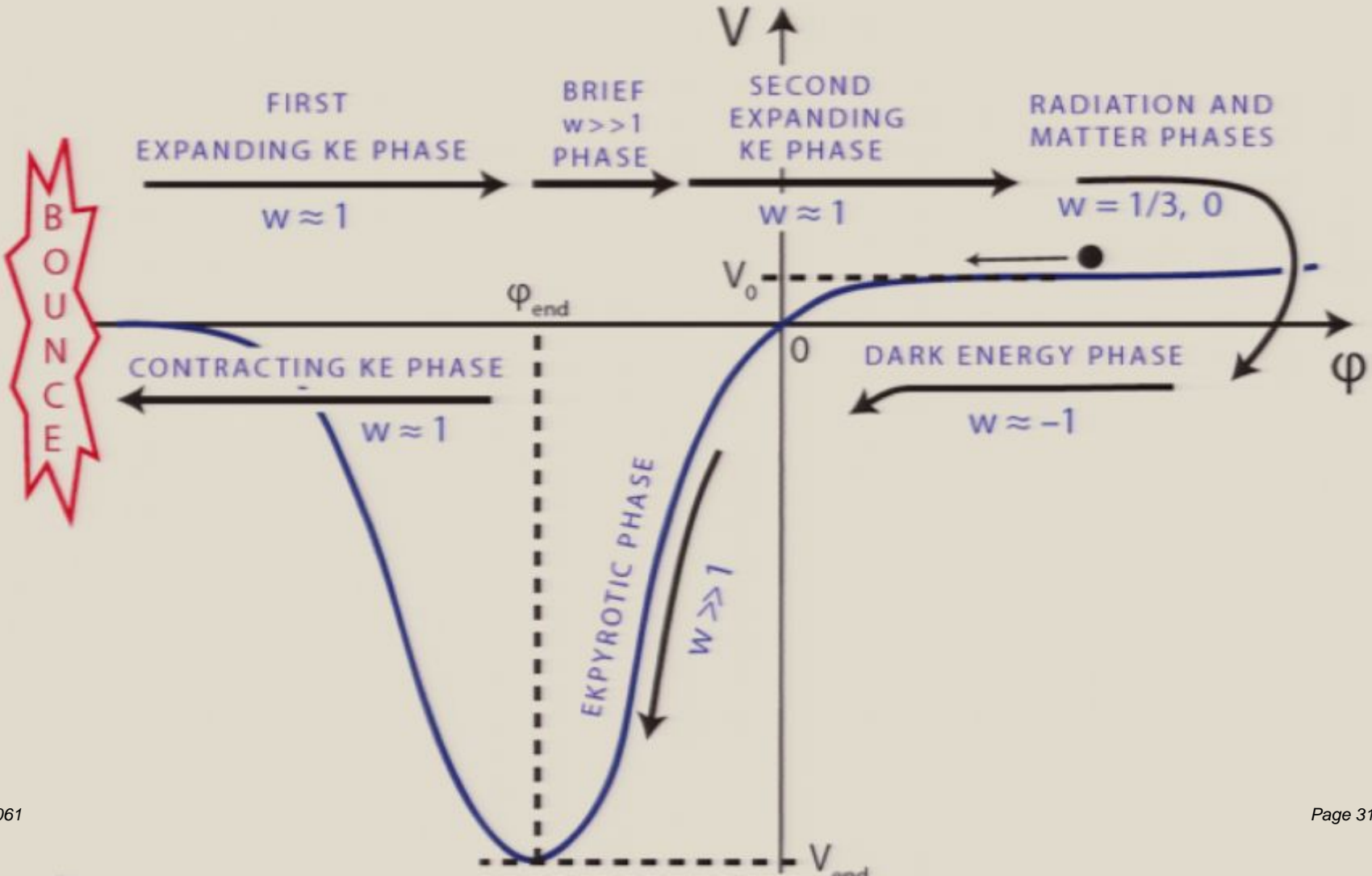
and to

$$\log \frac{k}{k_0} = N_{\text{rad}} + 4\gamma/3$$

$$\frac{1}{k^2} \sim \frac{1}{k^4}$$

$$10^{-5} e^{(-N_{\text{dark}} - 2\gamma/3 + N_{\text{rad}})4}$$

THE POTENTIAL



$$\log \frac{k}{k_0} = N_{\text{dark}} + 2\gamma/3 + N_{\text{rad}}$$

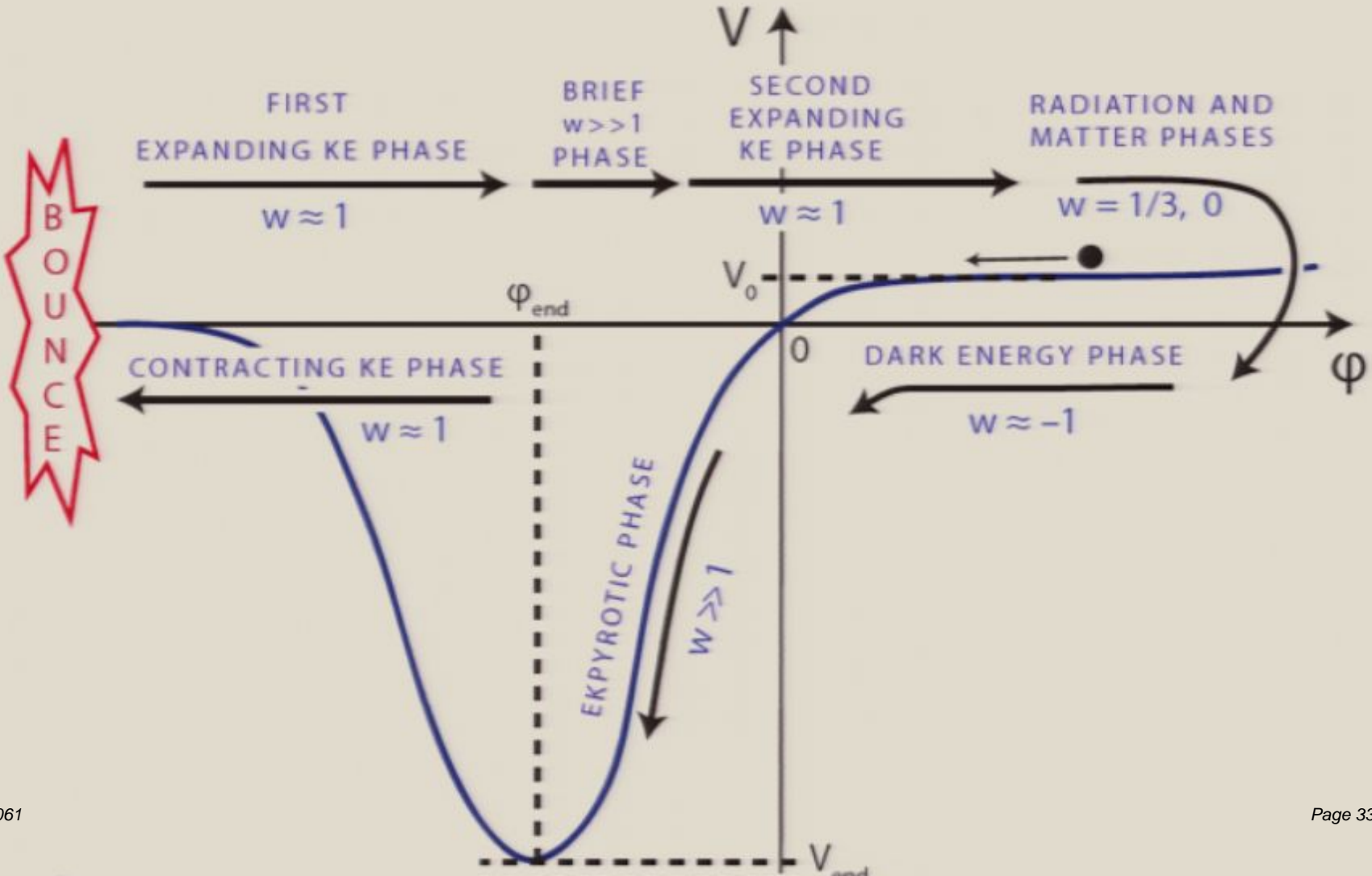
↖

$$\log \frac{k}{k_0} = N_{\text{rad}} + 4\gamma/3$$

$$\frac{1}{k^2} \sim \frac{1}{k^4}$$

$$10^{-5} e^{(-N_{\text{dark}} - 2\gamma/3 + N_{\text{rad}})4} < H_0$$

THE POTENTIAL



Flatness

$\rho_K \approx$



Flatness

$$\rho_K \sim \frac{1}{a^4}$$

$$\Omega_K \sim \frac{1}{(aH)^2}$$



Flatness

$$\rho_k \sim \frac{1}{a^2}$$

$$\Omega_k \sim \frac{1}{(aH)^2}$$

$$e_k: e^{-2N_{chp}}$$

Flatness

$$\rho_k \sim \frac{1}{a^2}$$

$$\Omega_k \sim \frac{1}{(aH)^2}$$

$$N_{\text{exp}} = 2(N_{\text{rad}} + \gamma)$$

$$e_k: e^{-2N_{\text{chp}}}$$

$$e^{-2N_{\text{rad}} + 8\gamma/3}$$

$$\Omega \propto e^{-N_{\text{chp}} + 2\gamma/3}$$

Flatness

$$\rho_k \sim \frac{1}{a^2}$$

$$\Omega_k \sim \frac{1}{(aH)^2}$$

$$N_{\text{exp}} = 2(N_{\text{rad}} + \gamma)$$

$$e_k: e^{-2N_{\text{chp}}}$$

$$= e^{2N_{\text{rad}} - 8\gamma/3}$$

$$\Omega \propto e^{-N_{\text{chp}} + 2\gamma/3} = e^{-2N_{\text{rad}} - 4\gamma/3}$$

Flatness

$$\rho_k \sim \frac{1}{a^2}$$

$$\Omega_k \sim \frac{1}{(aH)^2}$$

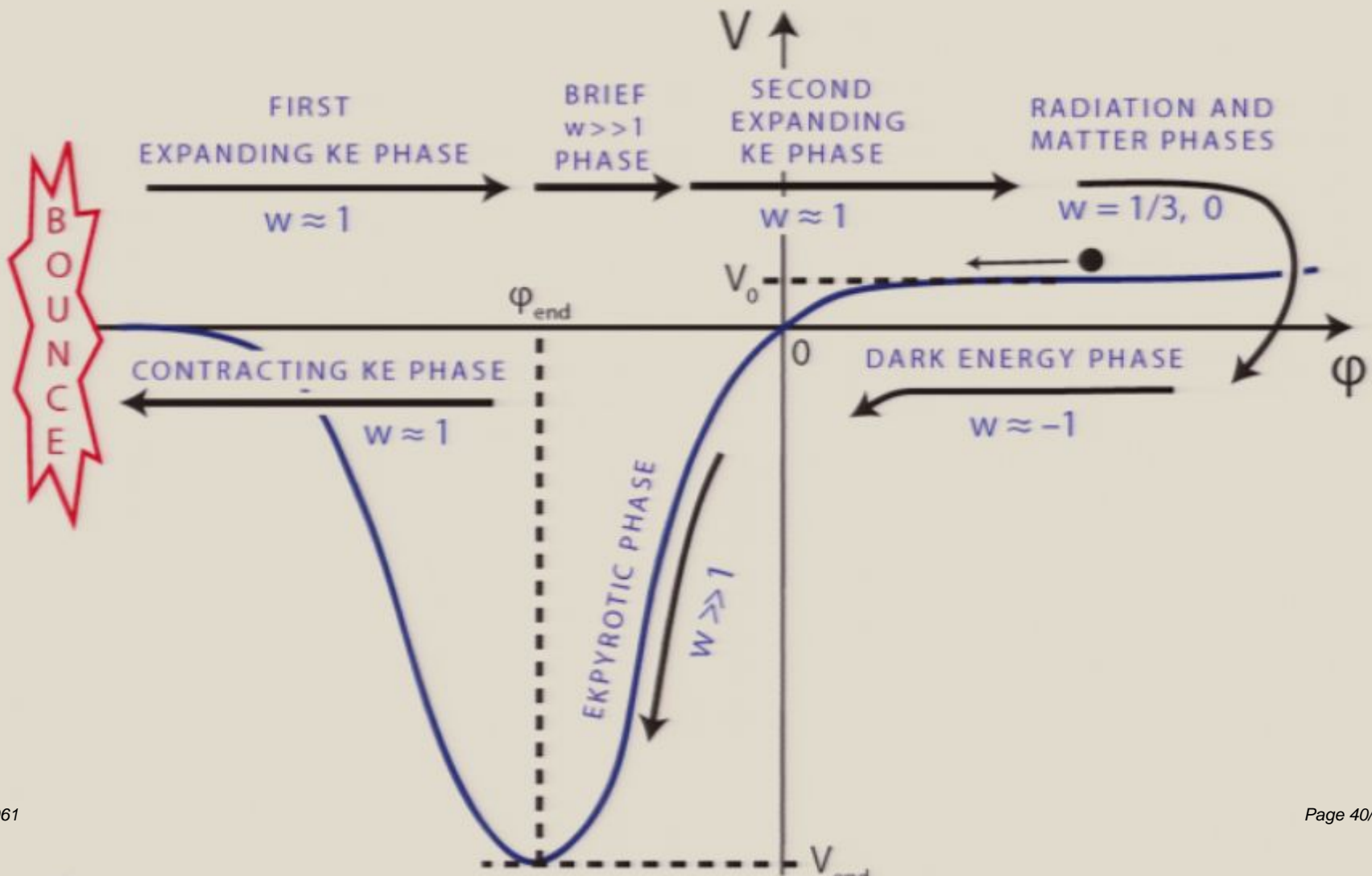
$$N_{\text{exp}} = 2(N_{\text{rad}} + \gamma)$$

$$e_k: e^{-2N_{\text{chp}}}$$

$$e^{2N_{\text{rad}} + 8\gamma/3}$$

$$\Omega \propto e^{-N_{\text{chp}} + 2\gamma/3}$$
$$= e^{-2N_{\text{rad}} - 4\gamma/3}$$

THE POTENTIAL



Flatness

$$\rho_K \sim \frac{1}{a^2}$$

$$\Omega_K \sim \frac{1}{(aH)^2}$$

$$\rho_\sigma \sim \frac{1}{a^6}$$

$$N_{\text{exp}} = 2(N_{\text{rad}} - \gamma)$$

$$e_k = e^{-2N_{\text{chp}}} \cdot e^{2N_{\text{rad}} - 8\gamma/3}$$

$$\Omega \propto e^{-N_{\text{chp}} + 2\gamma/3} = e^{-2N_{\text{rad}} - 4\gamma/3}$$

NO HAIR THEOREM

$$3\left(\frac{\dot{a}}{a}\right)^2 = \Lambda + \frac{\kappa}{a^2} + \frac{\rho_m}{a^3} + \frac{\rho_r}{a^4} + \frac{\sigma}{a^6} + \frac{\rho_\phi}{a^{-3(1+w)}}$$

cosmological constant, inflaton

dust, monopoles

anisotropy

curvature, gradients, domain walls

radiation

scalar field

In expansion, smallest w dominates (inflation)

In contraction, largest w dominates (ekpyrosis)

hep-th/0312009

CONCLUSIONS

- I. Large-scale structure in the cyclic universe is generated by quantum fluctuations from the previous cycle, without interference from older structure.
- II. The ekpyrotic phase is sufficient to make the universe homogeneous, isotropic and flat. Dark energy is superfluous.
- III. The global structure of the cyclic universe cannot be described by a uniform FRW picture.