

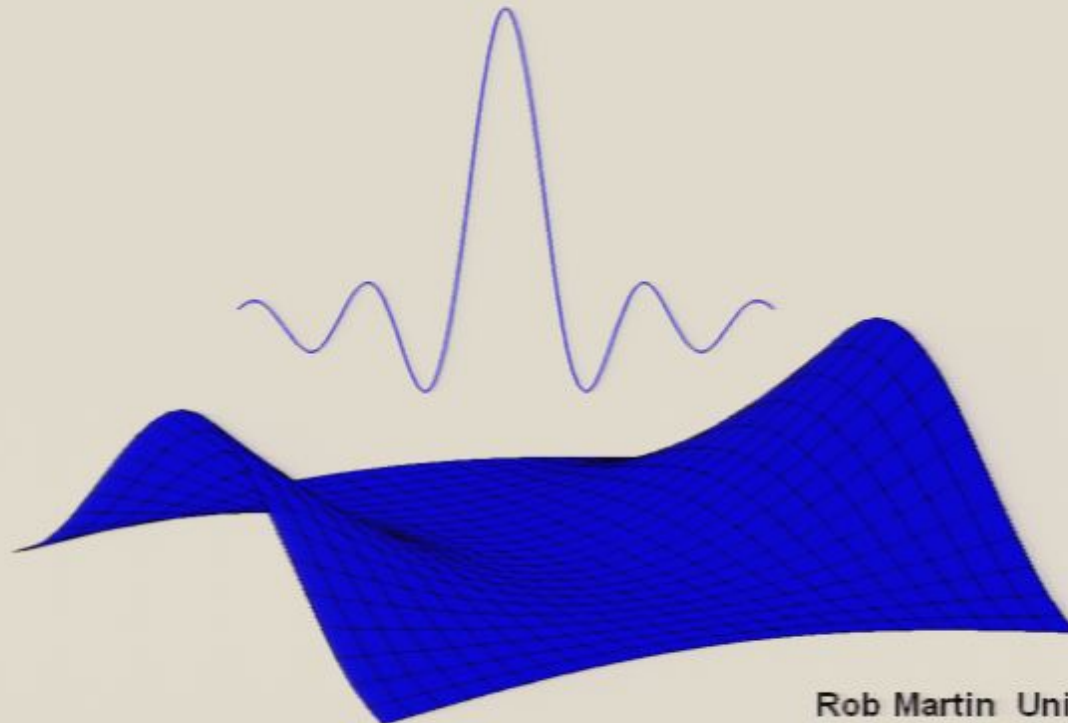
Title: Towards a generally covariant cutoff using sampling theory

Date: Sep 06, 2006 04:30 PM

URL: <http://pirsa.org/06090028>

Abstract:

A covariant UV cutoff: Sampling Theory on curved spacetime



Outline

1. **Motivation:** generally covariant UV cutoff using Sampling Theory
2. **Introduction to Sampling Theory:** Bandlimited functions
3. **Sampling Theory on Space-time:**
 - covariantly bandlimited functions
 - Minkowski spacetime
 - preliminary results on Sampling Theory in De Sitter spacetime
4. **Summary & Outlook**

1. Motivation: the problem

- Quantum gravity arguments \rightarrow Expect an UV cutoff
- UV divergencies \rightarrow QFT ill-defined on continuous curved spacetime
- QFT becomes well defined if spacetime is discrete
- GR lives on a smooth, continuous manifold, discreteness of spacetime violates manifold symmetries like Lorentz invariance
- Q: How to satisfy both?

1.

Motivation: A proposed solution

- **Idea:** Spacetime continuous but physical fields possess a finite spatial density of degrees of freedom.
 - path integral is over a countable number of spatial degrees of freedom
 - The field values $\{\phi(t, x_n)\}_{t \in R}$ on any 'sufficiently dense' discrete set of points $\Lambda = \{x_n\}_{n \in Z}$ determine ϕ uniquely & can reconstruct $\phi(t, x)$ perfectly from this info.
 - Any physical theory can be defined on the smooth manifold, or on any lattice of sufficient density. No preferred lattice → spacetime symmetries preserved
- Q:** Is there a theory that describes classes of fields/functions with these special properties?
- A:** Yes, its called Sampling Theory

Introduction to Sampling Theory

Sampling theory: study of classes of functions, reconstructible from their values taken on certain discrete sets of points

- **Example:** Ω -Bandlimited functions
 - Fourier transforms of elements of $L^2[-\Omega, \Omega]$
 - Ω called the bandlimit
- **Resources:** \exists large body of literature:
 - Mathematics
 - communication engineering
 - Information theory

Bandlimited functions

Example: $B(\Omega) :=$ Hilbert space of functions bandlimited by Ω :

$$f(x) = \int_{-\Omega}^{\Omega} F(w) e^{iwx} dw \quad F \in L^2(-\Omega, \Omega)$$

- Resolution of Identity: $1 = \sum_n (\cdot, b_{x_n}) b_{x_n}$

$$b_{x_n} := \frac{1}{\sqrt{2\Omega}} e^{ix_n w} \quad w \in [-\Omega, \Omega] \quad x_n := \frac{n\pi}{\Omega}$$



- The Shannon Sampling Formula:

$$F = \sum_n (F, b_{x_n}) b_{x_n}$$

$$(F, b_{x_n}) = \frac{1}{\sqrt{2\Omega}} \int_{-\Omega}^{\Omega} F(w) e^{iwx_n} dw = \frac{1}{\sqrt{2\Omega}} f(x_n)$$

$$f = \mathcal{F}^{-1} F = \sum_n (F, b_{x_n}) \mathcal{F}^{-1} b_{x_n}$$

$$\mathcal{F}^{-1} b_{x_n} = 2\Omega \frac{\sin \Omega(x - x_n)}{\Omega(x - x_n)}$$

$$\Rightarrow f(x) = \sum_n f(x_n) \frac{\sin \Omega(x - x_n)}{\Omega(x - x_n)}$$

Sampling Lattices for $B(\Omega)$

- **Sampling Lattice:** a set of points $\Lambda = \{y_n\}$, $|y_n - y_m| \geq \epsilon > 0$ s.t. there is a linear map $R : B(\Omega) \rightarrow l^2(Z)$, $Rf = \{f(y_n)\}_{n \in Z}$ which is bounded above and below.
 - E.g. $\Lambda = \{x_n = \frac{n\pi}{\Omega}\}$
- **Any lattice of sufficient density will do:**

$n^-(r) :=$ min number of y_n in any interval of length r

$$D^-(\Lambda) := \lim_{r \rightarrow \infty} \frac{n^-(r)}{r}$$

Theorem:

(a) if Λ is a set of sampling $D^-(\Lambda) \geq \frac{\Omega}{\pi}$ (Beurling)

(b) Λ is a set of sampling if $D^-(\Lambda) > \frac{\Omega}{\pi}$ (Duffin & Schaeffer / Beurling)

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Sampling Lattices for $B(S)$

- $B(S) :=$ Functions frequency limited by compact $S \subset \mathbb{R}^n$

$$f(x) = \int_S F(w) e^{iwx} dw \quad F \in L^2(S)$$

- **Theorem:** (*Landau*)

If Λ is a set of sampling for $B(S)$ then $D^-(\Lambda) \geq \frac{\mu(S)}{(2\pi)^n}$

minimum possible sample density \propto bandwidth volume

Sufficiency conditions not known in general

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A generally covariant bandlimit

- A bandlimit is a cutoff on the spectrum of the Laplacian in $L^2(M)$ on the manifold $M = \mathbb{R}$.

$$\Delta := \frac{d^2}{dx^2} \quad \frac{d^2}{dx^2} e^{\pm iwx} = w^2 e^{\pm iwx} \quad \forall w \in \mathbb{R}$$

- $B(\Omega) :=$ subspace spanned by $e^{\pm iwx}$ for $|w| \leq \Omega$

$$f(x) = \int_{-\Omega}^{\Omega} F(w) e^{iwx} dw \quad \Leftrightarrow \quad f \in B(\Omega)$$

- Definition:** $B(M, \Omega) :=$ subspace of $L^2(M)$ spanned by eigenfunctions to Δ (or \square) whose eigenvalues λ obey $|\lambda| \leq \Omega^2$

($\Omega :=$ the bandlimit)

A generally covariant UV cutoff

- **Idea:** Impose this generally covariant bandlimit / UV cutoff on physical fields on spacetime.
- **i.e.** Given spacetime M , we restrict the space of allowed physical fields to be $B(M, \Omega)$.
- This is a restriction on the set of fields that the path integral runs over
- **Investigate:** as in the simple case $M = \mathbb{R}$, does the subspace of fields $B(M, \Omega)$ have a countable (or even finite?) # of spatial degrees of freedom

3.

Sampling Theory on Spacetime

Sampling in Minkowski spacetime

- D'Alembertian: $-\frac{\partial^2}{\partial t^2} + \Delta$ Eigenfunctions: $e^{i(p^0 t - \mathbf{p} \cdot \mathbf{x})}$

Spectral cutoff: $|(p^0)^2 - \mathbf{p}^2| \leq \Omega^2$ Fix $p^0 = \Omega b$

$$\Omega\sqrt{1+b^2} \geq |\mathbf{p}| \geq \Omega\sqrt{b^2-1}$$

min. possible sample density \propto bandwidth volume (Landau's thm)

	$b \leq 1$	$b > 1$
1D	$2\Omega\sqrt{1+b^2}$	$2\Omega\left(\sqrt{1+b^2} - \sqrt{b^2-1}\right)$
2D	$\pi\Omega^2(1+b^2)$	$2\pi\Omega^2$
3D	$\frac{4\pi}{3}\Omega^3(1+b^2)^{\frac{3}{2}}$	$\frac{4\pi}{3}\Omega^3\left((1+b^2)^{\frac{3}{2}} - (b^2-1)^{\frac{3}{2}}\right)$



Reconstruction in flat space-time

- Note: in 1+1, 1+2, 1+3 D, each fixed temporal mode obeys a spatial sampling theorem and vice versa
- E.g. $\phi \in B(M, \Omega)$ $M := 1+1$ D Minkowski spacetime.
- $\Phi :=$ temporal Fourier transform of ϕ
- $p^0 = \Omega b$ fixed temporal frequency
- $\Phi(p^0, x)$ (spatial) frequency limited by

$$S_b = \left[-\Omega\sqrt{b^2 + 1}, \Omega\sqrt{b^2 + 1} \right] \quad b^2 \leq 1$$

$$S_b = \left[-\Omega\sqrt{b^2 + 1}, -\Omega\sqrt{b^2 - 1} \right] \cup \left[\Omega\sqrt{b^2 - 1}, \Omega\sqrt{b^2 + 1} \right] \quad b^2 \geq 1$$

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Reconstruction in flat space-time II

- Maximum bandwidth volume: $2\sqrt{2}\Omega$
 - There could be a set $\Lambda := \{y_n\}_{n=-\infty}^{\infty}$ satisfying Landau's thm and is sampling for every $B(S_b) \forall b$ i.e. for every fixed temporal mode
- One such a set is $\Lambda := \{y_n\}_{n=-\infty}^{\infty}$ where $y_n := \begin{cases} \frac{n\pi}{\sqrt{2}\Omega} & n \text{ even} \\ \frac{(n+\alpha)\pi}{\sqrt{2}\Omega} & n \text{ odd} \end{cases} \alpha \in (0, 1)$
- Every temporal mode $\Phi(p^0, x)$ can be reconstructed perfectly from $\{\Phi(p^0, y_n)\}_{y_n \in \Lambda}$
 - \Rightarrow every $\phi \in B(M, \Omega)$ can be reconstructed from $\{\phi(t, y_n)\}_{y_n \in \Lambda} t \in \mathcal{R}$
- Similarly every $\phi \in B(M, \Omega)$ can be reconstructed from $\{\phi(t_n, x)\}_{t_n \in \Lambda} x \in \mathcal{R}$

Reconstruction in flat space-time III

- In 1+3 dimensions, such a spatial sampling theorem is not possible as there is no upper bound on the bandwidth volume for a fixed temporal mode.
- Can show that \exists sets $\Lambda := \{x_n^1, x_m^2\}$ s.t. the information $\{\phi(t, x_n^1, x_m^2)\}$, $\{\phi(t, x_n^1, x_m^2, x^3)\}$ is sufficient to reconstruct $\phi(t, x^1, x^2)$, $\phi(t, x^1, x^2, x^3)$ everywhere.
- Shows that 2D surfaces have a finite density of degrees of freedom, good for holography
- Expect in more general 1+1 or 1+2 D spacetimes it may be possible to reconstruct $\phi \in B(M, \Omega)$ from its values taken on certain discrete sets of timelike curves.

Sampling theory in FRW Spacetimes

- For simplicity, consider 1+1 D. $ds^2 = -dt^2 + a^2(t)dx^2 = a^2(\eta) (-d\eta^2 + dx^2)$
 $\eta :=$ conformal time
- D'Alembertian: $\Delta = a^{-2}(\eta) (-\partial_\eta^2 + \partial_x^2) = -a^{-2}(\eta) (\partial_\eta^2 + k^2)$
- Expect as in Minkowski case, that for each fixed spatial mode 'k' there will be a temporal sampling theorem.
- E.g. De Sitter with a finite ending time: $a(t) = e^t \quad t \in (-\infty, 0]$

$$\eta = e^{-t} \quad \eta \in [1, \infty)$$

$$\Delta = -\eta^2 (\partial_\eta^2 + k^2) \quad L^2([1, \infty) \times (-\infty, \infty); \eta^{-2} d\eta dk)$$

De Sitter with finite end time

$$a(t) = e^t \quad t \in (-\infty, 0] \quad \eta = e^{-t} \quad \eta \in [1, \infty)$$

$$\Delta = -\eta^2 (\partial_\eta^2 + k^2) \quad L^2([1, \infty) \times (-\infty, \infty); \eta^{-2} d\eta dk)$$

- Spectral cutoff (bandlimit): $|\lambda| \leq \Omega^2 \quad \lambda \in \sigma(\Delta)$
- For each fixed 'k' this implies a cutoff on $\Delta_k = -\eta^2 (\partial_\eta^2 + k^2)$ in $L^2([1, \infty), \eta^{-2} d\eta)$
- **Preliminary results:**
 - The subspace of $L^2([1, \infty), \eta^{-2} d\eta)$ spanned by eigenfunctions $\{f_j^{(k)}\}_{j=1}^{N_k}$ to Δ_k obeying the bandlimit has finite dimension N_k for each fixed k.

\Rightarrow Any set $\{\eta_j\}_{j=1}^{N_k}$ of points s.t. the determinant $\det[S] \neq 0$ is a set of sampling for the fixed spatial mode k.

$$S_{jn} := f_j^{(k)}(\eta_n)$$

$$\phi(\eta, k) = \sum_{n=1}^{N_k} c_n f_n^{(k)}(\eta) \quad c_n = \sum_{j=1}^{N_k} \phi(\eta_j, k) S_{kn}^{-1}$$

De Sitter with finite end time

- **Further questions:**
 - How does the dimension N_k behave as a function of k , or of the end time of the manifold
 - We conjecture an overall temporal sampling formula that reconstructs $\phi(\eta, x)$ from $\{\phi(\eta_n, x)\}_{n \in Z}$
 - What about the full De Sitter spacetime $t \in (-\infty, \infty)$
- Some techniques used so far:
 - complex analysis
 - operator theory (self adjoint extensions of symmetric operators, spectral theory)

Outlook

Future Work:

- would like to develop a full understanding of reconstruction properties of $B(M, \Omega)$ in physically relevant spacetimes: De Sitter, Power law, Schwarzschild, etc.
- spatial sampling theorems for fixed temporal frequencies, higher dimensions
- given a general pseudo-Riemannian manifold, can we develop general sufficiency conditions on a discrete set of spatial hypersurfaces s.t. bandlimited fields can be reconstructed from their values on those surfaces.