

Title: Class 3 B

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Abstract: Graduate Course on Standard Model & Quantum Field Theory

Free Spin 1 Particles.
(Massive)

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M

Spin N : $|\vec{p}, h\rangle$

Spin N $|\vec{p}, \sigma\rangle$

$$h = \pm N$$

Free Spin 1 Particles. (Massive)

Massless Spin N : $|\vec{p}, h\rangle$
Massive Spin N : $|\vec{p}, \sigma\rangle$

$$h = \pm N$$
$$-N \leq \sigma \leq N$$

Free Spin 1 Particles. (Massive)

Massless Spin N : $|\vec{p}, h\rangle$ $h = \pm N$
 Massive Spin N : $|\vec{p}, \sigma\rangle$ $-N \leq \sigma \leq N$

$$V_\mu = \sum_{\sigma=0,1} \int d^3p \left[\epsilon_\mu(\vec{p}, \sigma) e^{ipx} a_{p\sigma} + \epsilon_\mu^\lambda(\vec{p}, \sigma) e^{-ipx} \bar{a}_p \right]$$

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$$p^\mu \epsilon_\mu(\vec{p}, \sigma) = 0$$

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$$\phi(x) = \sum_{\sigma=0,1} \int d^3p \left[\epsilon_\mu(\vec{p}, \sigma) e^{ipx} a_{p\sigma} + \epsilon_\mu^\dagger(\vec{p}, \sigma) e^{-ipx} a_{p\sigma}^\dagger \right]$$

$$p^\mu \epsilon_\mu(\vec{p}, \sigma) = 0 \quad \text{---} \quad \partial^\mu V_\mu = 0$$

Free Spin 1 Particles. (Massive)

$$E_p = \sqrt{\vec{p}^2 + m^2}$$

$$p^\mu p_\mu + m^2 = 0$$

$$(\partial^\mu \partial_\mu - m^2) \phi = 0$$

$$\psi$$

$$V_\mu$$

Massless Spin N : $|h\rangle$
Massive Spin N : $|\sigma\rangle$

$$h = \pm N$$

$$-N \leq \sigma \leq N$$

$$V_\mu(x) = \sum_{\sigma=0,1} \int d^3p \left[\epsilon_\mu(\vec{p}, \sigma) e^{ipx} a_{p\sigma} + \epsilon_\mu^\dagger(\vec{p}, \sigma) e^{-ipx} a_{p\sigma}^\dagger \right]$$

$$\partial^\mu V_\mu = 0$$

For 1 field:

$$\mathcal{L} = -A - B \vee V^a$$

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$$\mathcal{L} = -A - B \underline{V} V^{\mu} - C \partial_{\mu} V_{\nu} \partial^{\mu} V^{\nu} - D \partial_{\mu} V_{\nu} \partial^{\nu} V^{\mu}$$



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Require $C+D=0$ in order to ensure $\partial^{\mu} V_{\mu} = 0$ for all solutions.

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$$\mathcal{L} = -A - B V_\mu V^\mu - C \left(\partial_\mu V_\nu (\partial^\mu V^\nu - \partial^\nu V^\mu) \right)$$

$$= -A - \frac{B'}{2} V_\mu V^\mu - \frac{C'}{4} f_{\mu\nu} f^{\mu\nu}$$

$$f_{\mu\nu} = \partial_\mu V_\nu - \partial_\nu V_\mu$$

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$$\rightarrow (\Box - \mu^2) V_\mu = 0$$

Free Spin 1 Particles. (Massive)

$$E_p = \sqrt{\vec{p}^2 + m^2}$$

$$p^\mu p_\mu + m^2 = 0$$

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Massive Spin N : $|\vec{p}, \sigma\rangle$

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$$V_\mu(x) = \sum_{\sigma=0,1} \int d^3p \left[\epsilon_\mu(\vec{p}, \sigma) e^{ipx} a_{p\sigma} + \epsilon_\mu^\lambda(\vec{p}, \sigma) e^{-ipx} a_{p\sigma}^* \right]$$

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Massless spin

Massless spin

(A, B)
 (λ, λ)

$$\lambda = A - B$$

Massless spin 1

(i) It is impossible to represent a massless spin 1 particle by a vector field.

$$\lambda = A \cdot \theta$$

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simplest field which can represent massless spin 1 is $f_{\mu\nu}$

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The simplest field which can represent massless spin 1

is $f_{\mu\nu}$

$$f_{\mu\nu}(x) = \sum_{\lambda=\pm 1} \int \frac{d^3p}{p} \left[\underline{\epsilon_{\mu\nu}(\vec{p}, \lambda)} e^{ipx} a_{\lambda} + \epsilon_{\mu\nu}^*(\vec{p}, \lambda) e^{-ipx} \bar{a}_{\lambda} \right]$$

Massless spin 1

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The simplest field which can represent massless spin 1

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$\epsilon_{\mu\nu} = p_{\mu} \tilde{\epsilon}_{\nu} - p_{\nu} \tilde{\epsilon}_{\mu}$ where

Massless spin 1

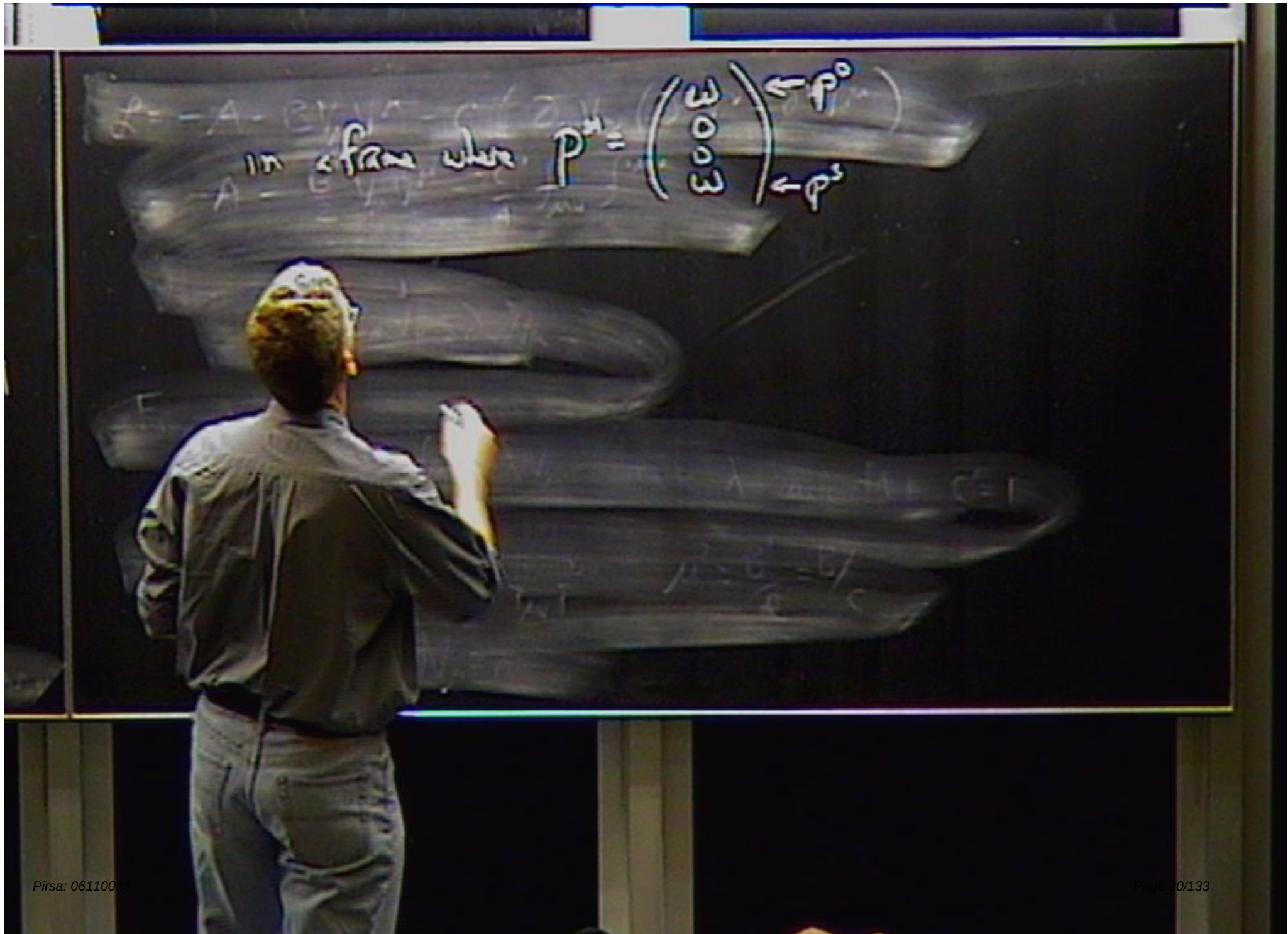
(i) It is impossible to represent a massless spin 1 particle by a vector field.

The simplest field which can represent massless spin 1

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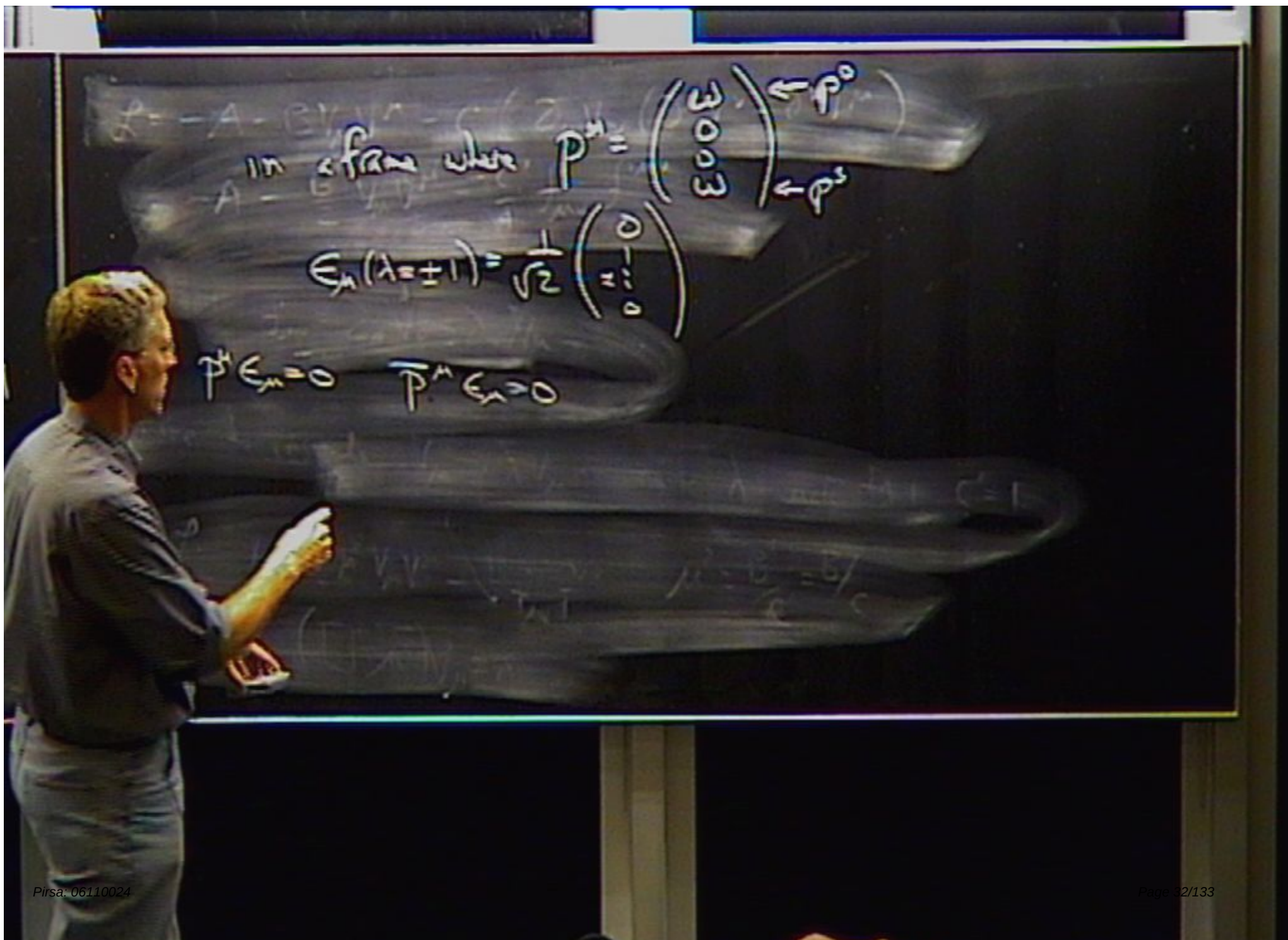
$$\epsilon_{\mu\nu} = p_\mu \tilde{\epsilon}_\nu - p_\nu \tilde{\epsilon}_\mu \quad \text{where}$$





in a frame where $P^\mu = \begin{pmatrix} E \\ 0 \\ 0 \\ 3 \end{pmatrix}$ $\leftarrow P^0$
 $\leftarrow P^3$

$$E_\lambda(\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$



in a frame where $p^\mu = \begin{pmatrix} E \\ 0 \\ 0 \\ 3 \end{pmatrix} \begin{matrix} \leftarrow p^0 \\ \\ \leftarrow p^3 \end{matrix}$

$$E_n(\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 1 \\ 0 \end{pmatrix}$$

$$p^\mu E_n = 0 \quad \bar{p}^\mu E_n = 0$$

in a frame where $p^\mu = \begin{pmatrix} E \\ 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow p^0$

$$\epsilon_\mu(\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \pm i \\ 0 \end{pmatrix} \leftarrow p^3$$

$$p^\mu \in \bar{p}^\mu \epsilon_\mu = 0$$

$$\bar{p}^\mu \bar{p}_\mu = 0 = p^\mu p_\mu$$

$$p^\mu \bar{p}_\mu = -1$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$S = \int d^4x$$

in a frame where $p^\mu = \begin{pmatrix} \omega \\ 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow p^0$

$\epsilon_\mu(\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \pm i \\ 0 \end{pmatrix} \leftarrow p^3$

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So $f_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$

$A_\mu = \sum_{\lambda=\pm 1} \int d^3p \left[\epsilon_\mu(p, \lambda) e^{ipx} a_{p\lambda} + \epsilon_\mu^*(p, \lambda) e^{-ipx} a_{p\lambda}^\dagger \right]$

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ϕ
 V_μ

$$U(\Lambda) f_{\mu\nu}(x) U^\dagger(\Lambda) = \Lambda_\mu^\alpha \Lambda_\nu^\beta f_{\alpha\beta}(\Lambda x)$$

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$$A: U(\Lambda) A_\mu(x) U^\dagger(\Lambda) = \Lambda_\mu^\alpha A_\alpha(\Lambda x) + \partial_\mu \Omega$$

for some Ω

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+
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$$\mathcal{L} = -A - \frac{1}{4} f_{\mu\nu} f^{\mu\nu}$$

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+
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$$\mathcal{L} = -A - \frac{1}{4} f_{\mu\nu} f^{\mu\nu}$$

in a frame where $p^\mu = \begin{pmatrix} \omega \\ 0 \\ 0 \\ \omega \end{pmatrix} \leftarrow p^0$

eg M

$$\vec{\epsilon}_\lambda(\lambda=\pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \pm i \\ 0 \end{pmatrix}$$

$$p^\mu \vec{\epsilon}_\lambda = 0$$

$$\bar{p}^\mu \vec{\epsilon}_\lambda = 0$$

$$\bar{p}^\mu \bar{p}_\mu = 0 = p^\mu p_\mu$$

$$p^\mu \bar{p}_\mu = -1$$

So $f_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$

$$A_\mu = \sum_{\lambda=\pm 1} \int d^3p \left[\vec{\epsilon}_\lambda(p, \lambda) e^{ipx} a_{p\lambda} + \vec{\epsilon}_\lambda^*(p, \lambda) e^{-ipx} a_{p\lambda}^\dagger \right]$$

in a frame where $p^\mu = \begin{pmatrix} \omega \\ 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow p^0$

eg M

$$\vec{\epsilon}_\lambda(\lambda=\pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \pm i \\ 0 \end{pmatrix}$$

$$\lambda \gtrsim \frac{4M}{\omega^2}$$

$$p^\mu \vec{\epsilon}_\lambda = 0$$

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$$\partial_\mu \Psi - \partial_\nu \Psi$$

frame where $P^\mu = \begin{pmatrix} \omega \\ 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow P^0$

$$\approx M$$

$$(\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \mp i \\ 0 \end{pmatrix}$$

$$\wedge \gtrsim \frac{4M}{\sigma^2}$$

$$P^\mu \tilde{\Sigma}_\mu = 0$$

$$\bar{P}^\mu \bar{P}_\mu = 0 = P^\mu P_\mu$$

$$P^\mu \bar{P}_\mu = -1$$

So

$$-\partial_\nu A_\mu$$

$$\left(\tilde{\Sigma}_\mu(\vec{p}, \lambda) e^{ipx} a_{p\lambda} + \tilde{\Sigma}_\mu^*(\vec{p}, \lambda) e^{-ipx} a_{p\lambda}^* \right)$$

$$\partial_\mu \psi - \partial_\mu \psi$$

in a frame where $P^\mu = \begin{pmatrix} \omega \\ 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow p^0$

eg M

$$\wedge \gtrsim \frac{4M}{\sigma^2}$$

$$\bar{\psi}_\lambda (\lambda = \pm 1) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ \pm i \\ 0 \end{pmatrix}$$

Chiral

(2,1)
(1,2)

$$\bar{P}^\mu \bar{\psi}_\lambda = 0$$

$$\bar{P}^\mu \bar{P}_\mu = 0 = P^\mu P_\mu$$

$$P^\mu \bar{P}_\mu = -1$$

So for $\nu - \partial_\nu A_\mu$

$$\left[\bar{\psi}_\lambda(\vec{p}, \lambda) e^{i\vec{p}\cdot\vec{x}} a_{\vec{p}, \lambda} + \bar{\psi}_\lambda^*(\vec{p}, \lambda) e^{-i\vec{p}\cdot\vec{x}} a_{\vec{p}, \lambda}^* \right]$$

Interactions

Residuos

$$E = \sqrt{p^2 + m^2}$$

$$F = \frac{1}{2} \frac{d^2 \phi}{dt^2} + V(\phi)$$

$$V(\phi) = \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{4!} \phi^4$$

$$H = \int d^3x \left(\frac{1}{2} \dot{\phi}^2 + \frac{1}{2} (\nabla \phi)^2 + V(\phi) \right)$$

$$H = H_0 + H_I$$

$$H_0 = \int d^3x \left(\frac{1}{2} \dot{\phi}^2 + \frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 \right)$$

$$H_I = \int d^3x \left(\frac{\lambda}{4!} \phi^4 \right)$$

Interactions

- Spins $D, 1/2$.

Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i

$N_{1/2}$ spin 1/2 : ψ^a

$\mathcal{L} =$

Interactions

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$$\mathcal{L} = -A - \frac{1}{2} \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + \mu_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a$$

Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i

$N_{1/2}$ spin 1/2 : ψ^a

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Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin 1/2 : ψ^a (majorana)

$$\mathcal{L} = -A - \frac{1}{2} \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + \mu_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a + \mathcal{L}_{int}$$

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Interactions

- Spins $0, 1/2$:

$\hbar=c=1$

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin $1/2$: ψ^a (majorana)

$$\mathcal{L} = \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + \mu_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a + \mathcal{L}_{int}$$

$$[\mathcal{L}] = 4$$

Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin 1/2 : ψ^a (majorana)

$$-A = \int d^4x \left(\sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + \mu_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a \right) + \mathcal{L}_{int}$$

$$[\alpha^\mu] = 4 \text{ i.e. dim of } \alpha^\mu \text{ is } (mass)^4$$

$$[\partial_\mu] = 1 \dots \partial \dots mass$$

Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin 1/2 : ψ^a (majorana)

$\hbar=c=1$

$$\mathcal{L} = -\mathcal{A} = -\frac{1}{2} \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + m_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a + \mathcal{L}_{int}$$

$$\mathcal{L} = \int d^4x \mathcal{L}$$

$$[\mathcal{L}] = 4 \text{ i.e. dim of } \mathcal{L} \text{ is } (mass)^4$$

$$[\partial_\mu] = 1 \quad \dots \quad \partial \sim mass$$

$$[\phi^i] = 1 \quad \dots \quad \phi^i \sim mass$$

Interactions

- Spins 0, 1/2:

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin 1/2 : ψ^a (majorana)

$$= -A - \frac{1}{2} \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + m_i^2 \phi^i \phi^i \right) - \frac{1}{2} \sum_a \bar{\psi}^a (\not{\partial} + m_a) \psi^a + \mathcal{L}_{int}$$

$$[A] = 4$$

$$[\phi] = 4 \quad \text{dim of } \phi \text{ is } (mass)^4$$

$$[\partial_\mu] = 1$$

$$[\phi^i] = 1$$

$$[\psi] = 3/2$$

Interactions

- Spins 0, 1/2:

$\hbar = c = 1$

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$$[\psi^i] = 3/2$$

L_{int}

= least squares (A) It is impossible to represent

A continuous space-time field.

continuous field which can represent a continuous space

is a continuous space-time field.

is a continuous space-time field.

is a continuous space-time field.

$$\mathcal{L}_{\text{int}} = g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l + K_{ij} \phi^i \phi^j + \dots$$

A dimensionless spin-1 particle (e.g. a vector field).

A field which can represent a spin-1 particle.

$$A_\mu = \sum_i A_\mu^i T_i$$

$$A_\mu = \sum_i A_\mu^i T_i = \sum_i A_\mu^i \frac{1}{2} \sigma_i$$

$$\mathcal{L}_{int} = g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l + K_{ij} \phi^i \phi^j + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$\mathcal{L}_{\text{int}} = g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l + \kappa_{ij} \phi^i \phi^j + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$[y_{iabc} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i + \text{c.c.}] + \tau_{ijab} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i \phi^j + \dots$$

$$\begin{aligned}
 \mathcal{L}_{\text{int}} = & g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l + \kappa_{ijk} \phi^i \phi^j \phi^k + \dots \\
 & + \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots \\
 & + \left[y_{iabc} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i + \text{c.c.} \right] + \tau_{ijkl} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i \phi^j + \dots \\
 & + w_{ijkl} \bar{\Psi}^a \gamma_\mu \gamma_\nu \Psi^b \partial^\mu \phi^i + \dots \\
 & + h_{ijkl} \bar{\Psi}^a \gamma_\mu \Psi^b \bar{\Psi}^c \gamma^\mu \Psi^d + \dots
 \end{aligned}$$

$$\begin{aligned}
\mathcal{L}_{\text{int}} = & g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l + \kappa_{ijk} \phi^i \phi^j \phi^k + \dots \\
& + \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots \\
& + \left[y_{iab} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i + \text{c.c.} \right] + \tau_{ijab} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i \phi^j + \dots \\
& + w_{iac} \bar{\Psi}^a \gamma_\mu \gamma_\nu \psi^b \partial^\mu \phi^c + \dots \\
& + h_{abcd} \bar{\Psi}^a \gamma_\mu \psi^b \bar{\Psi}^c \gamma_\nu \psi^d + \dots
\end{aligned}$$

$$\mathcal{L}_{int} = \overset{4}{g_{ijkl}} \overset{1}{\phi^i} \overset{1}{\phi^j} \overset{1}{\phi^k} \overset{0}{\phi^l} + \overset{0}{\lambda_{ijkl}} \overset{1}{\phi^i} \overset{1}{\phi^j} \overset{1}{\phi^k} \overset{1}{\phi^l} + \overset{0}{\kappa_{ij}} \overset{1}{\phi^i} \overset{1}{\phi^j} + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$[g_{ij}]^{-1}$$

$$[\lambda_{ijkl}]$$

$$+ [y_{iabc} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i + c.c.] + z_{ijkl} \bar{\Psi}^a \gamma_\mu \Psi^b \phi^i \phi^j + \dots$$

$$+ w_{iabc} \bar{\Psi}^a \gamma_\mu \gamma_\nu \Psi^b \partial^\mu \phi^i + \dots$$

$$+ l_{abcd} \bar{\Psi}^a \gamma_\mu \Psi^b \bar{\Psi}^c \gamma_\mu \Psi^d + \dots$$

$$\mathcal{L}_{int} = \overset{4}{g_{ijk}} \overset{1}{\phi^i} \overset{1}{\phi^j} \overset{1}{\phi^k} + \overset{0}{\lambda_{jke}} \overset{1}{\phi^i} \overset{1}{\phi^k} \overset{1}{\phi^e} \overset{1}{\phi^j} + \overset{4-n}{K_{i_1 \dots i_n}} \overset{1}{\phi^{i_1}} \dots \overset{1}{\phi^{i_n}} + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$[g_{ijk}] = 1$$

$$[\lambda_{jke}] = 0$$

$$[K_{i_1 \dots i_n}] = 4-n$$

$$+ [y_{iab} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i + c.c.] + \tau_{ijab} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i \phi^j + \dots$$

$$+ w_{iab} \bar{\Psi}^a \gamma_\mu \gamma^\nu \psi^b \partial_\nu \phi^i + \dots$$

$$+ l_{abcd} \bar{\Psi}^a \gamma^\mu \psi^b \bar{\Psi}^c \gamma_\mu \psi^d + \dots$$

$$\mathcal{L}_{\text{int}} = \overset{4}{g_{ijk}} \overset{1}{\phi^i} \overset{1}{\phi^j} \overset{1}{\phi^k} + \overset{0}{\lambda_{jke}} \overset{1}{\phi^i} \overset{1}{\phi^k} \overset{1}{\phi^e} \overset{1}{\phi^j} + \overset{4-n}{K_{i_1 \dots i_n}} \overset{1}{\phi^{i_1}} \dots \overset{1}{\phi^{i_n}} + \dots$$

$$+ \alpha_{ijk} \overset{1}{\phi^i} \overset{1}{\partial_\mu \phi^j} \overset{1}{\partial^\mu \phi^k} + \dots$$

$$[g_{ijk}] = 1$$

$$[\lambda_{jke}] = 0$$

$$[K_{i_1 \dots i_n}] = 4-n$$

$$[\alpha_{ijk}] = -1$$

$$[y_{iab}] = 0$$

$$+ \left[y_{iab} \bar{\Psi}^a \gamma_\mu \psi^b \overset{1}{\phi^i} + \text{c.c.} \right] + \tau_{ijab} \bar{\Psi}^a \gamma_\mu \psi^b \overset{1}{\phi^i} \overset{1}{\phi^j} + \dots$$

$$+ w_{iac} \bar{\Psi}^a \gamma_\mu \gamma_\nu \psi^c \partial^\mu \phi^i + \dots$$

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$$\mathcal{L}_{int} = \overset{4}{g_{ijk}} \overset{1}{\phi^i} \overset{1}{\phi^j} \overset{1}{\phi^k} + \overset{0}{\lambda_{jke}} \overset{1}{\phi^i} \overset{1}{\phi^k} \overset{1}{\phi^e} \overset{1}{\phi^j} + \overset{4-n}{K_{i_1 \dots i_n}} \overset{1}{\phi^{i_1}} \dots \overset{1}{\phi^{i_n}} + \dots$$

$$+ \alpha_{ijk} \overset{1}{\phi^i} \overset{1}{\partial_\mu \phi^j} \overset{1}{\partial^\mu \phi^k} + \dots$$

$$[g_{ijk}] = 1$$

$$[\lambda_{jke}] = 0$$

$$[K_{i_1 \dots i_n}] = 4-n$$

$$[\alpha_{ijk}] = -1$$

$$[y_{iabc}] = 0$$

$$[\tilde{z}_{ijab}] = -1$$

$$[w] = -1$$

$$[l] = -2$$

$$+ \left[y_{iabc} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\Psi^c} \overset{1}{\phi^i} + c.c \right] + \tilde{z}_{ijab} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\Psi^c} \overset{1}{\phi^i} \overset{1}{\phi^j} + \dots$$

$$+ w_{iabc} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\gamma_c} \overset{1}{\Psi^d} \overset{1}{\partial_\mu \phi^i} + \dots$$

$$+ l_{abcd} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\Psi^c} \overset{1}{\bar{\Psi}^d} \overset{1}{\gamma_e} \overset{1}{\Psi^e} + \dots$$

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$$+ \alpha_{ijk} \overset{1}{\phi^i} \partial_\mu \overset{1}{\phi^j} \partial^\mu \overset{1}{\phi^k} + \dots$$

$$* [g_{ijk}] = 1$$

$$* [\lambda_{jke}] = 0$$

$$[K_{i_1 \dots i_n}] = 4-n$$

$$[\alpha_{ijk}] = -1$$

$$* [y_{iab}] = 0$$

$$[\bar{z}_{ijab}] = -1$$

$$[w] = -1 \quad [l] = -2$$

$$+ \left[y_{iab} \bar{\Psi}^a \gamma_\mu \psi^b \overset{1}{\phi^i} + c.c. \right] + \bar{z}_{ijab} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i \phi^j + \dots$$

$$+ w_{iab} \bar{\Psi}^a \gamma_\mu \gamma^\nu \psi^b \partial_\nu \phi^i + \dots$$

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$$+ \alpha_{ijk} \overset{1}{\phi^i} \overset{1}{\partial_\mu \phi^j} \overset{1}{\partial^\mu \phi^k} + \dots$$

$$* [g_{ijk}] = 1$$

$$* [\lambda_{jke}] = 0$$

$$[K_{i_1 \dots i_n}] = 4-n$$

$$[\alpha_{ijk}] = -1$$

$$* [y_{iabc}] = 0$$

$$[\varepsilon_{ijkl}] = -1$$

$$[w] = -1$$

$$[l] = -2$$

$$+ \left[y_{iabc} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i + c.c. \right] + \varepsilon_{ijkl} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i \phi^j + \dots$$

$$+ w_{ijkl} \bar{\Psi}^a \gamma_\mu \gamma_\nu \psi^b \partial^\mu \phi^i + \dots$$

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$$[\tau_{ijab}] = -1$$

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$$+ w_{iabc} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\gamma_c} \overset{1}{\psi^d} \overset{1}{\partial_\mu \phi^i} + \dots$$

$$+ l_{abcd} \overset{1}{\bar{\Psi}^a} \overset{1}{\gamma_b} \overset{1}{\psi^c} \overset{1}{\bar{\Psi}^d} \overset{1}{\gamma_e} \overset{1}{\psi^f} + \dots$$

Most general renormalizable interactions:

$\mathcal{L}_{int} =$

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^* \phi + \dots$$

$$\mathcal{L}_{int} = -g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijke} \phi^i \phi^j \phi^k \phi^e + K_{ij\dots n} \phi^i \phi^j \dots \phi^n + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$* [g_{ijk}] = 1$$

$$* [\lambda_{ijke}] = 0$$

$$[K_{ij\dots n}] = 4-n$$

$$[\alpha_{ijk}] = -1$$

$$* [y_{i\mu}] = 0$$

$$[\bar{z}_{ijab}] = -1$$

$$[w] = -1$$

$$[l] = -2$$

$$+ \left[y_{i\mu} \bar{\Psi}^{\frac{1}{2}} \gamma_\mu \psi^{\frac{1}{2}} \phi^i + c.c. \right] + \bar{z}_{ijab} \bar{\Psi}^{\frac{1}{2}} \gamma_\mu \psi^{\frac{1}{2}} \phi^i \phi^j + \dots$$

$$+ w_{icd} \bar{\Psi}^{\frac{1}{2}} \gamma_\mu \gamma_\nu \psi^{\frac{1}{2}} \partial^\mu \phi^i + \dots$$

$$+ l_{abcd} \bar{\Psi}^{\frac{1}{2}} \gamma_\mu \psi^{\frac{1}{2}} \bar{\Psi}^{\frac{1}{2}} \gamma_\nu \psi^{\frac{1}{2}} + \dots$$

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^* \phi + g_{ijk} \phi^i \phi^j \phi^k$$

cubic.

$$+ \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l$$

$$\mathcal{L}_{int} = -g_{ijk} \phi^i \phi^j \phi^k + \lambda_{ijk} \phi^i \phi^j \phi^k + \kappa_{ijk} \phi^i \phi^j \phi^k + \dots$$

$$+ \alpha_{ijk} \phi^i \partial_\mu \phi^j \partial^\mu \phi^k + \dots$$

$$* [g_{ij}] = 1$$

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$$[\tilde{e}_{ijk}] = -1$$

$$[w] = -1$$

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$$- \left[y_{ijk} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i + c.c. \right] + \tilde{e}_{ijk} \bar{\Psi}^a \gamma_\mu \psi^b \phi^i \phi^j + \dots$$

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Most general renormalizable interactions:

$$\mathcal{L} = V(\phi)$$

$$\left[y_{ial} \bar{\psi}^a \gamma_L \psi^b \phi^i + \text{cc} \right]$$

Yukawa

$$V = + \frac{1}{2} \mu^2 \phi^a \phi^a + g_{ijk} \phi^i \phi^j \phi^k$$

cubic.

$$+ \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l$$

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^* \phi + g \phi^* \phi^2 + \lambda \phi^* \phi^3$$

cubic.

$$- \left[y_{ij} \bar{\psi}_i \gamma_5 \psi_j \phi^* + \text{cc} \right]$$

$$+ \lambda_{ijk} \bar{\psi}_i \psi_j \phi^* \phi^k$$

quartic interactions

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^2 + g \phi^3 + \lambda \phi^4$$

cubic.

$$- \left[y_{\text{int}} \bar{\psi} \gamma_5 \psi \phi + \text{cc} \right]$$

$$+ \lambda_{ij} \phi^i \phi^j \phi^k \phi^l$$

"Yukawa" interactions

Most general renormalizable interactions:

$$\mathcal{L}_{int} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^2 + g \phi^3 + \dots$$

cubic.

$$- \left[y_{int} \bar{\psi} \gamma_5 \psi \phi + cc \right]$$

$$+ \lambda_{ij} \phi_i \phi_j \phi_k$$

'Yukawa' interactions

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$$\mathcal{L} = -A - \frac{1}{2} \sum_i (\partial_\mu \phi^i \partial^\mu \phi^i + m_i^2 \phi^i \phi^i) - \frac{1}{2} \sum_i \bar{\psi}^i (\not{\partial} + m_i) \psi^i + \mathcal{L}_{int}$$

$$S = \int d^4x \mathcal{L}$$

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$$[\mathcal{L}] = 4 \text{ = dim of } \mathcal{L} \text{ is (mass)}^4$$

$$[\partial_\mu] = 1$$

$$[\phi^i] = 1$$

$$[\psi^i] = 3/2$$

Spin 10/10/05

Spin Interactions

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^2$$

Spin 1/2 Dirac:

$\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^{\alpha} f^{\mu\nu}_{\alpha}$$

$$f_{\mu\nu}^{\alpha} = \partial_{\mu} A_{\nu}^{\alpha} - \partial_{\nu} A_{\mu}^{\alpha}$$

Spin 1 in $\mathbb{R}^D$:

$\alpha = 1, \dots, D$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}^\alpha} f_{\alpha}^{\mu\nu}$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

Spin in $U(1)$:

$\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^\alpha f_{\mu\nu}^\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

Most general renormalizable interactions:

$$\mathcal{L}_{int} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^2 + g \phi^3 + \frac{\lambda}{4!} \phi^4$$

$$- \left[y_{int} \bar{\psi} \gamma_L \psi \phi + c.c. \right]$$

$$+ \lambda_{\mu} \phi^2 \Box \phi$$

'Yukawa' interactions

$$\phi \phi f_{\mu\nu}$$

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = + \frac{1}{2} \mu^2 \phi^* \phi + g_{ijk} \phi^i \phi^j \phi^k$$

cubic.

$$- \left[y_{i\alpha\beta} \bar{\psi}^\alpha \gamma_L^\mu \psi^\beta \phi^i + \text{cc} \right]$$

'Yukawa' interactions

$$+ \lambda_{jkl} \phi^j \phi^k \phi^l$$

$$c_{ijkl} \partial^\mu \phi^i \partial_\mu \phi^j \phi^k \phi^l$$

Most general renormalizable interactions:

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$V = +\frac{1}{2}\mu^2 \phi^* \phi + g \phi^\dagger \phi^2$$

$$- \left[y_{ij} \bar{\psi}^i \gamma_L \psi^j \phi + \text{c.c.} \right]$$

$$+ \lambda_{ijk} \phi^i \phi^j \phi^k$$

"Yukawa" interactions

Spin interaction: $\alpha = 1, \dots, N_s$

$$\mathcal{L} = -\frac{1}{4} \underline{f}_{\mu\nu}^\alpha \underline{f}_{\mu\nu}^\alpha$$

$$\underline{f}_{\mu\nu}^\alpha = \partial_\mu \underline{A}_\nu^\alpha - \partial_\nu \underline{A}_\mu^\alpha$$

there are no interactions built using only $\underline{f}_{\mu\nu}^\alpha$
for

Spin interaction: $\alpha = 1, \dots, N_s$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}^\alpha} f_{\mu\nu}^\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

there are no interactions built using only $f_{\mu\nu}^\alpha$
for which the couplings have non-neg. dimension.

$$f_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a$$

no terms built using only $f_{\mu\nu}^a$
 the terms have non-zero dimension.

Most general vector

$$\mathcal{L}_{\text{int}} = -V(\phi)$$

$$-[\bar{\psi}_L \gamma^\mu \psi_L]$$

Yukawa

$$c_{\mu\nu} \bar{\psi} \gamma^{\mu\nu} \psi f_{\mu\nu}^a$$

Spin Interp: $\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^\alpha f_{\mu\nu}^\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- are no interactions built using only $f_{\mu\nu}^\alpha$
which the couplings have non-neg dimension.

Ψ_X^μ

Spin in $\psi \bar{\psi} \gamma^\mu \psi$ $\alpha = 1, \dots, 4$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^\alpha f_{\alpha}^{\mu\nu}$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- the no interactions built using only $f_{\mu\nu}^\alpha$
the couplings have non-neg dimension.

$$\bar{\psi} \gamma_\mu \gamma_5 \psi A_\mu^\alpha \text{ has } + \dots$$

Spin Interactions: $\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}^\alpha} f_{\mu\nu}^\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the couplings have non-neg. dimension.

Use $A_{\mu i}$

$$\mathcal{L}_{int} = \underline{\bar{\Psi} \gamma_\mu \gamma_5 \Psi} A_\mu^\alpha h_{\alpha i} + \dots$$

$$[h_{\alpha i}] = 0 \quad [f_{\mu\nu}^\alpha] = 0$$

$$+ \phi^i \partial^\mu \phi^j A_\mu^\alpha f_{\alpha ij} + \dots$$

Spin Interactions: $\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^\alpha f^{\mu\nu}_\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the α 's have non- ∞ dimension.

Use $A_{\mu ij}$

$h_{\mu\nu} + \dots$

$$[h_{\mu\nu}] = 0 \quad [f_{\mu\nu}] = 0$$

A_μ^α

$f_{\mu\nu}^{\alpha\beta}$

Spin in $\gamma_{\mu\nu}$: $\alpha=1, \dots, 4$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}}^\alpha \underline{f_{\mu\nu}}_\alpha \quad f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the couplings have non- neg dimension.

Use $A_{\mu i}$ $\mathcal{L}_{\text{int}} = \bar{\psi} \gamma_\mu \gamma_5 \psi A_\mu^\alpha h_{\alpha i} + \dots$ $[h_{\alpha i}] = 0$ $[f_{\mu\nu}] = 0$

$$+ \phi^i \partial^\mu \phi^j A_\mu^\alpha f_{\alpha ij} + \dots$$

$$+ A_\mu^\alpha A_\nu^\beta A_\rho^\gamma A_\sigma^\delta c_{\alpha\beta\gamma\delta} + AA\partial A$$

Spin 1 in 4D: $\alpha=1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}}^\alpha \underline{f^{\mu\nu}}_\alpha \quad f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the couplings have non-neg dimension.

Use A_μ : $\mathcal{L}_{int} = \bar{\psi} \gamma_\mu \psi A^\mu h_{int} + \dots$

to lorentz invariant

regular by $A \rightarrow A + \partial\omega$

$$+ \phi^i \partial^\mu \phi^j A_\mu^\alpha f_{\alpha ij} + A_\mu^\alpha A_\nu^\beta A_\rho^\gamma A_\sigma^\delta c_{\alpha\beta\gamma\delta} + A A \partial A$$

$$[h_{int}] = 0 \quad [f_{\alpha ij}] = 0$$

Spin 1 in 4D: $\alpha=1, \dots, 4$

$$\mathcal{L}_0 = -\frac{1}{4} f_{\mu\nu}^\alpha f^{\mu\nu}_\alpha$$

$$f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the couplings have non-neg. dimension.

Use A_μ : $\mathcal{L}_{int} = \bar{\psi} \gamma_\mu \gamma_5 \psi A^\mu + \dots$

to Lorentz invariant

for massless fe.

$$+ \phi^i \partial^\mu \phi^j A_\mu^\alpha f_{\alpha ij} + \dots$$

$$+ A_\mu^\alpha A_\nu^\beta A_\rho^\gamma A_\sigma^\delta \epsilon_{\alpha\beta\gamma\delta} C_{\mu\nu\rho\sigma} + A A \partial A$$

$$[h_{\mu\nu}] = 0 \quad [f_{\mu\nu}] = 0$$

Spin 1 fields: $\alpha = 1, \dots, N$

$$\mathcal{L}_0 = -\frac{1}{4} \underline{f_{\mu\nu}^\alpha} f_{\mu\nu}^\alpha \quad f_{\mu\nu}^\alpha = \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha$$

- there are no interactions built using only $f_{\mu\nu}^\alpha$ for which the couplings have non-neg dimension.

Use $A_{\mu i}$ $\mathcal{L}_{int} = \underline{\bar{\Psi} \gamma_\mu \gamma_5 \Psi} A_\mu^\alpha h_{\alpha i} + \dots$

to lorentz invariant

require $\dim A = 1 + \frac{1}{2} \Omega$
for massless case.

$$+ \phi^i \partial^\mu \phi^j A_\mu^\alpha f_{\alpha ij} + A_\mu^\alpha A_\nu^\beta A_\rho^\gamma A_\sigma^\delta c_{\alpha\beta\gamma\delta} + A A \partial A$$

$$[h_{\alpha i}] = 0 \quad [f_{\alpha ij}] = 0$$

Interactions

$\hbar=c=1$
- Spins 0, 1/2:

N_0 spin 0 : ϕ^i
 $N_{1/2}$ spin 1/2 : ψ

$$\mathcal{L} = -A - \frac{1}{2} \sum_i \left(\partial_\mu \phi^i \partial^\mu \phi^i + M_i^2 \phi^i \phi^i \right) - \frac{1}{2}$$

$$S = \int d^4x \mathcal{L}$$

$$[A] = 4$$

$$[\mathcal{L}] = 4 \text{ is dim of } \mathcal{L} \text{ is}$$

$$[\partial_\mu] = 1$$

$$[A_\mu] = [\phi^i] = 1$$

$$\Lambda \approx 10^{16} \text{ GeV}$$

$$+ \mathcal{L}_{int}$$

Interactions

- Spins 0, 1/2:

$\hbar = c = 1$

N_0 spin 0 : ϕ^i (real)

$N_{1/2}$ spin 1/2 : ψ^a (fermions)

$$\Lambda \geq \frac{4M}{\pi^2}$$

$$\mathcal{L} = -A - \frac{1}{2} \sum_i (\partial_\mu \phi^i \partial^\mu \phi^i + \mu_i^2 \phi^i \phi^i) -$$

$$S = \int d^4x \mathcal{L}$$

$$[\mathcal{L}] = 4 \text{ = dim of } \mathcal{L}$$

$$[A] = 4$$

$$[\partial_\mu] = 1$$

$$[A_\mu^\nu] = [\phi^i] = 1$$

$$(\not{\partial} + m_a) \psi^a + \mathcal{L}_{int}$$

$$[\psi^a] = 3/2$$

for 1 meson spin 1: $\vec{S} = \hbar \vec{e}_z$ interaction

$$L_{int} = -V(\vec{r}) \quad V = +\frac{1}{2} \mu_m \vec{\phi} \cdot \vec{b} + \frac{1}{2} \mu_m \vec{\phi} \cdot \vec{e}_z$$

$$\vec{b} = \frac{1}{2} \mu_m \vec{\phi} \cdot \vec{e}_z$$

$$\vec{b} = \frac{1}{2} \mu_m \vec{\phi} \cdot \vec{e}_z$$

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$$\vec{b} = \frac{1}{2} \mu_m \vec{\phi} \cdot \vec{e}_z$$

$$\vec{b} = \frac{1}{2} \mu_m \vec{\phi} \cdot \vec{e}_z$$

for massless spin 1/2 particles in interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

$$V = + \frac{1}{2} \mu^2 \phi^2 + \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \dots$$

for massless spin 1 interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

for massless spin 1: interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{int} = J^\mu(\psi, \phi) A_\mu$

↓ Aside:

for massless spin 1 interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

inside: Suppose we have a system of spin 0 ($\pm 1/2$) fields ϕ , described by $\mathcal{L}(\phi)$

for massless spin 1/2

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

↓ Aside: Suppose we have a system of spin 0 ($\pm 1/2$) fields ϕ^i , described by $\mathcal{L}(\phi^i)$

Suppose $\mathcal{L}$ is unchanged by $\delta \phi^i = f^i(\phi)$

then: $\delta x = 0$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} f^i(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu f^i = 0$ for all ϕ^i

for 1 massless spin 1: massless interactions

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interactions $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

Aside:

we have a system of spin 0 ($\rightarrow 1/2$)
described by $\mathcal{L}(\phi^i)$

$$\delta \phi^i = \epsilon f^i(\phi)$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_\alpha(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i_\alpha \equiv 0$ for all ϕ^i

for massless spin 1: gauge interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

↓ Aside: Suppose we have a system of spin 0 ($\pm \frac{1}{2}$) fields ϕ^i , described by $\mathcal{L}(\phi^i)$.

Suppose $\mathcal{L}$ is unchanged by $\delta \phi^i = \epsilon^a f_a^i(\phi) \quad \partial_\mu \epsilon^a = 0$

$$\text{then: } \delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_\mu(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i_\mu = 0 \text{ for all } \phi^i$$

for all masses, spin 1: negligible interaction

$$A_\mu \rightarrow A_\mu + \partial_\mu \Omega$$

interested in $\mathcal{L}_{\text{int}} = J^\mu(\psi, \phi) A_\mu$

↓ Aside: Suppose we have a system of spin 0 ($\pm 1/2$) fields ϕ^i , described by $\mathcal{L}(\phi^i)$

Suppose $\mathcal{L}$ is unchanged by $\delta \phi^i = \epsilon^i f^i(\phi)$ $\partial_\mu \epsilon^\mu = 0$

$$\mathcal{L} = \partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi^\dagger \phi) \quad \delta \phi = i \epsilon \phi$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_0(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i_0 = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^a is not constant?

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} f_i'(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu f_i = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^a is not constant?

$$\begin{aligned} & \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \\ &= i (\phi^* \partial_\mu \phi - \partial_\mu \phi^* \phi) \partial^\mu \epsilon \end{aligned}$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^a is not constant?

$$\begin{aligned} \mathcal{L} &= \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \\ &= i (\partial_\mu \phi^* \phi - \partial_\mu \phi \phi^*) \partial^\mu \epsilon \end{aligned}$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^a is not constant?

$$\left[\begin{aligned} \delta \mathcal{L} &= \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \\ &= i (\partial_\mu \phi^* \phi - \partial_\mu \phi \phi^*) \partial^\mu \epsilon \end{aligned} \right]$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu \epsilon f^i_\nu = 0$ for all ϵ

Q: What is $\delta \mathcal{L}$ if ϵ^μ is not constant?

$$\left[\delta \mathcal{L} = \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \right. \\ \left. = i (\partial_\mu \phi^* \phi - \partial_\mu \phi \phi^*) \partial^\mu \epsilon \right]$$

$$\delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon^\nu f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu (\epsilon^\nu f^i_\nu) = \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu \epsilon^\nu f^i_\nu$$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i_\nu = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^μ is not constant?

$$\begin{aligned} \delta \mathcal{L} &= \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \\ &= i (\partial_\mu \phi^* \phi - \partial_\mu \phi \phi^*) \partial^\mu \epsilon \end{aligned}$$

$$\begin{aligned} \delta \mathcal{L} &= \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon^\nu f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu (\epsilon^\nu f^i_\nu) = \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu \epsilon^\nu f^i_\nu \\ &= J^\mu_\nu(x) \partial_\mu \epsilon^\nu \quad \text{Noether current} \end{aligned}$$

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

Use A_μ $\mathcal{L}_{int} = \bar{\psi} \gamma_\mu \gamma_5 \psi A^\mu h_{int} + \dots$

$[h_{int}] = 0$ $[f_{ij}] = 0$

to lorentz invariant
require $A \rightarrow A + \partial \Lambda$
for massless case.

$+ \phi^i \partial^\mu \phi^j A_\mu f_{ij} + \dots$
 $+ A_\mu^\alpha A_\nu^\beta A_\rho^\gamma A_\sigma^\delta C_{\alpha\beta\gamma\delta} + A A \partial A$

$$\frac{\partial \mathcal{L}}{\partial \phi^i} \partial_\mu \phi^i = 0 \text{ for all } \phi^i$$

not constant?

$\phi \in \phi^*$

$$\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \partial_\mu \phi^i = f^i$$

$T^\mu_\nu(x) \partial_\mu \phi^i$ Noether current

$$\text{eq. of motion: } \frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0$$

$$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \phi^i \right] = 0$$

Use A_μ $\mathcal{L}_{int} = \bar{\psi} \gamma_\mu \gamma_5 \psi A^\mu + \dots$
 to lorentz invariant
 require sym $A \cdot A + \partial A$
 for massless case.
 $+ \phi^i \partial_\mu \phi^j A^\mu_{ij} + \dots$
 $+ A^\mu_\lambda A^\nu_\lambda A^\rho_\lambda A^\sigma_\lambda C_{\mu\nu\rho\sigma} + A A \partial$

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f_v^i(\phi) + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f_v^i = 0$ for all ϕ^i

Q: What is ϵ if ϵ^μ is not constant?

$$\left[\begin{aligned} \delta \mathcal{L} &= \partial_\mu \phi^i \epsilon^\mu \frac{\partial \mathcal{L}}{\partial \phi^i} \\ &= i \left(\frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon^\mu \partial_\mu \phi^i \right) = \epsilon^\mu \left(\frac{\partial \mathcal{L}}{\partial \phi^i} \partial_\mu \phi^i \right) \end{aligned} \right]$$

$$\delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon^\mu f_v^i$$

$$f_v^i = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \partial_\mu \epsilon^\nu f_v^i$$

$$= J_\nu^\mu \partial_\mu \epsilon^\nu \quad \text{Noether current}$$

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i_\nu \right] = 0$

$\textcircled{J^\mu_\nu}$

Use A_μ $\mathcal{L}_{int} = \bar{\psi} \gamma_\mu \psi A^\mu h_{int} + \dots$ $[h_{int}] = 0$ $[f_{int}] = 0$

to lorentz invariant $+ \phi^i \partial_\mu \phi^j A^\mu_{ij} f_{int} + \dots$

require by $A \rightarrow A + \partial \Lambda$ $+ A^\mu_\lambda A^\lambda_\nu A^\nu_\rho A^\rho_\mu c_{\mu\nu\rho} + A A \partial A$

for massless case.

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi'} = \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi')} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi')} f'_\nu \right] = 0$

$\textcircled{J^\mu_\nu}$

Note: if we write $\mathcal{L} = \mathcal{L}(\phi) - \frac{1}{4} f^\mu_\nu f^\nu_\mu$

$- J^\mu_\nu A_\mu$ Noether current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

$\textcircled{J^\mu_\nu}$

Note: if we write

Claim:

$\mathcal{L}$ is invariant

wrt: $\delta A_\mu = \partial_\mu \Omega(x)$

$\delta \phi^i = \Omega f^i(x)$

$\delta \mathcal{L} = 0$

$J^\mu_\nu A_\mu$

Noether current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

Claim:

$\mathcal{L}$ is invariant

wrt: $\delta A_\mu = \partial_\mu \Omega(x)$

$\delta \phi^i = \Omega f^i(x)$

provided $\delta J_\mu = 0.$

Not! we write $\mathcal{L} = \mathcal{L}(x) - \frac{1}{4} f_{\mu\nu}^a f^{\mu\nu a}$

$- J_\mu^a A_\mu^a$

Worther current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

$\textcircled{J^\mu_\nu}$

Claim:

$\mathcal{L}$ is invariant

$J_\mu = \partial_\mu \Omega_{ij}$

$\phi^i = \Omega f^i(x)$

total $\delta J_\mu = 0.$

Note: if we write $\mathcal{L} = \mathcal{L}$

rather current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

$\textcircled{J^\mu_\nu}$

Claim:

$\mathcal{L}$ is invariant

wrt: $\delta A_\mu = \partial_\mu \Omega(x)$

$\delta \phi^i = \Omega f^i(x)$

provided $\delta J_\mu = 0.$

Notice: if we write $\mathcal{L} = \underline{\mathcal{L}(x)} - \frac{1}{4} f^{\mu\nu} f_{\mu\nu}$

$- J^\mu_\nu \textcircled{A_\mu}$

Wortler current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

$\textcircled{J^\mu_\nu}$

Claim:

$\mathcal{L}$ is invariant "Gauging the symmetry"

wt: $\delta A_\mu = \partial_\mu \Omega$

$\delta \phi^i = \Omega f^i(x)$

provided $\delta J_\mu = 0$.

Notice: if we write $\mathcal{L} = \mathcal{L}(x) - \frac{1}{4} f^\mu_\nu f^\nu_\mu$

$- J^\mu_\nu \textcircled{A_\mu}$

Weyler current

eq. of motion: $\frac{\partial \mathcal{L}}{\partial \phi^i} - \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \right] = 0.$

$\rightarrow \partial_\mu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} f^i \right] = 0$

$\textcircled{J^\mu_\alpha}$

Claim:

$\mathcal{L}$ is invariant Gauging the symmetry

wt: $\delta A_\mu = \partial_\mu \Omega$

$\delta \phi^i = \Omega f^i(x)$

provided $\delta J_\mu = 0.$

Note: if we write $\mathcal{L} = \mathcal{L}(x) - \frac{1}{4} f^{\mu\nu} f_{\mu\nu}$

$- J^\mu_\alpha \textcircled{A_\mu}$

Wether current

then: $\delta \mathcal{L} = 0 = \frac{\partial \mathcal{L}}{\partial \phi^i} \epsilon f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \epsilon \partial_\mu f^i_\nu = 0$ for all ϕ^i

Q: What is $\delta \mathcal{L}$ if ϵ^μ is not constant?

$$\begin{aligned} \delta \mathcal{L} &= \partial_\mu \phi^* [i \partial_\mu \epsilon \phi] - \partial_\mu \phi [i \partial_\mu \epsilon \phi^*] \\ &= i (\partial_\mu \phi^* \phi - \partial_\mu \phi \phi^*) \partial^\mu \epsilon \end{aligned}$$

$$\begin{aligned} \epsilon^\nu f^i_\nu + \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^i} \partial_\mu (\epsilon^\nu f^i_\nu) &= \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^i)} \partial_\mu \epsilon^\nu f^i_\nu \\ &= J^\mu_\nu(x) \partial_\mu \epsilon^\nu \quad \text{Noether current} \end{aligned}$$