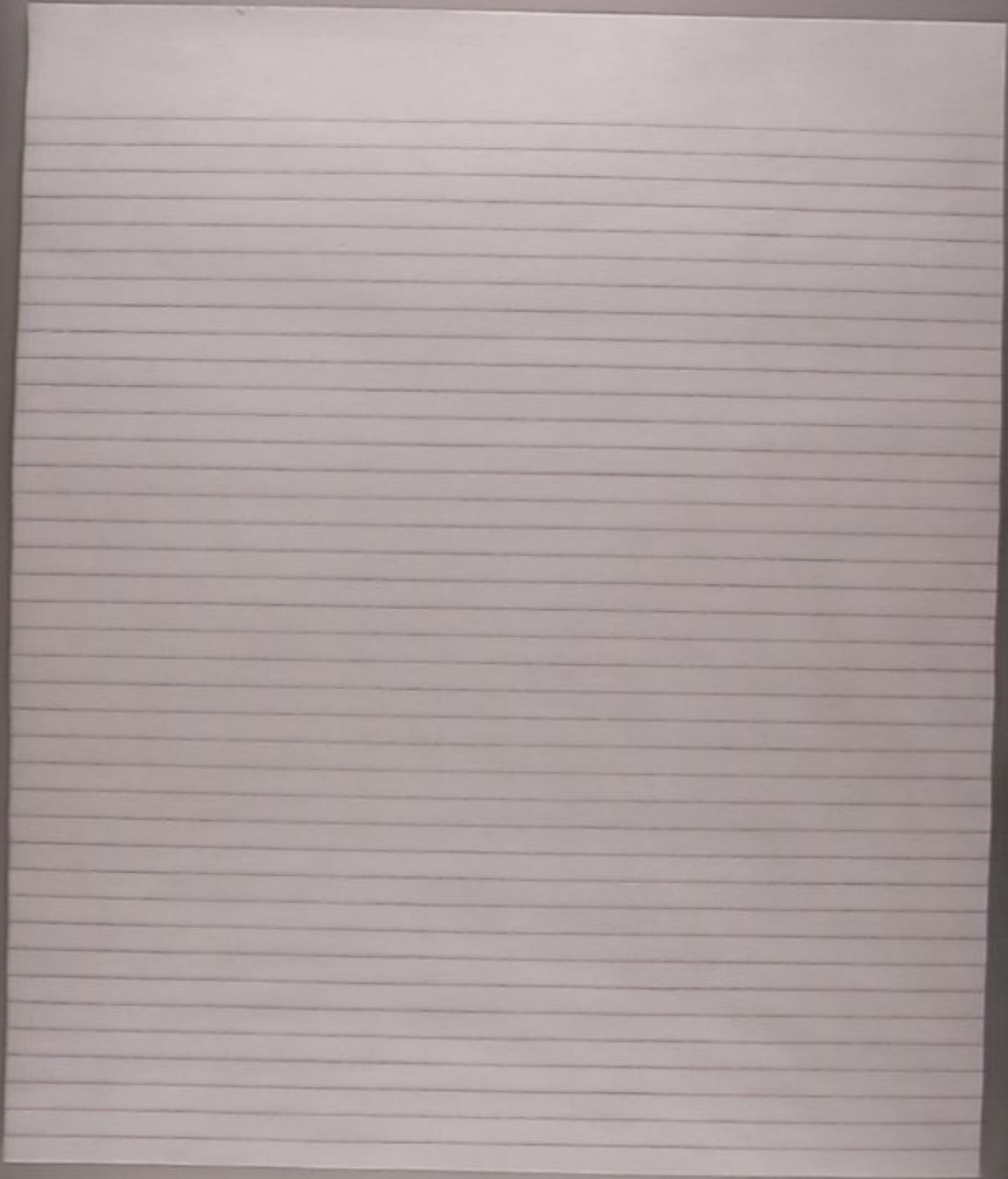


Title: A discrete, Lorentz-invariant wave equation and its continuum limit

Date: Sep 05, 2006 02:00 PM

URL: <http://pirsa.org/06090022>

Abstract:



disclosures

disciplines

cutoff (



disorder

cutoff (Lor. Invar.)

fluctuations

disorder

cutoff (Lor. Invar.)

fluctuations

scattering

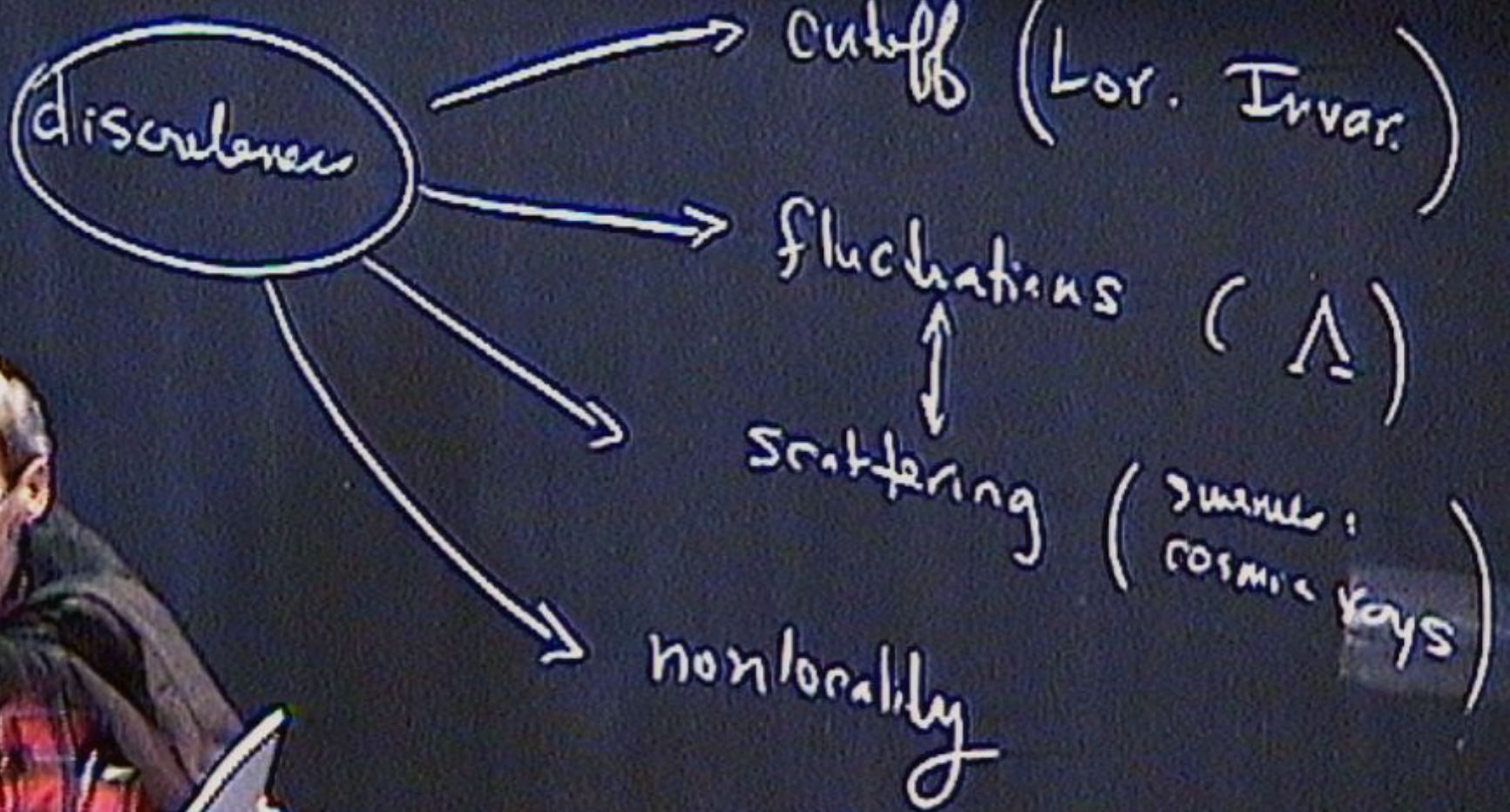


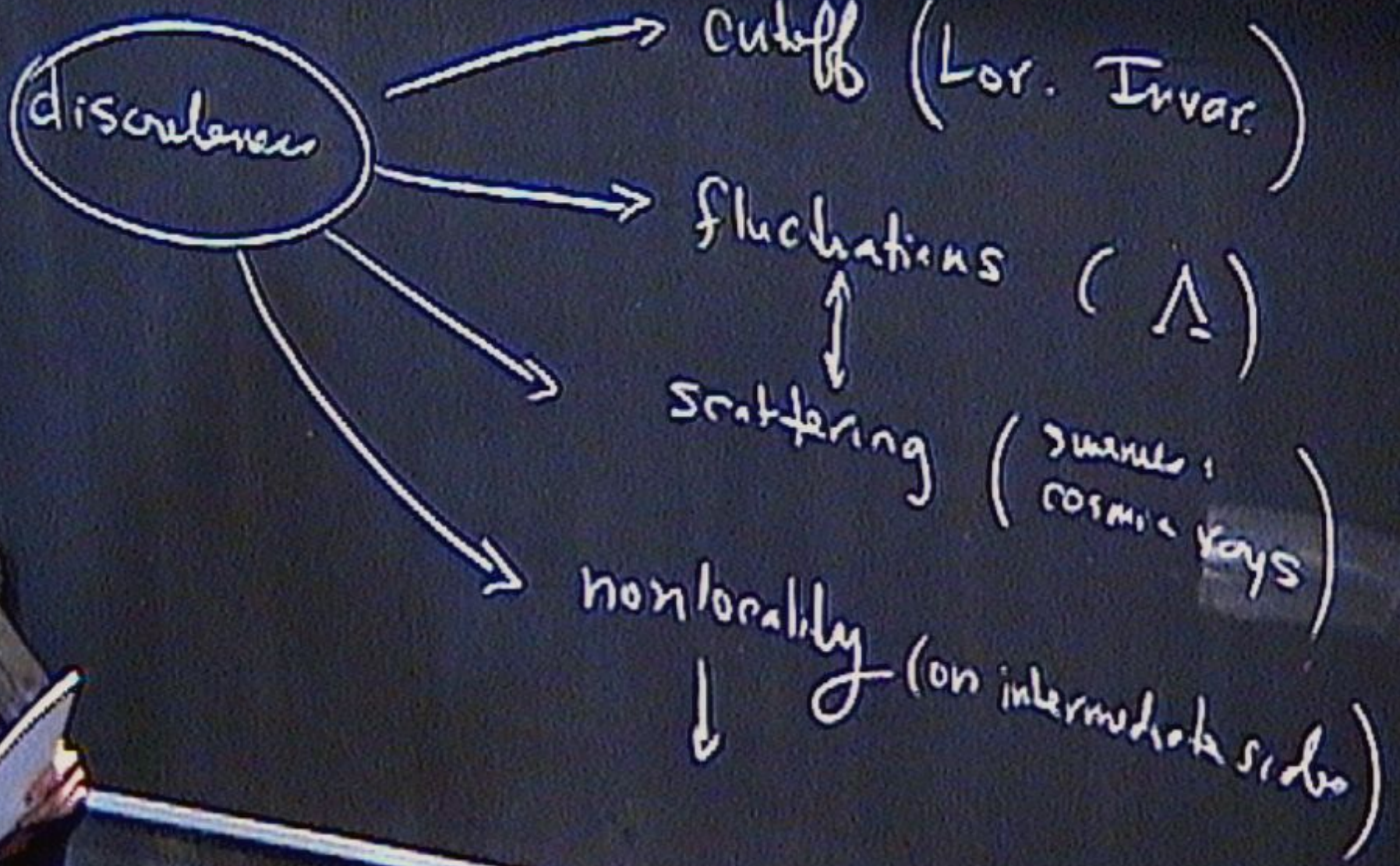
discorlance

cutoff (Lor. Invar.)

fluctuations (Λ)

scattering (small: cosmic)





disorder

cutoff (Lor. Invar)

fluctuations (Λ)

scattering (source: cosmic rays)

nonlocality (on intermediate scales)
↓ sources curvature
effect on ϕ in expanding cosmos

1. Discreteness *can respect* Lorentz-transformations
(Kinematic randomness plays a role – Poisson processes)
2. But locality must be abandoned
Implies radical nonlocality at fundamental level (*micro-scale* l)
3. One can recover locality approximately at large scales (*macro-scale*)
4. But residual nonlocality survives at *intermediate* length-scales
(*meso-scale*, below λ_0)
5. An effective meso-theory would be *continuous* but *nonlocal*

Illustrate these claims with scalar field ϕ on a *fixed* causet C :

Recovery of $\square\phi$.

($\delta\Lambda$ is also a nonlocal effect of discreteness; I'll not discuss it)

$V = N \xrightarrow{\text{Poisson}} \text{Lor Invar}$

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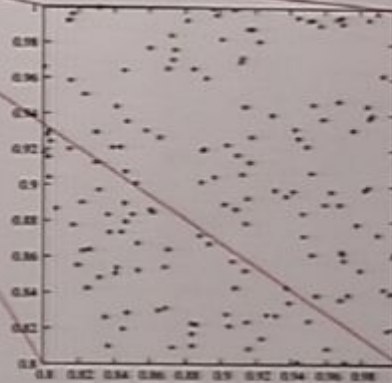
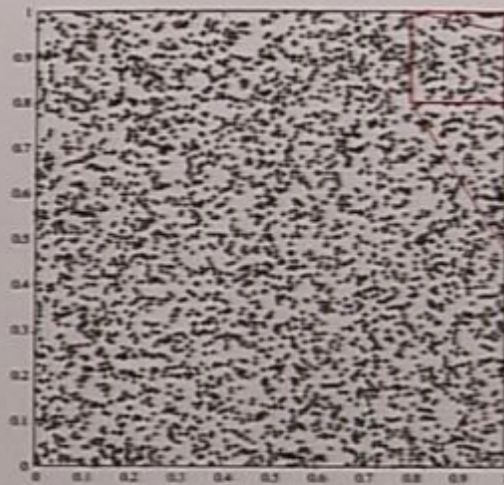
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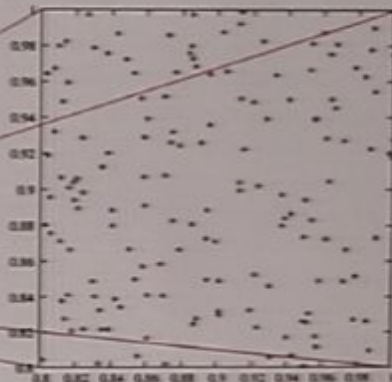
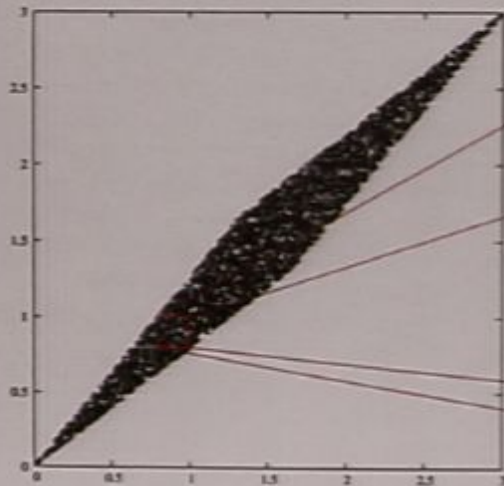
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Red squares are in same coordinate location.
 (Boost is an active transformation.)

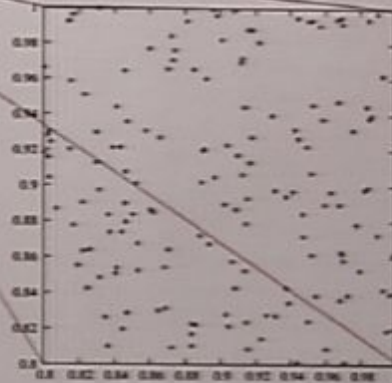
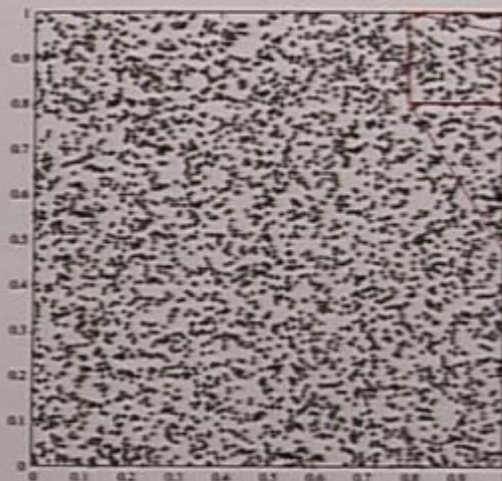


$$\mathbb{D}_K$$

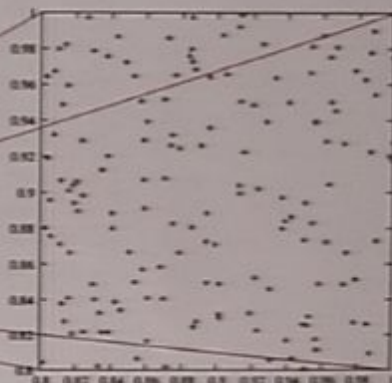
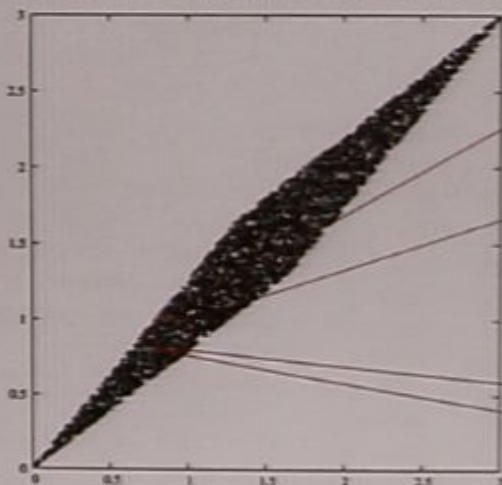
$$\sim q^{-2}$$

$$V = N \xrightarrow{\text{Poisson}} \text{Lor Invar}$$

$$C \xrightarrow{\text{faithful embedding}} M(1)$$



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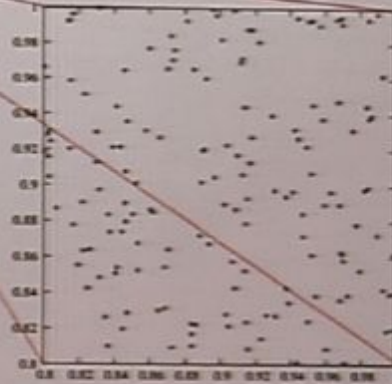
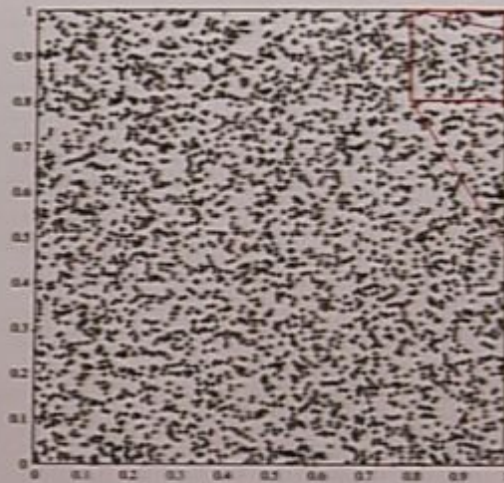


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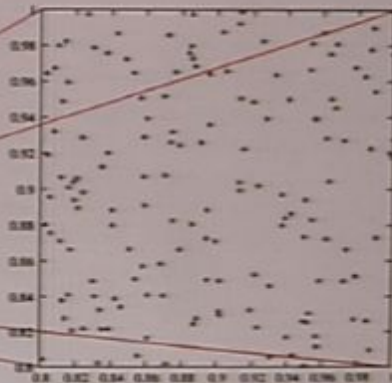
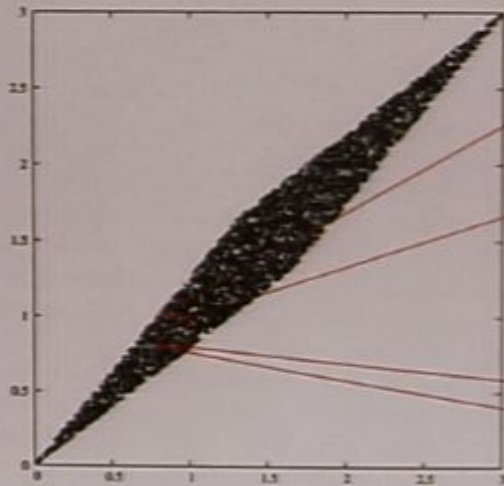
$$V = N \xrightarrow{\text{Poisson}} \text{Lor Inver}$$

$$C \xrightarrow{\text{faithful embedding}} M(3)$$

$$\frac{R}{\binom{N}{2}} \rightarrow d$$



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A theorem on Poisson processes

Ω = space of all sprinklings of M^d (sample space)

Poisson process induces a measure μ on Ω

Let f be a rule for deducing a direction from a sprinkling
 $f : \Omega \rightarrow H = \text{unit vectors in } M^d$

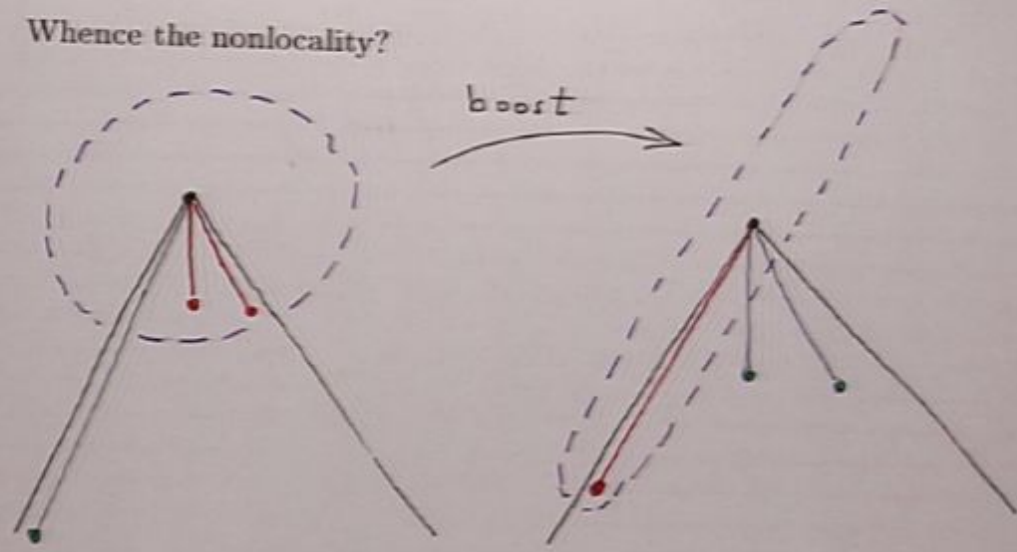
Require f *equivariant* ($f\Lambda = \Lambda f$, $\Lambda \in \text{Lorentz}$)

Assume that f is measurable (hardly an assumption)

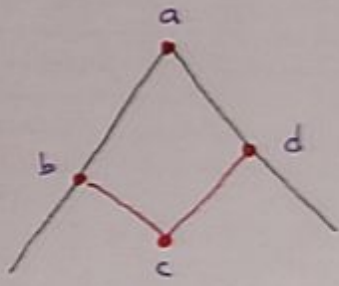
THEOREM *No such f exists*
(not even on a partial domain of positive measure)

(So with probability 1, a sprinkling will not determine a frame.)

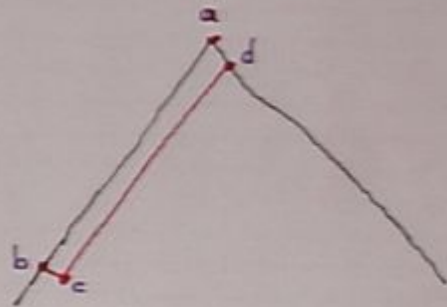
Whence the nonlocality?



Needs a miracle. (consider eg $\phi = t^2 - x^2$, invariance $\Rightarrow \infty$?)



$$(a+c) - (b+d) \rightarrow \square \phi$$

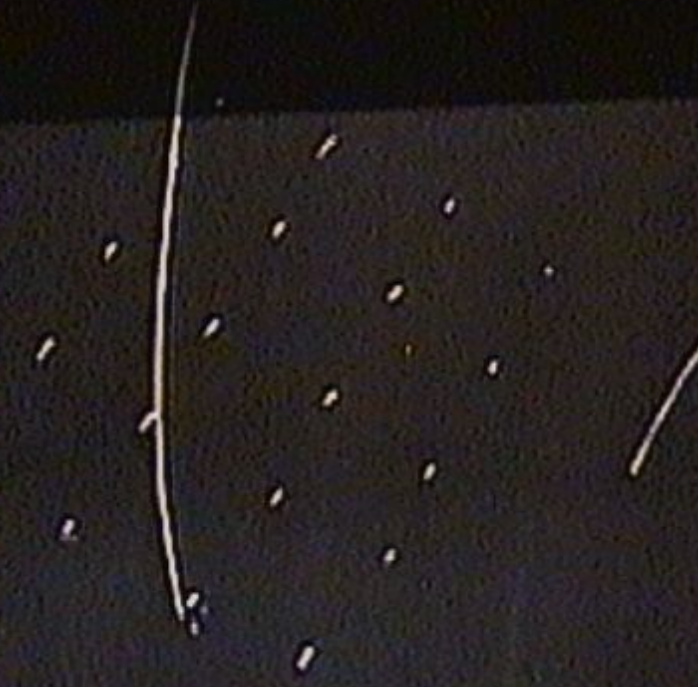


$$(a+c) - (b+d) = (a-d) - (b-c) \rightarrow \bigcirc$$

$$x < y$$

↑
independent

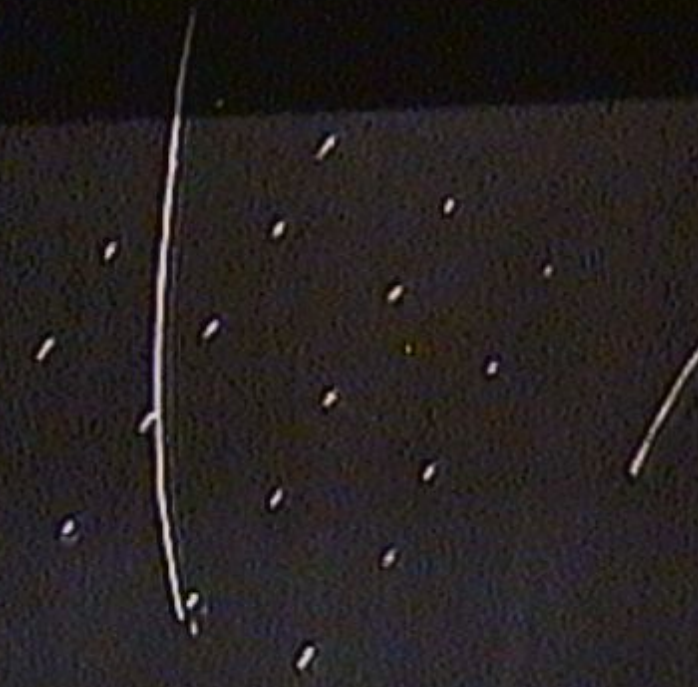
↑
dependent



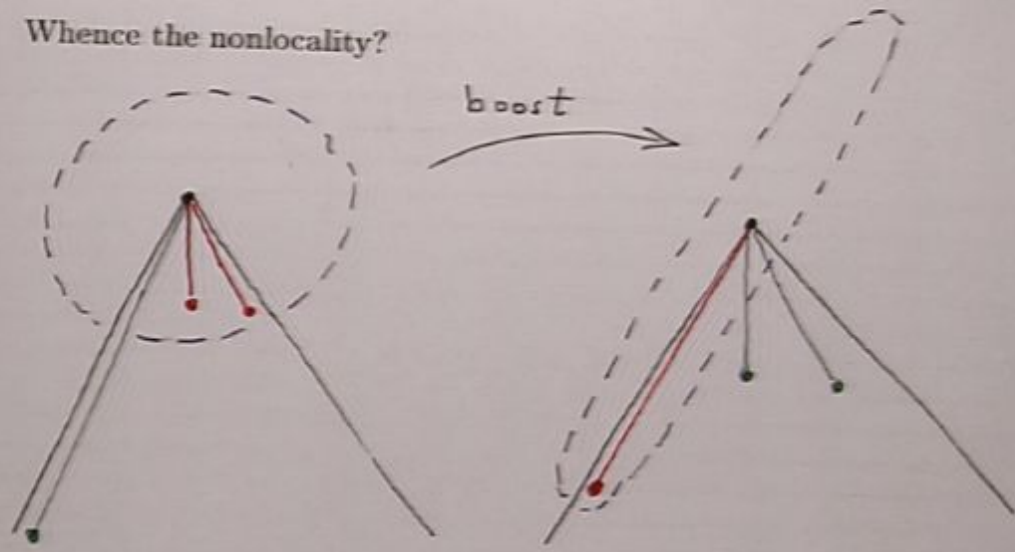
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↑
descender

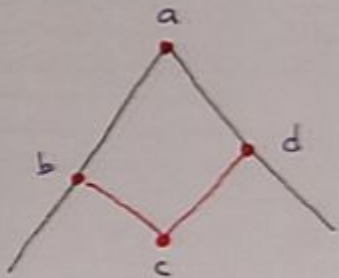
↑
descendant



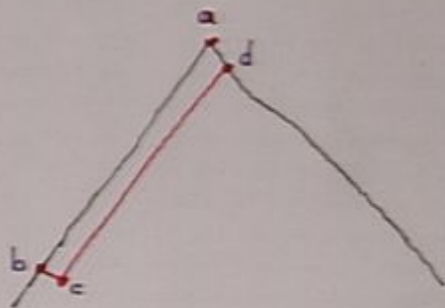
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$$\begin{aligned} (a+c) - (b+d) \\ = (a-d) - (b-c) \rightarrow \bigcirc \end{aligned}$$

These ideas lead to expressions like

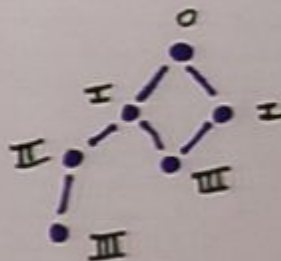
$$\frac{4}{l^2} \left(-\frac{1}{2} \phi(0) + \sum_{x \in I} \phi(x) - 2 \sum_{x \in II} \phi(x) + \sum_{x \in III} \phi(x) \right)$$

i.e.

$$\square \phi(i) \leftrightarrow \sum_k B(i, k) \phi(k)$$

where

$$\frac{l^2}{4} B(i, k) = \begin{cases} -\frac{1}{2} & \text{if } i = k \\ 1 & \text{if } i \prec k \text{ is a link (NN) } |\langle i, k \rangle| = 0 \\ -2 & \text{if } i \prec k \text{ and (NNN) } |\langle i, k \rangle| = 1 \\ 1 & \text{if } i \prec k \text{ and (NNNN) } |\langle i, k \rangle| = 2 \end{cases}$$



Can prove that, as $l \rightarrow 0$

$$S \equiv \mathbf{E} \sum_k B_{ik} \phi_k \rightarrow \square \phi(x_i)$$

using e.g.

$$\mathbf{E} \sum_{x \in I} \phi(x) = \int \frac{dudv}{l^2} \exp\{-uv/l^2\} \phi(u, v)$$

Problem: $\Delta S \rightarrow \infty$ (fluctuations) as $l \rightarrow 0$!

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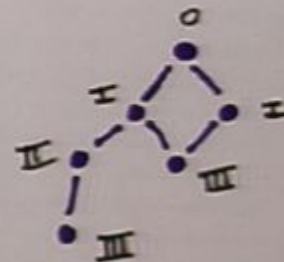
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IDEA: Our averaged sum is a *continuum* expression,

$$\int B(x-x') \phi(x') d^2x' ,$$

where

$$B(x) = \theta(x) (-2K\delta(x) + 4K^2 p(\xi) e^{-\xi}) ,$$

with $p(\xi) = 1 - 2\xi + \frac{1}{2}\xi^2$, $\xi = Kuv$, and $K = 1/l^2$.

But can *decouple* K from l^2 . We get a nonlocal continuum analog of the D'Alembertian! Call it \square_K .

Umkehren: approximate \int by \sum over sprinkled points!

This produces the causet expression,

$$\frac{4\varepsilon}{l^2} \left(-\frac{1}{2}\phi(y) + \varepsilon \sum_{x \prec y} p(\xi) e^{-\xi} \phi(x) \right) ,$$

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The "trick" drives down the fluctuations, but nonlocality survives at the intermediate scale $\lambda_0 = 1/\sqrt{K}$.

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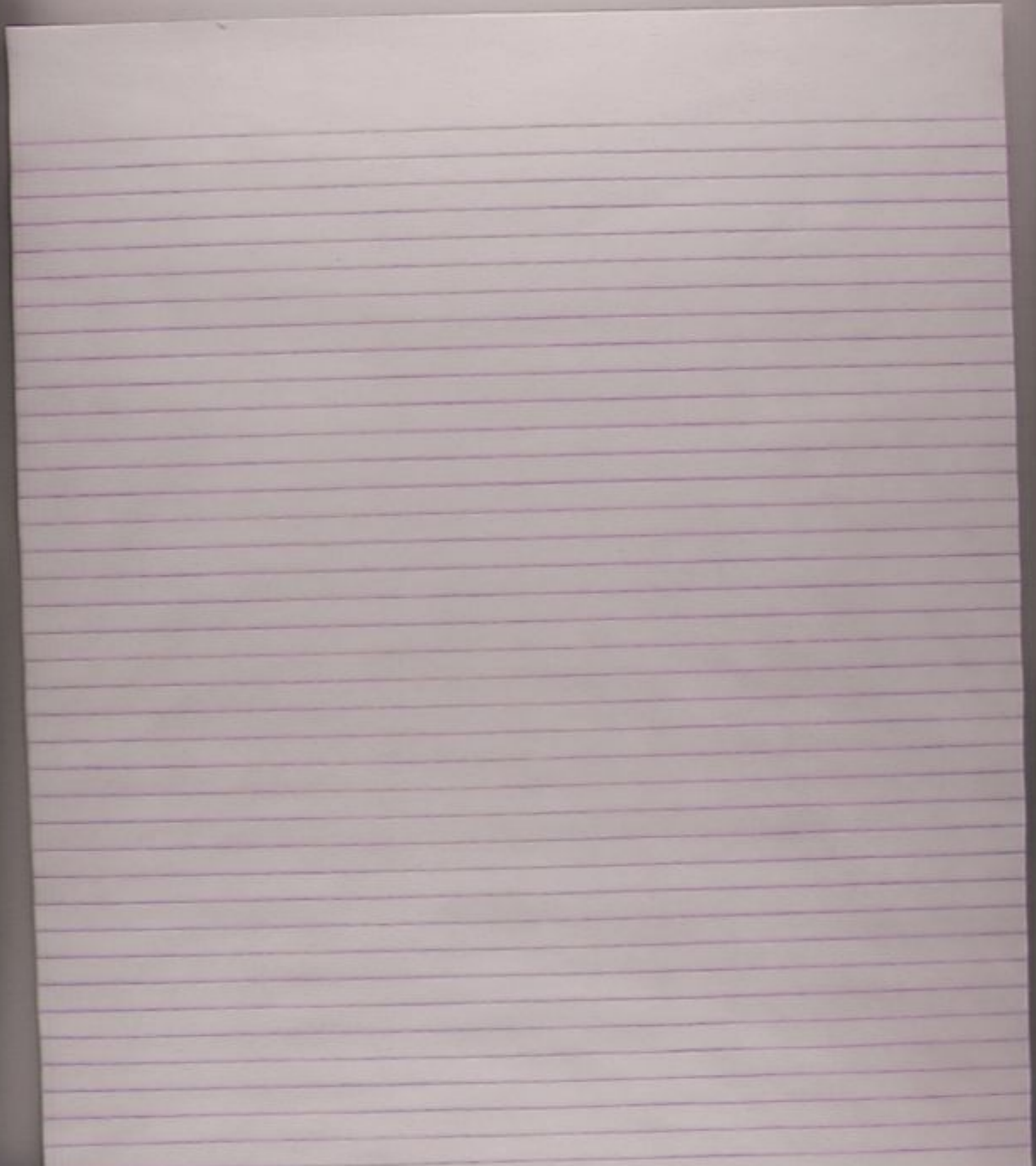
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Remarks and applications

- Analogous expressions exist in other dimensions. In $d = 4$

$$p(\xi) = 1 - 3\xi + (3/2)\xi^2 - (1/6)\xi^3$$

$$\leftrightarrow \sum_I -3 \sum_{II} +3 \sum_{III} - \sum_{IV}$$

- Can now study propagation on sprinkled causet (Rideout)
cf. swerves
- The continuum theory's free field is *stable*: ($\ker \square_K = \ker \square$)
But response to sources differs
- Quantum Field Theory version? New approach to renormalization?
Our nonlocality does *not* remove ∞ 's, but perhaps it will allow an
invariant (Lorentzian) cutoff.
- How big is λ_0 ? Must balance fluctuations vs. nonlocality.
 $L = \text{Hubble}^{-1}$, $l = \text{Planck length}$.

$$\lambda_0 \gtrsim (l^2 L)^{1/3}$$

if want \square_K pointwise accurate. \Rightarrow nuclear size!!