

Title: Little String Theory From a Double Scaled Matrix Model

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URL: <http://pirsa.org/06100049>

Abstract:

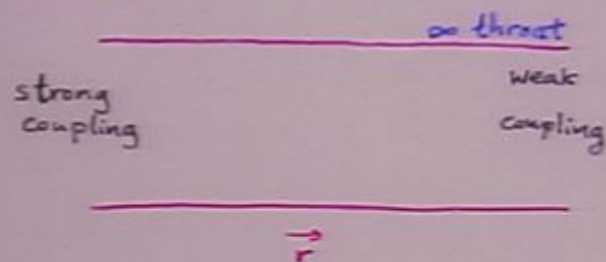


N_5 IA NS5-branes

Type IIA LST:

- defined as worldvolume dynamics of IIA NS5-branes for $g_s \rightarrow 0$ or fixed
- Hagedorn density of states, flows to $(0,2)$ CFT in IR
- No direct Lagrangian description (nonlocal)

Gravity Dual: near-horizon NS5-brane solution



$$ds^2 = N_5 \alpha' [-dt^2 + d\vec{x}_5^2 + dr^2 + d\Omega_3^2]$$

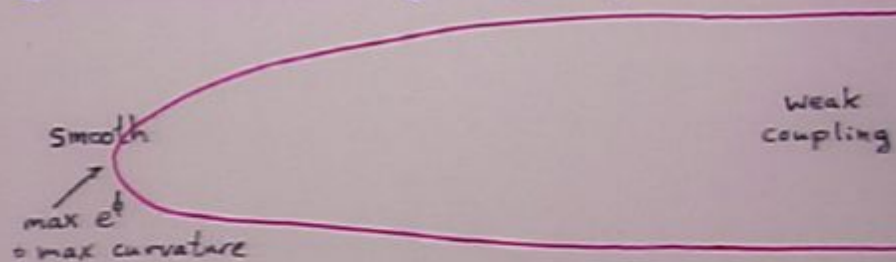
$$e^\phi = g_s e^{-r} \quad \text{linear dilaton, strongly coupled}$$

+ H flux for $r \rightarrow -\infty$

BOTH PICTURES DIFFICULT.

THIS TALK: LST on S^5

Lin + Maldacena: gave gravity solution for LST on S^5



Asymptotically:

$$ds^2 = N_5 \alpha' [2r(-dt^2 + d\Omega_3^2) + dr^2 + d\Omega_2^2]$$
$$e^\phi = g_s e^{-r}$$

+ H flux

- tunable maximum curvature, $R_{max} \alpha' \sim \frac{1}{N_5}$
- tunable maximum dilaton

MORE TRACTABLE

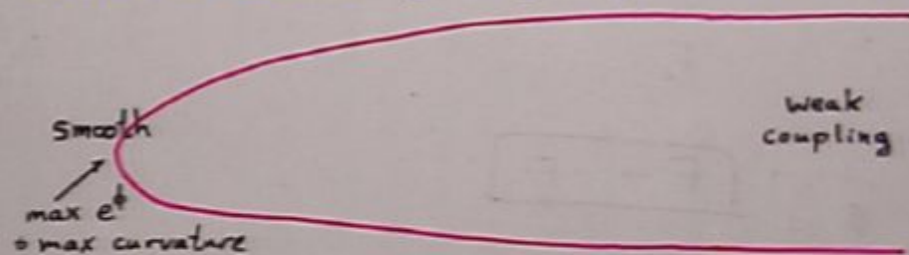
FIELD THEORY: $\text{LST on } S^5$

GOAL: Determine field theory dual

- Explicit Lagrangian description of LST on S^5

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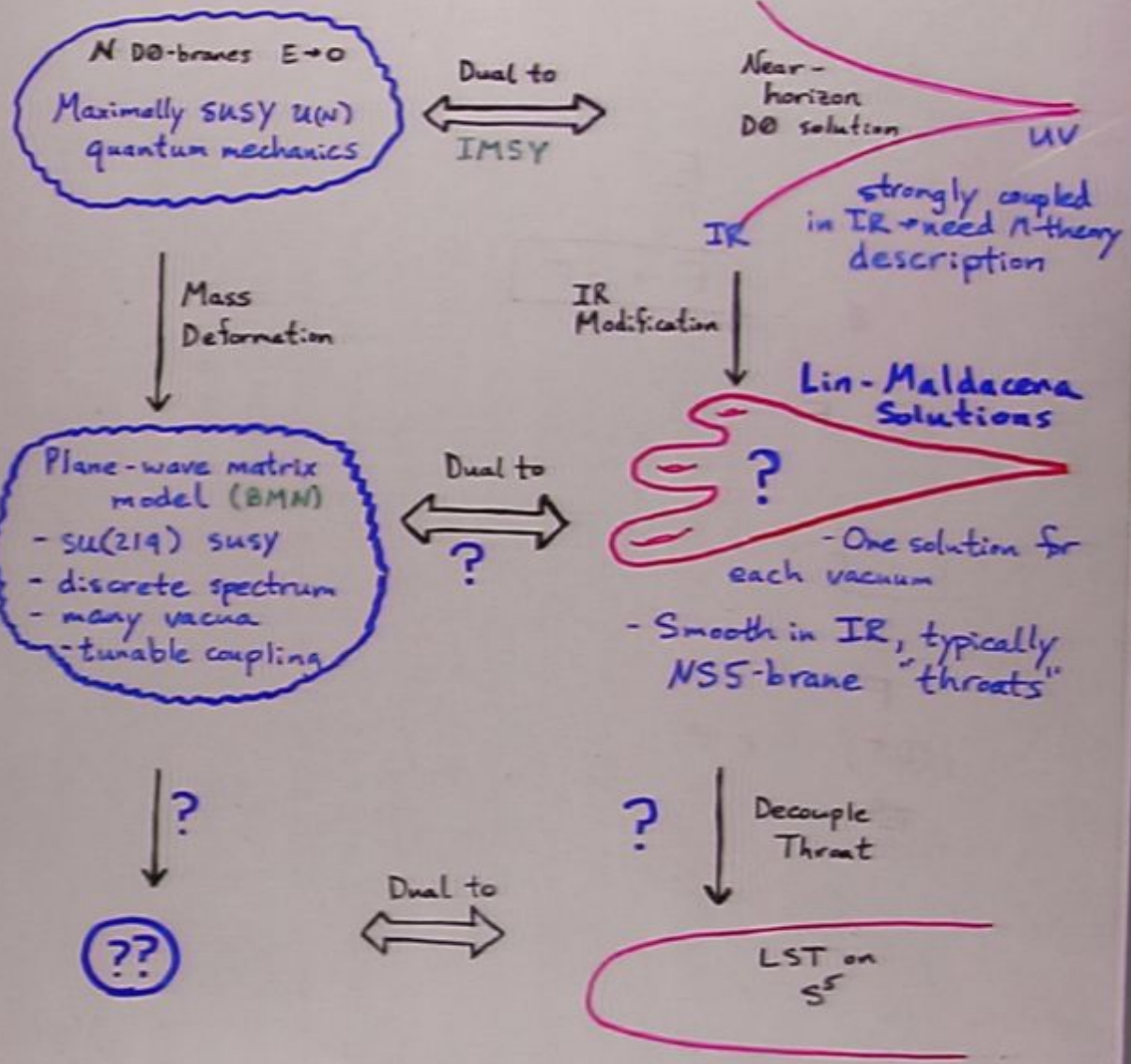
MORE TRACTABLE

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ASPECT FOR D0-BRANE



AdS/CFT FOR D0-BRANES

N D0-branes $E \rightarrow 0$
 Maximally SUSY $U(N)$
 quantum mechanics

Dual to
 \longleftrightarrow
 IMSY

Near-horizon
 D0 solution \rightarrow UV

strongly coupled
 in IR \rightarrow need M-theory
 description

Mass
 Deformation
 \downarrow

IR
 Modification
 \downarrow

Plane-wave matrix
 model (BMN)
 - $SU(2|1|4)$ susy
 - discrete spectrum
 - many vacua
 - tunable coupling

Dual to
 \longleftrightarrow
 ?

Lin-Maldacena
 Solutions
 ?
 - One solution for
 each vacuum

- Smooth in IR, typically
 NS5-brane "throats"

?
 \downarrow

?
 \downarrow Decouple
 Throat

??

Dual to
 \longleftrightarrow

LST on
 S^5

PLANE WAVE MATRIX MODEL

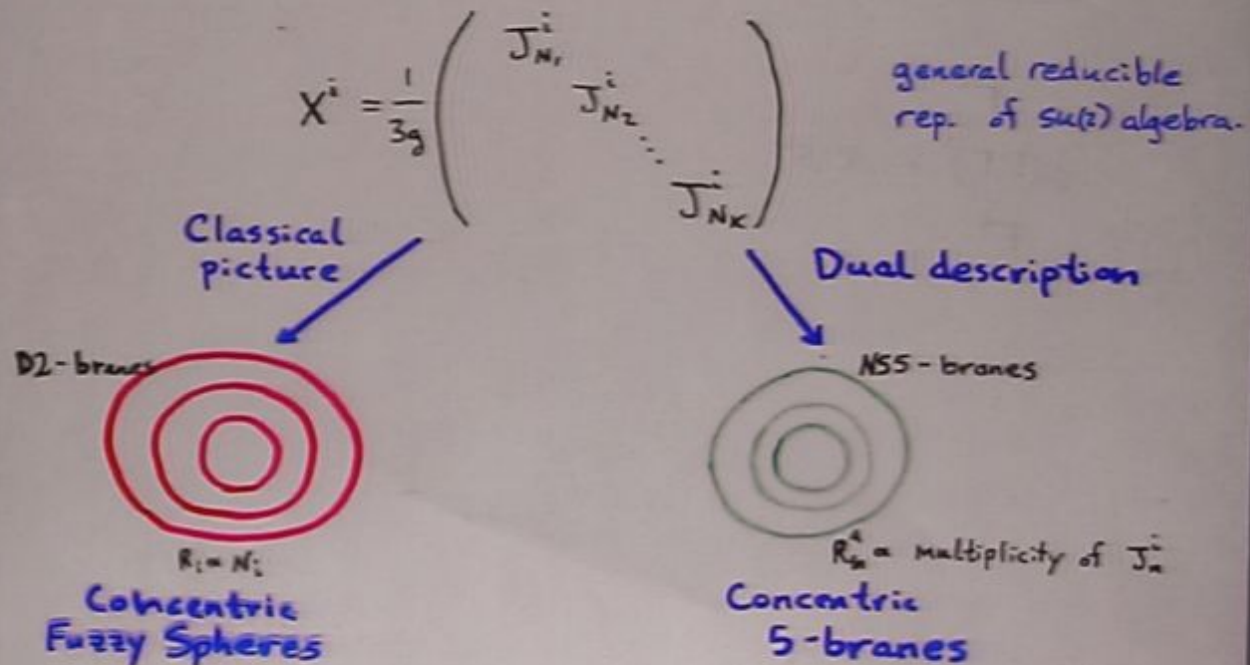
- Mass deformation of 0+1 maximal SYM preserving
32 susys \supset $SU(2|4)$:

$$H = \text{Tr} \left(\frac{1}{2} P_A^2 + \frac{1}{2} (X_i/3)^2 + \frac{1}{2} (X_a/6)^2 + \frac{i}{8} \Psi^T \gamma^{123} \Psi \right. \\ \left. + \frac{i}{3} g \epsilon^{ijk} X^i X^j X^k - \frac{g}{2} \Psi^T \gamma^4 [\chi_a, \Psi] - \frac{g^2}{4} [X_a, X_b]^2 \right)$$

$A=1-9$ $i=1-3$ $a=4-9$

Parameters: N (size of matrices)
 g tunable coupling

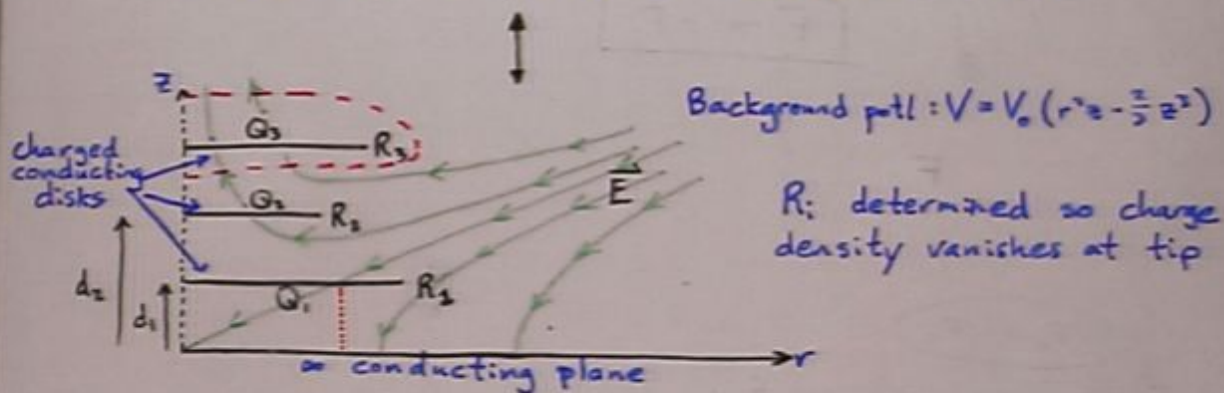
Vacua: require $[X^i, X^j] = \frac{i}{3g} \epsilon^{ijk} X^k$ $X^a = 0$



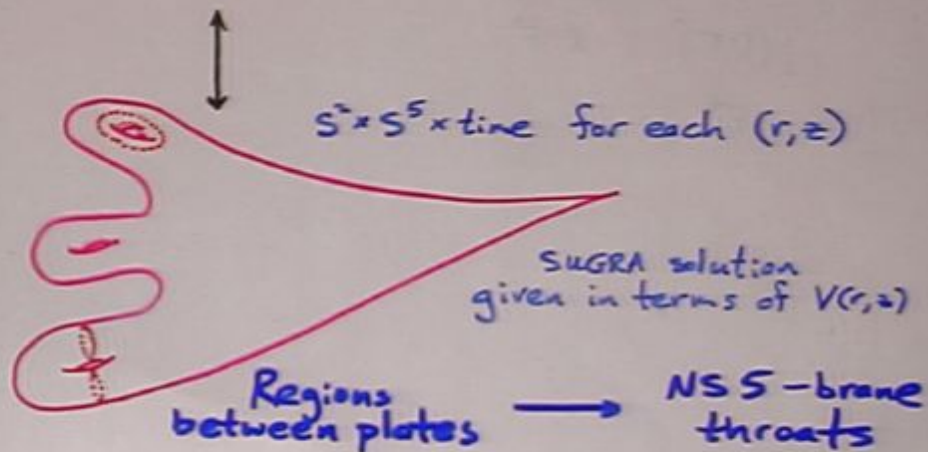
LIN-MALDALENA GEOMETRIES

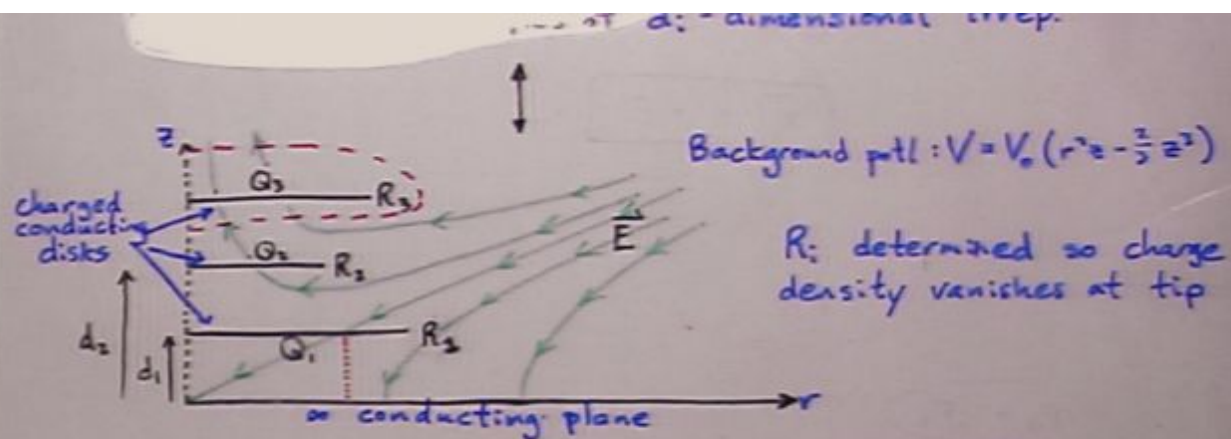
- Supergravity (IIA) solution dual to each vacuum given in terms of solution to auxiliary electrostatics problem:

vacuum with Q_i copies of d_i -dimensional irrep.

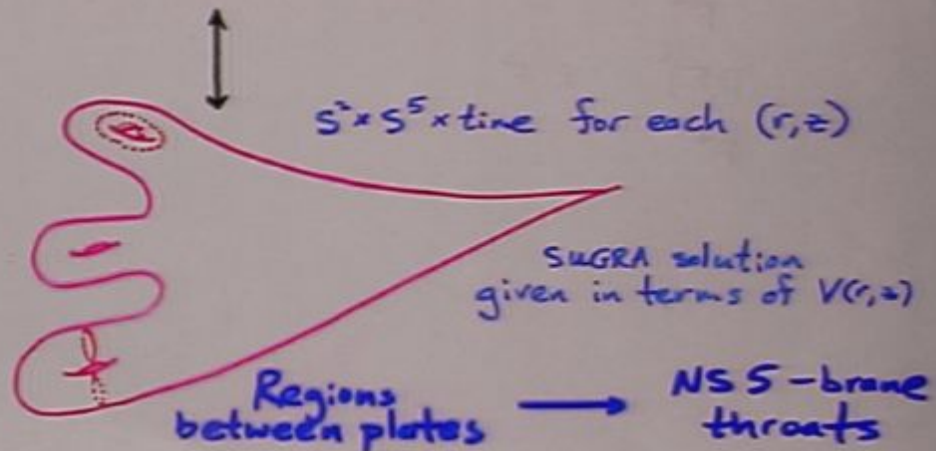


Find $V(r, z)$ for given Q_i, d_i, V_0

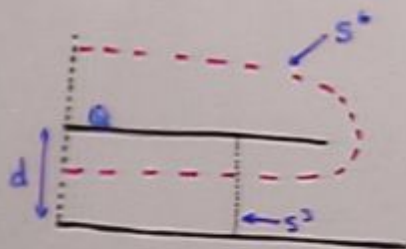




Find $V(r, z)$ for given Q_i, d_i, V_0



Lin-Maldacena: Single throat geometry has d/Q units of NS5/D2 flux through S^3/S^6 .



recall: $X^i = \frac{1}{3g} \begin{pmatrix} J_d \\ \uparrow Q \\ J_d \end{pmatrix}$ associated with $\begin{cases} Q \text{ D2-branes} \\ d \text{ NS5-branes} \end{cases}$

Remaining parameter: V_0

scaling this doesn't change metric, takes

$$e^\phi \rightarrow \frac{V_0}{\sqrt{V_0}} e^\phi$$

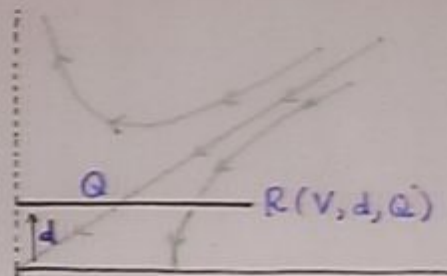
Associate $V_0 \sim \frac{1}{g_s} \sim \frac{1}{g^2}$

Simplest case: single throat



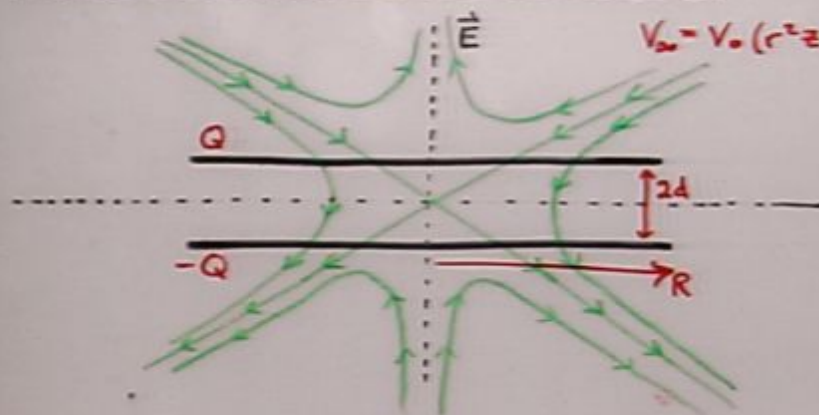
$$X^i = \frac{1}{3g} \begin{pmatrix} J_d^i & & & \\ & \ddots & & \\ & & Q & \\ & & & J_d^i \end{pmatrix}$$

Q copies
of
 d dimensional
irrep.



Single charged
disk

THE ELECTROSTATICS PROBLEM: SIMPLEST CASE



$$V_0 = V_0(r^2 z - \frac{z^3}{3})$$

$V_0 = 0$: Maxwell
 Clausius
 Helmholtz
 Kirchhoff
 Polya-Seego
 Nicholson
 Love 1949

Solution: $V = V_0 R^3 \phi_{\frac{d}{R}}(\frac{r}{R}, \frac{z}{R})$

solve for ϕ with fixed potential Δ on plates.

$$\phi_k(r, z) = \frac{1}{\sigma} (\Delta + \frac{2}{3} k^2) \int_{-1}^1 dt f(t) \left\{ \frac{1}{\sqrt{r^2 + (z + k + kt)^2}} - \frac{1}{\sqrt{r^2 + (z - d + kt)^2}} \right\}$$

$$f(t) - \int_{-1}^1 dx K(t, x) f(x) = 1 - \frac{2k}{\Delta + \frac{2}{3} k^2} t^2$$

$$K(t, x) = \frac{k/\pi}{k^2 + (x - t)^2}$$

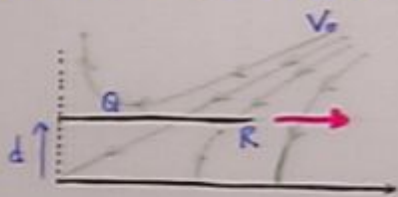
CONSTRAINT:

$\sigma \rightarrow 0$ at edge $\Rightarrow f(1) = 0$

FIXES $\Delta(k)$ (potl on plates)
 \downarrow
 FIXES $R(V_0, Q, d)$

DECOUPLING LIMIT

Want infinite throat in geometry



Need $R \rightarrow \infty$, d fixed

Scale V_0 to get non-trivial solution between plates

$$V = \frac{1}{g_s} \sin\left(\frac{\pi z}{d}\right) I_0\left(\frac{\pi r}{d}\right)$$

Requires: $Q \rightarrow \infty$ d fixed $V_0 \rightarrow \frac{1}{g_s} (Rd)^{-\frac{3}{2}} e^{\frac{\pi R}{d}}$
 $R \rightarrow \infty$

Field theory: $N = Qd \rightarrow \infty$ d fixed = # fivebranes
 $\frac{1}{g^2} \lambda^{\frac{3}{8}} e^{-\alpha \lambda^{\frac{1}{4}} / d^{\frac{5}{4}}}$ fixed

★ LST for d NS5-branes on S^5 described by large N limit with $\lambda \sim d^5 \ln^4(N)$ ★

THE DOUBLE-SCALING LIMIT

Usual 'tHooft limit: $N \rightarrow \infty$ λ fixed

- only planar diagrams remain
- free strings in bulk

Our limit: $N \rightarrow \infty$ $\lambda \sim d^2 \ln^2(N)$

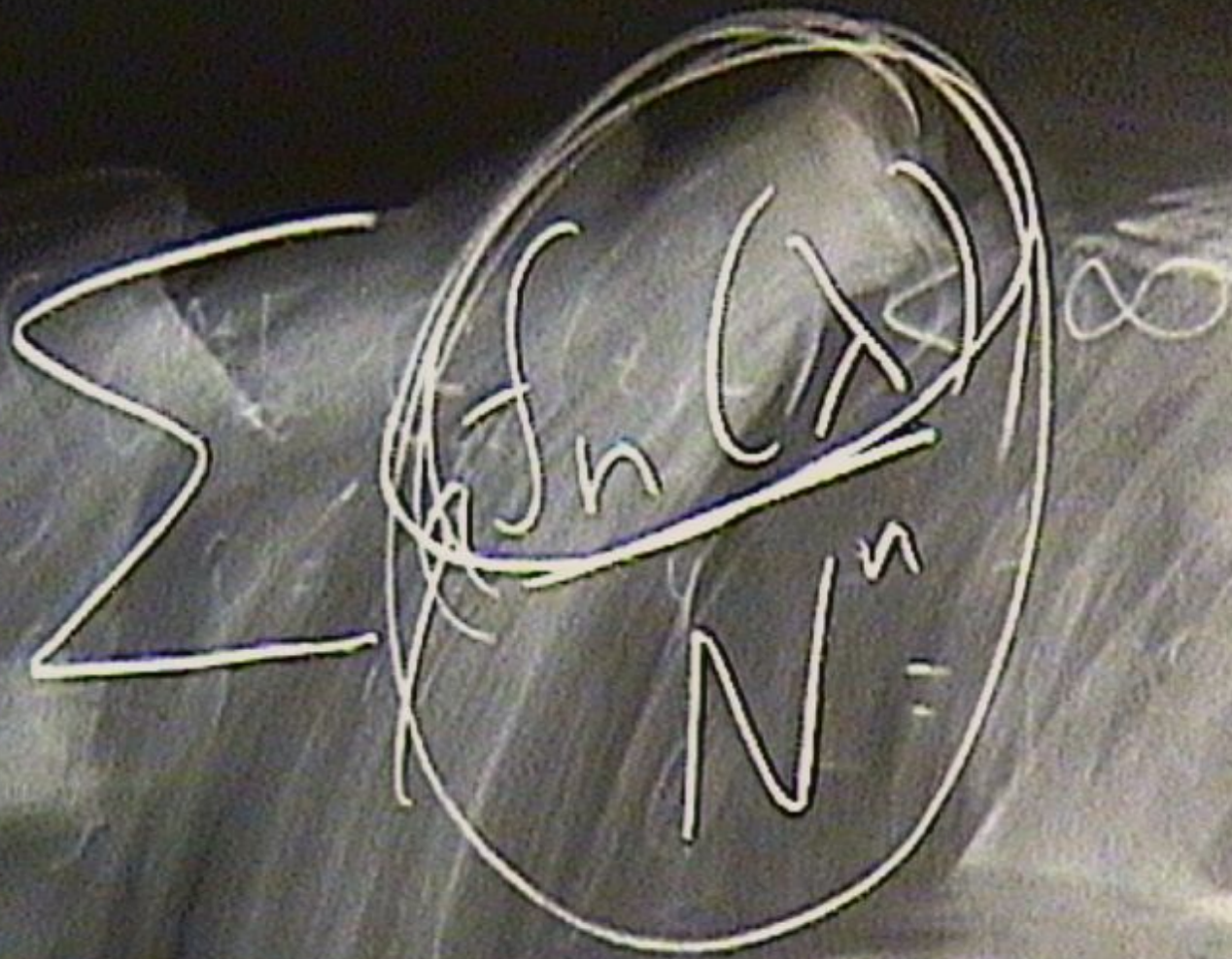
- $g_s \rightarrow 0$ asymptotically but not at $r=0$
- \therefore interacting strings \rightarrow all terms in genus expansion must contribute

PREDICTS

$$f_n \sim c_n \left[\lambda^{\frac{d-2}{2}} e^{\frac{a\lambda^{\frac{d-2}{2}}}{d}} \right]^n \quad \text{as } \lambda \rightarrow \infty$$

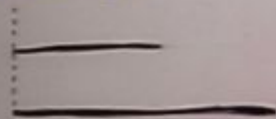
\uparrow
sum of all genus n diagrams for a physical observable

$$\sum_{n=1}^{\infty} \frac{1}{n^2} = \frac{\pi^2}{6}$$



LIMITS OF PWMM

① Finite g, N



IIA strings / M theory on LM geometries



② M-theory limit

$$N \rightarrow \infty$$

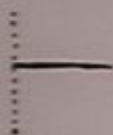
$$\lambda \sim N^4$$

M-theory on max. susy plane-wave

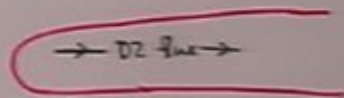
③ D2-brane limit

$$N \rightarrow \infty$$

$$\lambda \sim N^2 \quad Q \text{ fixed}$$



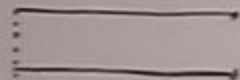
Max. susy $U(Q)$ gauge theory on S^2



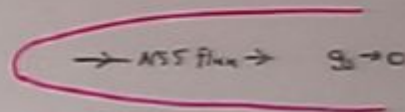
④ NS5-brane limit

$$N \rightarrow \infty$$

$$\lambda \sim \ln^4(N) \quad d \text{ fixed}$$



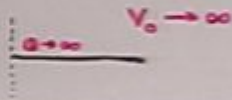
LST on S^5



⑤ 't Hooft limit

$$N \rightarrow \infty$$

$$\lambda \text{ fixed}$$



Free strings on LM geometry