

Title: Decoding the fundamental laws of gravity

Date: Nov 11, 2006 03:00 PM

URL: <http://pirsa.org/06110055>

Abstract:

Spin-2



Spin-2

$$\underline{h_{\mu\nu}} \quad \nabla^\kappa \bar{T}_{\mu\nu} = 0, \quad h \equiv h_\mu^\mu$$

$$(\epsilon h_{\mu\nu} + m^2 \sigma) (h_{\mu\nu} - \eta_{\mu\nu} h) = T_{\mu\nu}$$



Spin-2  $\underline{h}_{\mu\nu} \quad \nabla^\kappa \bar{T}_{\mu\nu} = 0, h \equiv h_\mu^\mu$

$$(\epsilon h)_{\mu\nu} + k^2(\sigma)(h_{\mu\nu} - \eta_{\mu\nu}h) = T_{\mu\nu}$$

$$\Box h = \Box h_{\mu\nu} - \eta_{\mu\nu} \Box h - 2\partial_\mu \partial_\nu h + \eta_{\mu\nu} (\partial_\mu \partial_\nu h) - 2\lambda h$$



Spin-2  $\underline{h_{\mu\nu}}$   $\nabla^\lambda \underline{T_{\mu\nu}} = 0, h \equiv h^\mu_\mu$

$$(\mathcal{E} h)_{\mu\nu} + \text{Im}(\sigma) (\underline{h_{\mu\nu}} - \eta_{\mu\nu} h) = T_{\mu\nu}$$

BdL-Fres.

$$T_{\mu\nu} = \square h_{\mu\nu} - \eta_{\mu\nu} \square h - 2\partial_\mu \partial_\nu h_{\mu\nu} + \eta_{\mu\nu} \partial_\mu \partial_\nu h + 2\partial_\mu \partial_\nu h_{\mu\nu}$$

Spin-2

$$\underline{h_{\mu\nu}}$$

$$\nabla^\mu \underline{T_{\mu\nu}} = 0$$

$$h = h^{\mu\nu}$$

$$\underline{\zeta_c} = H_0^{-1}$$

$$(Eh)_{\mu\nu} + k^2(\sigma)(h_{\mu\nu} - \eta_{\mu\nu}h) = T_{\mu\nu}$$

Bulk-Fried.

$$E_{\text{action}} = D h_{\mu\nu} - \eta_{\mu\nu} D h - 2\lambda \delta^{ab} h_{ab} + \eta_{\mu\nu} \partial_\mu \partial_\nu h$$

Spin-2

$$h_{\mu\nu} \quad \nabla^\kappa T_{\mu\nu} = 0, \quad h \equiv h^{\mu\nu}$$

$$\underbrace{(\mathcal{E} h)_{\mu\nu} + m^2(\sigma) (h_{\mu\nu} - \eta_{\mu\nu} h)}_{\text{Ricci - Fuchs}} = T_{\mu\nu} \quad \boxed{\gamma_c = H_0'}$$

$$E_{\text{action}} = D h_{\mu\nu} - h_{\mu\nu} D h - 2 \partial_\mu \partial_\nu h + \eta_{\mu\nu} (\partial_\mu \partial_\nu h)^2 + 2 \partial_\mu \partial_\nu h$$

$$\lim_{D \rightarrow 0} m^2(\sigma) = \sqrt{c}^{2(k-1)} \Delta^k$$

Spin-2  $\underline{h_{\mu\nu}}$   $\nabla^\mu \bar{T}_{\mu\nu} = 0, h \equiv h^{\mu\nu}$

$$(\epsilon h_{\mu\nu} + m^2(\sigma) (\underline{h_{\mu\nu}} - \eta_{\mu\nu} h)) = T_{\mu\nu}$$

$\boxed{\epsilon_c = H_0'}$

Ricci - Frans.

$$Einstein = \square h_{\mu\nu} - \eta_{\mu\nu} \square h - 2\lambda \square^2 h_{\mu\nu} + 2\lambda \square^2 h + 2\lambda h$$

$$\frac{m^2(0)}{l \rightarrow 0} = \sqrt{c^{2(k-1)}} \quad \square^k$$

$$\alpha = 0 \\ \alpha = \frac{1}{2}$$

Spin-2

$$h_{\mu\nu}$$

$$\nabla^\mu \tilde{T}_{\mu\nu} = 0$$

$$h = h^{\mu}_{\mu}$$

$$\gamma_c = H_0^{-1}$$

$$(Eh_{\mu\nu} + m^2(\sigma) \underbrace{(h_{\mu\nu} - \eta_{\mu\nu} h)}_{\text{Ricci-Fres.}}) = T_{\mu\nu}$$

$$E_{\text{action}} = \square h_{\mu\nu} - \eta_{\mu\nu} \square h - 2\partial_\mu \partial_\nu h + 2\partial_\mu \partial_\nu h_{\mu\nu} + 2\lambda h$$

$$\frac{m^2(\sigma)}{D \rightarrow 0} = \sqrt{c^{2(k-1)}} \quad \square^k$$

$$0 < \alpha < 1$$

$$h_{\mu\nu} = \int d^4x \delta^{(4)}(x) h_{\mu\nu}^{(4)}$$



$$\rho_{\mu\nu} = \int d^4x S^{\mu\nu} h_{\mu\nu}$$

$$\frac{1}{D + \omega(0)} = \frac{g(r)}{r^2}$$

$$h_{\mu\nu} = \int_0^\infty d\mu s(\mu) h_{\mu\nu}^{(n)}$$

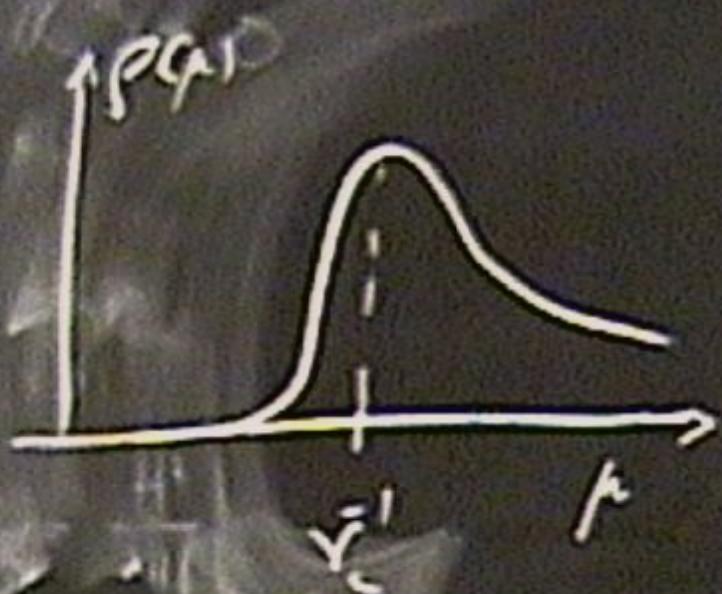
$$\frac{1}{D+\mu^2(0)} = \int_0^\infty \frac{d\mu s(\mu)}{D+\mu^2}$$

$$h_{\mu\nu} = \int_0^\infty d\mu \, S(\mu) h_{\mu\nu}^{(0)}$$

$$\frac{1}{D + \mu^2(0)} = \int_0^\infty \frac{d\mu \, S(\mu)}{D + \mu^2}$$

$$h_{\mu\nu} = \int_0^\infty d\mu \, g(\mu) h_{\mu\nu}$$

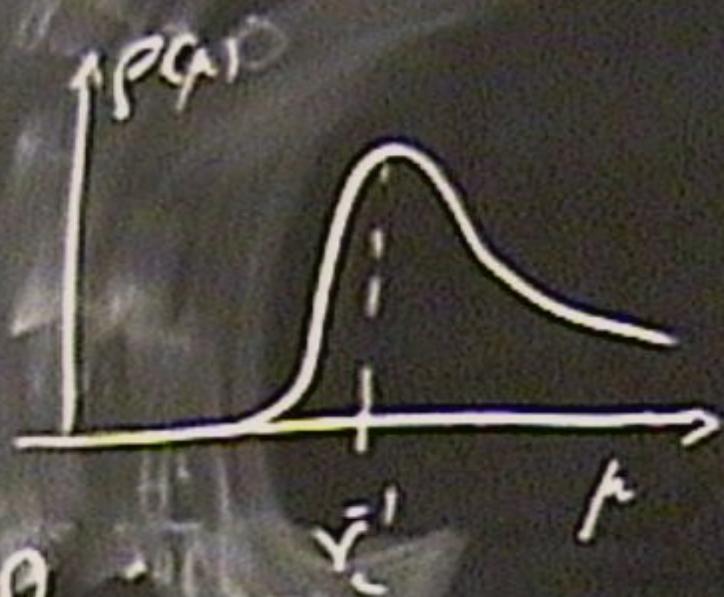
$$\frac{1}{D + \mu^2(0)} = \int_0^\infty \frac{d\mu \, g(\mu)}{D + \mu^2}$$



$$h_{\mu\nu} = \int_0^\infty d\mu \, g(\mu) h_{\mu\nu}^{(0)}$$

$$\frac{1}{D + \mu^2} = \int_0^\infty \frac{g(\mu)}{D + \mu^2}$$

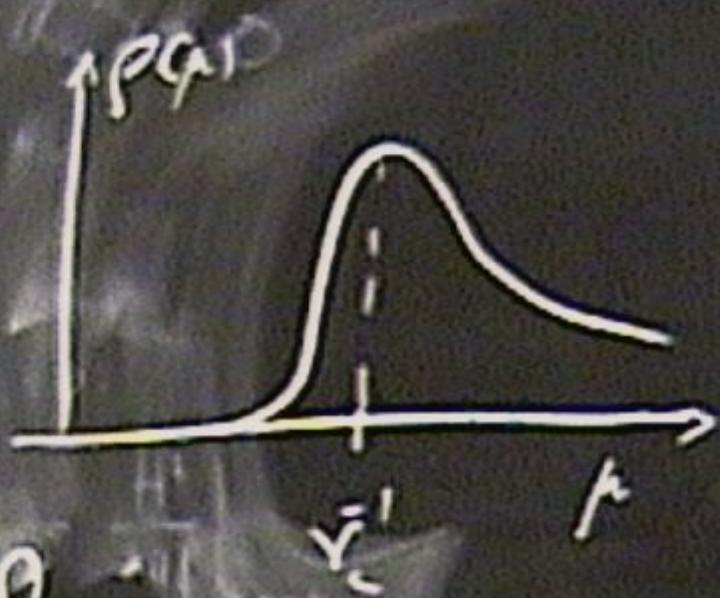
$$g(\mu) \geq 0 \Rightarrow \alpha > 0$$



$$h_{\mu\nu} = \int_0^\infty d\mu S(\mu) h_{\mu\nu}^{(n)}$$

$$\frac{1}{D + \mu^2} = \int_0^\infty \frac{d\mu S(\mu)}{D + \mu^2}$$

$$S(\mu) \geq 0 \Rightarrow \mu > 0$$



$$\text{Spin-2} \quad \underline{h_{\mu\nu}} \quad \underline{\nabla^\mu \Gamma_{\mu\nu} = 0, h = h_\mu^{\mu\nu}} \quad \underline{\gamma_c = H_0^{-1}} \\ (\epsilon h)_{\mu\nu} + m^{(0)}(h_{\mu\nu} - \eta_{\mu\nu} h) = T_{\mu\nu} \quad \sqrt{0}$$

Ricci-Fres.

$$E_{\text{action}} = D h_{\mu\nu} - \eta_{\mu\nu} D h - \frac{1}{2} \epsilon D^2 h_{\mu\nu} + g^{(0)} R_{\mu\nu} - \frac{1}{2} \lambda h$$

$$\frac{m^{(0)}}{D \rightarrow 0} = \sqrt{c^{2(k-1)}} \quad \square$$

$$D \leq \rho < 1$$

$$d = \frac{1}{2}$$

$$f(w) \geq 0 \iff \alpha > 0$$

DVZ-distribution



$$f(w) \geq 0 \Leftrightarrow \alpha > 0$$

### DVZ-discontinuity

$$\mu^2 = 0 \rightarrow 2 \\ \mu^2(\text{off} \neq 0) \rightarrow 5 = 3 + 2\lambda$$

$$\begin{aligned} & \text{on R} & & \text{on L} \\ & \eta = 0 & & \eta = 0 \\ & 1 - \frac{1}{2}\eta & & \eta < 0 \\ & 1 - \frac{1}{2}\eta & & \eta < 0 \end{aligned}$$

$$F(3-\eta) \\ 14 - 1 = 3 \rightarrow \eta = 5$$

$$f(u) \geq 0 \Leftrightarrow \alpha > 0$$

DVZ - discontinuity 1-2:

$$\begin{aligned} k^2 < 0 &\rightarrow 2 \\ k^2 (k \neq 0) &\rightarrow 5 = 3 + 2 - 2 + 2 + 1 \end{aligned}$$

$$f(w) \geq 0 \iff \alpha > 0$$

DVZ-dicretinuity

$$\begin{aligned} k=0 &\rightarrow 2 \\ k \neq 0 &\rightarrow 5 = 3 + 2 = 2 + 2 + 1 \end{aligned}$$

Spin-2

$h_{\mu\nu}$

$\nabla^\lambda T_{\mu\nu} = 0$

$h = h^{\mu\nu}_{\mu\nu}$

$\gamma_c = H_0^{-1}$

$$(\epsilon h)_{\mu\nu} + M^{(0)}(h_{\mu\nu} - \eta_{\mu\nu} h) = T_{\mu\nu}$$

$B_{\mu\nu} - F_{\mu\nu}$

$\sqrt{0}$

$$E_{action} = D h_{\mu\nu} - h_{\mu\nu} D h - 2\alpha D^2 h_{\mu\nu} + 2(D)^\mu \rho_{\mu\nu} + 2\omega h$$

$$\begin{aligned} M^{(0)} &= \sqrt{c^{2(k-1)}} \\ D &\rightarrow 0 \end{aligned}$$

$$D < q < 1$$

$$\omega = \frac{1}{2}$$

$$f(n) > 0 \Leftrightarrow \alpha > 0$$

$\sqrt{DVZ}$  - discontinuity

$k=0 \rightarrow 2$

$k^2 \neq 0 \rightarrow 5 = 3 + 2 \neq 2 + 2 + 1$

$$T_{00} =$$

$$1 = 3 - 1$$

$$14 - 1 = 3 - 2\gamma - 5$$

$$f(n) > 0 \Rightarrow \alpha > 0$$

$\sqrt{D}V^2$  - discontinuity  $\begin{matrix} h-2 \\ \downarrow \\ h-1 \\ \downarrow \\ h-0 \end{matrix}$   $\left| \begin{matrix} \sqrt{g} = 2GM \\ \eta < 0 \end{matrix} \right.$

$m=0 \rightarrow 2$

$m^2(p) \neq 0 \rightarrow 5 = 3+2 = 2+2+1$

$T_{00} = M\delta(r)$

$f_{\mu\nu} = \frac{\delta_\mu^\circ \delta_\nu^\circ - \frac{1}{3}(\eta_{\mu\nu} + \frac{2m^2}{r})}{\Box + m^2(p)} \quad \boxed{M\delta(r)}$

$$\beta(n) > 0 \iff \alpha > 0$$

$\sqrt{D}V^2$  - discontinuity  $\left. \begin{matrix} L_2 \\ L_1 \\ L_0 \end{matrix} \right| \quad \left. \begin{matrix} \sqrt{g} = 2GM \end{matrix} \right|$

$$k^2 = 0 \rightarrow 2$$

$$k^2(\rho \neq 0) \rightarrow 5 = 3 + 2 \quad \left. \begin{matrix} 2+2+1 \\ \vdots \end{matrix} \right| \quad \left. \begin{matrix} 2 \\ 1 \\ 1 \end{matrix} \right| \quad \left. \begin{matrix} 0 \\ 0 \\ 0 \end{matrix} \right| = 0$$

$$T_{00} = M \delta(r) \quad \left. \begin{matrix} 1 \\ 1 \\ 1 \end{matrix} \right| \quad \left. \begin{matrix} 0 \\ 0 \\ 0 \end{matrix} \right| \quad \left. \begin{matrix} 0 \\ 0 \\ 0 \end{matrix} \right|$$

$$f_{\mu\nu} = \frac{\delta_\mu^\circ \delta_\nu^\circ - \frac{1}{3}(\eta_{\mu\nu} + \frac{2m}{m^2(0)})}{D + m^2(0)} \quad \left. \begin{matrix} 3-1 \\ 3-1 \\ 3-1 \end{matrix} \right|$$

$$T^{\mu\nu} \rho_{\mu\nu} = (T_{00} - \frac{1}{3}T) \left[ \frac{Yg}{r} + O(m^2) \right] \rightarrow 0$$

$\downarrow$

$$(T_{00} - \frac{1}{2}T) \frac{Yg}{r}$$
$$\frac{1}{2} = \frac{1}{3}$$



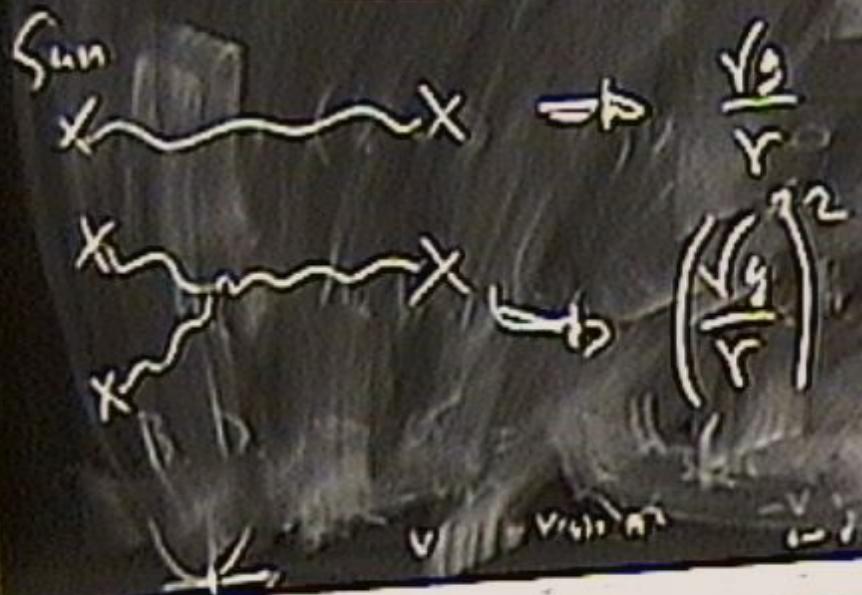
$$T^{\mu\nu}g_{\mu\nu} = \left(T_{00} - \frac{1}{3}T\right) \left[ \frac{G}{r} + O(m^2) \right] \quad n=10$$
$$\left(T_{00} - \frac{1}{2}T\right) \frac{G}{r}$$
$$\frac{1}{2} \neq \frac{1}{3}$$

Sun  
X ~ X



$$T^{\mu\nu} \delta_{\mu\nu} = \left( T_{00} - \frac{1}{3} T \right) \int \frac{\sqrt{g}}{r} + \mathcal{O}(m^2) \xrightarrow{n \rightarrow 0}$$

$$\left( T_{00} - \frac{1}{2} T \right) \frac{\sqrt{g}}{r} \quad \frac{1}{2} \neq \frac{1}{3}$$



$$T^{\mu\nu}g_{\mu\nu} = \left(T_{00} - \frac{1}{3}T\right) \left[ \frac{\sqrt{g}}{r} + \mathcal{O}(m^2) \right] + \frac{1}{r} \left( T_{00} - \frac{1}{2}T \right)$$

$m^2 \rightarrow 0$

$$\frac{1}{2} = \frac{1}{3}$$

Sun

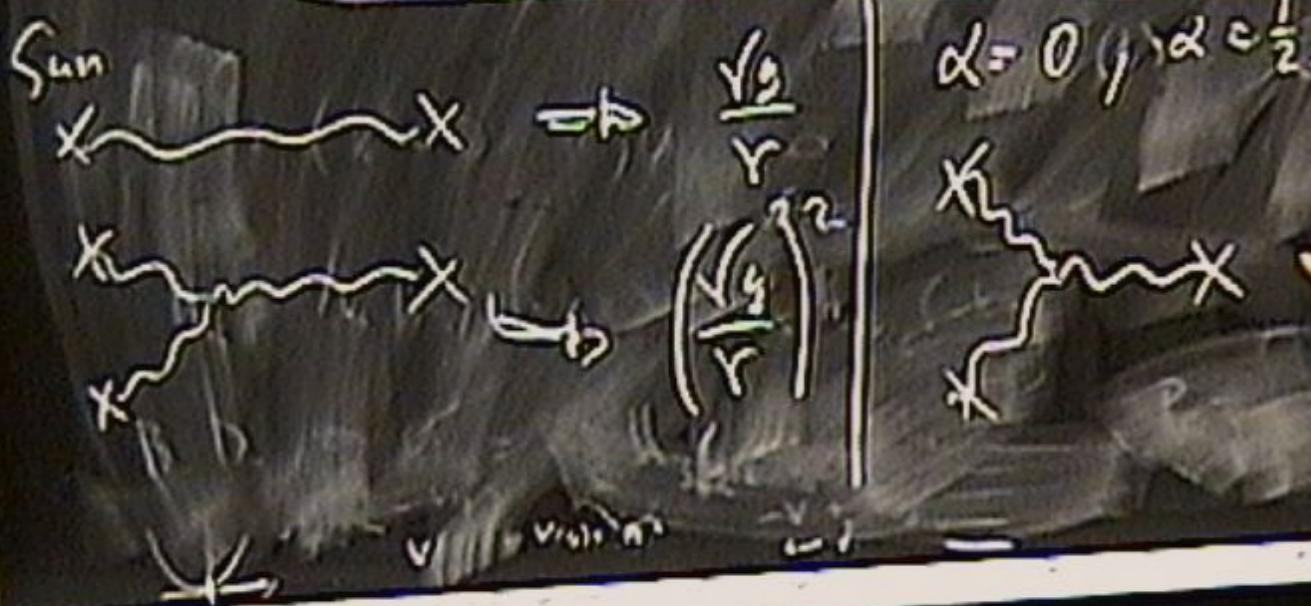
$$\frac{\sqrt{g}}{r}$$

$$\left(\frac{\sqrt{g}}{r}\right)^2$$



$$T^{\mu\nu} \rho_{\mu\nu} = \left( T_{00} - \frac{1}{3} T \right) \left[ \frac{\sqrt{g}}{r} + \mathcal{O}(m^2) \right]_{n=0}$$

$$\left( T_{00} - \frac{1}{2} T \right) \frac{G}{r} \quad \frac{1}{2} = \frac{1}{3}$$



$$T^{\mu\nu}g_{\mu\nu} = \left(T_{00} - \frac{1}{3}T\right) \left[ \frac{\sqrt{g}}{r} + \mathcal{O}(m^2) \right]_{n=1=0} + \left(T_{00} - \frac{1}{2}T\right) \frac{G}{r}$$

$\frac{1}{2} \pm \frac{1}{3}$

Sun

$$\xrightarrow{\frac{\sqrt{g}}{r}}$$

$$\xrightarrow{\left(\frac{\sqrt{g}}{r}\right)^2}$$

$$\alpha = 0, 1, \alpha < \frac{1}{2}$$

$$\xrightarrow{\left(\frac{c}{r}\right)^{4.42} \frac{\sqrt{g}}{r}}$$

$$\kappa_{\text{max}} = 10^{-32}$$



$$x_{\text{max}} = 10^{-32}$$



$\Delta \alpha \approx 10^{-10}$



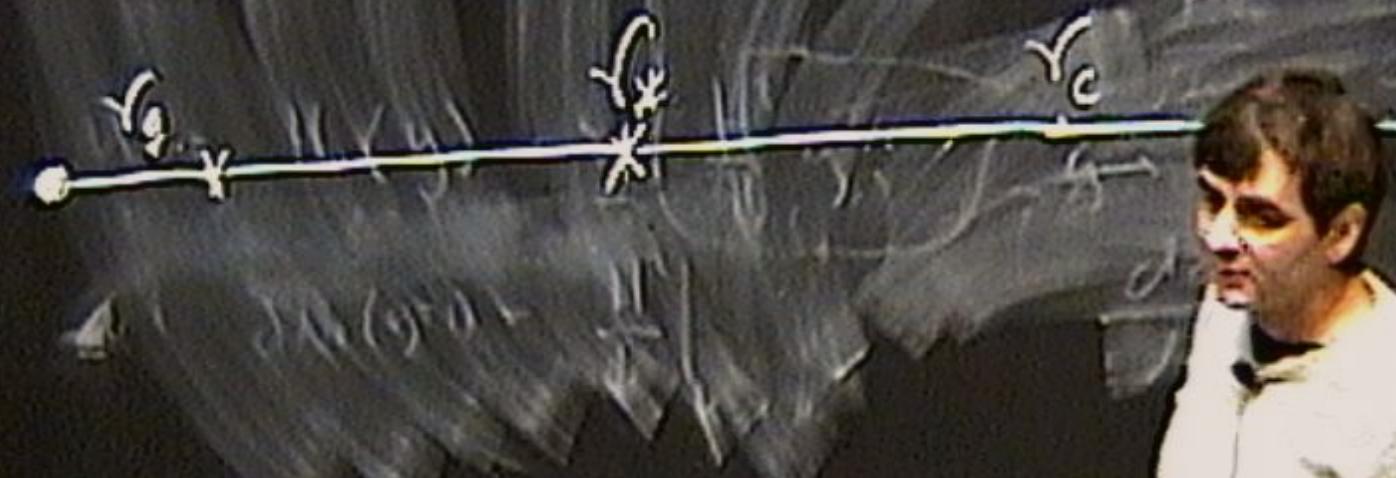
$$x_{\text{max}} = 10^{-32}$$

$$\gamma_* = \left( \sqrt{\zeta} \sqrt{\alpha} \right)^{\frac{1}{5-4\alpha}}$$



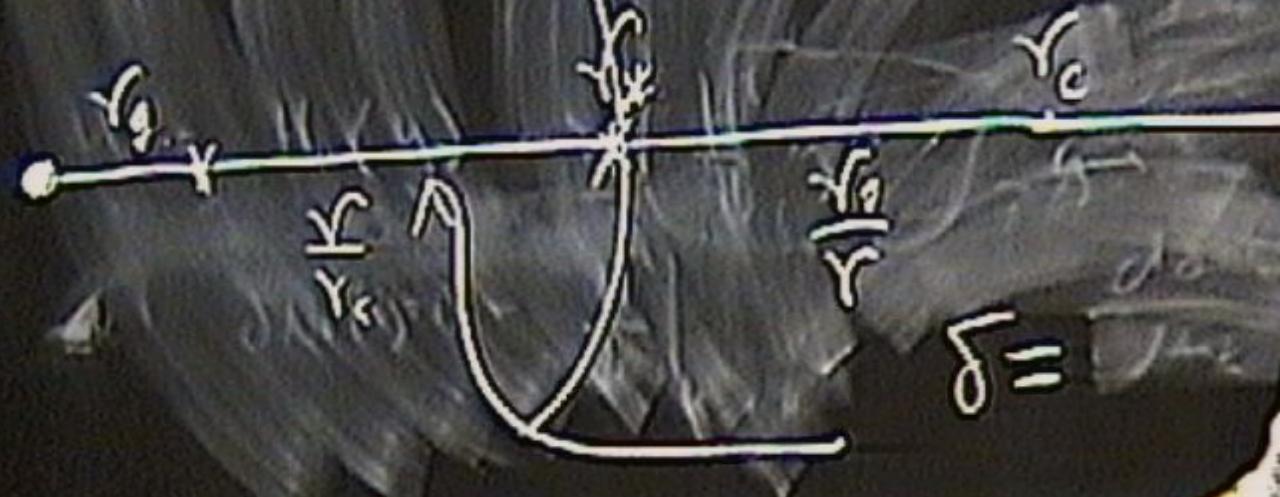
$$x_{\text{max}} = 10^{-32}$$

$$\gamma_* = \frac{(\sqrt{4-4\alpha} - \sqrt{\alpha})}{\sqrt{5-4\alpha}}$$



$$x_{\text{max}} = 10^{-32}$$

$$\gamma_* = \left( \sqrt{\frac{4-4\alpha}{\alpha}} \sqrt{\frac{5-4\alpha}{1-\alpha}} \right)^{\frac{1}{1-\alpha}}$$



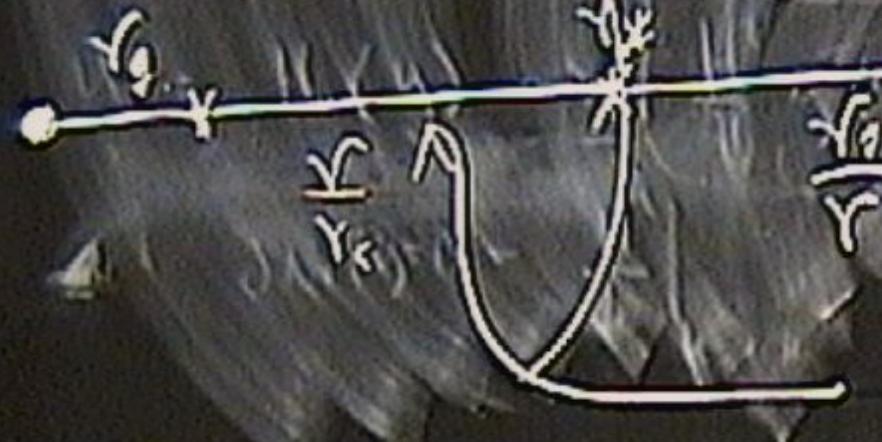
$$\delta =$$



$$x_{\text{max}} = 10^{-32}$$

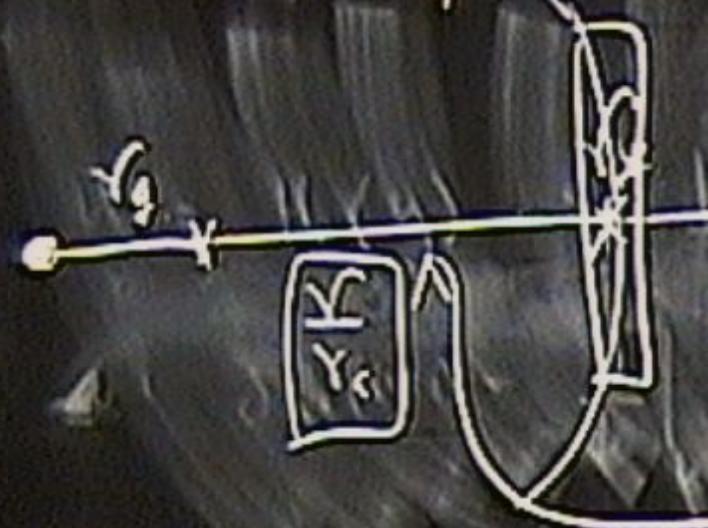
$$\gamma_* = \left( \sqrt{\frac{4-4\alpha}{c}} \sqrt{\frac{5-4\alpha}{g}} \right)^{\frac{1}{5-4\alpha}}$$

$$\gamma_c \equiv H_0$$



$$\delta = (\sqrt{H_0})^{2(1-\alpha)} \sqrt{\frac{r}{r_g}}$$

$$x_{max} = 10^{-32}$$



$$\gamma^* = \left( \sqrt{4-4\alpha} \right)^{\frac{1}{5-4\alpha}}$$

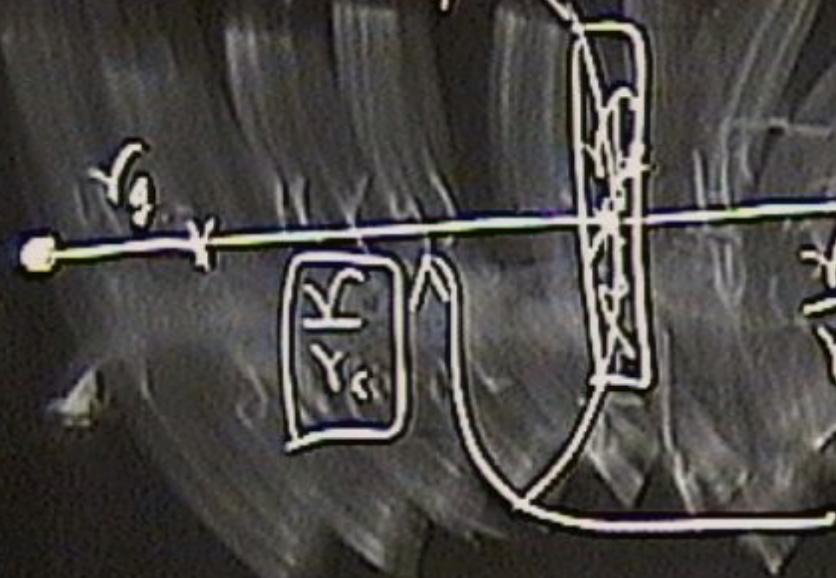
$$\gamma_c \equiv H_0'$$

$$\frac{\gamma^*}{\gamma}$$

$$(H_0)^{2(1-\alpha)} \sqrt{\frac{r}{rg}}$$



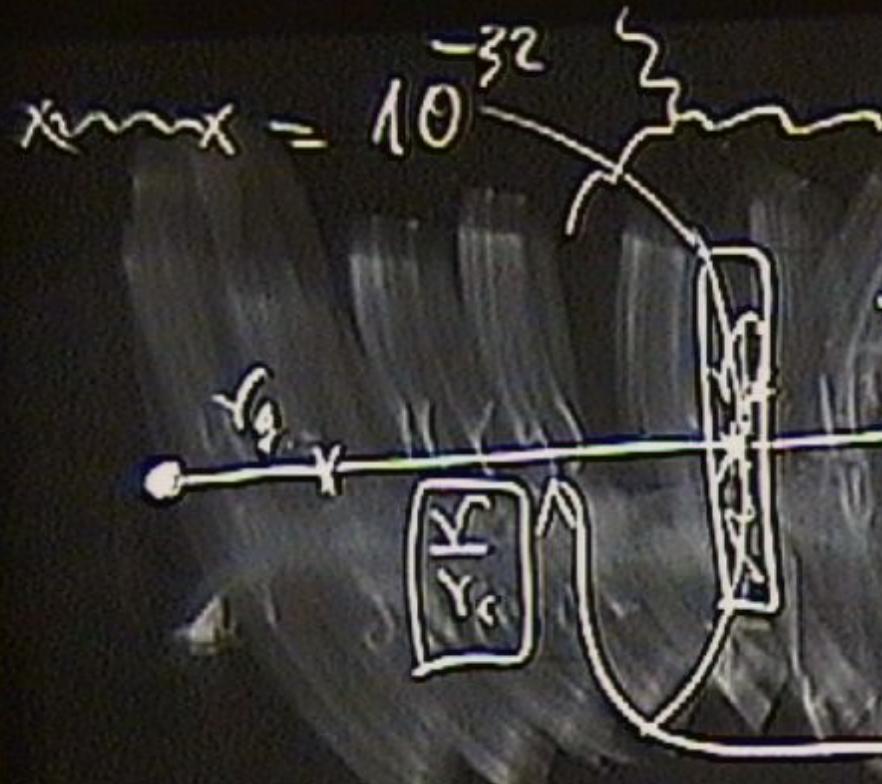
$$x_{\text{max}} = 10^{-32}$$



$$\gamma_* = \left( \sqrt{\gamma_c} \sqrt{\gamma_g} \right)^{\frac{1}{S - 4\alpha}}$$

$\gamma_c \equiv H_0'$

$$\delta = (\sqrt{H_0})^{2(1-\alpha)} \sqrt{\frac{r}{r_g}}$$



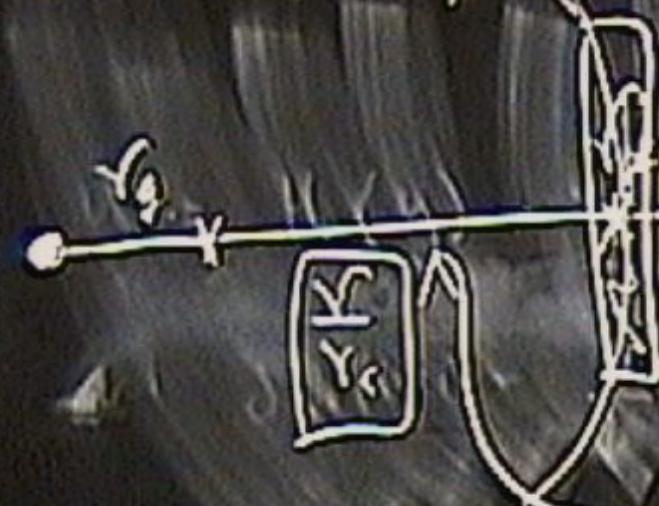
$$\gamma_* = \left( \gamma_c^{4-4\alpha} \gamma_q^{5-4\alpha} \right)^{\frac{1}{5-4\alpha}}$$

$\gamma_c \equiv H_0$

$$\delta = (CH_0)^{2(1-\alpha)}$$



$$x_{\text{max}} = 10^{-32}$$



$$\gamma_* = \left( \gamma_c^{4-4\alpha} \gamma_g \right)^{\frac{1}{5-4\alpha}}$$

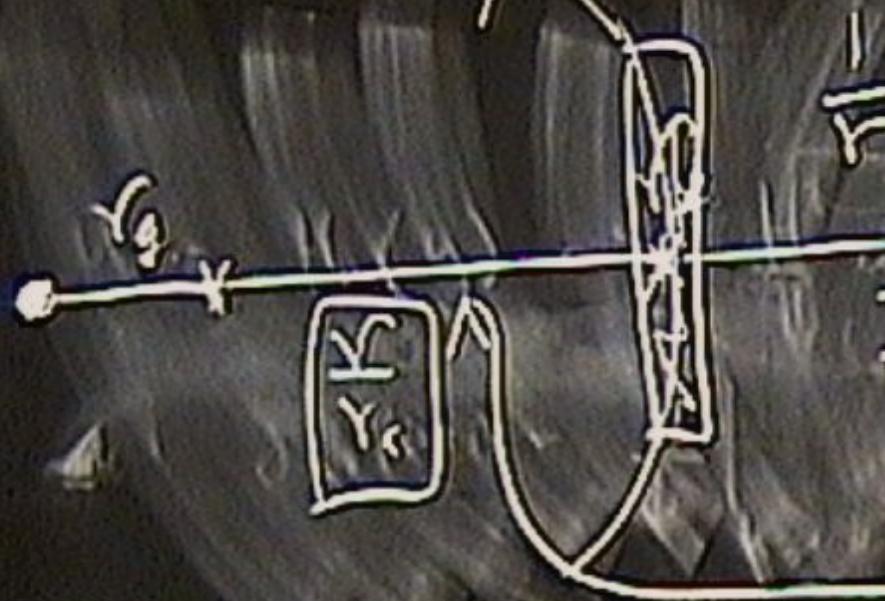
$$\gamma_c \equiv H_0'$$

$$\delta = \left( \gamma H_0 \right)^{2(1-\alpha)} \sqrt{\frac{r}{r_g}}$$

$\underline{\text{Spin-2}}$        $\underline{h_{\mu\nu}}$        $\underline{\nabla^\mu \Gamma_{\mu\nu} = 0, h = h_\mu^\mu}$   
 $(\mathcal{E}h)_{\mu\nu} + M^2(0)(h_{\mu\nu} - \eta_{\mu\nu}h) = T_{\mu\nu}$        $\gamma_c = H_0^{-1}$   
 $E_{\text{matter}} = D h_{\mu\nu} - h_{\mu\nu} D h -$   
 $M^2(0) = \sqrt{c^2(k-1)}$        $\alpha = \frac{1}{2}$

Lunay  $\Rightarrow 0 \leq \alpha \leq 0.5$

$$x_{\text{max}} = 10^{-32}$$



$$\frac{1}{F}$$

$$r_* = \left( \sqrt{\gamma_c} - \sqrt{g} \right)^{\frac{1}{5-4\alpha}}$$

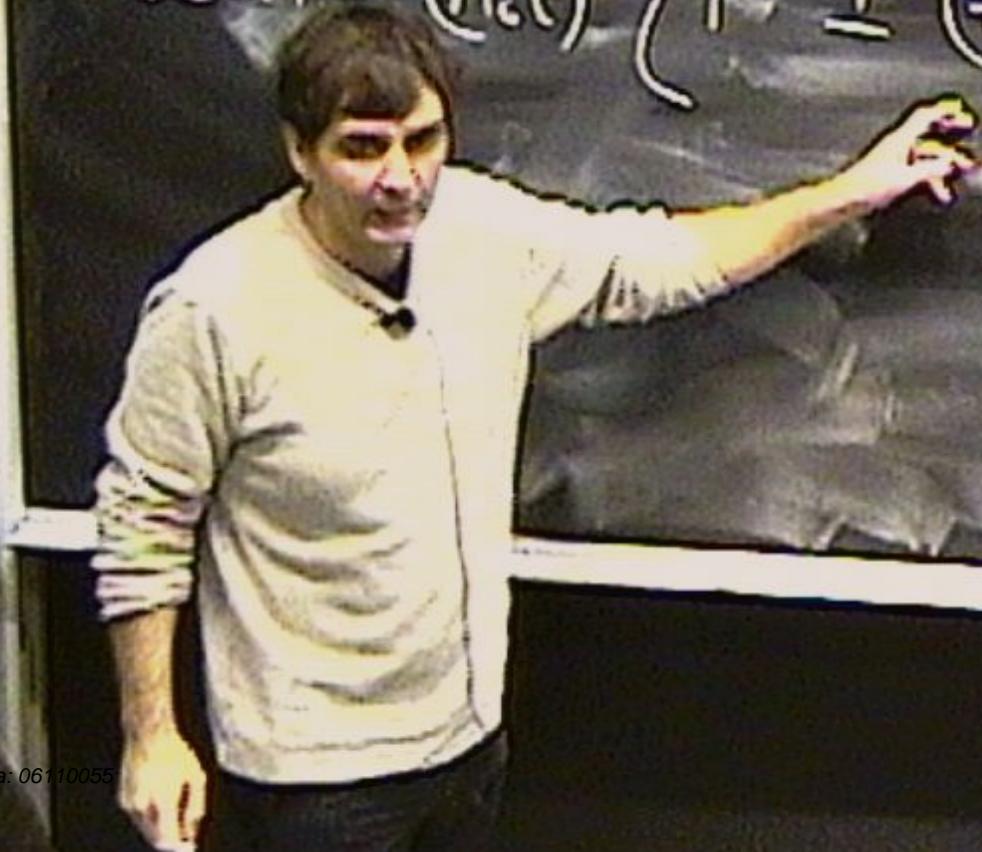
$$\gamma_c = H_0$$

$$\frac{\gamma_1}{\gamma}$$

$$\delta = (\gamma H_0)^{2(1-\alpha)} \sqrt{\frac{F}{rg}}$$

Lunary  $\Rightarrow 0 \leq \alpha \leq 0.5$

$$\vec{a}^3(t) = (H_0 t)^2 \left[ 1 \pm (t H_0)^{2(1-\alpha)} + \dots \right]$$



Lunary  $\Rightarrow 0 \leq \alpha \leq 0.5$

$$\vec{a}(t) = (H_0 t)^2 \left[ 1 \pm (t H_0)^{2(1-\alpha)} + \dots \right]$$

$$t = r_g$$



Lunar  $\Rightarrow 0 \leq \alpha \leq 0.5$

$$\vec{a}(t) = (H(t))^2 \left[ I \pm (tH_0)^{2(1-\alpha)} + \dots \right]$$

$$t = \sqrt{g} \cdot \sqrt{\frac{t}{r_g} - 1}$$

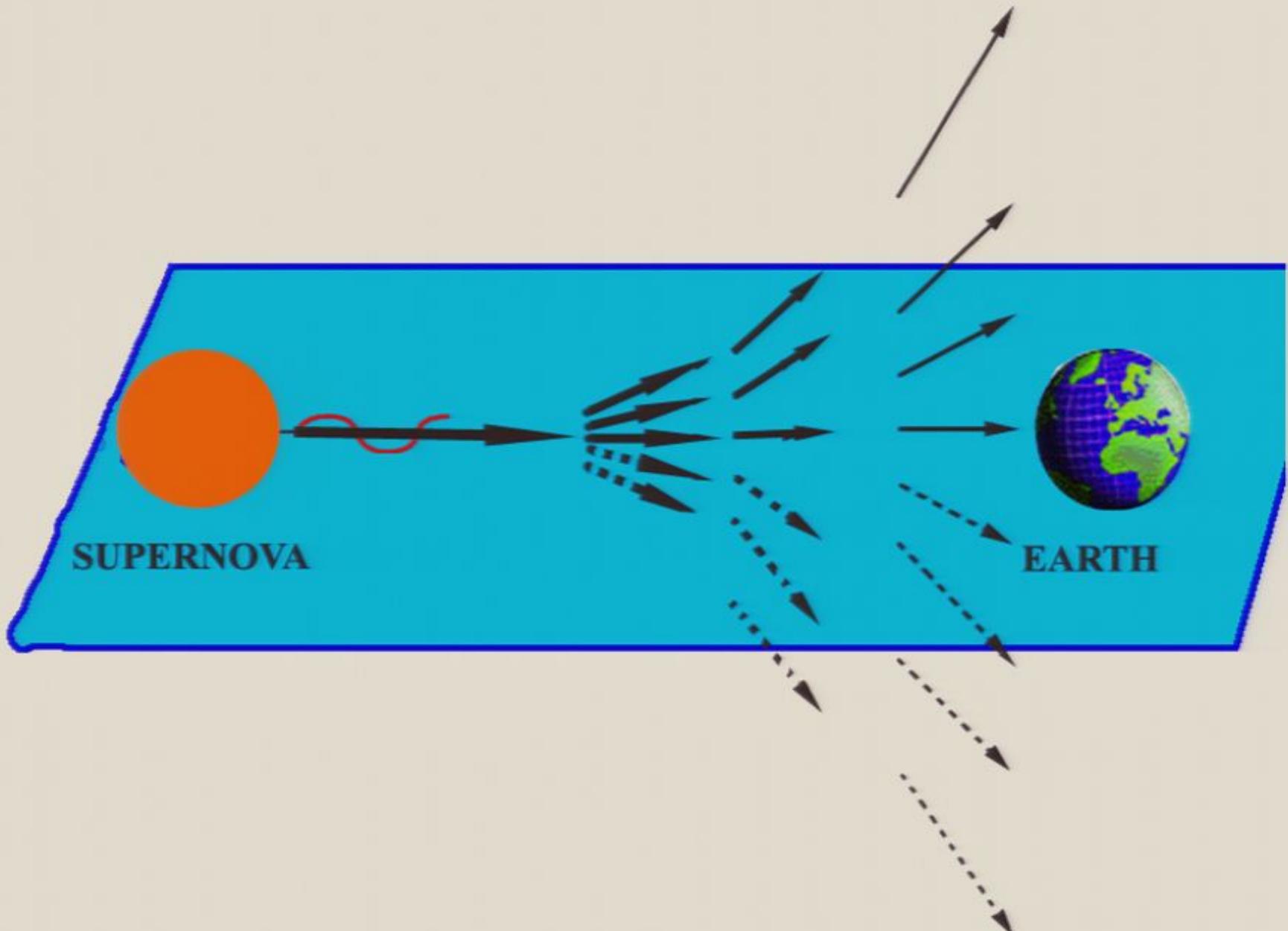


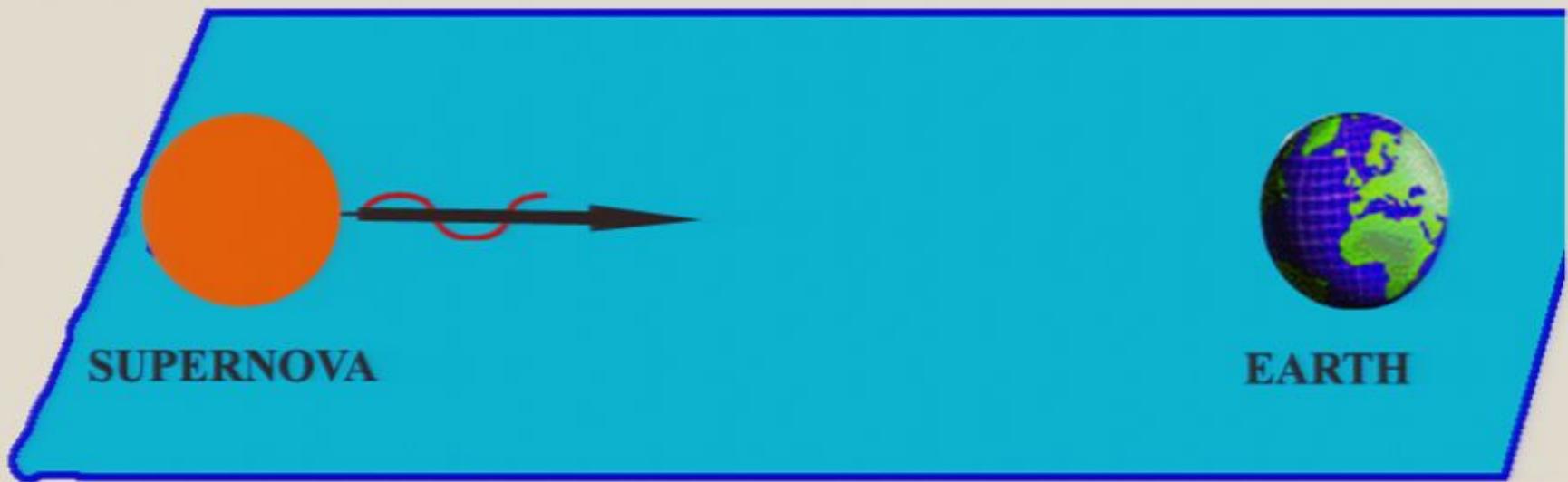
unary  $\Rightarrow 0 \leq \alpha \leq 0.5$

$$\vec{a}(t) = (A_0 t)^2 \left[ 1 \pm (2H_0)^{2(1-\alpha)} + \dots \right]$$

$$\alpha = \frac{1}{2}$$

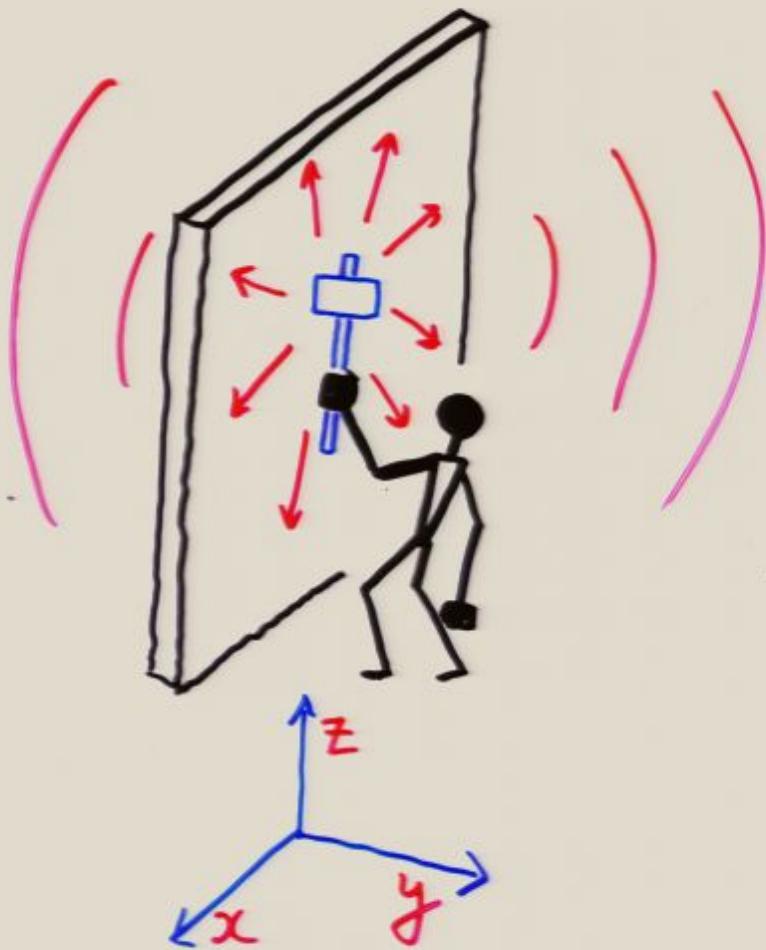
$$t = \sqrt{g} \cdot \sqrt{\frac{t}{g} - 1}$$



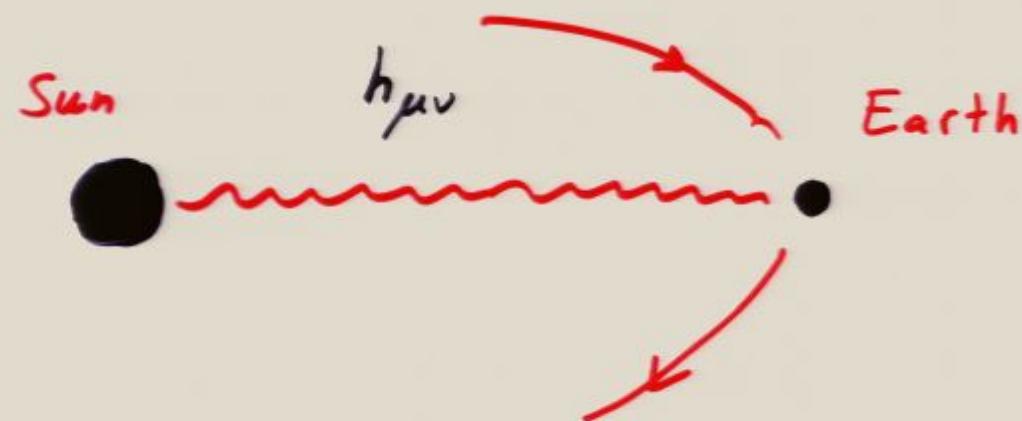
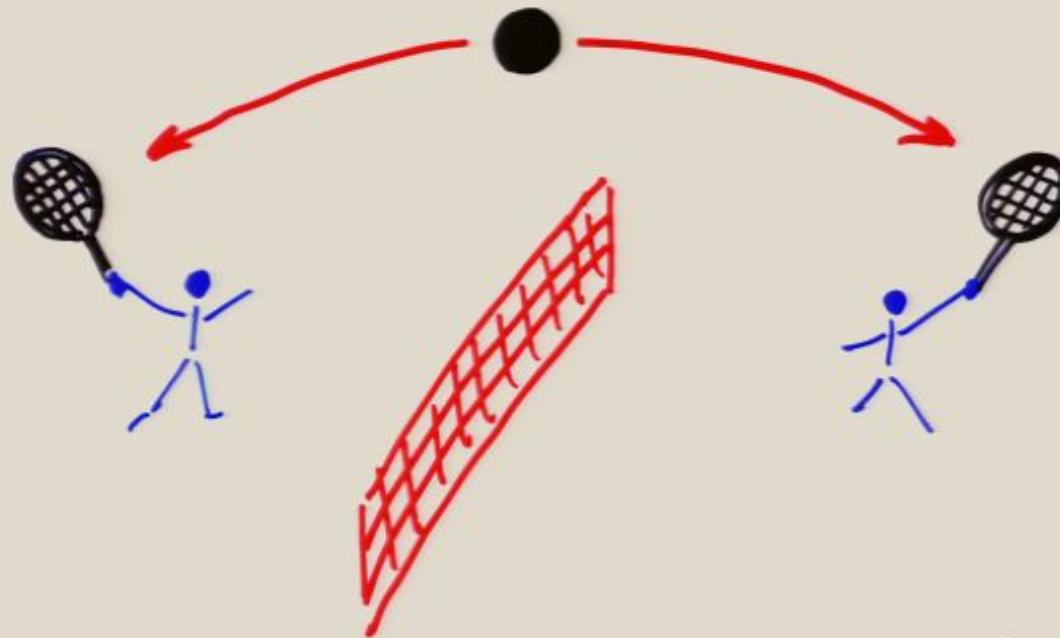


# SOUND WAVE

DGP '00



$$\left\{ \delta(y) \square_{2+1} + \frac{1}{c} \square_{3+1} \right\} \psi = \delta(z)$$



FRW Equation is modified

$$H^2 - H/r_c = \frac{8\pi}{3} G_N \rho$$

Early cosmology in normal  $H \gg r_c^{-1}$

Late cosmology  $H \rightarrow H = r_c^{-1}$

At late times Universe is self-accelerating!

NO NEED IN DARK ENERGY.

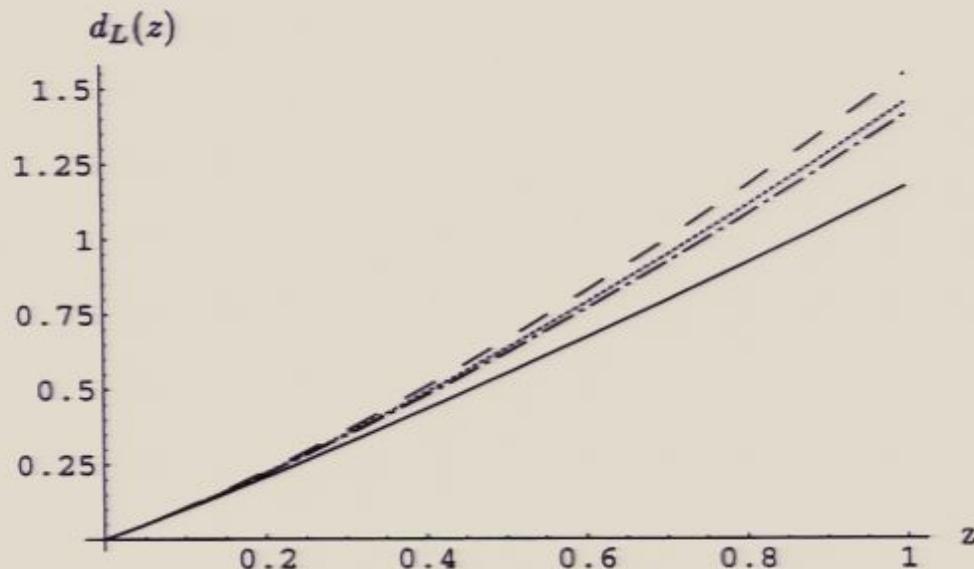
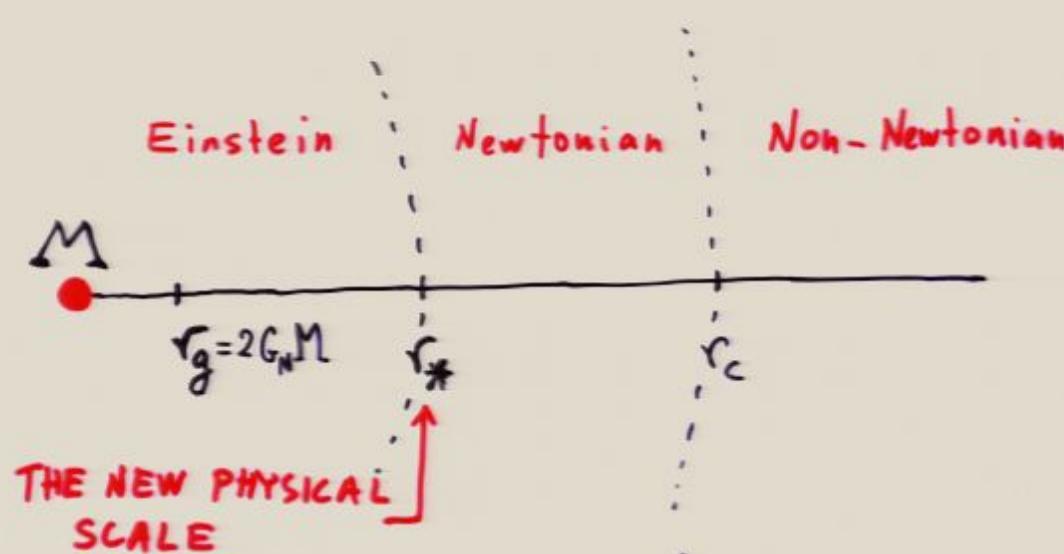


Figure 2: Luminosity distance as a function of red-shift for ordinary cosmology with  $\Omega_\Lambda = 0.7, \Omega_M = 0.3, k = 0$  (dashed line),  $\Omega_\Lambda = 0, \Omega_M = 1, k = 0$  (solid line), and dark energy with  $\Omega_X = 0.7, w_X = -0.6, \Omega_M = 0.3, k = 0$  (dotted-dashed line) and in our model (dotted line) with  $\Omega_M = 0.3$  and a flat universe (for which one gets from equation (28)  $\Omega_{r_c} = 0.12$  and  $r_c = 1.4H_0^{-1}$ ).

Corrections to Einstein are  
Source-dependent

Corrections to Einstein

$$\left\{ \begin{array}{l} \neq \frac{r}{r_c} \\ = \left( \frac{r}{r_*} \right)^{\frac{3}{2}} \quad r_* \ll r_c \end{array} \right.$$



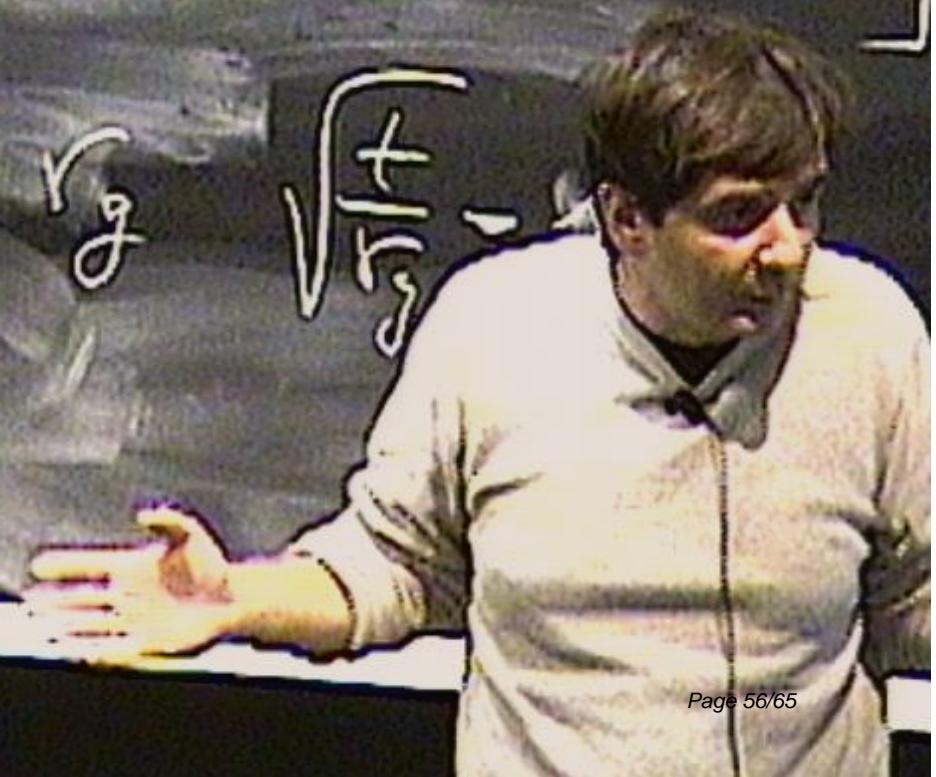
$$r_* = (r_c^2 r_g)^{\frac{1}{3}}$$

Unstable  $\Rightarrow 0 \leq \alpha \leq 0.5$  |  $\alpha \approx 0$

$$\tilde{a}^3(t) = (H_0 t)^2 \left[ 1 \pm (t H_0)^{2(1-\alpha)} + \dots \right]$$

$$\alpha = \frac{1}{2}$$

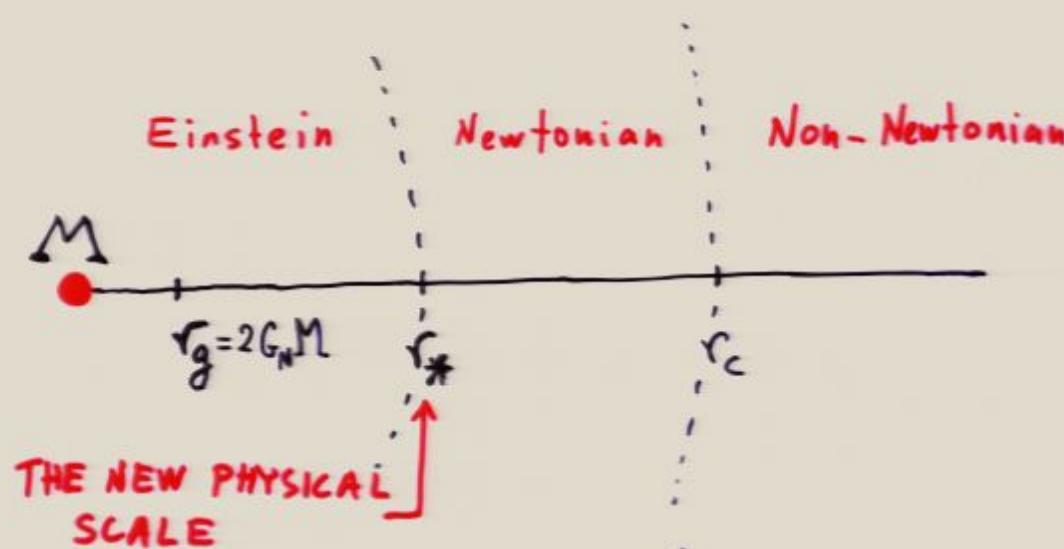
$$t = r_g - \sqrt{\frac{t}{r_g}}$$



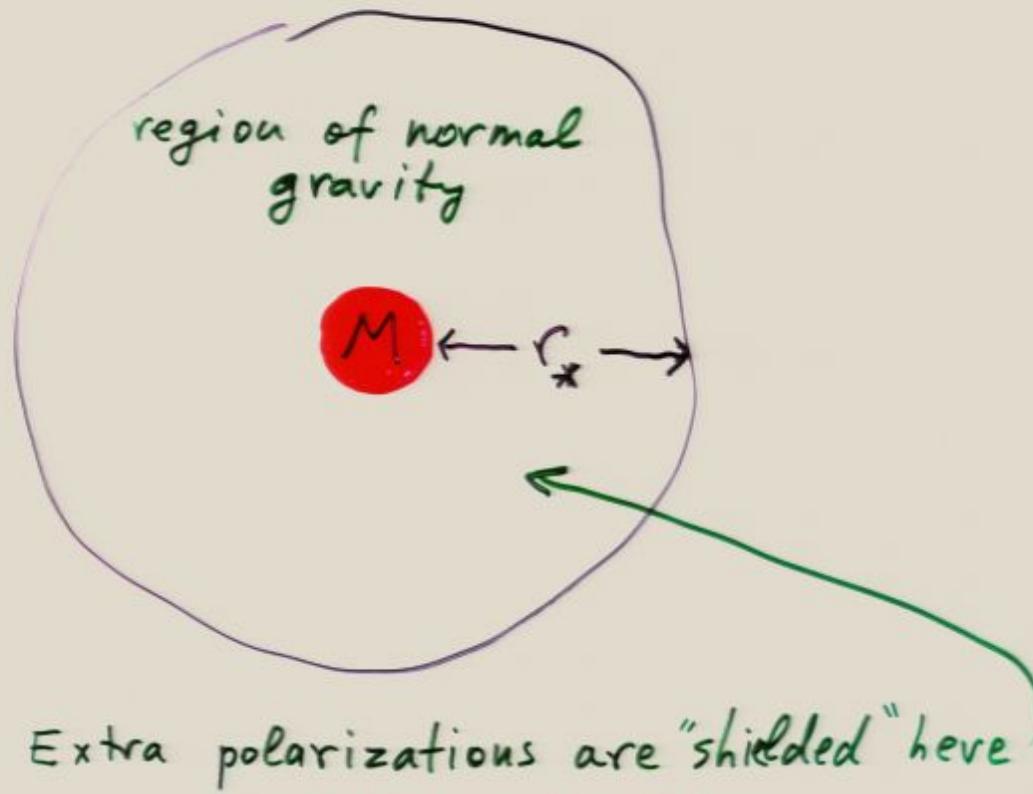
Corrections to Einstein are  
Source-dependent

Corrections to Einstein

$$\left\{ \begin{array}{l} \neq \frac{r}{r_c} \\ = \left( \frac{r}{r_*} \right)^{\frac{3}{2}} \quad r_* \ll r_c \end{array} \right.$$



$$r_* = (r_c^2 r_g)^{\frac{1}{3}}$$



$$r_* \sim (r_c^2 r_g)^{\frac{1}{3}} \ll r_c !$$

$$r_g = 2 G_N M$$

# LUNAR RANGING TEST OF MODIFIED GRAVITY

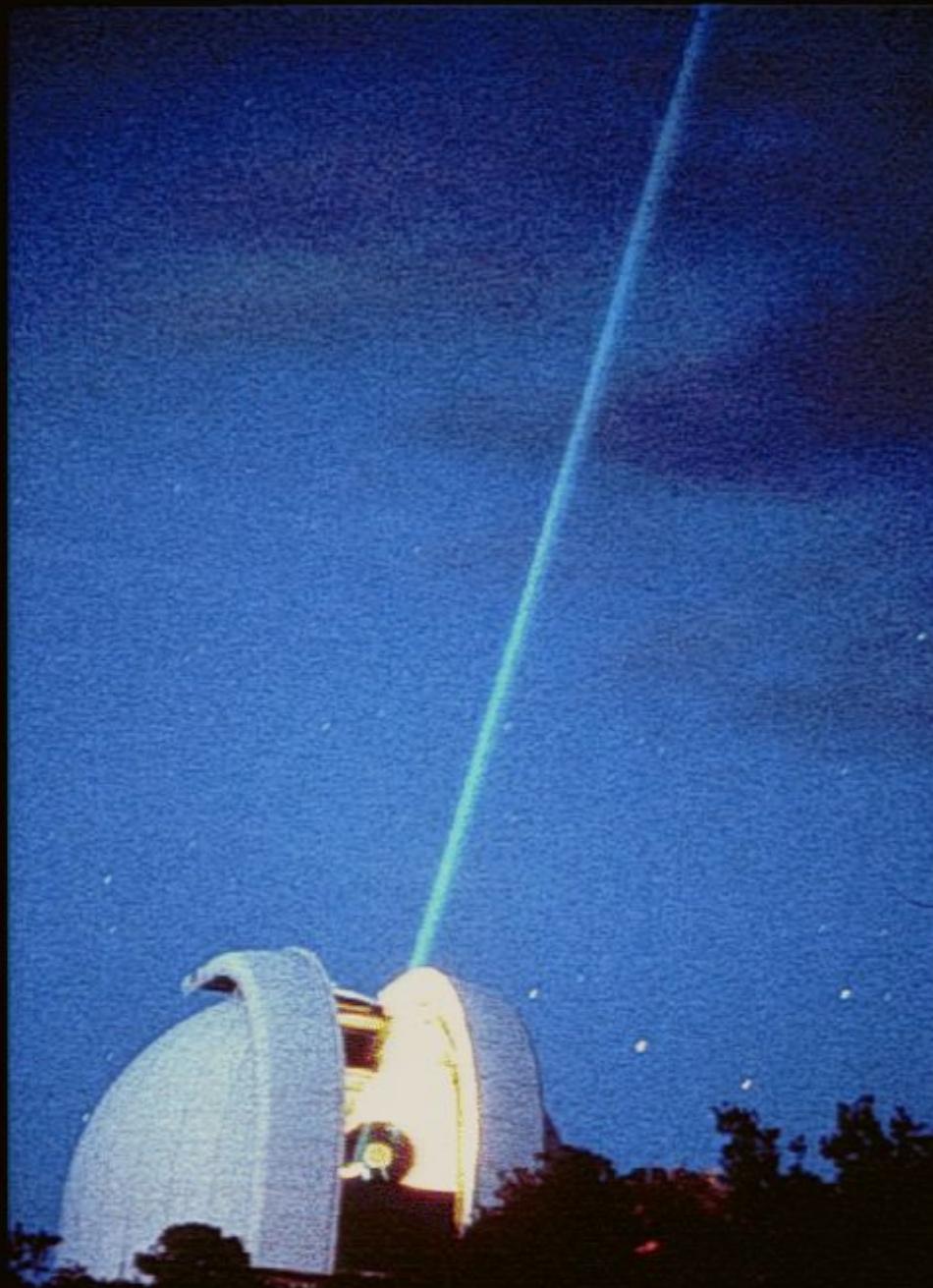
Predicted anomalous perihelion precession of the lunar orbit

$$\delta\phi = 1.4 \times 10^{-12}$$

Todays accuracy:

$$\sigma_\phi = 2.4 \times 10^{-11}$$

10-fold improvement is expected



# LUNAR RANGING TEST OF MODIFIED GRAVITY

Predicted anomalous perihelion precession of the lunar orbit

$$\delta\phi = 1.4 \times 10^{-12}$$

Todays accuracy:

$$\sigma_\phi = 2.4 \times 10^{-11}$$

10-fold improvement is expected

