

Title: ISSYP 2006a - Student Presentations

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URL: <http://pirsa.org/06080006>

Abstract:

# Thermodynamics

How a Black Hole is formed



- Cloud contracts due to gravity
- ~ Star is formed and burns, increasing the temperature
- ~ Star cools
- ~ Black Hole is formed



The curved spacetime of a Black Hole.



Schwarzschild Radius

$$R = \frac{2GM}{c^2} \quad (\text{where } G \text{ and } c \text{ are constants})$$

From this equation we can see that  
The distance from the singularity to the  
Event horizon is directly proportional to

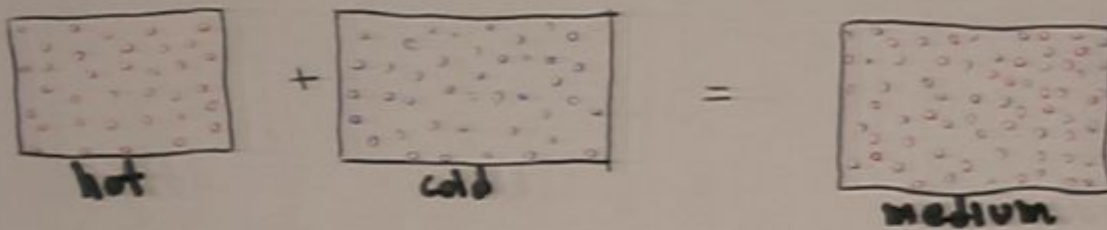
# 1st Law: Conservation of Energy

$$\Delta u = q - w$$

↑ change in internal energy  
↑ heat added to system  
work done by system

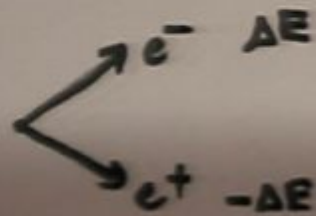
# 2nd Law: Entropy

Heat flows from hot to cold

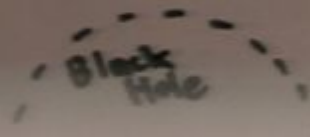


Hawking hypothesized that Black Holes may radiate particles.

According to Quantum Mechanics, particles and their corresponding antiparticles (eg. electrons and positrons) may spontaneously appear in a vacuum.



Total change in energy in system  
= sum of energies formed  
=  $\Delta E + (-\Delta E)$   
= 0



$$M_{BH} = M - \Delta E$$

Hawking suggested that black holes may be 'leaking' particles as the total mass decreases ( $M_{\text{BH}} = M - \Delta E$ ) and a particle is emitted.

From this assumption, and knowing that black holes radiate matter/energy, we came up with an equation for the lifespan of a black hole (how long it takes for it to radiate all its energy/mass away).

$$t_i = m_0^3 \left( \frac{2^3 \pi^3 k^4 G^2}{3\sigma k^4 c^6} \right)$$

where  $\pi \approx 3.14$

$5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$   
 $2.998 \times 10^8 \text{ m s}^{-1}$   
 $6.674 \times 10^{-11} \text{ N m}^2 \text{ kg}^{-2}$

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$$t_l = m_0^3 \left( \frac{2^3 \pi^3 k^4 G^2}{3\sigma k^4 c^6} \right)$$

where  $\pi \approx 3.14$

$$k = 1.38 \times 10^{-23} \text{ J/K}$$

$$\sigma = 5.67 \times 10^{-8} \text{ W m}^2 \text{ K}^{-4}$$

Entropy

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From this assumption, and knowing that black holes radiate matter/energy, we came up with an equation for the lifespan of a black hole (how long it takes for it to radiate all its energy/mass away).

$$t_l = m_0^3 \left( \frac{2^9 \pi^3 k^4 G^2}{30 \hbar^4 c^6} \right)$$

where  $\pi \approx 3.14$

$$k = 1.38 \times 10^{-23} \text{ J K}^{-1}$$

$$G = 6.67 \times 10^{-11} \text{ N m}^2 \text{ kg}^{-2}$$

$$\sigma = 5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$$

$$c = 3.00 \times 10^8 \text{ m s}^{-1}$$

$$\hbar = 6.63 \times 10^{-34} \text{ J s}$$

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## \* Entropy

- Measure of the order/disorder of a system
- Given as S.

$$\Delta S = \frac{Q}{T}$$

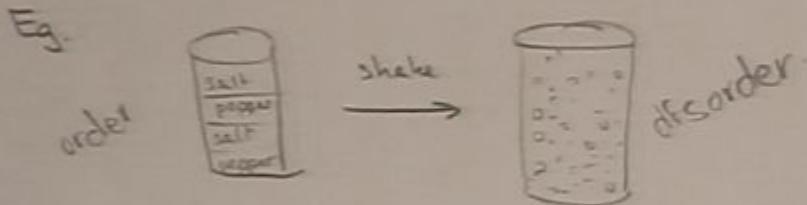
← heat added  
← temperature

Dimensions of entropy are J per K

According to the second law of thermodynamics

$$\Delta S \geq 0$$

Higher Entropy = Greater Disorder



Entropy of systems increases over time

Macrostate

Possible

1

3H

2

2H1T

3

2T1H

3



Macrostate	Possible
1	
2	
3	
4	

Macrostate

Possible

1

3H

2

2H1T

3

2T1H

4

3T



Macrostate	Possible Microstates	Number of Microstates
1 3H	HHH	1
2 2HT		
3 HTT		
4 THT		
5 TTH		
6 TTT		



Macrostate	Possible Microstates	Number of microstates
1 3H	HHH	1
2 2HT	HHT	
3 1H		
4 T		



	Macrostate	Possible Microstates	Number of microstates
1	3H	HHH	1
2	2H1T	HHT HTH THH	3
3	2T1H		
4	3T		

Macrostate	Possible Microstates	Number of Microstates
1 3H	HHH	1
2 2H1T	HHT HTH THH	3
RT 1H	TTH THT HTT	
RT		



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4	TTT	1

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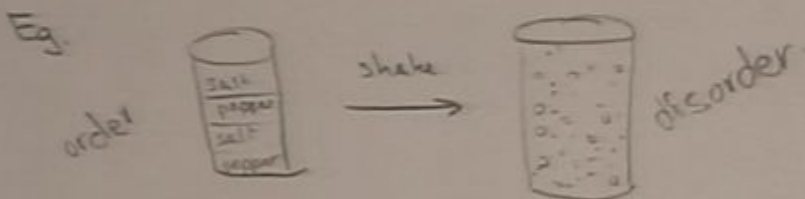
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Entropy of systems increases over time

A system is made up of microstates and macrostates.

Microstates - Specific information about systems like velocity & velocity

Macro states - General information about systems like temperature or # of molecules

Consider a case where 3 coins are tossed.

Macrostate	Possible microstates	Number of microstates
3H	HHH	1
2H 1T	HHT, HTH, THH	3
1H 2T	HTT, THT, TTH	3
3T	TTT	1

low entropy →

(3)

The least probable states are the most orderly and have lower entropy.

The reverse also applies.



$$\Delta S = \frac{Q}{T}$$

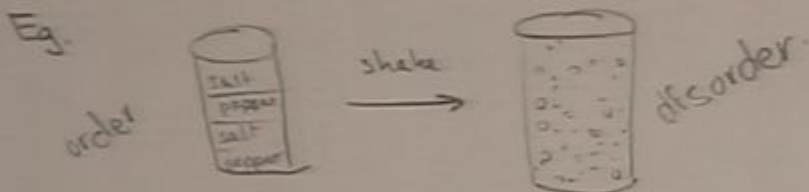
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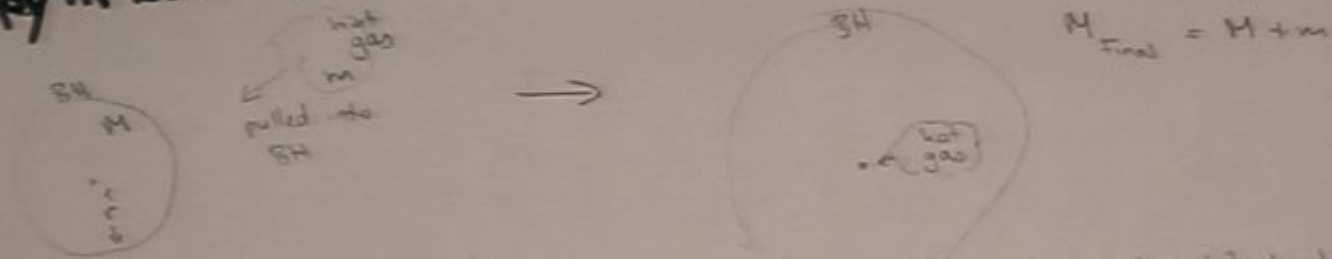
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# Entropy in Black Holes



We know the hot gas has entropy, and it combines with the black hole. Since entropy can't decrease, a black hole must have entropy. ... and surface area of the

We know the hot gas has entropy, and it combines with the black hole.  
Since entropy can't decrease, a black hole must have entropy.  
We derived equations for entropy in terms of mass and surface area of the event horizon (A).

$$S(M) = \left( \frac{4\pi Gk}{\pi c} \right) M^2$$

all constants

$$S(A) = \left( \frac{k c^3}{4G\hbar} \right) A$$

### Summary

- Black holes have a lifespan
- Black holes have entropy
  - They have microstates

↓  
What are these microstates?

The answer to this question, as well as many others will be contained in a complete Quantum Theory of Gravity.

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all constants

$$S(A) = \left( \frac{kc^3}{4G\hbar} \right) A$$

### Summary

- Black holes have a lifespan
- Black holes have entropy
  - They have microstates

↓  
What are these microstates?

all constants

$$S(A) = \left( \frac{Kc^3}{46\pi} \right) A$$

## Summary

- Black holes have a lifespan
- Black holes have entropy
  - They have microstates
  - ↓
  - What are these microstates?

The answer to this question, as well as many others will be contained in a complete Quantum Theory of Gravity.

(4)

$E \sim \hbar \omega$

$$\frac{\hbar}{2\pi}$$

SCHRODINGER

wave function  
partial differentiation

VS

HEISENBERG.

matrices  
eigenvalue/vectors



START: E of SHO

END: E of N-particle  
Boson / Fermion systems

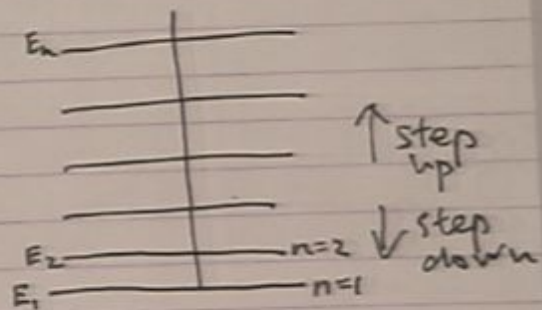
HOW? : FOLLOW THE E TRAIL!

$$\bar{z}z = \frac{m\omega^2}{2} \hat{x}^2 + \frac{\hat{p}^2}{2m}$$

$$z = \sqrt{\frac{m\omega^2}{2}} \hat{x} + \frac{i\hat{p}}{2m}$$

$$\frac{1}{\sqrt{\hbar m \omega}} z \Rightarrow \hat{a} \quad \text{annihilation operator}$$

$$\frac{1}{\sqrt{\hbar m \omega}} \bar{z} \Rightarrow \hat{a}^\dagger \quad \text{creation operator}$$



$$\hat{a}^\dagger = \sqrt{\frac{m\omega}{\hbar}} |n+1\rangle$$

$$\overline{z}z = \frac{m\omega^2}{2} \hat{x}^2 + \frac{\hat{p}^2}{2m}$$

$$z = \sqrt{\frac{m\omega^2}{2}} \hat{x} + \frac{i\hat{p}}{2m}$$

START: E

END: E

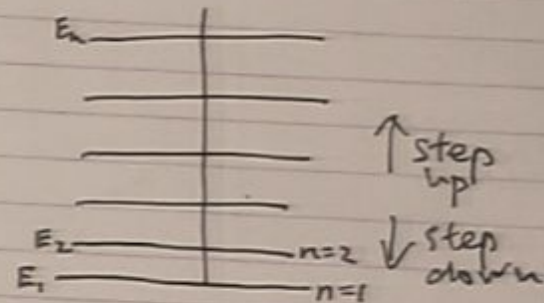
FROM: E

$$\bar{z}z = \frac{m\omega}{2} \hat{x}^2 + \frac{\hat{p}}{2m}$$

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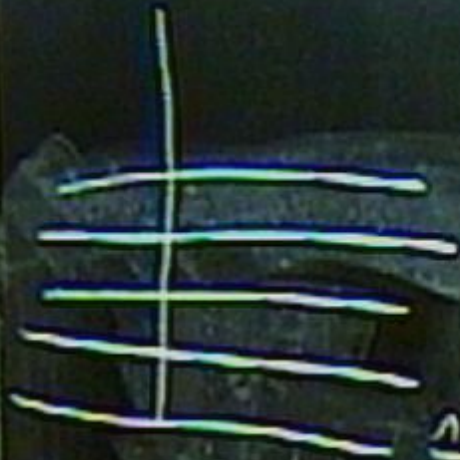


$$\hat{a}|n\rangle = \sqrt{n}|n-1\rangle, \quad \hat{a}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle.$$

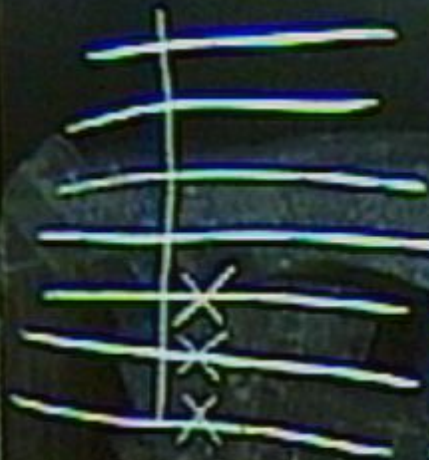
START: E

END: E

HOM: E

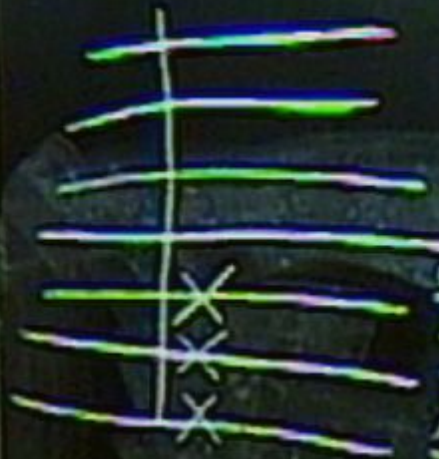


$$\frac{h}{2} \omega$$



2h  
 3h  
 2h  
 2h

6  
 5  
 4  
 3  
 2  
 1  
 0



$\frac{5}{2}h\omega$   
 $\frac{3}{2}h\omega$   
 $\frac{1}{2}h\omega$

6  
 5  
 4  
 3  
 2  
 1  
 0



ing in the harmonic oscillator. These are called into two categories: Bosons and Fermions. Fermions are called 'Fermi particles'. What is Pauli exclusion principle? occupy a single quantum state.

to you! Recall that in our context of harmonic oscillator the energy eigenstates have unique energy  $E_n$  - that is, the energy eigenstates have unique energy. The  $k^{th}$  boson. The  $N$ -particle quantum state is

not Fermions. That is, there is no exclusion principle. We can take them to be  $|s_1, s_2, \dots, s_N\rangle$  with  $s_1 \geq s_2 \geq \dots \geq 0$ .

specified by  $N$ -tuples  $\vec{f} = (f_1, f_2, \dots, f_N)$  level number. This means all  $f_k$ 's are  $f_1, \dots, f_j, \dots, f_i, \dots, f_N$  are not

implies

you have to use the formula  $\sum_{k=1}^p k = \frac{p(p+1)}{2}$ . The energy measured

$$H_f = E_f - E_0 = \hbar\omega \left[ \frac{N(N-1)}{2} + \sum_{k=1}^N k \right]$$

## 2 N Bosons in SHO

Next let us consider  $N$  identical free bosons in a harmonic oscillator. The  $N$ -particle quantum state is

$$|s_1, s_2, \dots, s_N\rangle$$

Again we consider identical particles and so we have to get rid of the labels. We can take them to be  $|s_1, s_2, \dots, s_N\rangle$  with  $s_1 \geq s_2 \geq \dots \geq 0$ .

$$E_{\vec{s}_1} = \hbar\omega \sum_{k=1}^N (s_k + \frac{1}{2}) = \frac{N\hbar\omega}{2} + \hbar\omega \sum_{k=1}^N s_k$$

The ground state is when all the bosons are in the ground state  $|n=0\rangle$ . We again choose to measure the energy of a given state above the ground state.

$$H_{\vec{s}_1} = E_{\vec{s}_1} - E_0 = \hbar\omega \sum_{k=1}^N s_k$$

## 4.3 Bosons in SHO

The  $N$ -particle quantum state is  $|s_1, s_2, \dots, s_N\rangle$  with  $s_1 \geq s_2 \geq \dots \geq 0$ . We can take them to be  $|s_1, s_2, \dots, s_N\rangle$  with  $s_1 \geq s_2 \geq \dots \geq 0$ . Then a  $N$ -particle quantum state can be specified by  $N$  bosons and the corresponding energy is

$$\sum_{k=1}^N k\hbar\omega$$



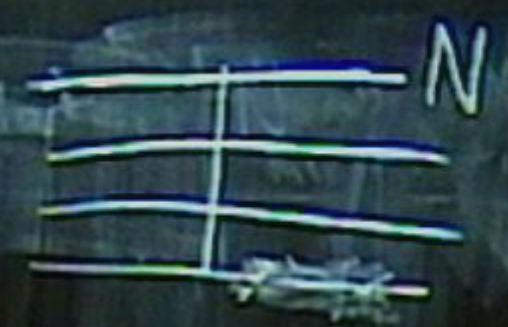
N



2hw  
 3hw  
 2hw  
 1hw

6  
 5  
 4  
 3  
 2  
 1  
 0



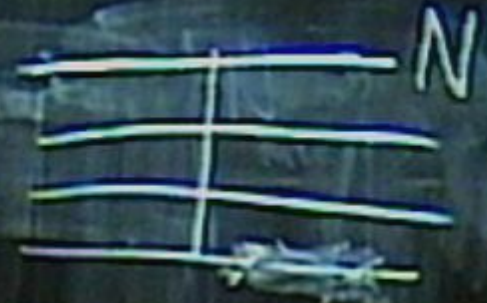


N



5 h w  
 3 h w  
 2 h w

6  
 5  
 4  
 3  
 2  
 1  
 0



N

$$|f_1, f_2, \dots, f_n\rangle$$

$$E_n = (n + \frac{1}{2}) \hbar \omega$$

$$E_{n+1} =$$

$$E_0 = \frac{N^2}{2} \hbar \omega$$

$$|f_1, f_2, \dots, f_n\rangle$$

$$E_n = (n + \frac{1}{2})\hbar\omega$$

$$E_{\pm} = \frac{N}{2}\hbar\omega + \hbar\omega \sum_{k=1}^N f_k$$

$$E_0 = \frac{N^2}{2}\hbar\omega$$

$$|f_1, f_2, \dots, f_n\rangle$$

$$E_n = (n + \frac{1}{2})\hbar\omega$$

$$E_f = \frac{N}{2}\hbar\omega + \hbar\omega \sum_{k=1}^N f_k$$

$$H_f \rightarrow -\frac{N(N-1)}{2} + \sum_{k=1}^N f_k$$

$$E_0 = \frac{N^2}{2}\hbar\omega$$

$$|f_1, f_2, \dots, f_n\rangle$$

$$E_n = (n + \frac{1}{2})\hbar\omega$$

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$$\underline{H_f} = -\frac{N(N-1)}{2} + \sum_{k=1}^N f_k$$

$$E_0 = \frac{N^2}{2}\hbar\omega$$

$|f_1, f_2, \dots, f_n\rangle$

degeneracy

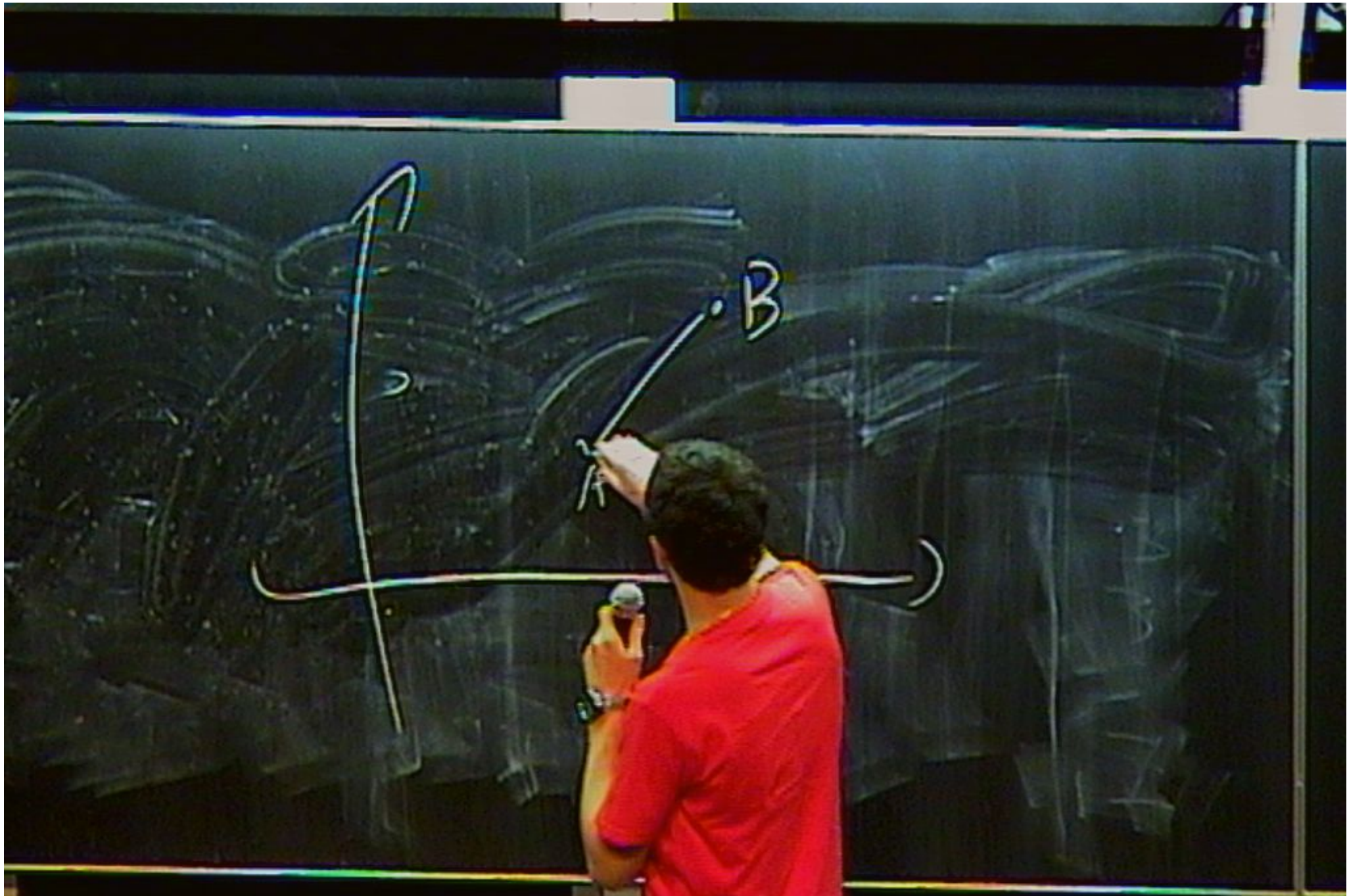
$$E_n = (n + \frac{1}{2})\hbar\omega$$

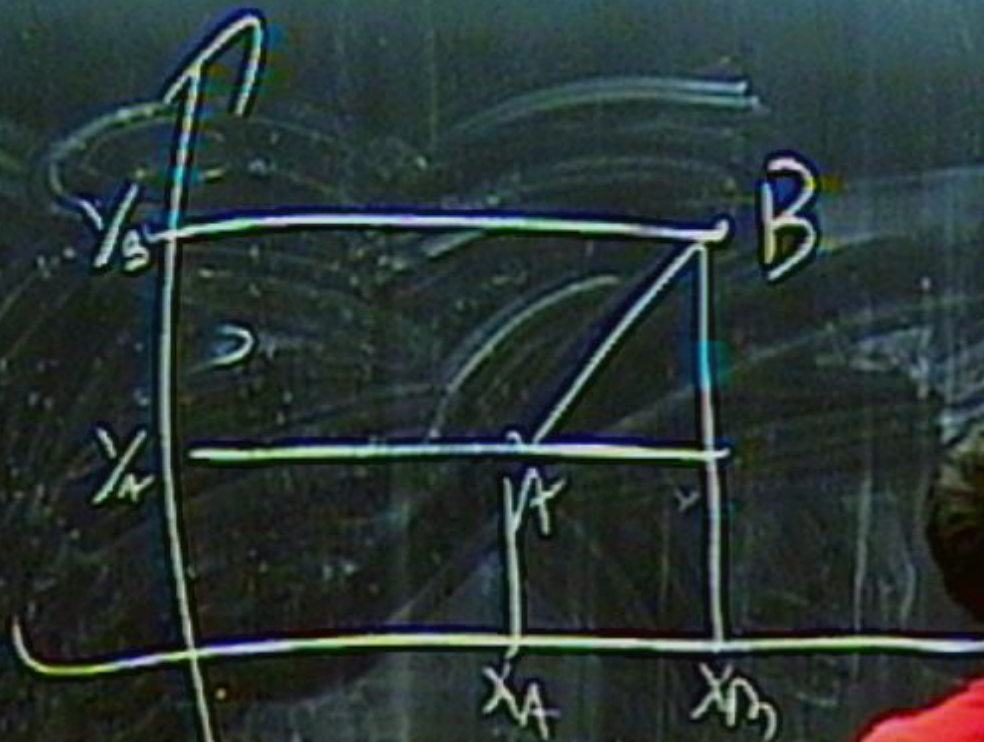
$$E_f = \frac{N}{2}\hbar\omega + \hbar\omega \sum_{k=1}^N f_k$$

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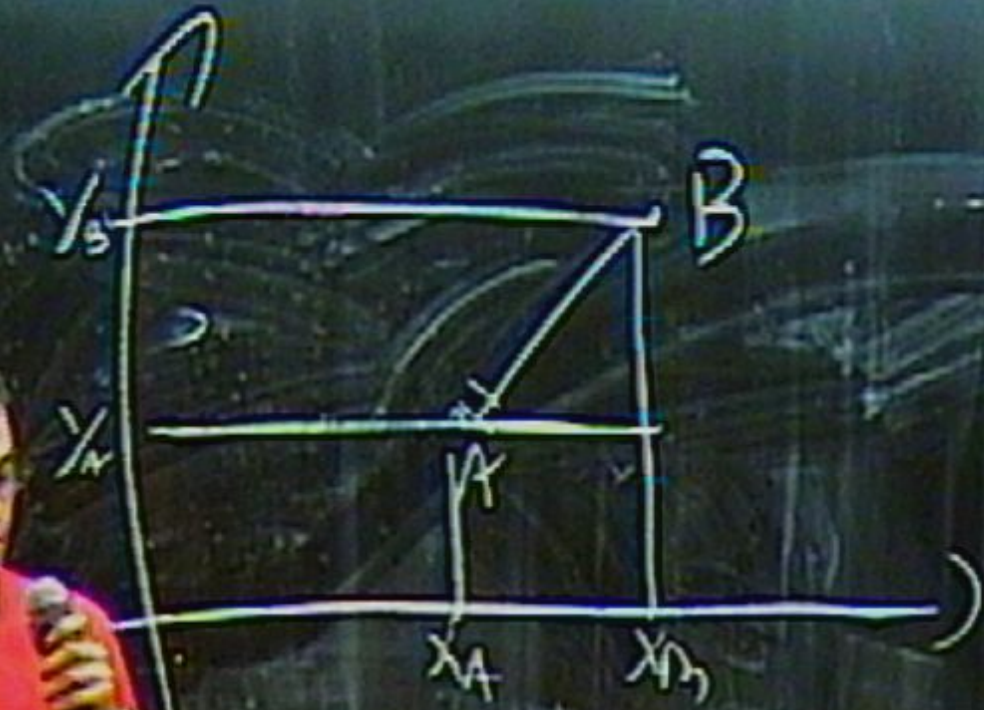
$$E_0 = \frac{N^2}{2}\hbar\omega$$



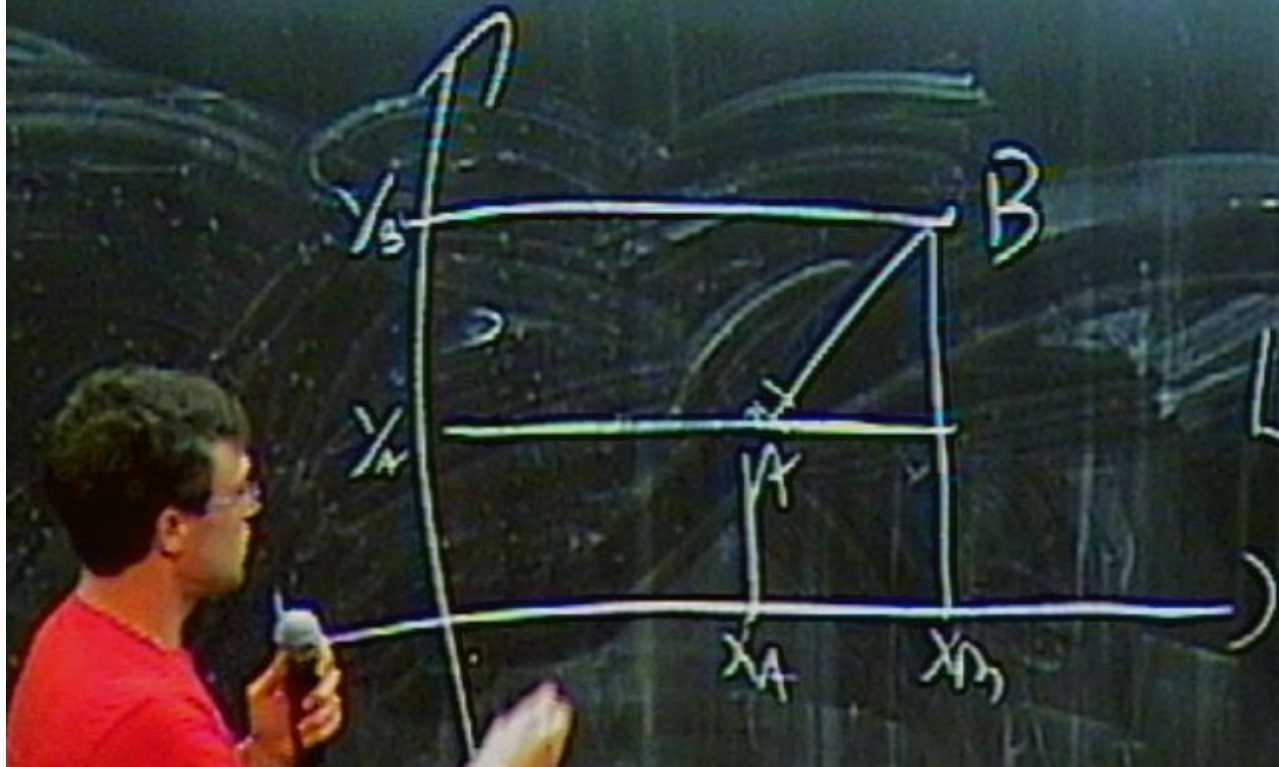




$$\Delta S^2 \equiv \Delta$$



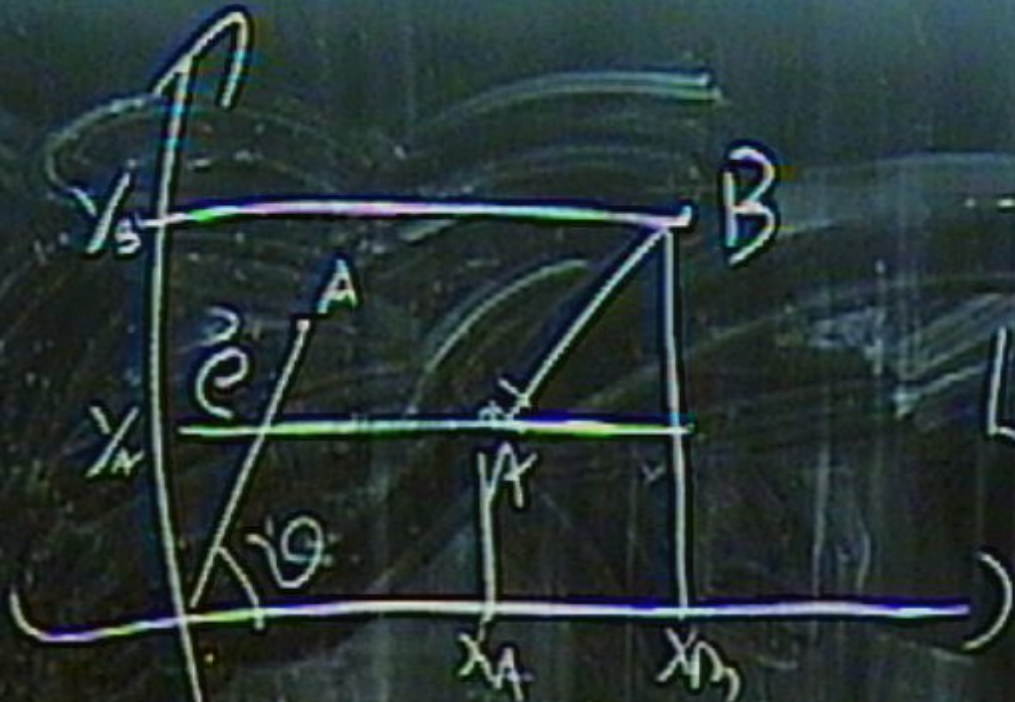
$$\Delta s^2 = \Delta x^2 + \Delta y^2$$



$$\Delta s^2 = \Delta x^2 + \Delta y^2$$

$$\boxed{ds^2 = dx^2 + dy^2}$$





$$\Delta s^2 = \Delta x^2 + \Delta y^2$$

$$\boxed{ds^2 = dx^2 + dy^2}$$

$$ds^2 = e^2 d\theta^2 + dp^2$$

$E \sim \hbar \omega$

$$\frac{\hbar}{2\pi}$$

$$+ \Delta y^2$$

$$dy^2$$

$$ds^2 = e^2 d\theta^2 + dp^2$$

$$E \sim \hbar \omega$$

$$\frac{\hbar}{2\pi}$$

$$r^2 + \Delta y^2$$

$$dy^2$$

$$ds^2 = e^2 d\theta^2 + dc^2$$



$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

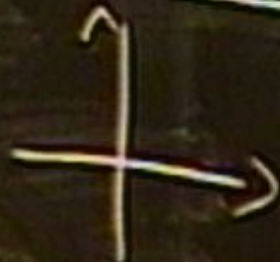
$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + \left( r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2 \right)$$

$$c=1$$
$$G=1$$

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + \frac{r^2}{1 - \frac{2M}{r}} (d\theta^2 + \sin^2\theta d\phi^2)$$



$$ds^2 = e^2 d\Omega^2 + de^2$$

BH



$$ds^2 = e^2 d\theta^2 + de^2$$

BH



$$ds^2 = e^2 dt^2 + dr^2$$

BH

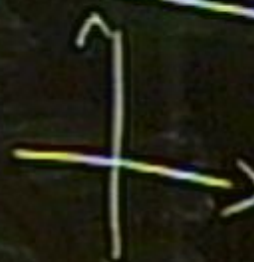


$r$	$1 - \frac{2M}{r}$
10 M	0,8
8 M	0,75
6 M	$\frac{2}{3}$
4 M	0,5
3 M	$\frac{1}{3}$
2 M	0

$$c=1$$
$$G=1$$

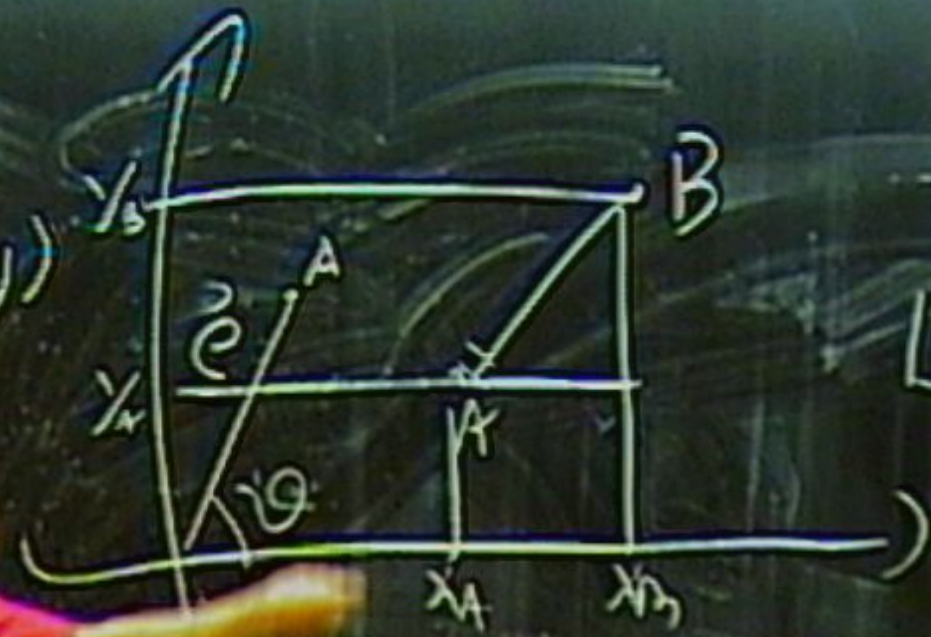
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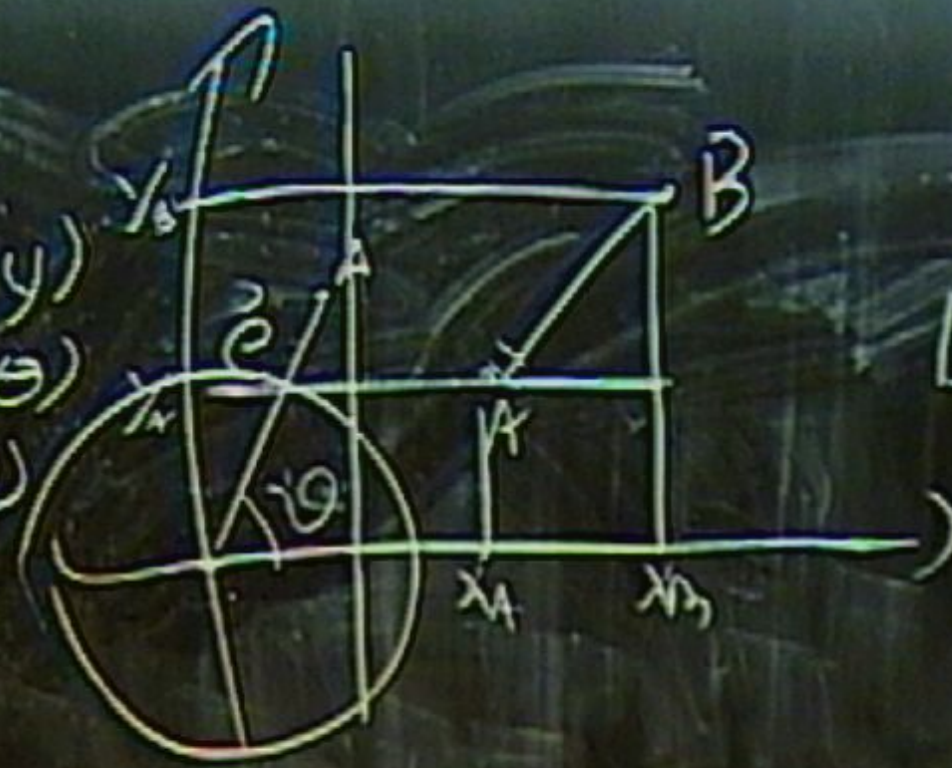


$(x, y)$   
 $(\Delta x, \Delta y)$   
 $(\Delta x, \Delta y)$



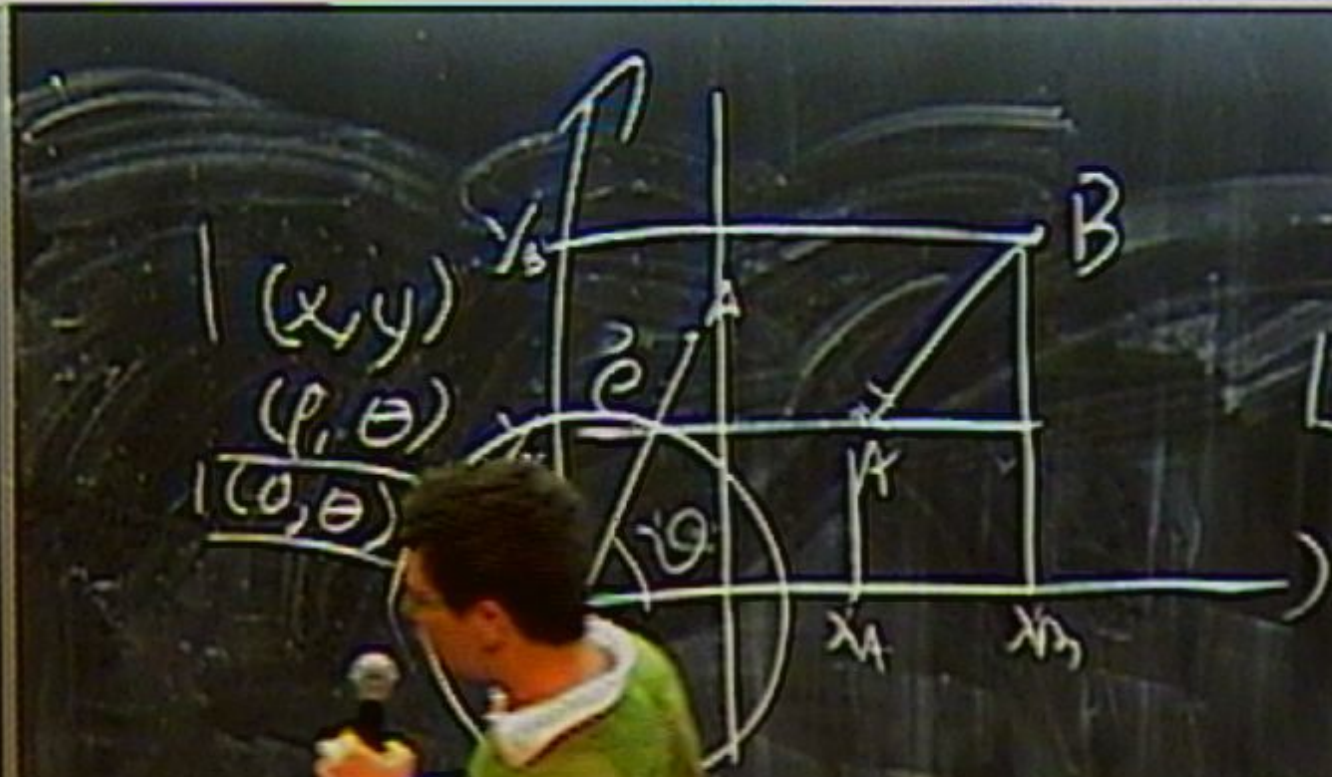
$$\Delta s^2 = \Delta x^2 + \Delta y^2$$
$$\boxed{ds^2 = dx^2 + dy^2}$$

$(x, y)$   
 $(r, \theta)$   
 $(0, \theta)$



$$\Delta s^2 = \Delta x^2 + \Delta y^2$$

$$\boxed{ds^2 = dx^2 + dy^2}$$



$(x, y)$   
 $(r, \theta)$   
 $(0, 0)$

$$\Delta s^2 = \Delta x^2 + \Delta y^2$$

$$\boxed{ds^2 = dx^2 + dy^2}$$



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$$G=1$$

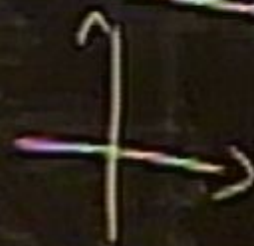
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$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + \frac{d\theta^2 + \sin^2\theta d\phi^2}{\left(\frac{2M}{2M+E}\right)}$$

$$2M+E$$

$$\frac{E}{2M} \ll 1$$

$$\left(\frac{2M}{2M+E}\right)$$



$$c=1$$

$$G=1$$

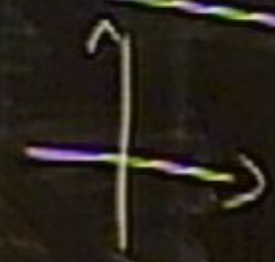
$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

$$ds^2 = \left( \frac{2M}{r} \right) dt^2 + \left( 1 - \frac{2M}{r} \right)^{-1} dr^2 + \left( r^2 + \frac{2M}{r} \right) d\phi^2$$

$$M + E$$

$$\frac{v}{2M} \ll 1$$

$$\left( \frac{2M}{r} \right) \approx \frac{2M}{r}$$



$$c=1$$

$$G=1$$

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 +$$

$$r = 2M + \epsilon$$

$$\left(1 - \frac{2M}{2M + \epsilon}\right)$$

$$dr = d\epsilon$$

$$ds^2 = \frac{\epsilon}{2M} dr^2 + \frac{2M}{\epsilon} d\phi^2$$

$$\left(1 - \frac{2M}{r}\right) dr^2 + \frac{2M}{r} d\phi^2$$

$$d\phi = \frac{2M}{\sqrt{\epsilon}} d\epsilon$$

$$c=1$$

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$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + \underbrace{r^2 (d\theta^2 + \sin^2\theta d\phi^2)}_{\text{angular part}}$$

$$r = 2r_s$$

$$\frac{2M}{r} \ll 1$$

$$\int dp = \frac{2M}{\sqrt{E}} dE$$

$$p = 2\sqrt{2M} \sqrt{E}$$

$$ds^2 =$$

$$\frac{2M}{E} dE^2 = dp^2$$

$$c=1$$

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$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

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$$r = 2M + \epsilon \quad \frac{\epsilon}{2M} \ll 1$$

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$$dr = d\epsilon$$

$$ds^2 = \frac{\epsilon}{2M} dt^2 + \frac{2M}{\epsilon} d\epsilon^2$$

$$\int d\phi = \sqrt{\frac{2M}{r}} d\epsilon$$

$$r = 2M + \epsilon$$

$$\epsilon = \frac{r - 2M}{3}$$

$$ds^2 =$$



$$c=1$$

$$G=1$$

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$$ds^2 = \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + \frac{r^2}{1 - \frac{2M}{r}} (d\theta^2 + \sin^2\theta d\phi^2)$$

$$r = 2M + \epsilon$$

$$\left(1 - \frac{2M}{2M + \epsilon}\right)$$

$$dr = d\epsilon$$

$$ds^2 = \frac{\epsilon}{2M} dr^2$$

$$\int dp = \sqrt{\frac{2M}{r}} d\epsilon$$

$$p = 2\sqrt{2M}\sqrt{\epsilon}$$

$$\epsilon = \frac{p^2}{8M}$$

$$\frac{r^2}{16M^2} = \theta^2$$

$$ds^2 = \frac{p^2}{16M^2} d\tau^2 + dp^2 \frac{r^2}{4M} = \theta$$

$$c=1$$

$$G=1$$

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

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$$ds^2 = \frac{\epsilon}{2M} dt^2 + \frac{2M}{\epsilon} d\epsilon^2 = dp^2$$

$$\int dp = \sqrt{\frac{2M}{\epsilon}} d\epsilon$$

$$p = 2\sqrt{2M\epsilon} \quad \frac{\epsilon^2}{16M^2} = \theta^2$$

$$ds^2 = \frac{p^2}{16M^2} dt^2 + dp^2 \frac{1}{4M} = \theta$$

$$ds^2 = e^2 d\theta^2 + de^2$$



# Quantum Memory

# Quantum Memory

Julian Haw Far Chin  
Kristin Flowers

Franziskuss Benedict Moritz Konstantin Hillerbrand

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Julian Haw Far Chin

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Tim Horton

Rachel Mahnke

Richard McNamara

David Nissimoff

Michael Perssons

# O u t l i n e

# Outline



# Outline

- Classical Memory
  - Energy and Spin
  - Dimensions and the Ising Model
  - Decay
  - Canonical Ensemble

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- Quantum Memory
  - Energy, Spin, Qubits
  - Entanglement
  - Major Problems
  - Self-Healing Memory
  - Possible Solutions
  - Quest for 3-D and Control of Entanglement Patterns
  - Future of Quantum Memory

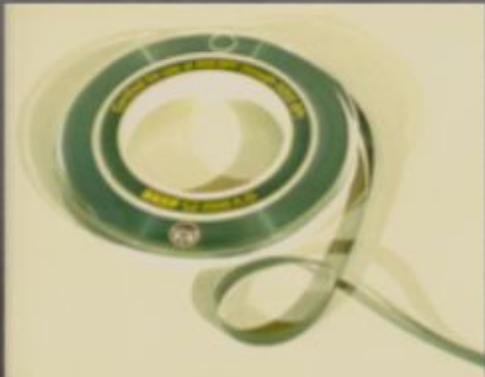
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- Lifetime: roughly 50 years
- Examples include magnetic tape, hard disk drives, CDs, DVDs

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n e r g y a n d S p

# Energy and Spin

# Energy and Spin

- Hamiltonian Equation for the energy of a state:

$$E = E_0 \cdot (\# \text{ of nearest neighbors})$$

- Spin states: up or down
- Uses millions of spins to represent one bit.
- Spins tend to follow their neighbors.





# Dimensions and the Ising Model

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- 0-D: spins isolated

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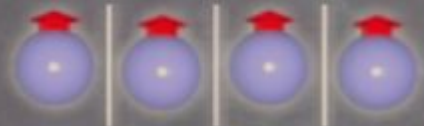
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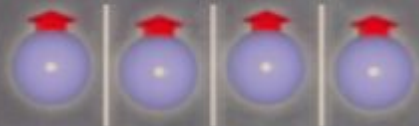


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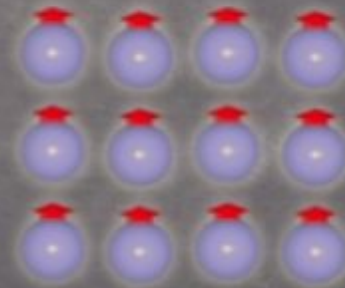
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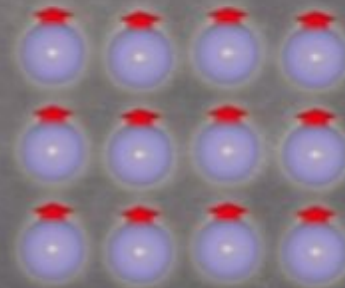
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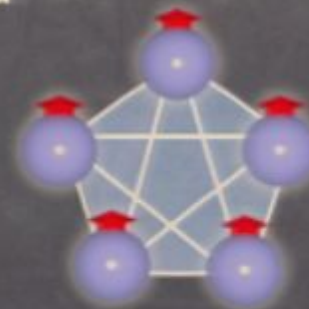
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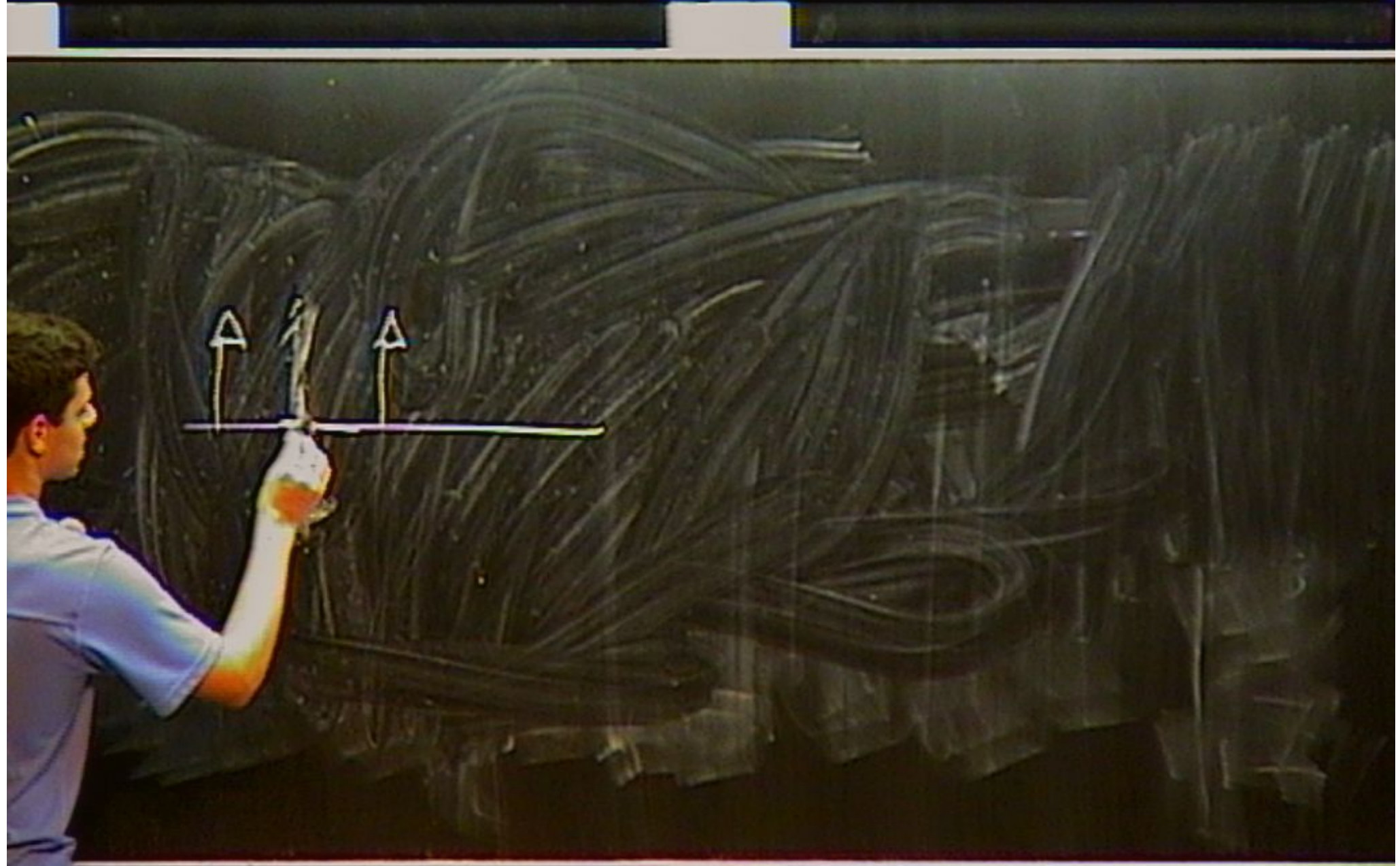


- 2-D: spins aligned in a grid



- $\infty$ -D: all spins interconnected

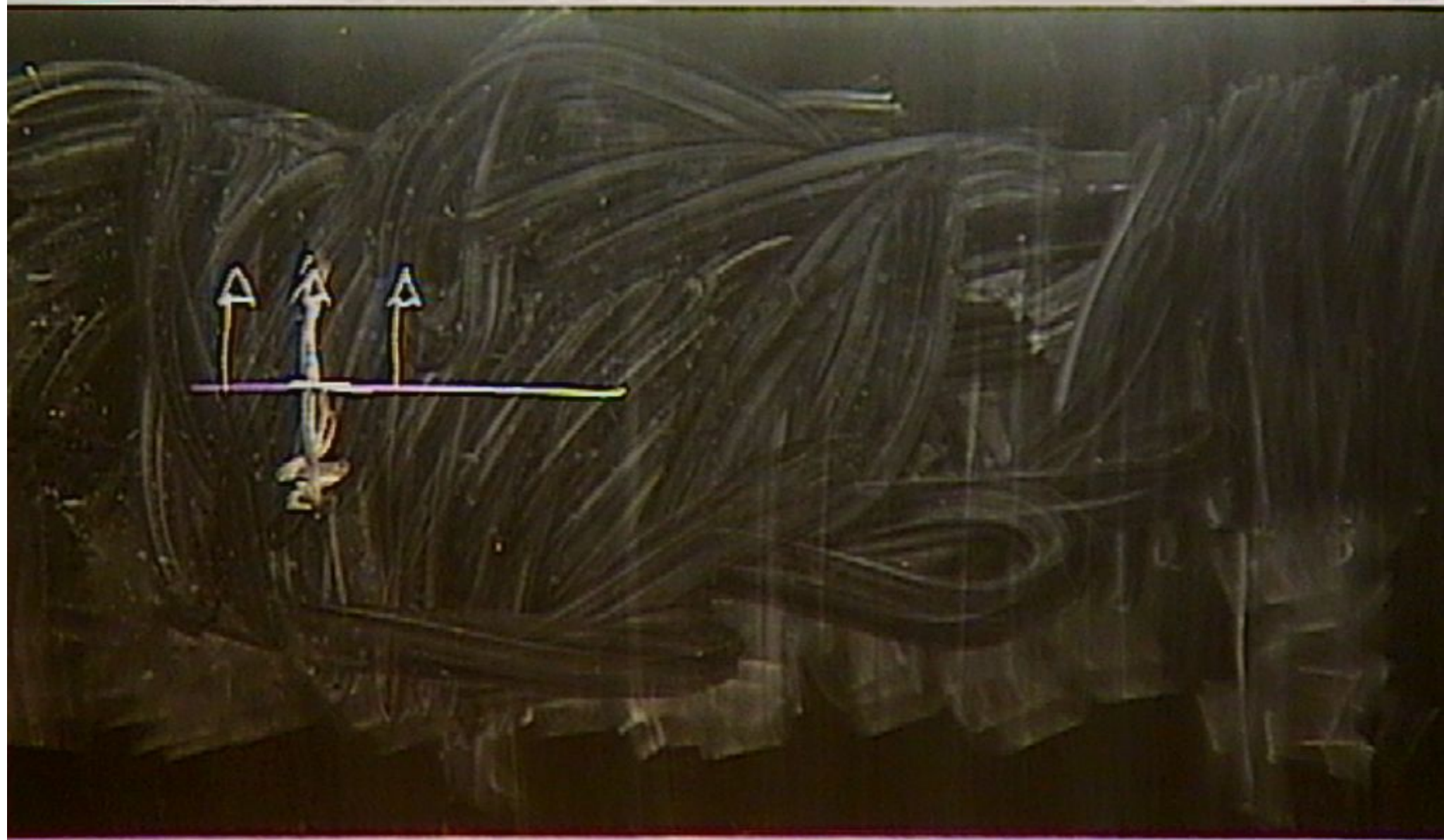












# Dimensions and the Ising Model

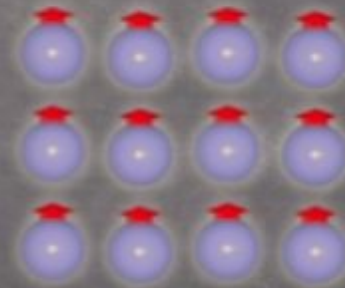
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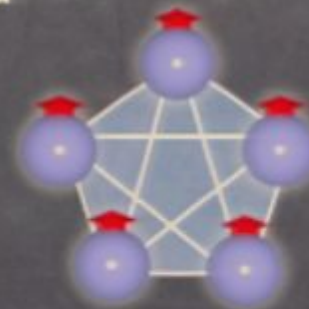
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# Decay

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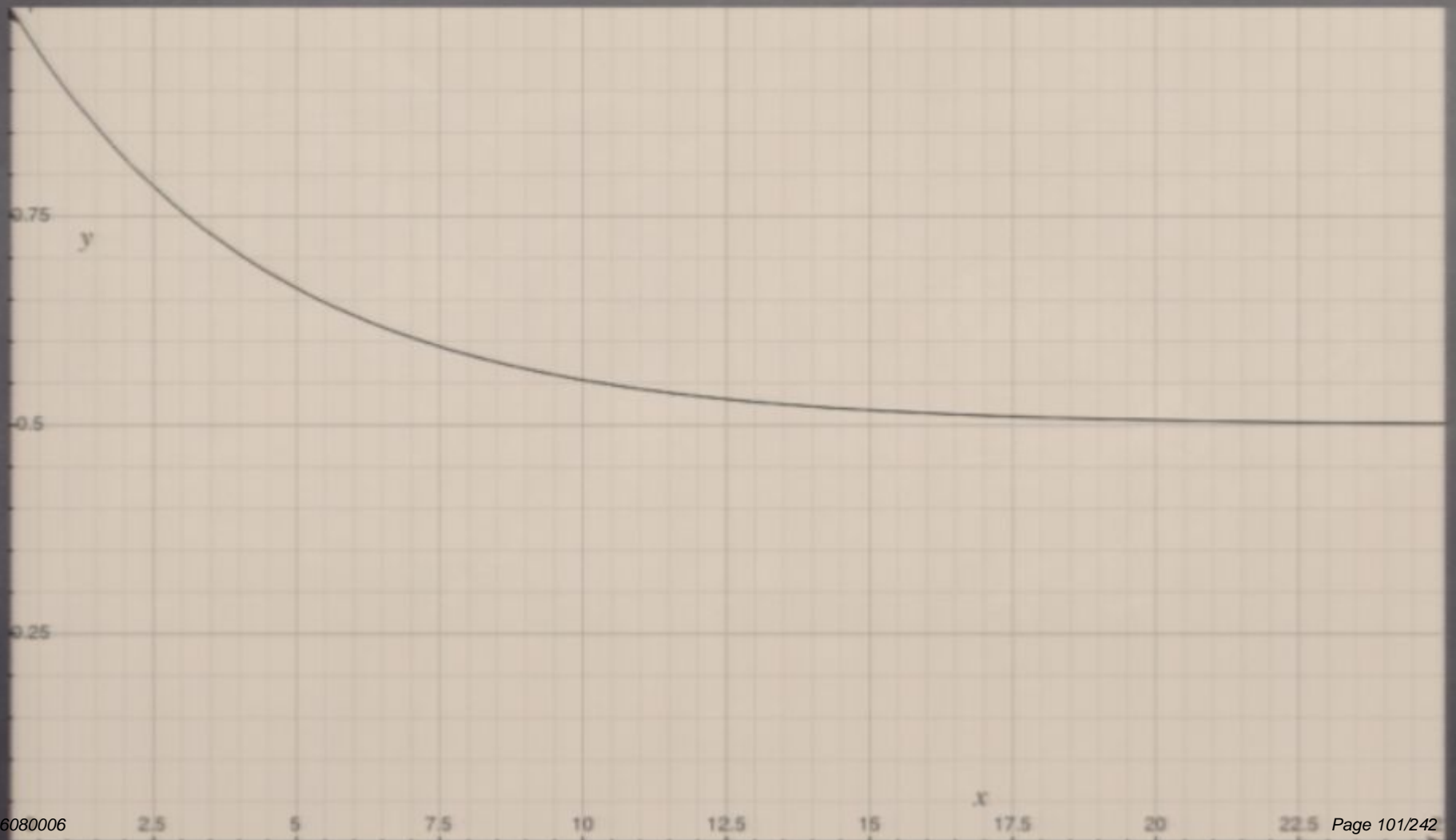
# Decay

- All memory systems undergo decay over time, mainly due to collisions with external particles.
- This decay can be modeled mathematically, as will be demonstrated momentarily.

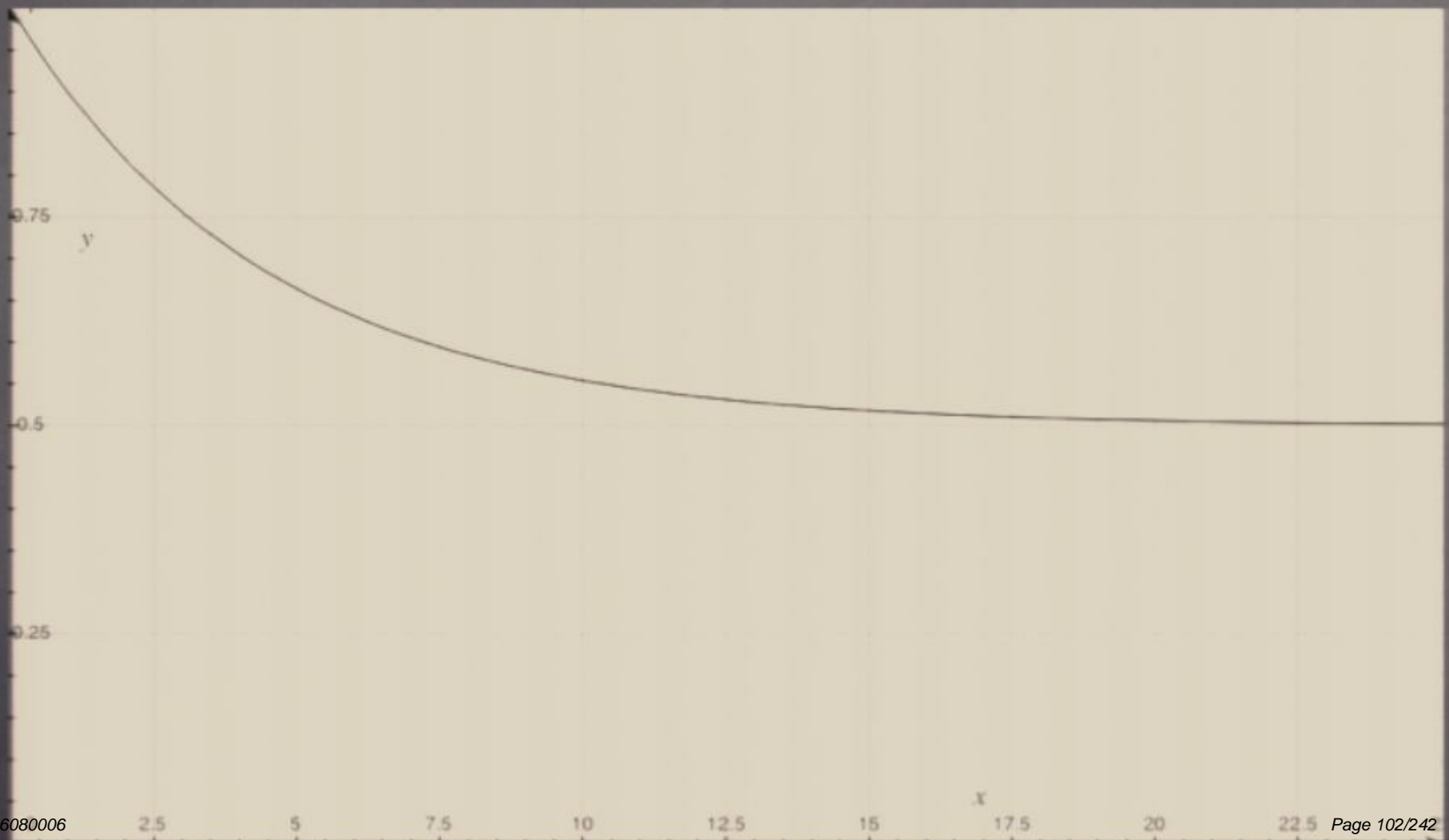


# Chance of a Single Bit Persisting in its Original State over Time

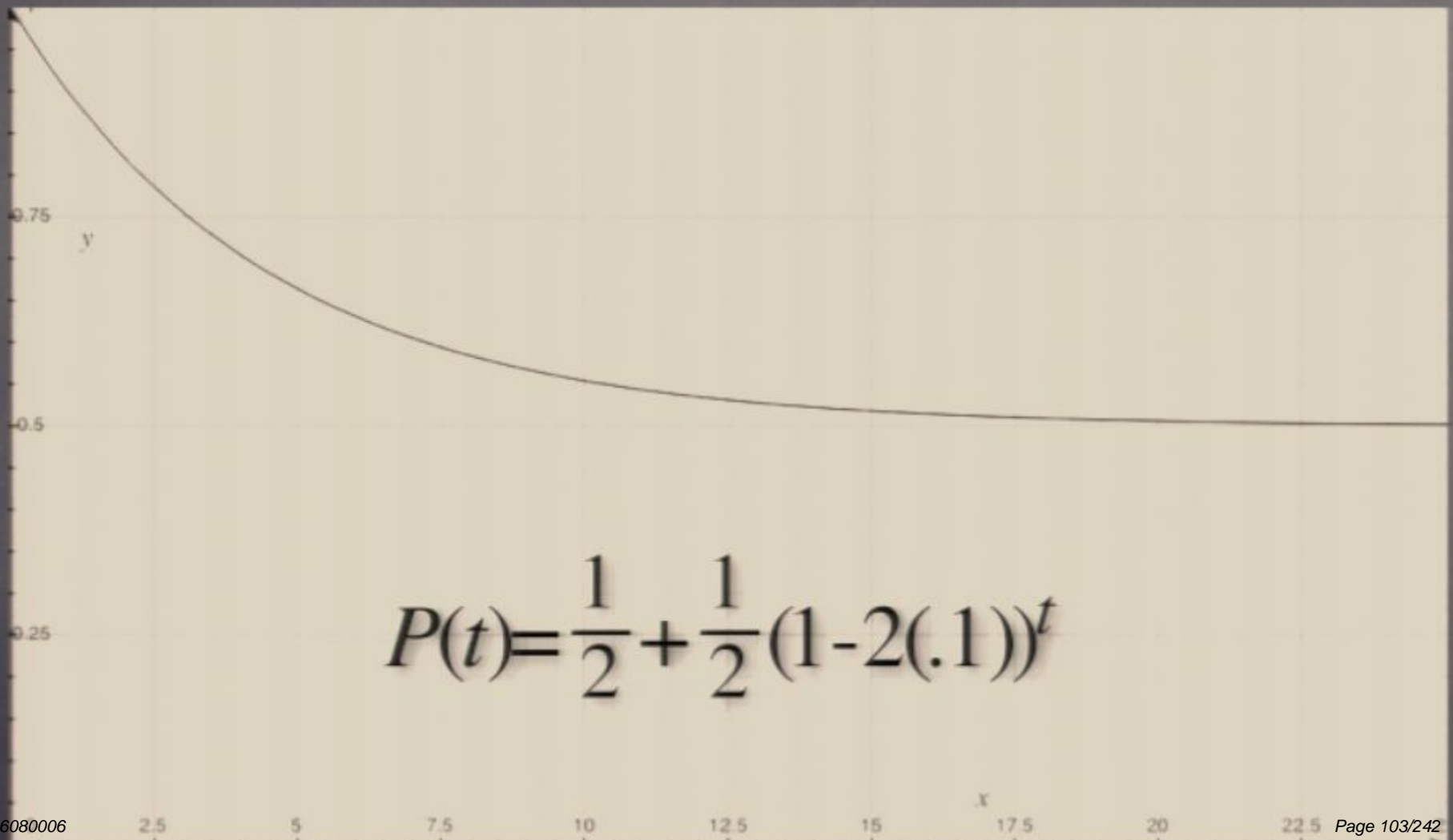
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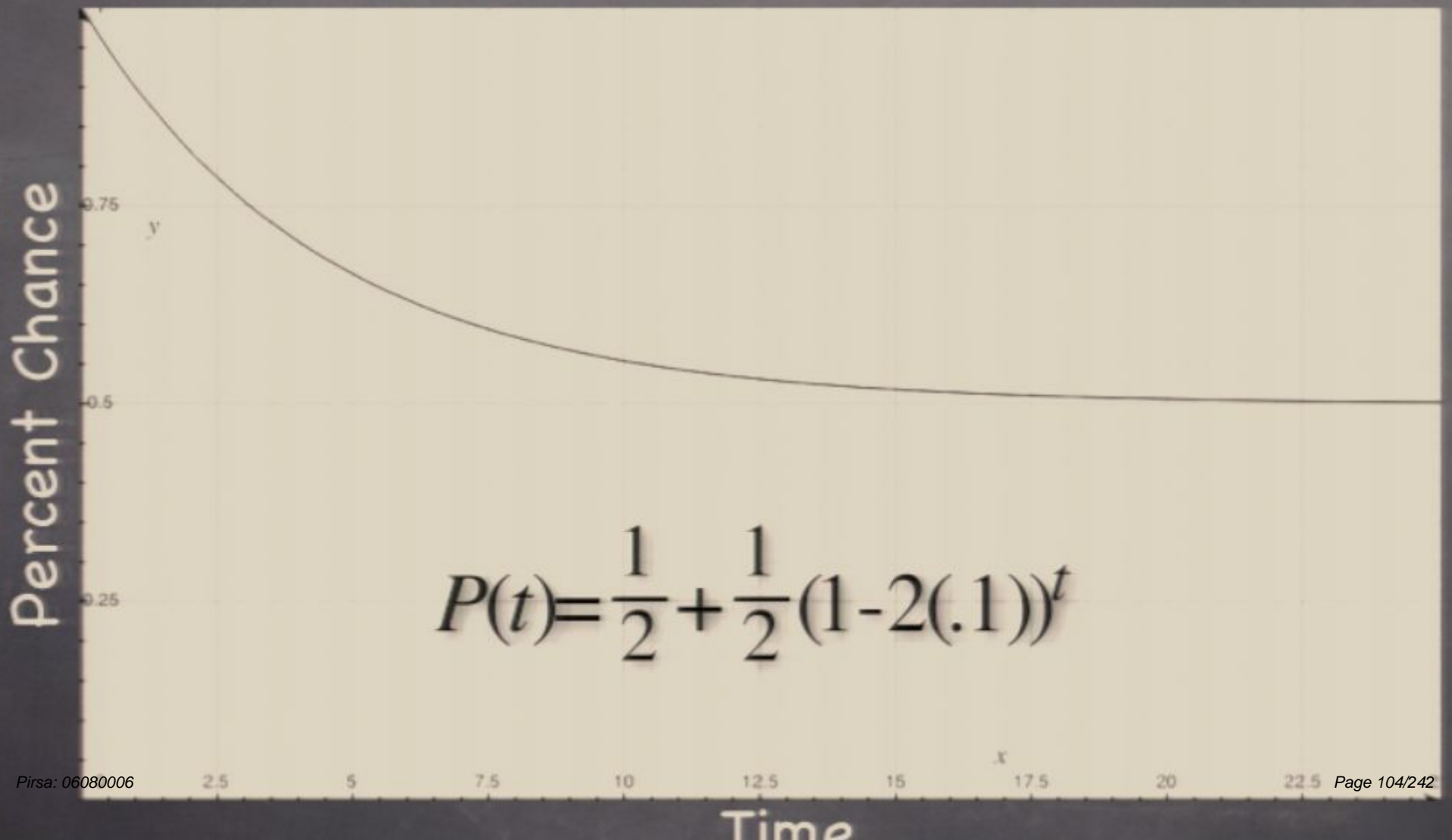
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- Determines the probability of a system being in a certain state at a certain temperature and energy.

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$$Prob(state) = \frac{1}{Z} e^{\frac{-E(state)}{kT}}$$

- Determines the probability of a system being in a certain state at a certain temperature and energy.
- Therefore, the lower the energy of a state, the more likely it is to exist in that state, potentially causing an error.

a n t u m M e m o

# Quantum Memory

# Quantum Memory

- Standard unit of memory: qubit.
- Qubit is based on the quantum state of one particle, e.g. the polarization of a photon, the spin of an electron, etc.
- Current quantum memory systems hold information for about 1 second.
- One goal of quantum information is to achieve reliable and high capacity memory systems.





# Entanglement

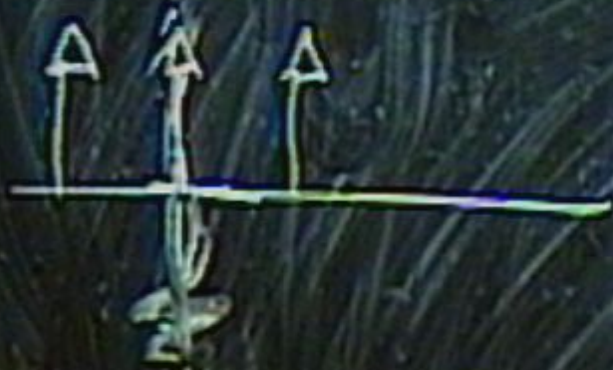
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- The key to faster computations with quantum memory.
- Allows for teleportation, superdense coding, solving “hard” mathematical problems (factoring), etc.
- Without entanglement, the qubit would be equivalent to classical bits.



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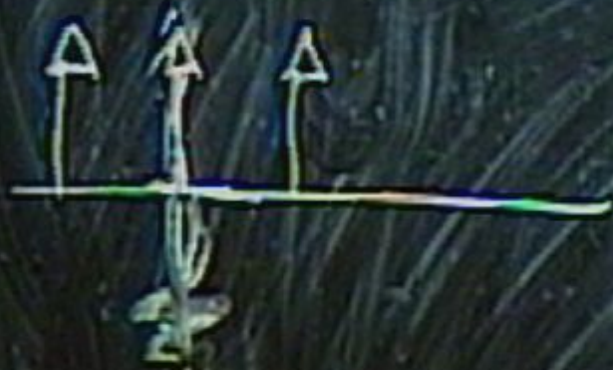
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$$a \uparrow \uparrow \uparrow + b \downarrow \downarrow \downarrow$$



$$a | \uparrow \rangle + b | \downarrow \rangle$$



$$a | \uparrow \rangle + b | \downarrow \rangle$$



$$a \uparrow \uparrow \uparrow + b \downarrow \downarrow \downarrow$$



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- The genuinely random nature of subatomic particles creates many possibilities for error.
- X Error
- Y Error
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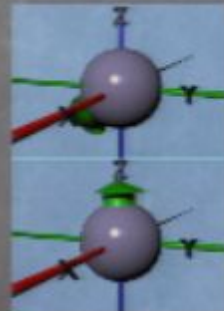
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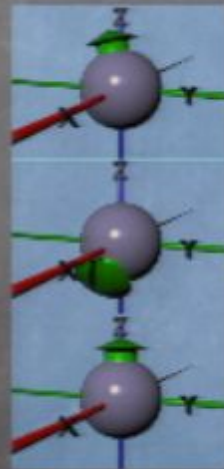
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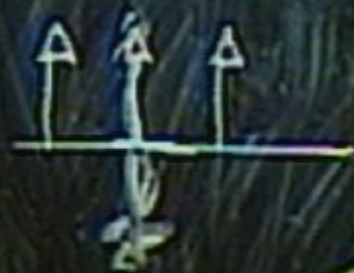
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$$a \uparrow \uparrow \uparrow + b \downarrow \downarrow \downarrow$$

xor  
 zero

$$a \downarrow \downarrow \downarrow +$$

$$a \downarrow \downarrow \downarrow -$$

$$b \uparrow \uparrow \uparrow$$

$$b \downarrow \downarrow \downarrow$$



$$a \uparrow \uparrow \uparrow + b \downarrow \downarrow \downarrow$$

$x \text{ error}$   $a \downarrow \downarrow \downarrow + b \uparrow \uparrow \uparrow$   
 $z \text{ error}$   $a \downarrow \downarrow \downarrow + b \downarrow \downarrow \downarrow$   
 $y \text{ error}$   $a \times \neq$



$$a \uparrow \uparrow \uparrow + b \downarrow \downarrow \downarrow$$

$x \text{ error}$   
 $z \text{ error}$   
 $y \text{ error}$

$a \downarrow \downarrow \downarrow +$   
 $a \downarrow \downarrow \downarrow -$   
 $a \neq z$

$$b \uparrow \uparrow \uparrow$$

$$b \downarrow \downarrow \downarrow$$



# Self-Healing Memory

- Passive vs. Active error correction: using nature or man-made algorithms
- Current passive healing rate is much slower than the rate of errors
- If we can raise the energy required to maintain erratic states, the probability of errors decreases.

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# 9-Quarter Current Progress

# Current Progress

- 9-Qubit Code

$$a \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right) \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right) \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right) +$$

# Current Progress

## • 9-Qubit Code

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$$b \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle - |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right) \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle - |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right) \left( \frac{|\uparrow\uparrow\uparrow\uparrow\rangle - |\downarrow\downarrow\downarrow\downarrow\rangle}{\sqrt{2}} \right)$$

## • Toric Code

# Current Progress

## • 9-Qubit Code

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## • Toric Code















Energy =  $E_0$



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Energy = 0



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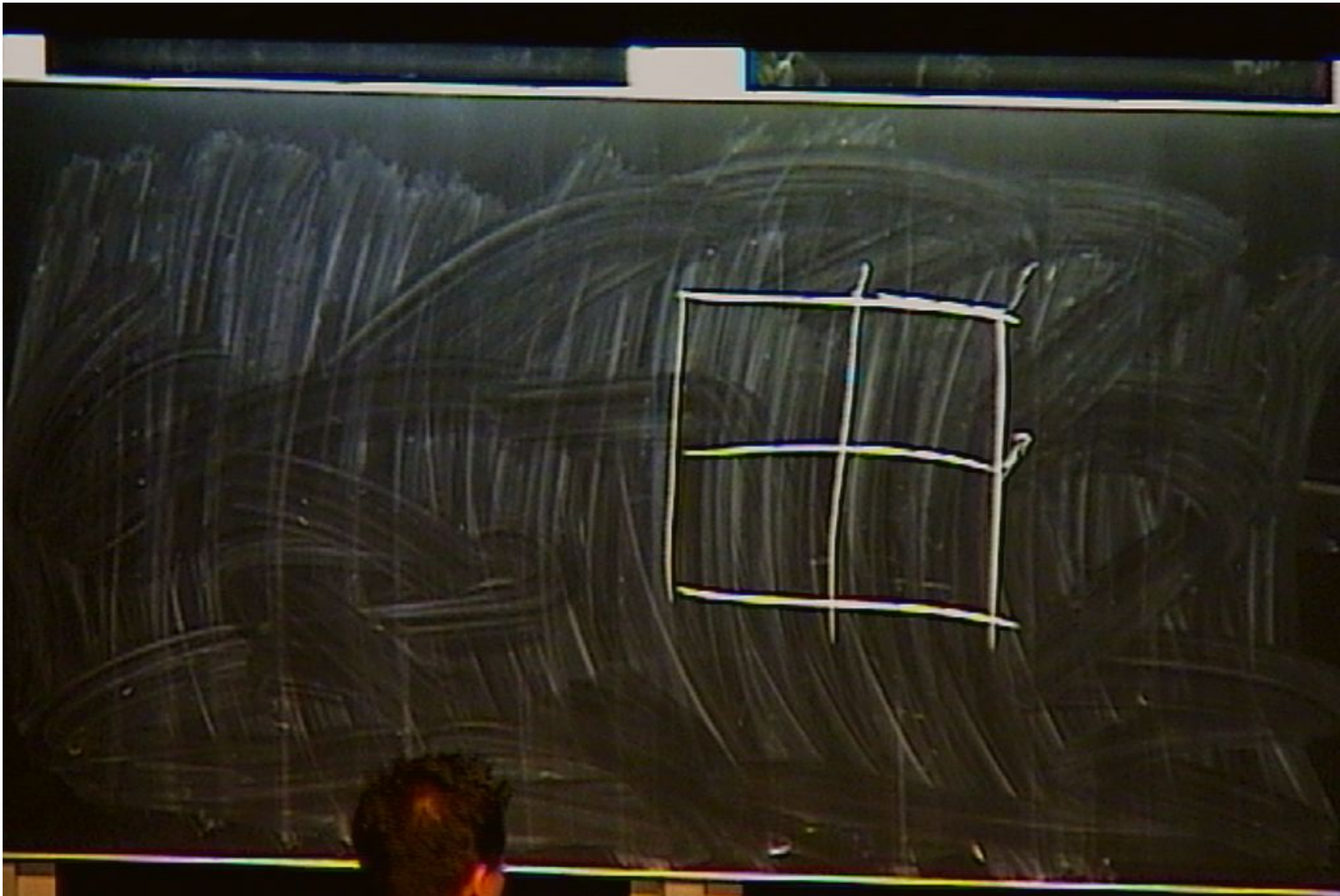
(Count of UP Spins) MOD 2



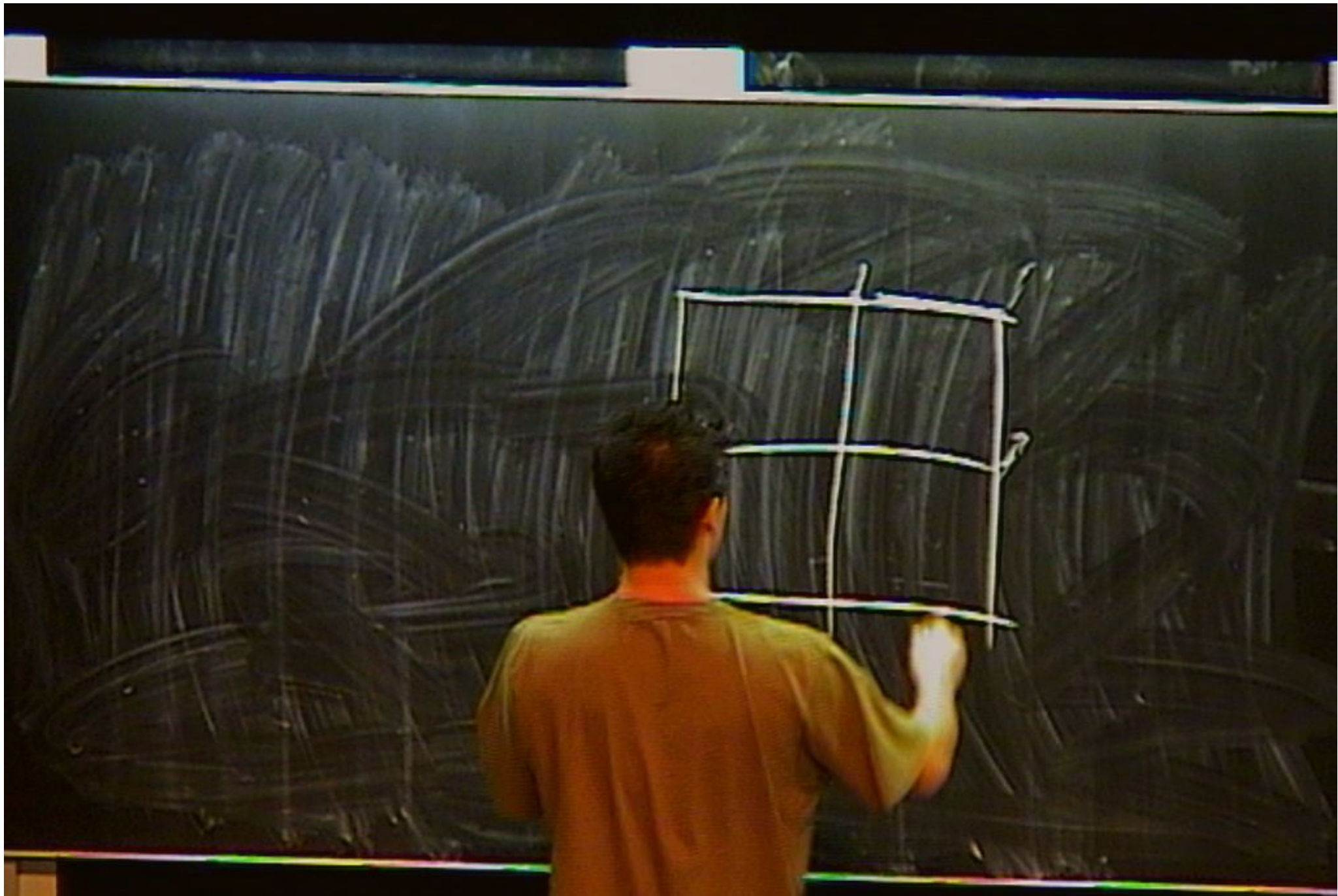
Energy = 0

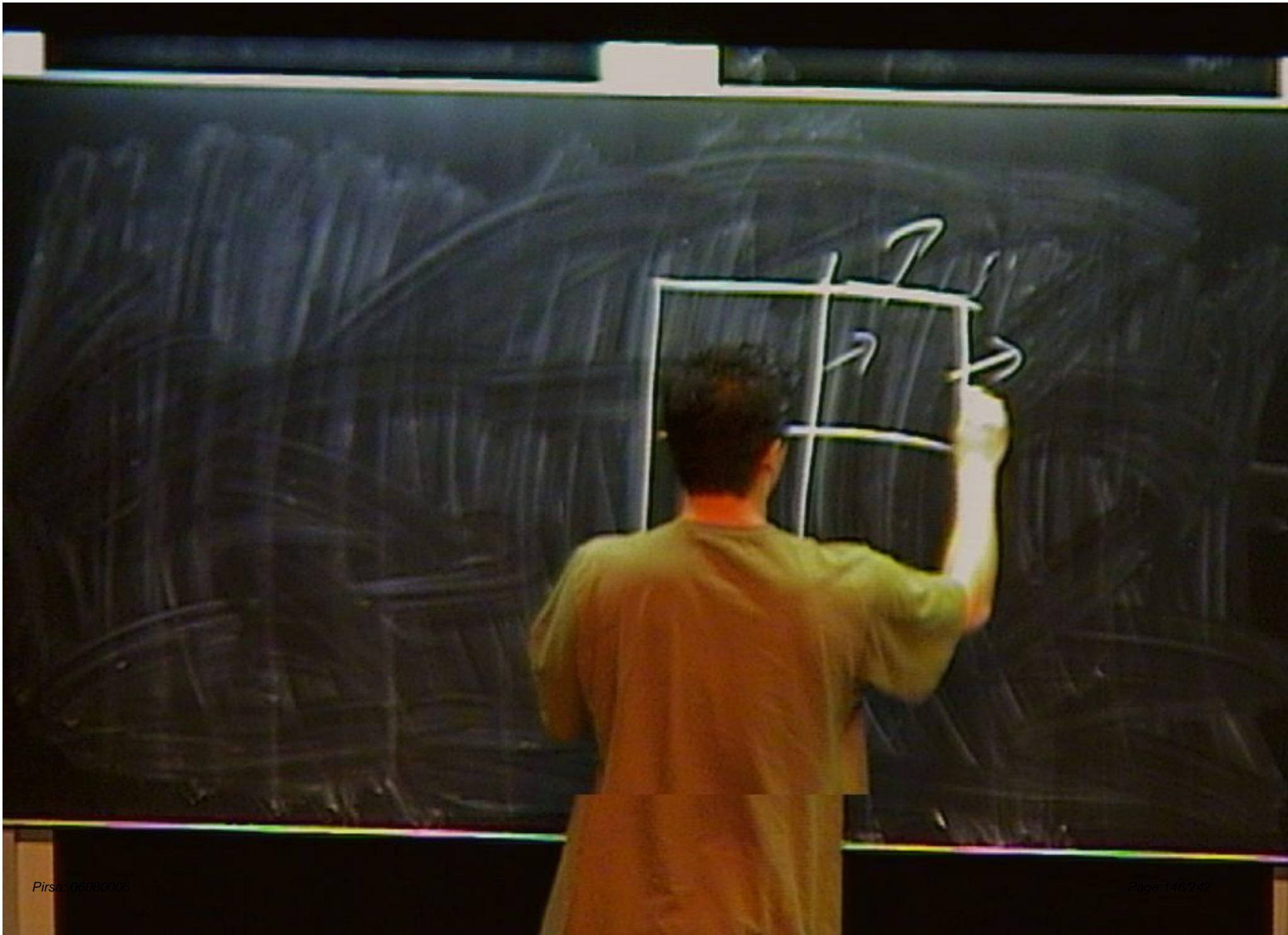


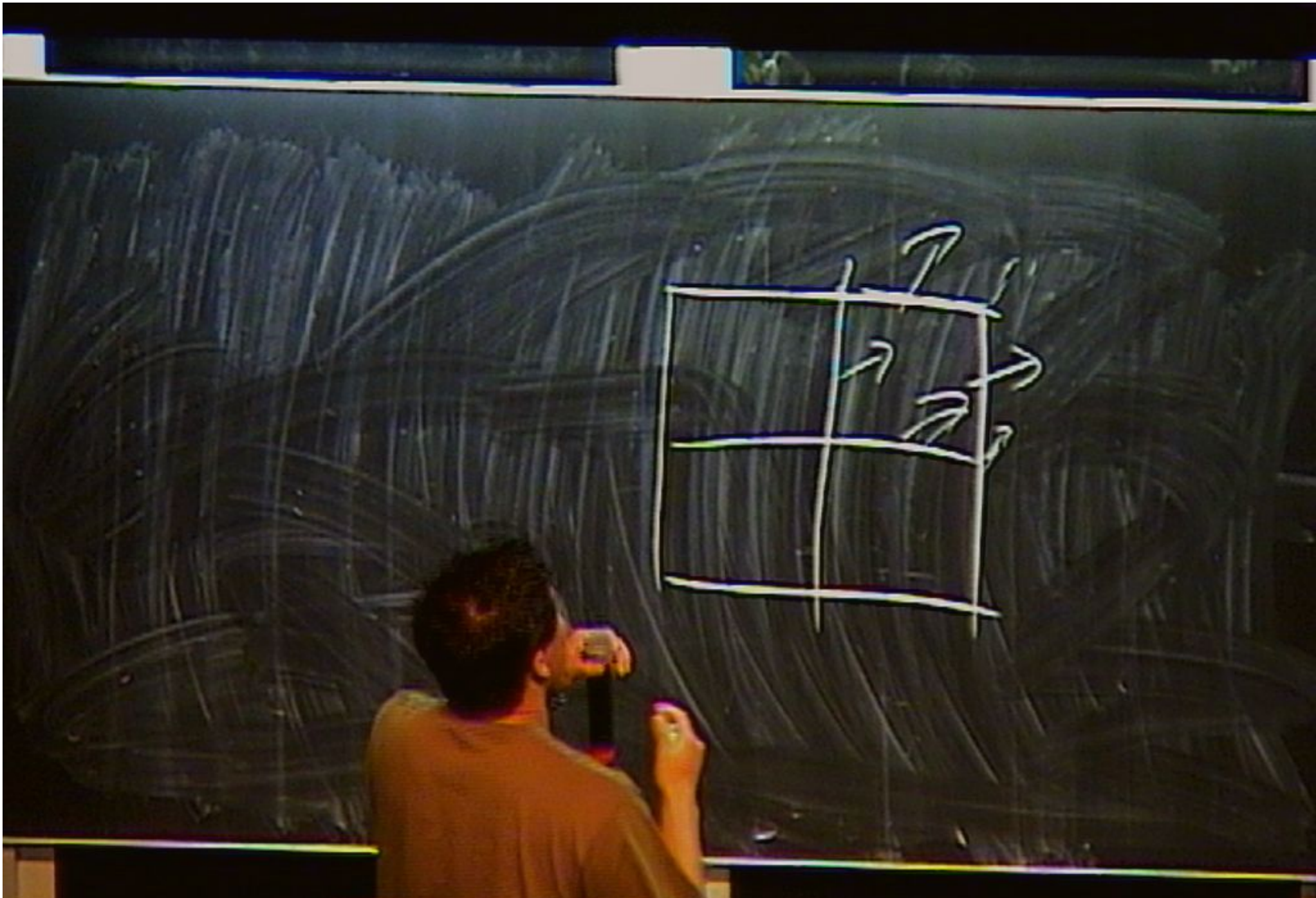
Energy = 0



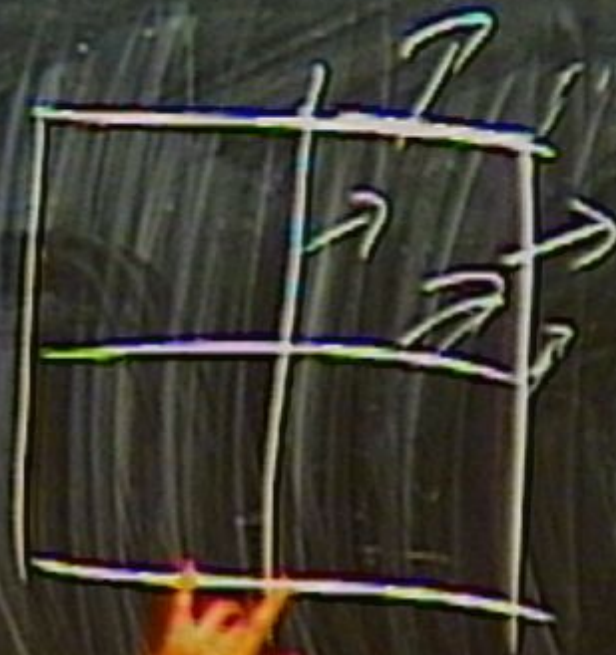












Energy =  $E_0$



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(Count of UP Spins) MOD 2

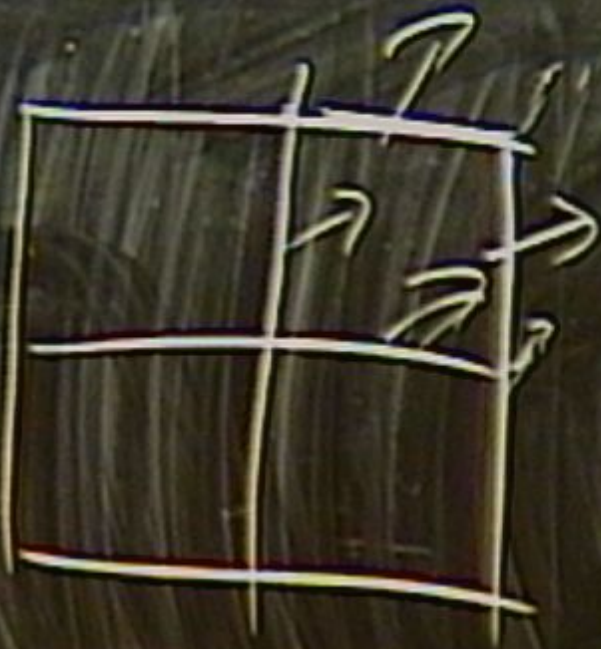


Energy = 0



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# Quest for 3-D Toric Code and Deterministic Entanglement

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- Flaws in 2-D Toric code

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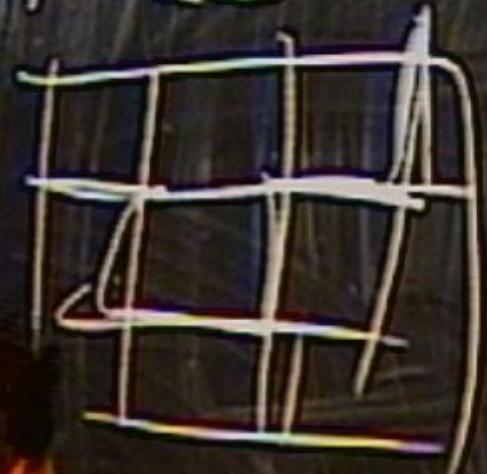
- $C$
  - Flaws in 2-D Toric code
  - 4-D Toric code is  $?$  equivalent to 2-D Ising Model
- $3$        $r_c$   
 $t$     $e a$        $i$     $d$

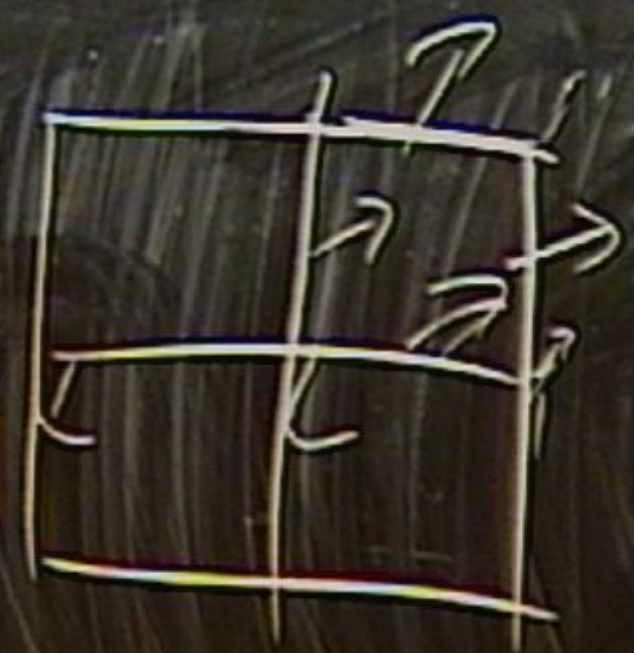
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# Quest for 3-D Toric Code and Deterministic Entanglement

- Flaws in 2-D Toric code
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- Is there a 3-D Toric Code?
- Can we create and maintain useful entanglement patterns?









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**Many thanks to  
Dr. Carlos Mochon**

**Major props to  
STIMAC**







aqua cool \*

No Signal

VGA-1

No Signal

VGA-1

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Schrodinger's equation:

$$i\hbar \frac{\partial}{\partial t} \psi(x,t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x,t) + V(x) \psi(x,t)$$

take:  $\hbar = 1 \Rightarrow$

$\Rightarrow$  wave function:

$$i \frac{\partial}{\partial t} \psi(x,t) = -\frac{1}{2m} \frac{\partial^2}{\partial x^2} \psi(x,t) + V(x) \psi(x,t)$$

$+0 \Rightarrow$  general formula for waves:

$$\psi(x,t) = R(x,t) e^{iS(x,t)}$$

Movement formula:

$$\frac{\partial x}{\partial t}(t) = \frac{1}{m} \frac{\partial S}{\partial x}(x,t) \quad (\text{by Lagrange})$$

$$\Rightarrow m \cdot \frac{\partial^2 x}{\partial t^2} = -\frac{\partial}{\partial x} (V(x) + Q)$$

Initial momentum is determined by the incident wave function,

$$p = \partial s / \partial z = p_0$$

In practice, however, we do not control the initial location (so, though it goes through a definite slit, we cannot predict which slit it will be.)

Acted on by "Quantum-mechanical" potential; or  $U_0$

$$U = (-\hbar^2/2m) \nabla^2 R/R$$

← case-simplified equations

Because  $U$  becomes infinite when  $R$  becomes zero (making a repelling force) we have now got a simple and accurate model of why particles are never found where the wave-function vanishes.

If one slit is closed, the  $\psi$ -field is changed and so the particle can now reach some places where it could not go when both slits were open.

THEREFORE the slit can only affect the motion of the particle indirectly, through an effect on the Schrodinger  $\psi$ -field.

$U(x)$  acting on a particle depends on a wave intensity,  $P(x)$ ; also numerically equal to a probability density. So the wave function can be interpreted both as a force and a probability density.

Source



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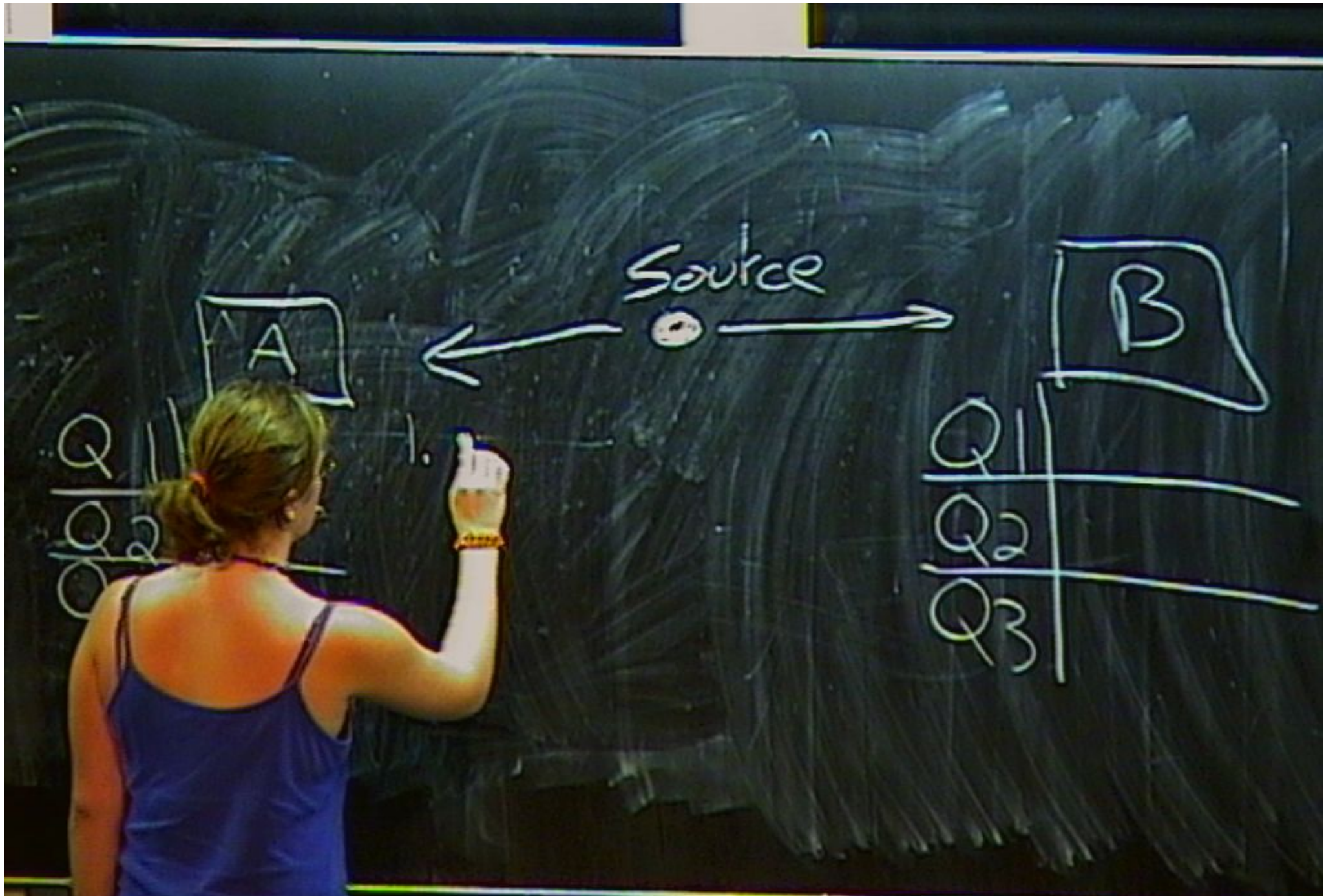
Because  $U$  becomes infinite when  $R$  becomes zero (making a repelling force) we have now got a simple and accurate model of why particles are never found where the wave-function vanishes.

If one slit is closed, the  $\psi$ -field is changed and so the particle can now reach some places where it could not go when both slits were open.

THEREFORE the slit can only affect the motion of the particle indirectly, through an effect on the Schrodinger  $\psi$ -field.

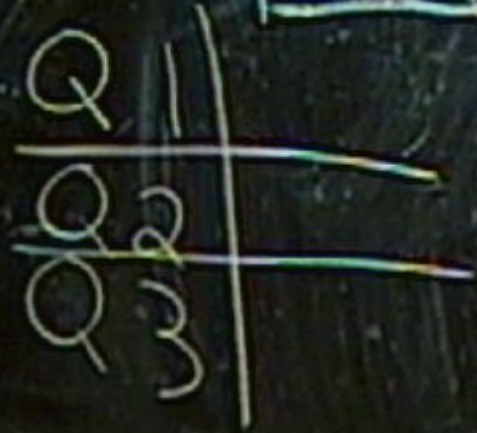
$U(x)$  acting on a particle depends on a wave intensity  $I$   $P(x)$ ; also numerically equal to a probability density. So the wave function can be interpreted both as a force and a probability density.







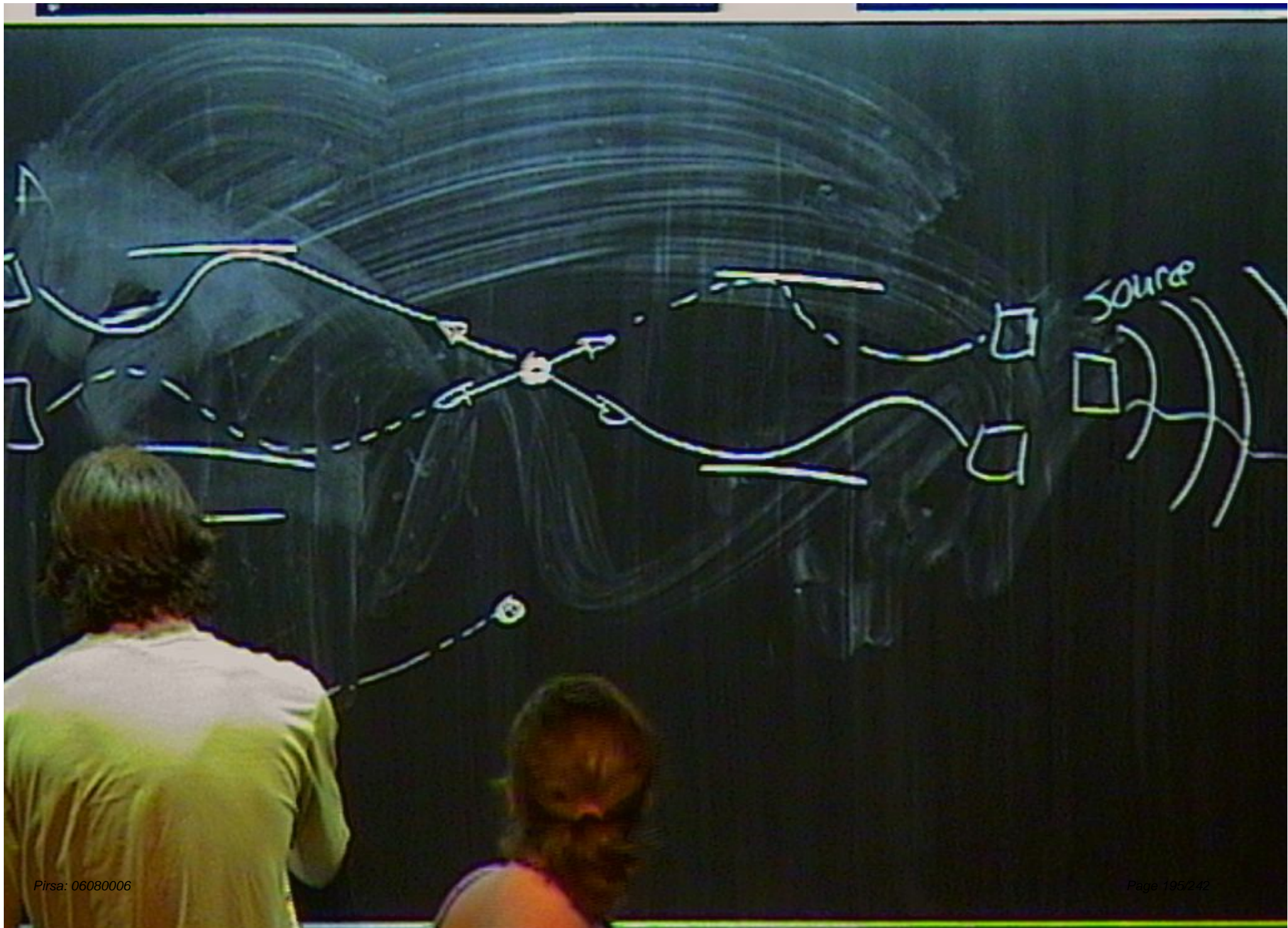
Source



1. Same questions  $\Rightarrow$  Same answer



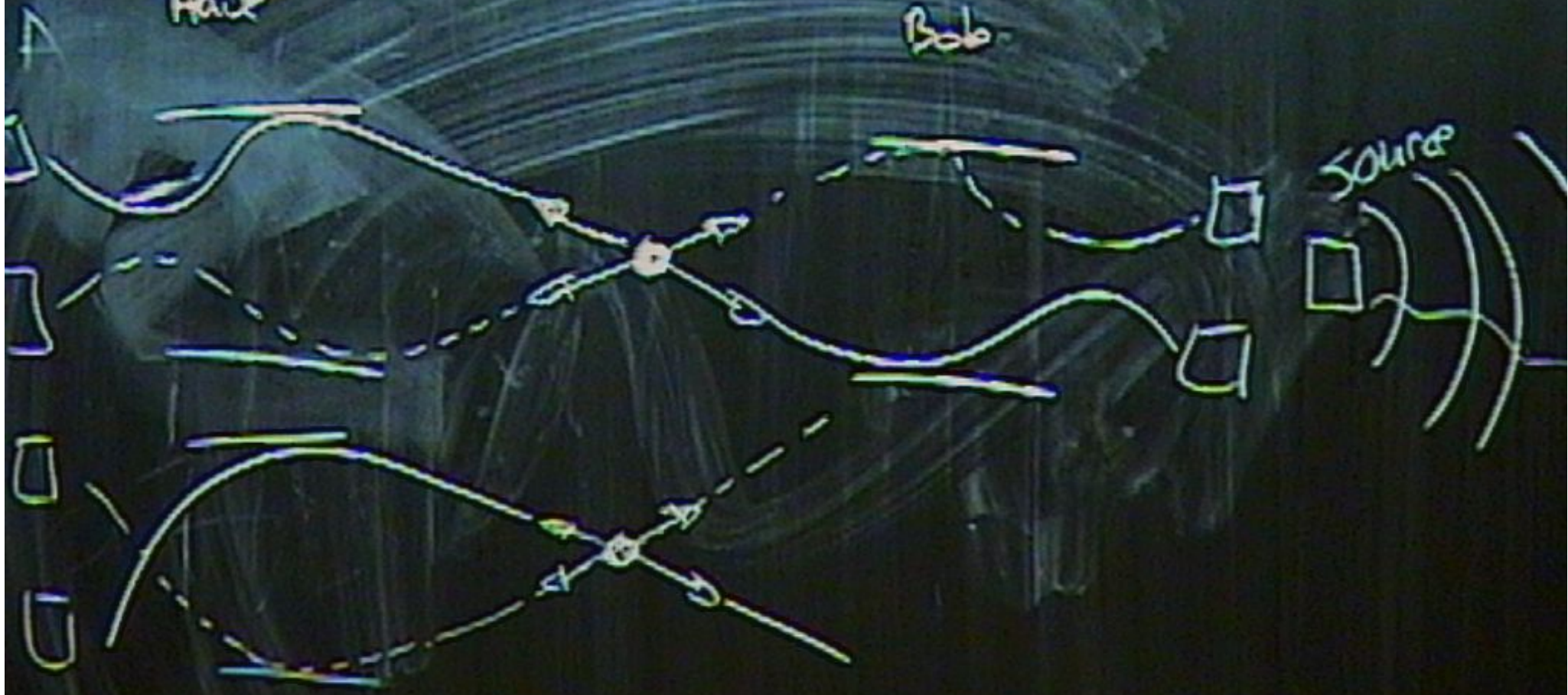
2. # Agree = # Disagree

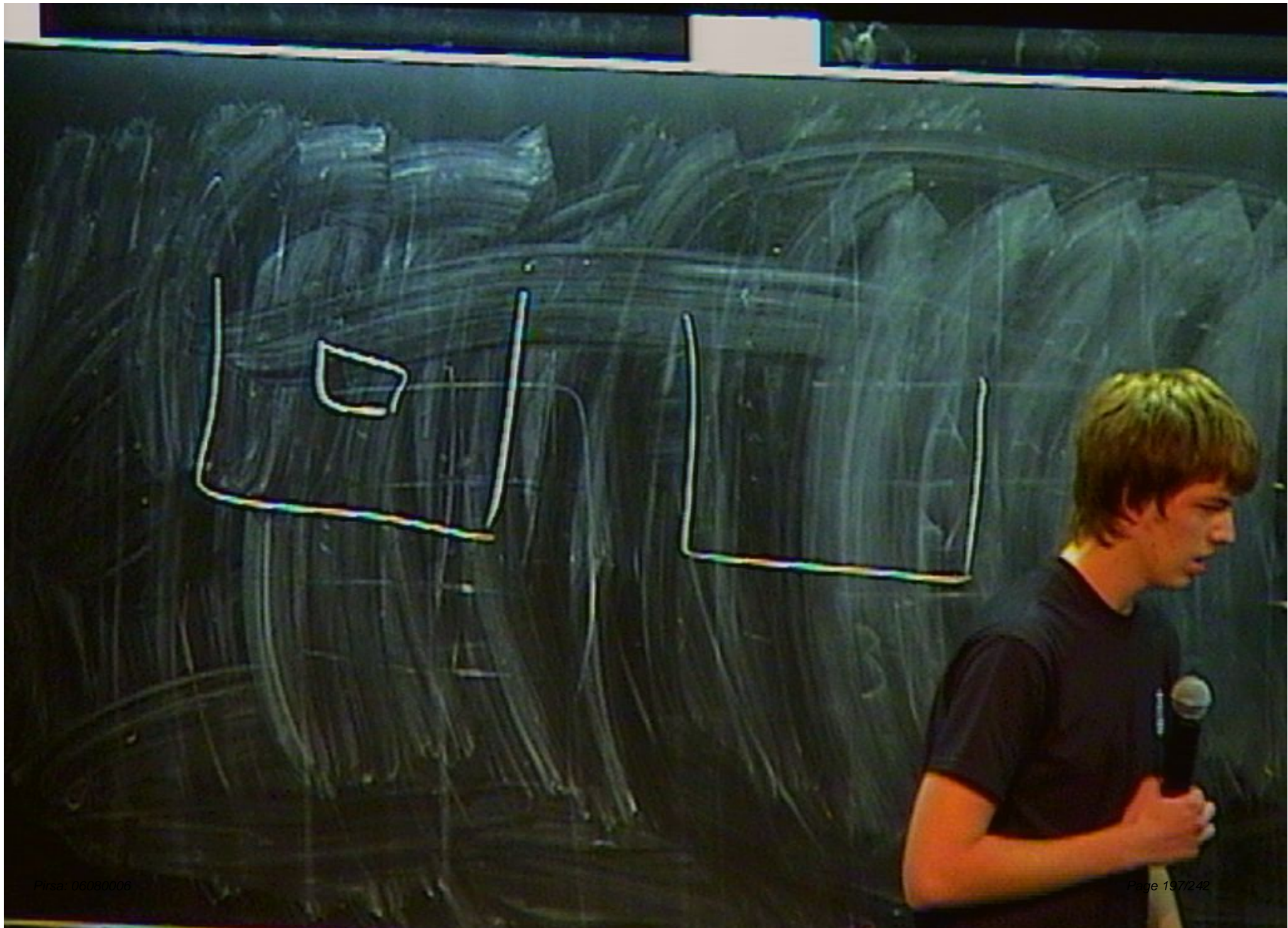


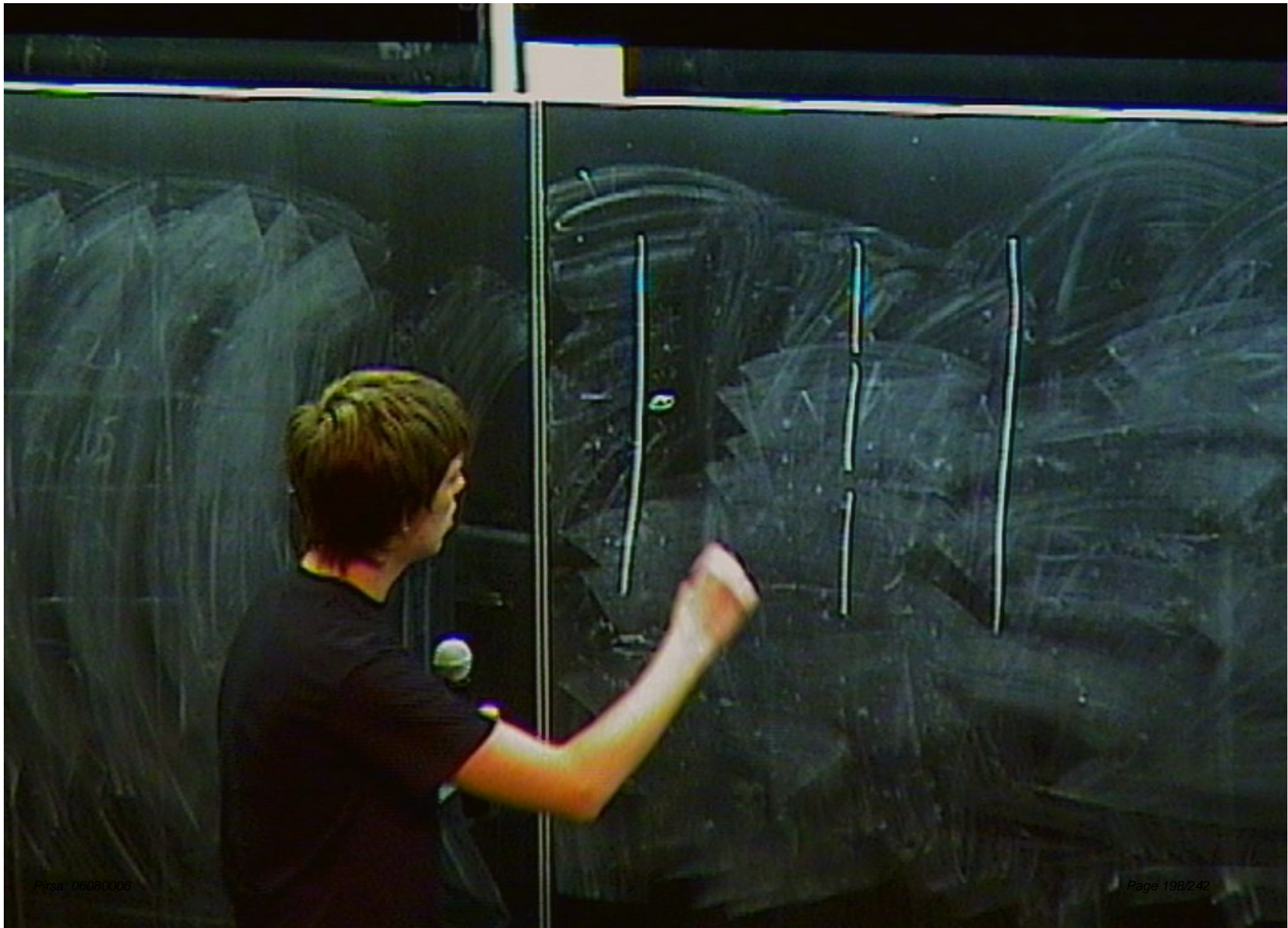
Alice

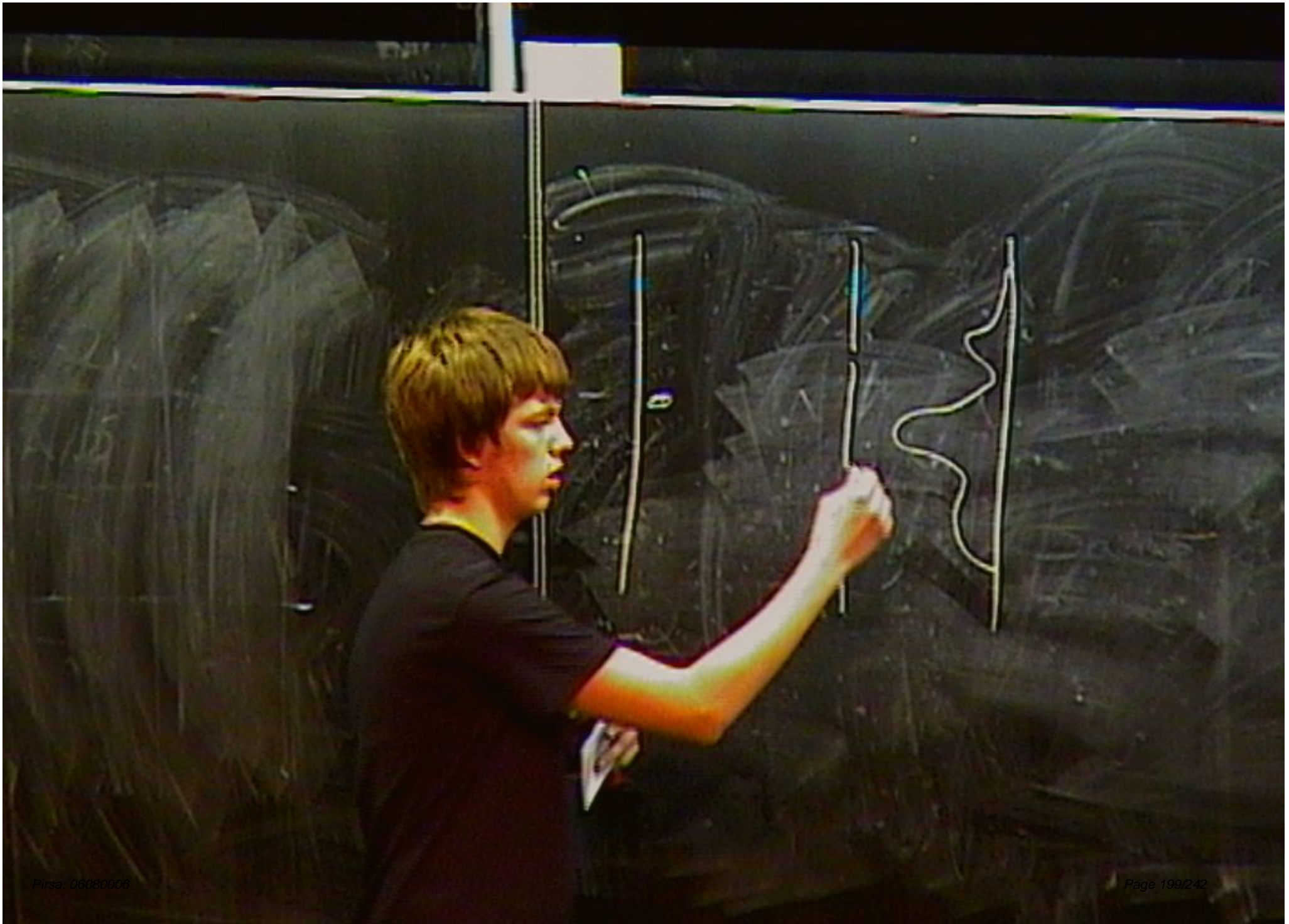
Bob

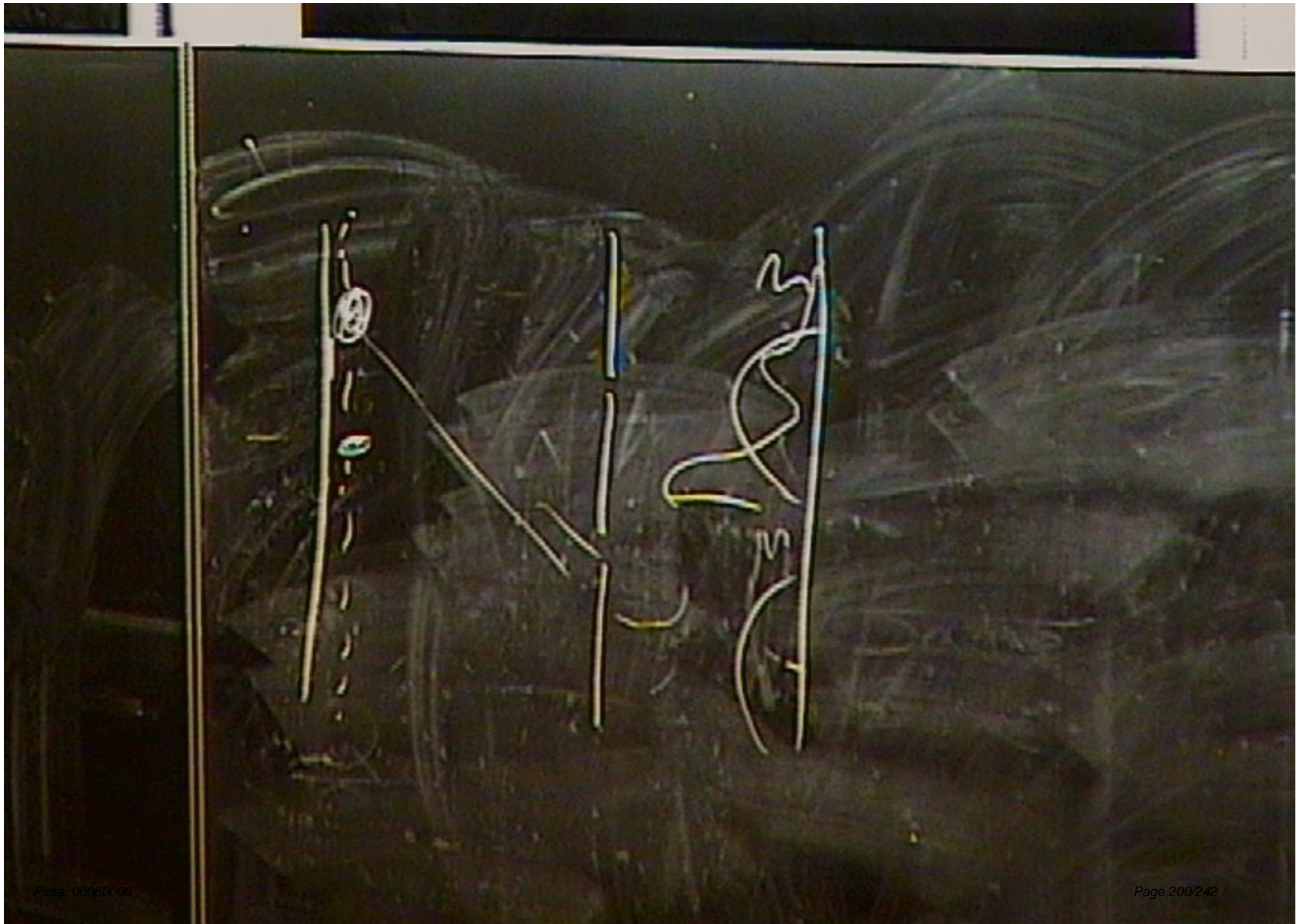
Source



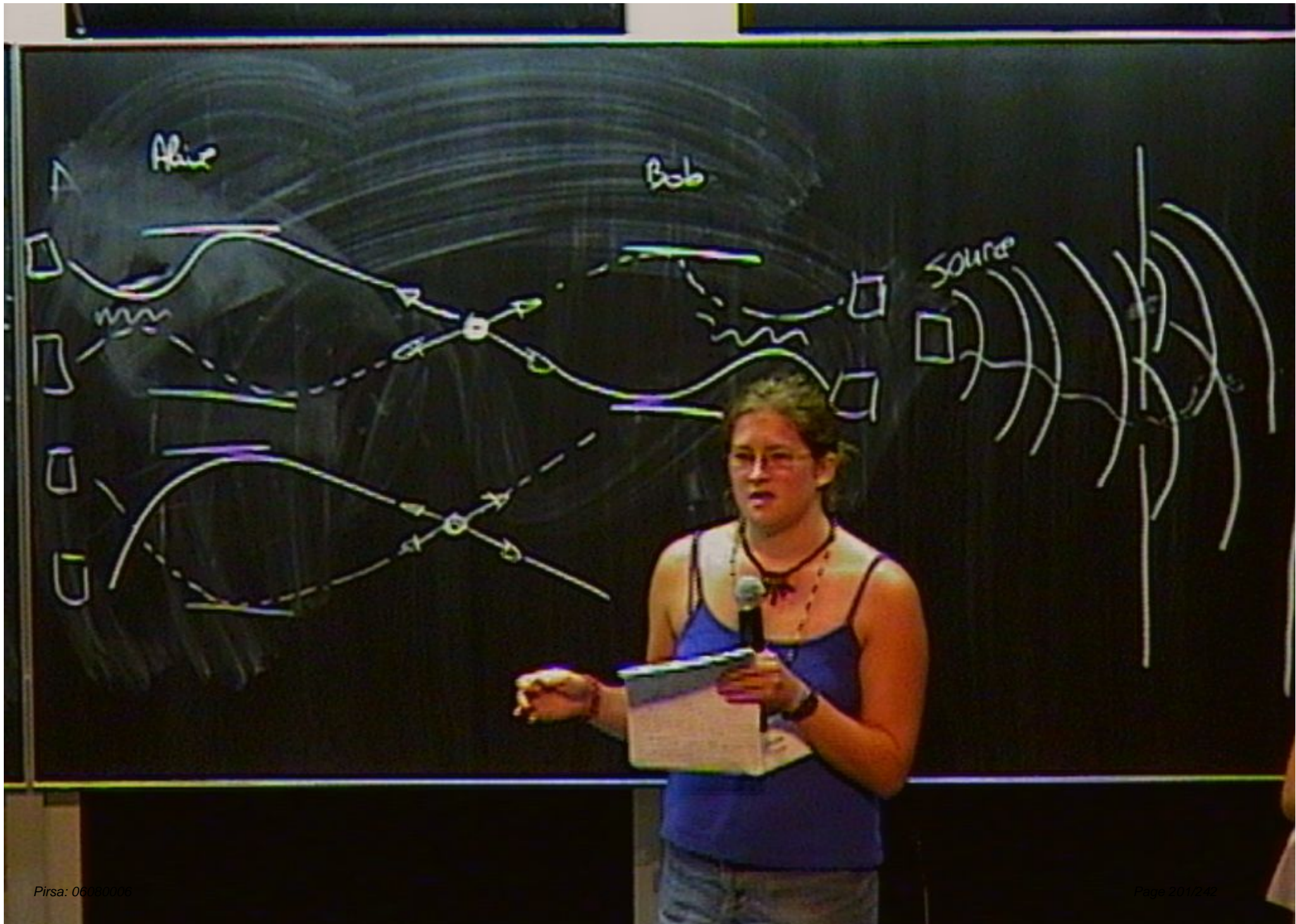


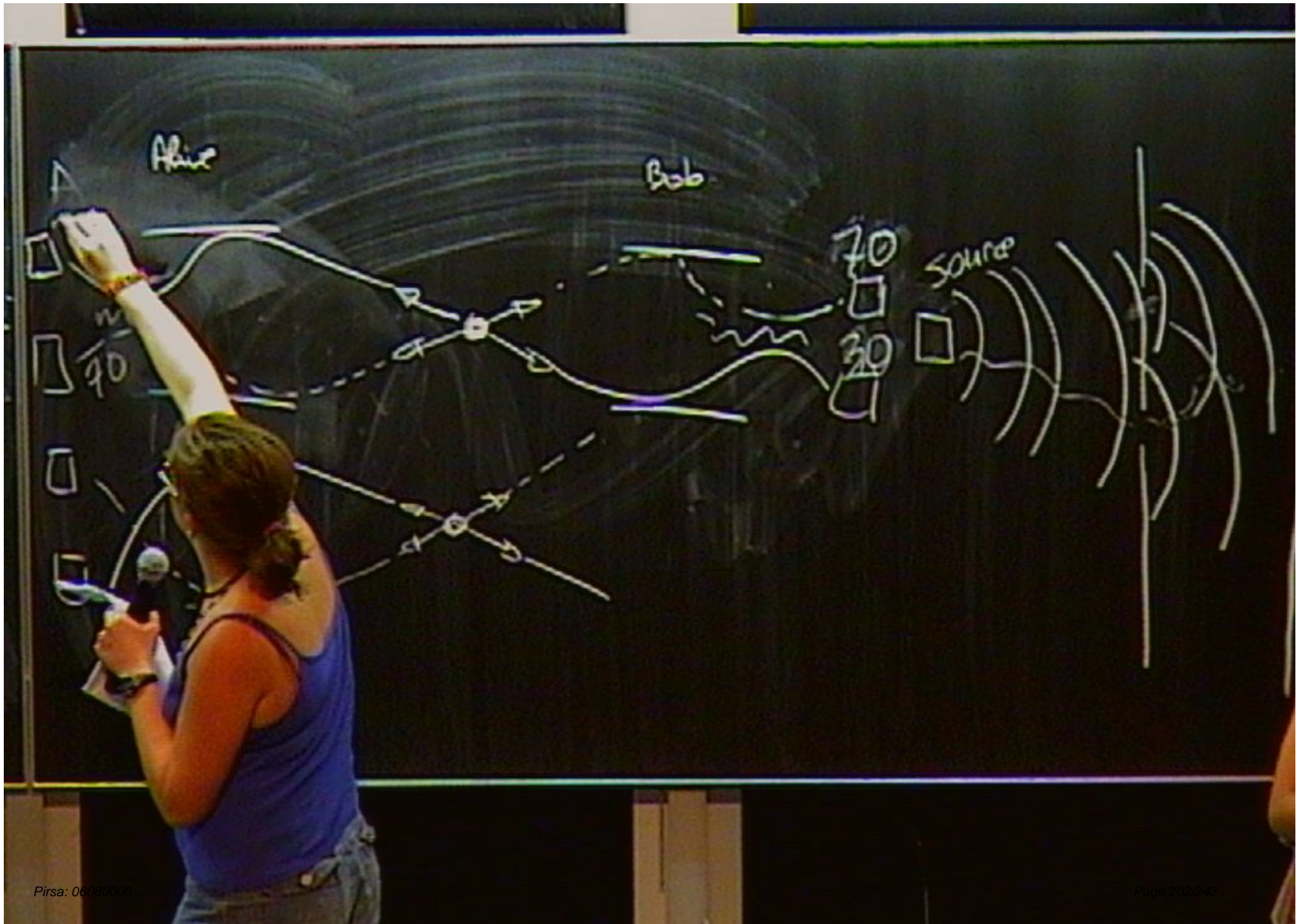


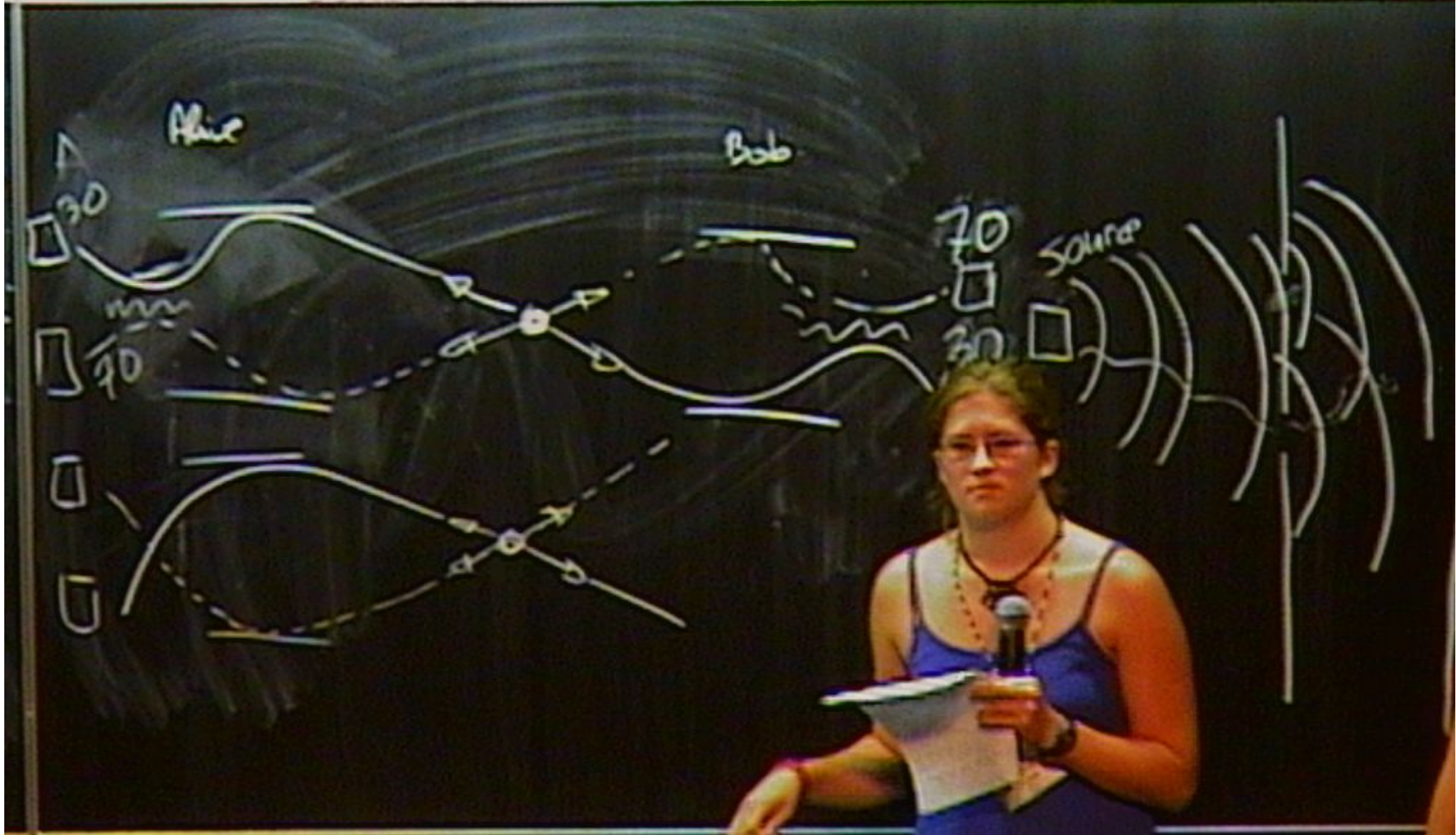




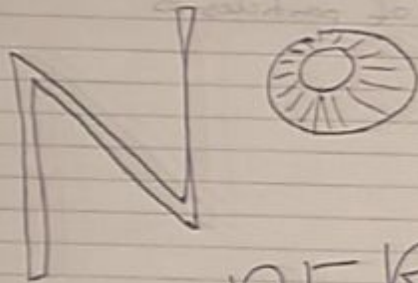








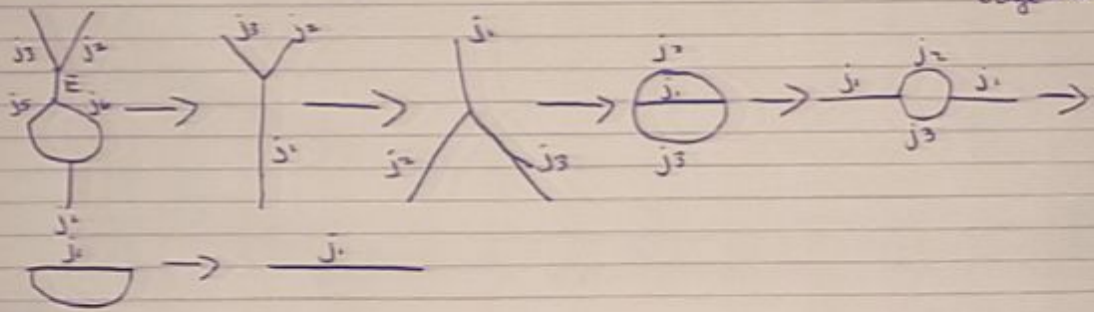
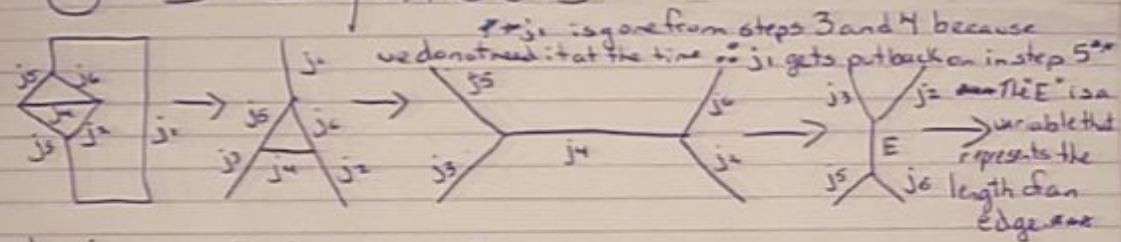
most importantly,



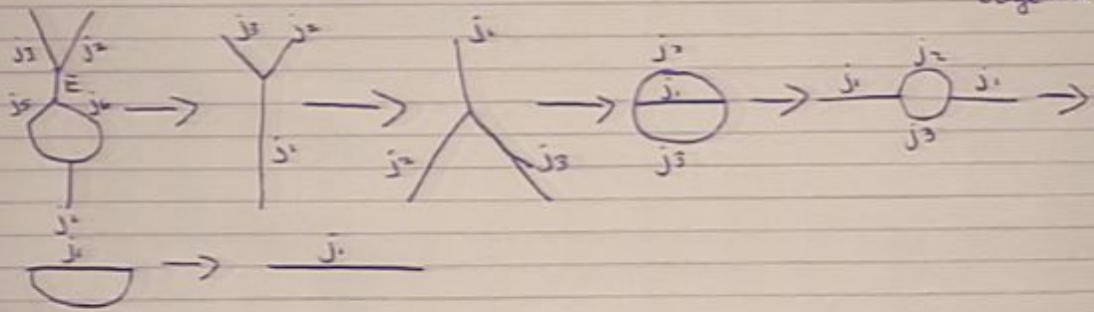
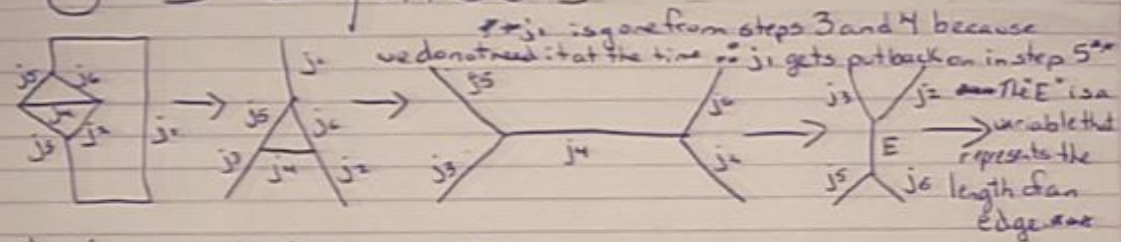
SUPERPOSITION!

yay!

# 6 J Symbols

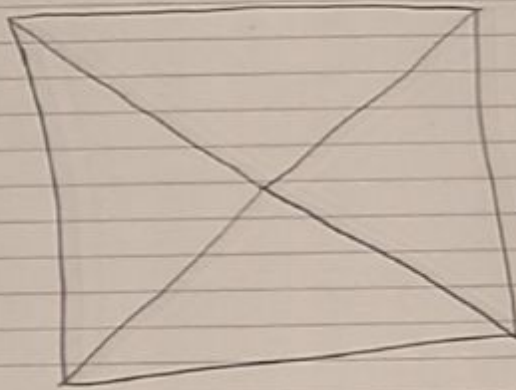


# 6 J Symbols

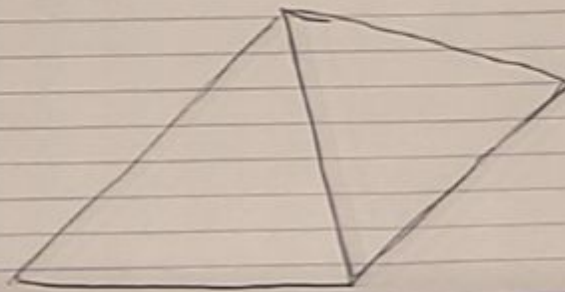


Tetrahedron

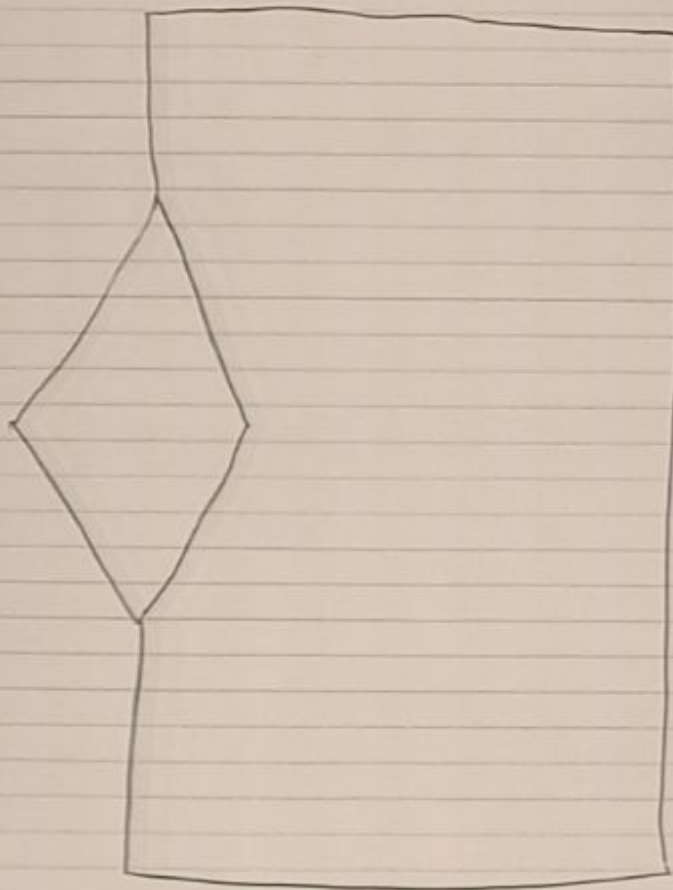
Top



Side

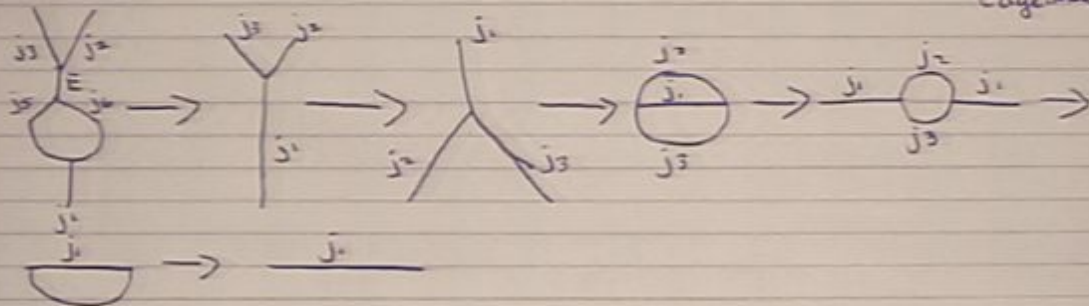
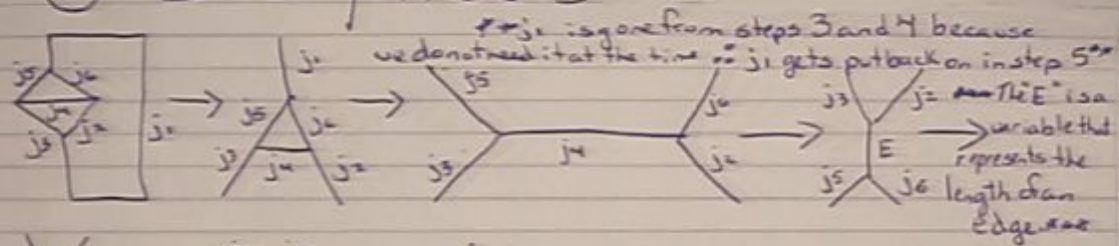


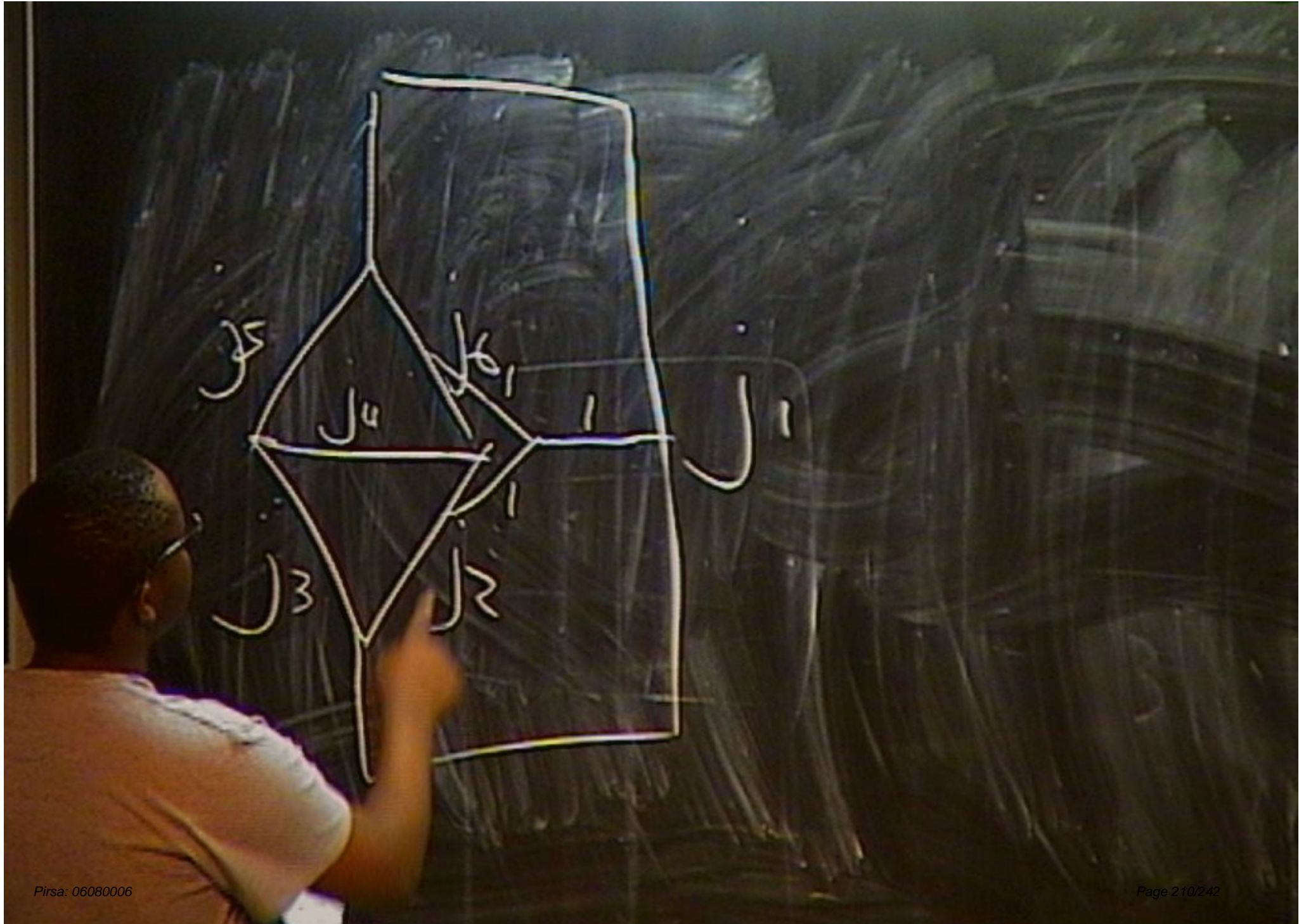
# Projection of a tetrahedron

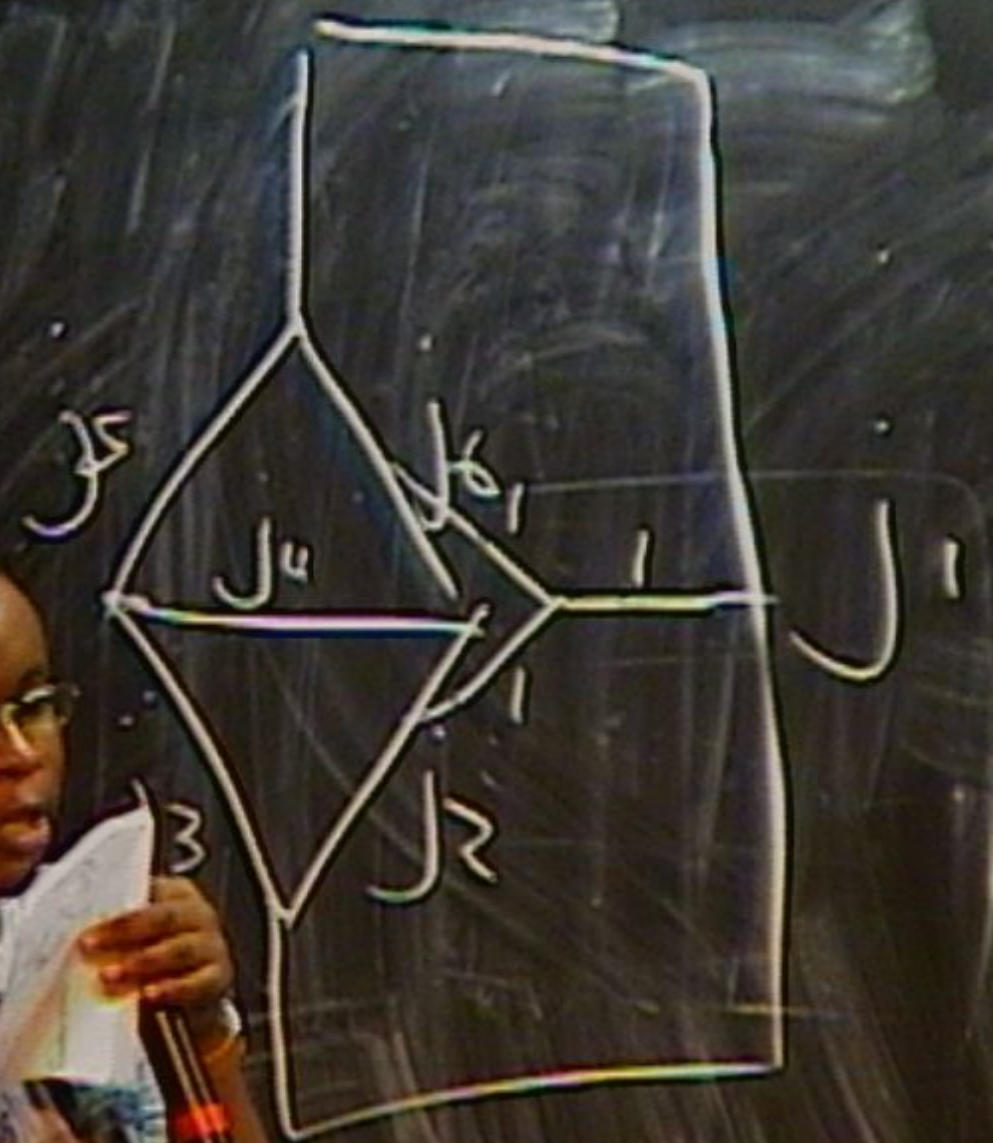




# 6 J Symbols

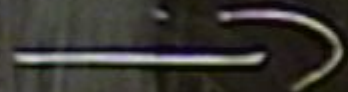
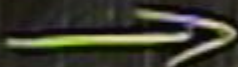


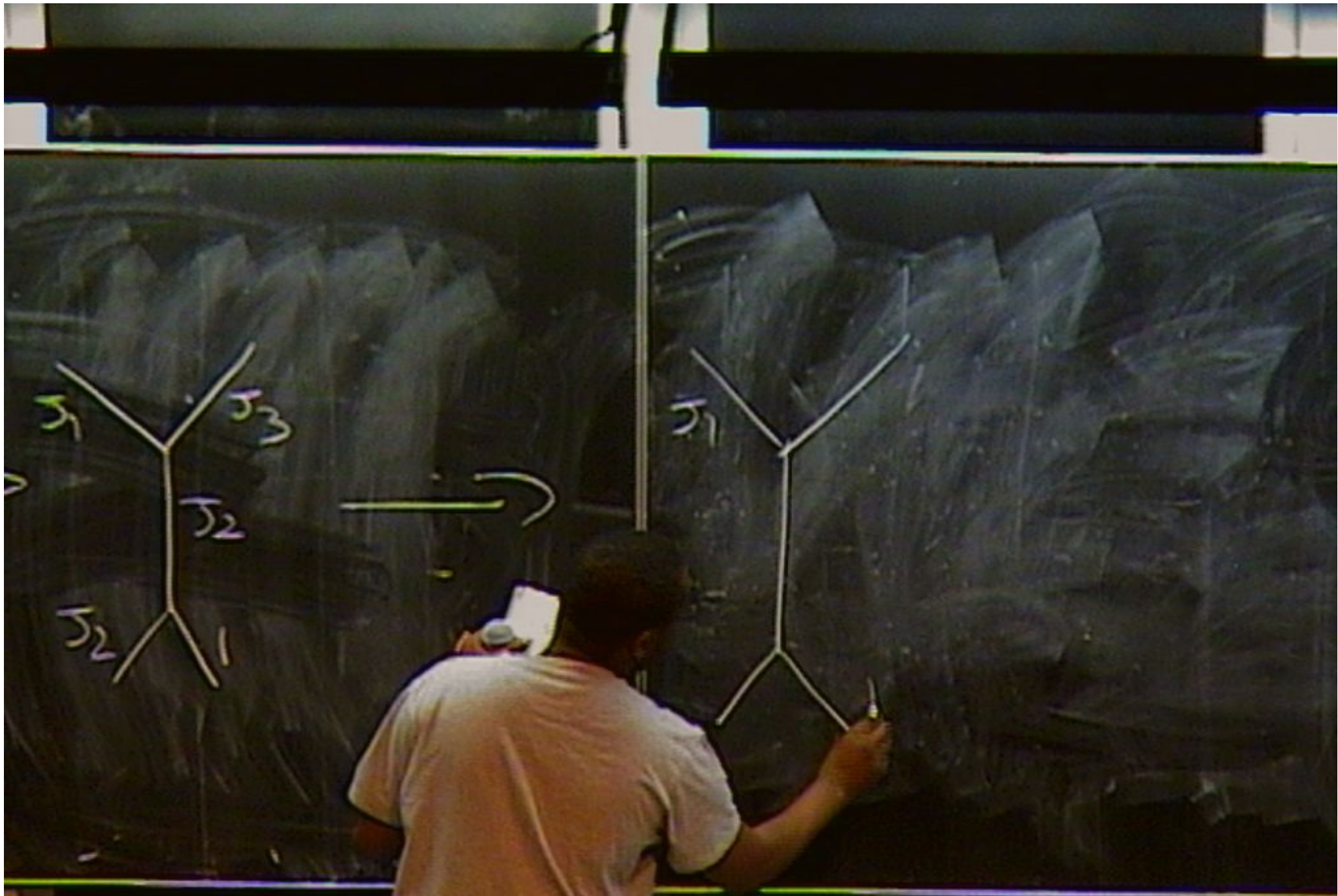


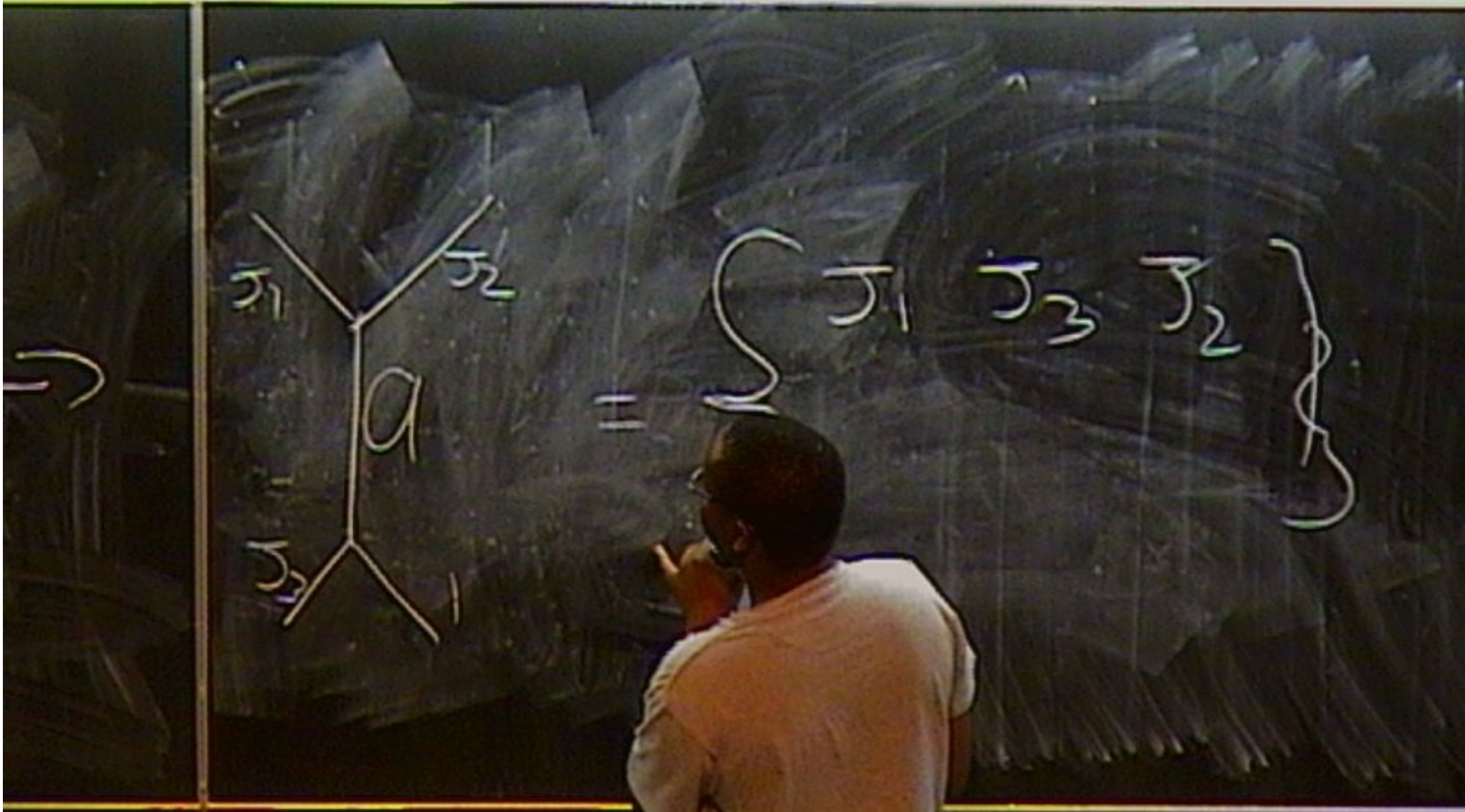


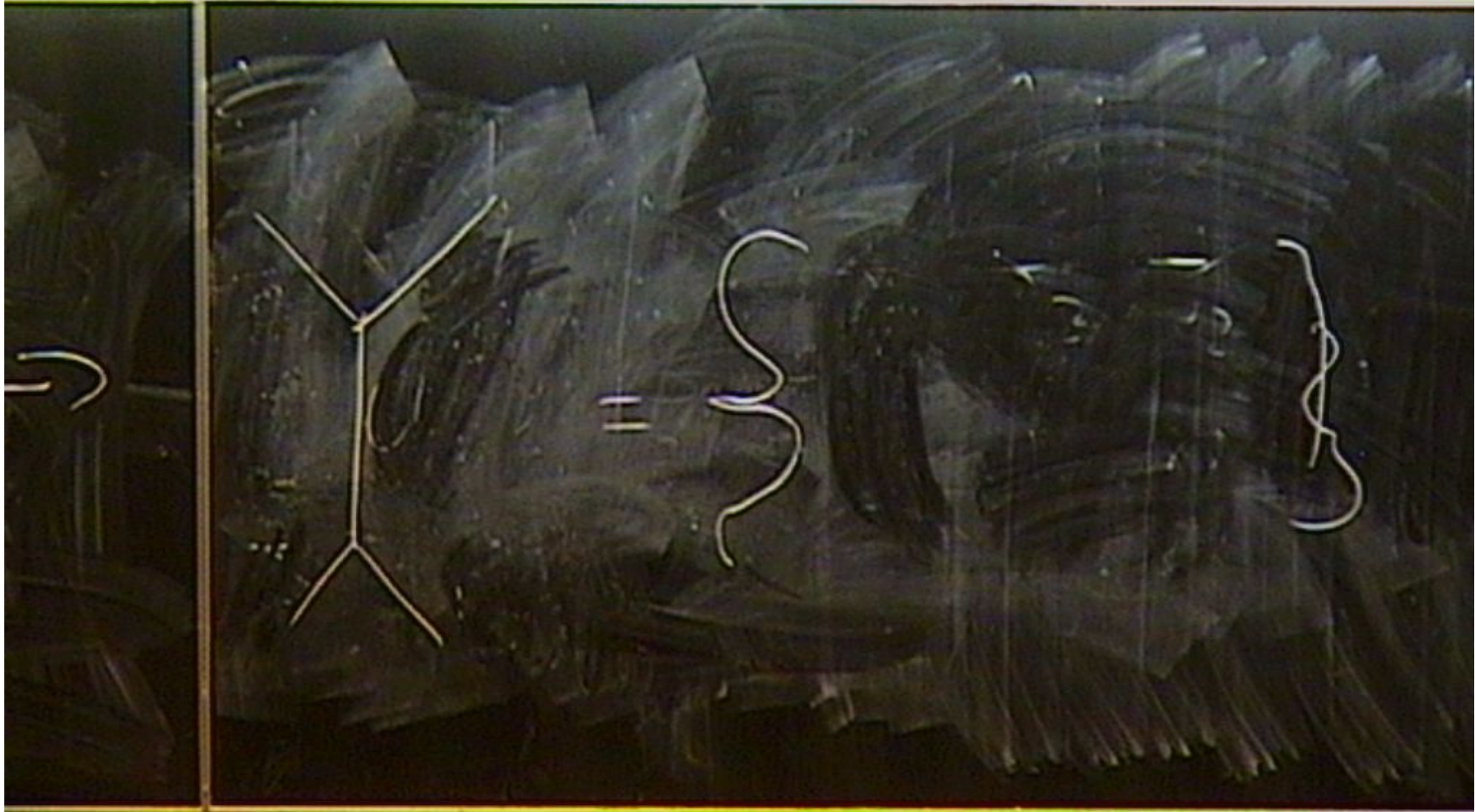


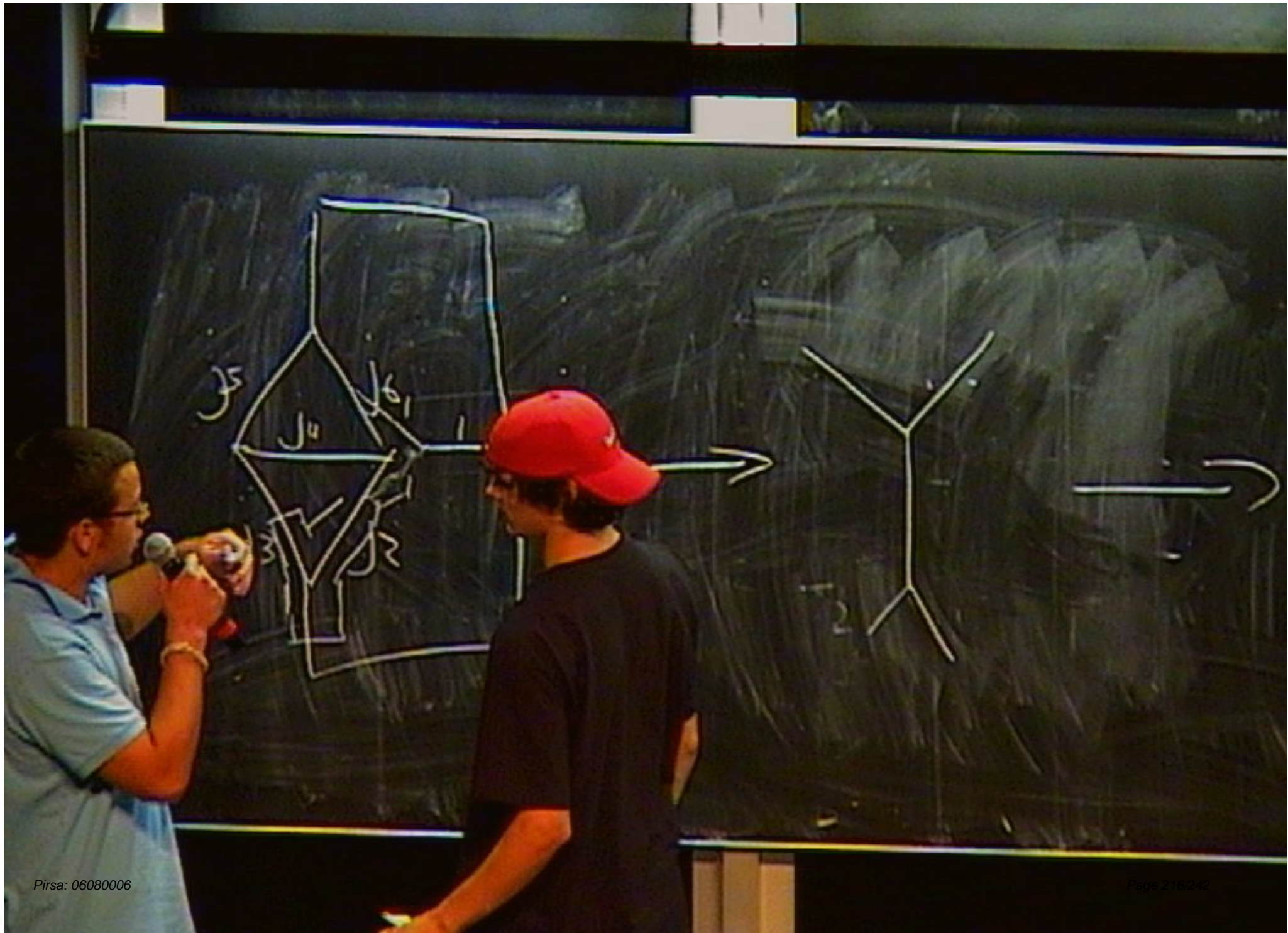
1





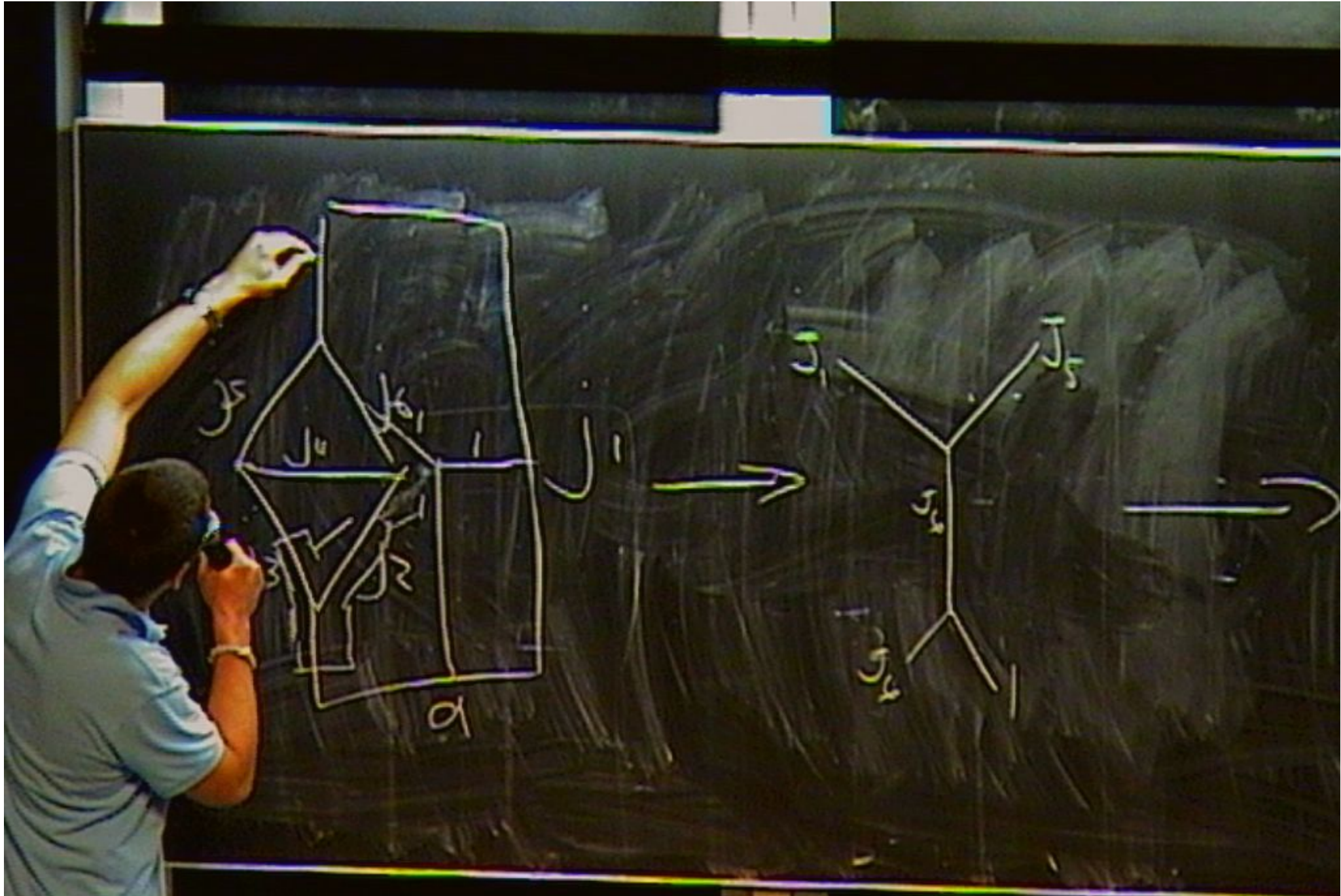


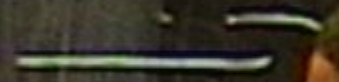
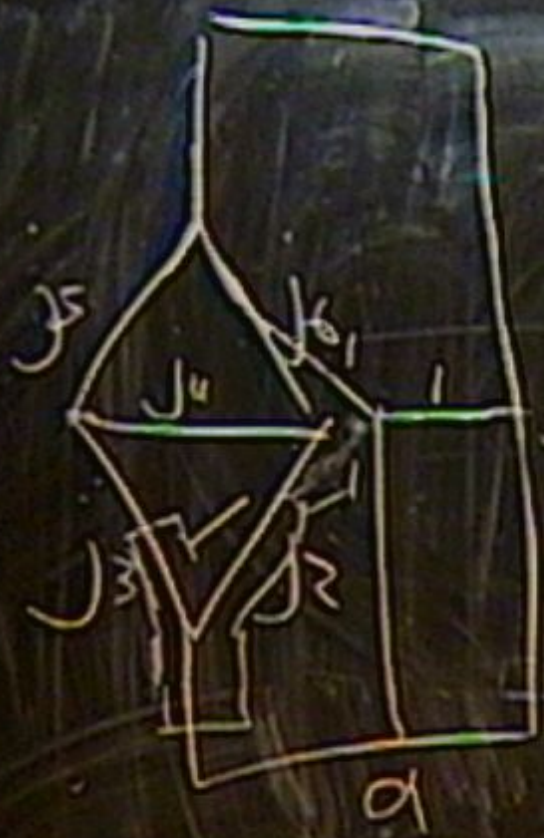


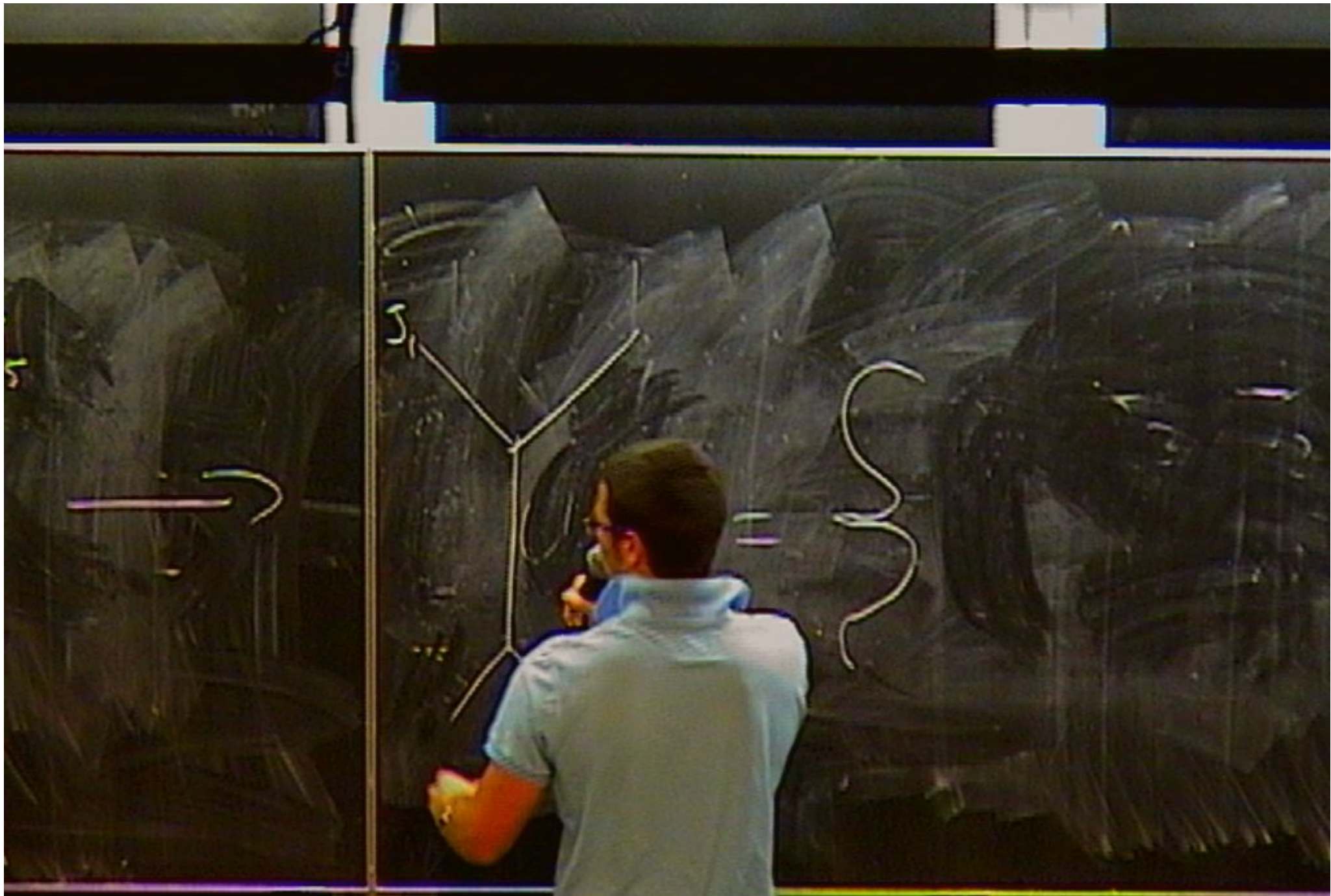




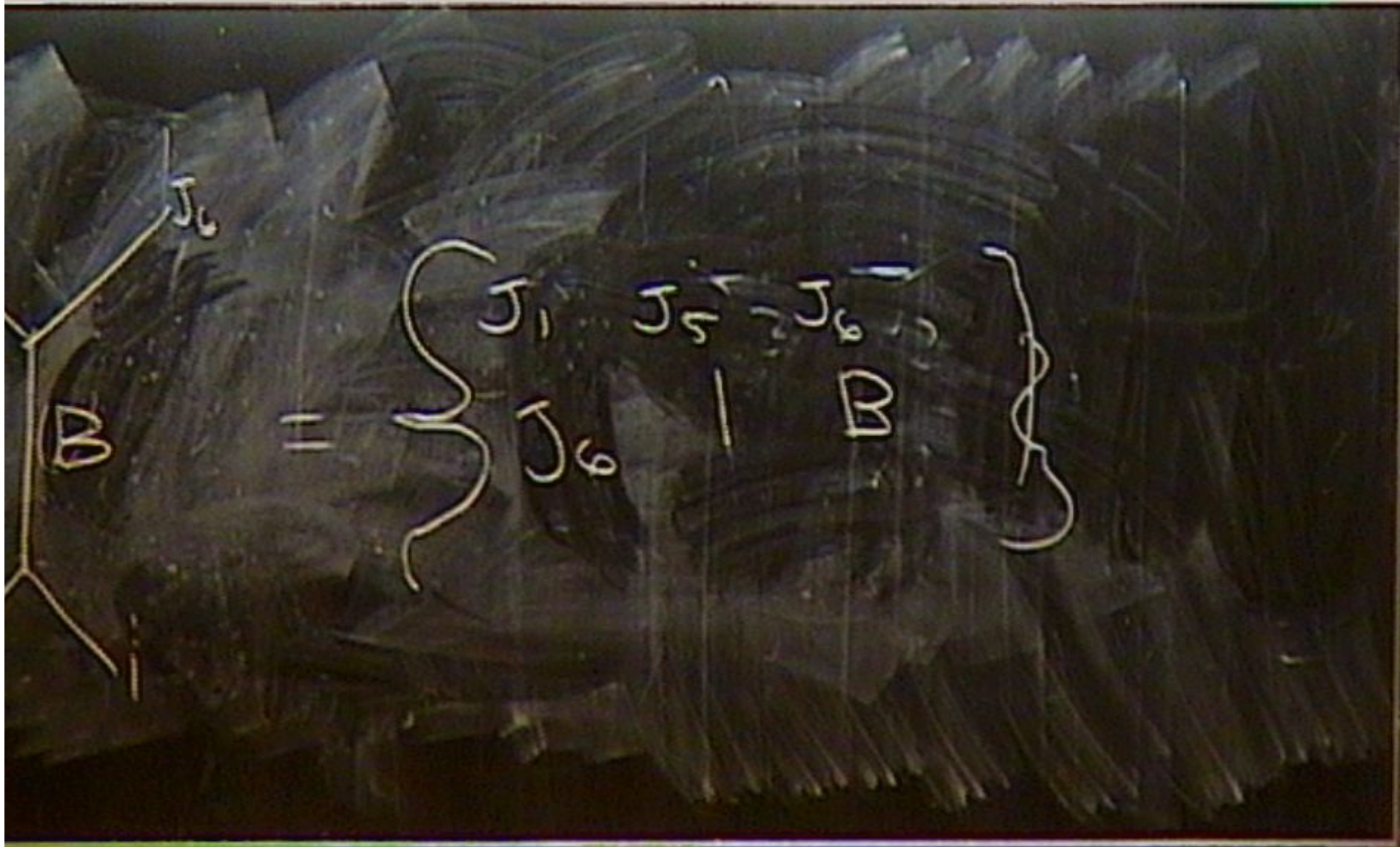


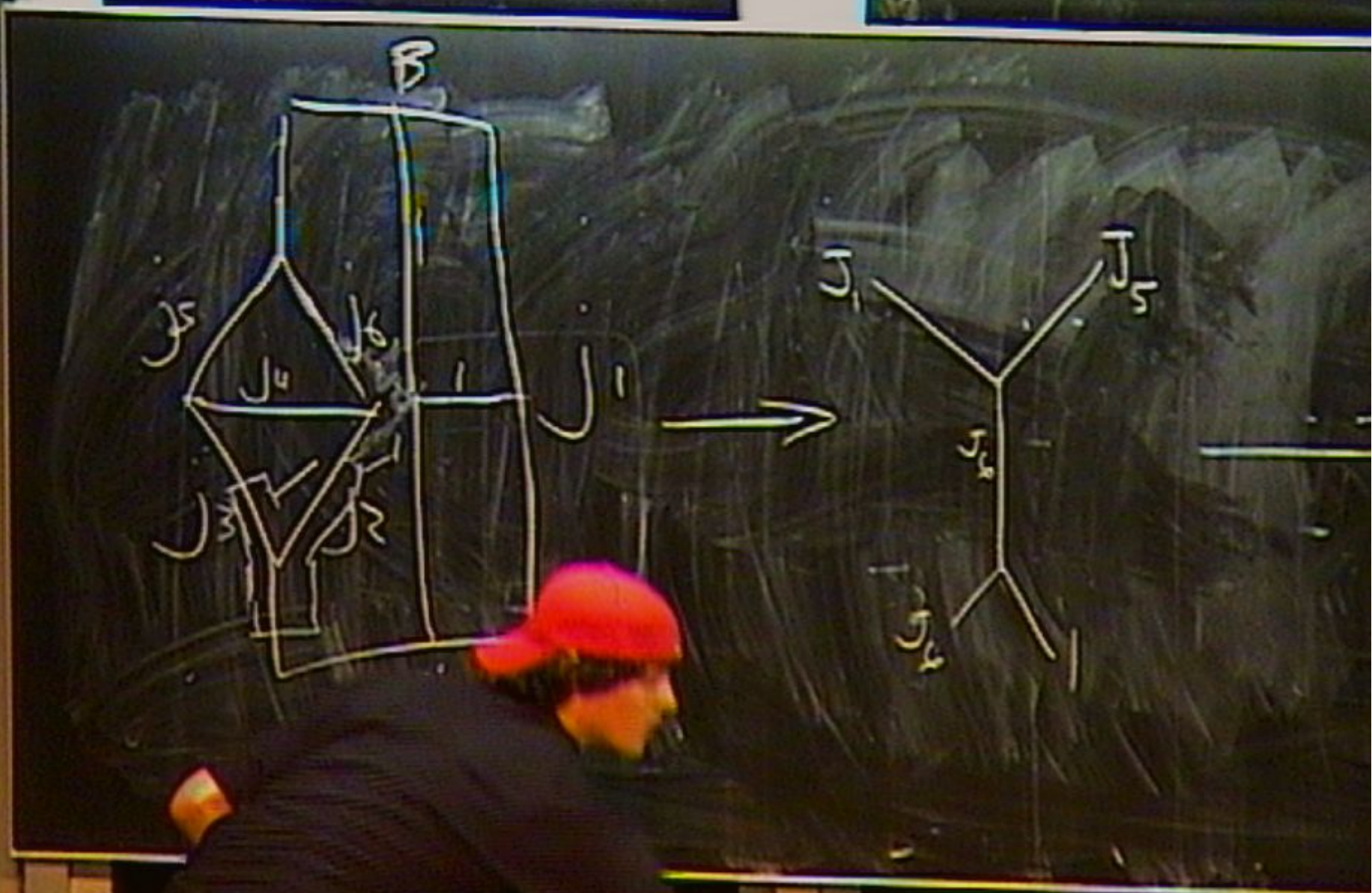


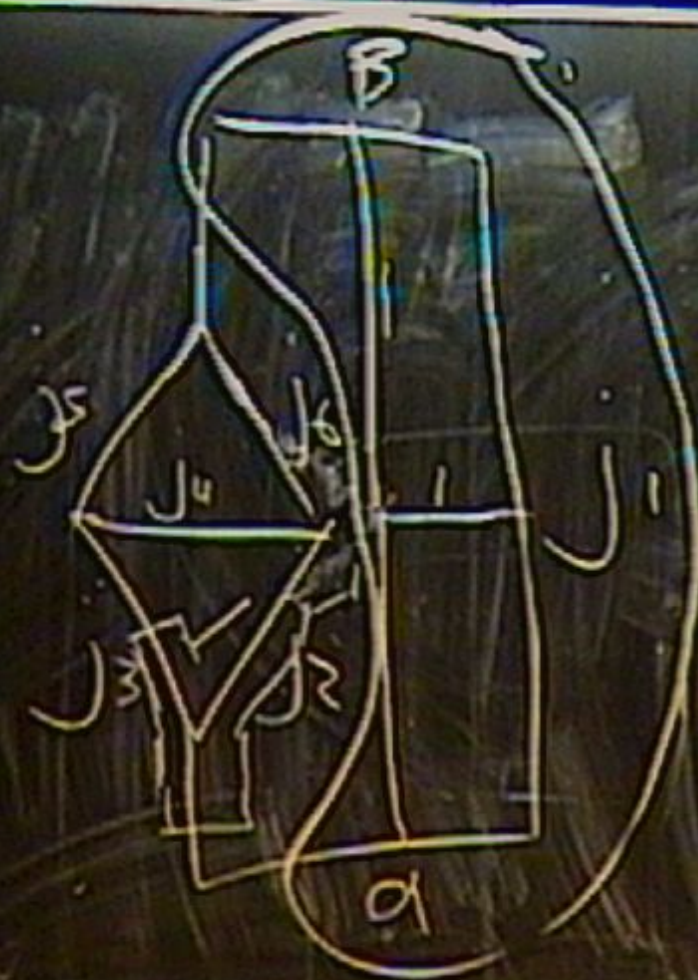




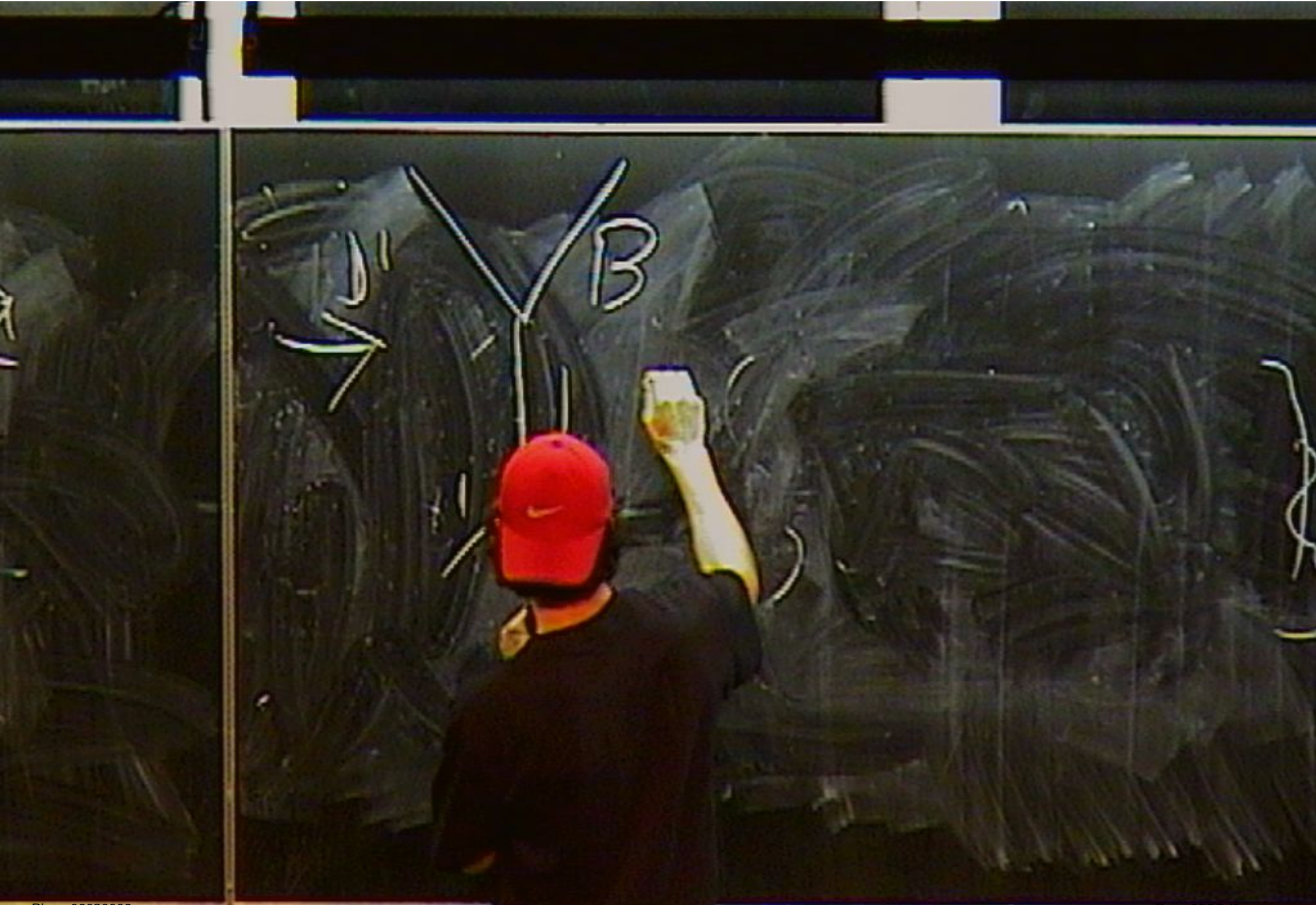




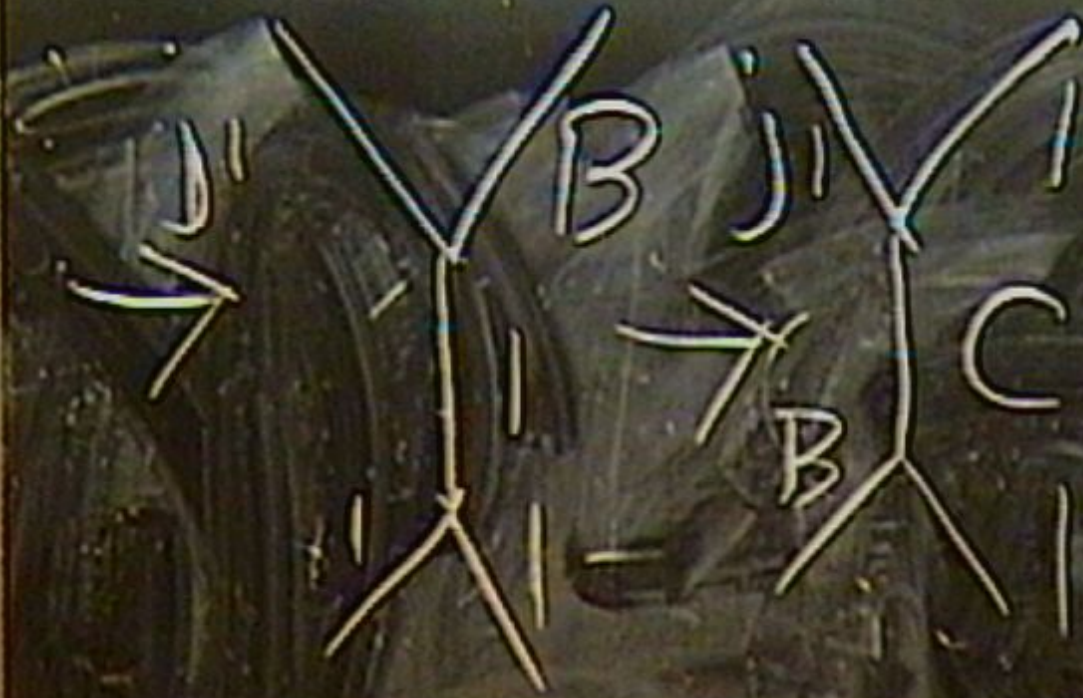


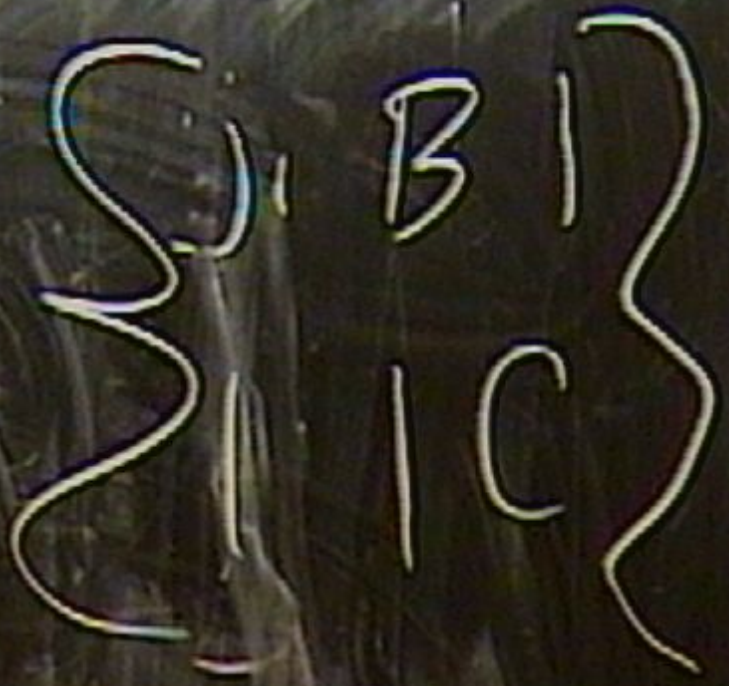


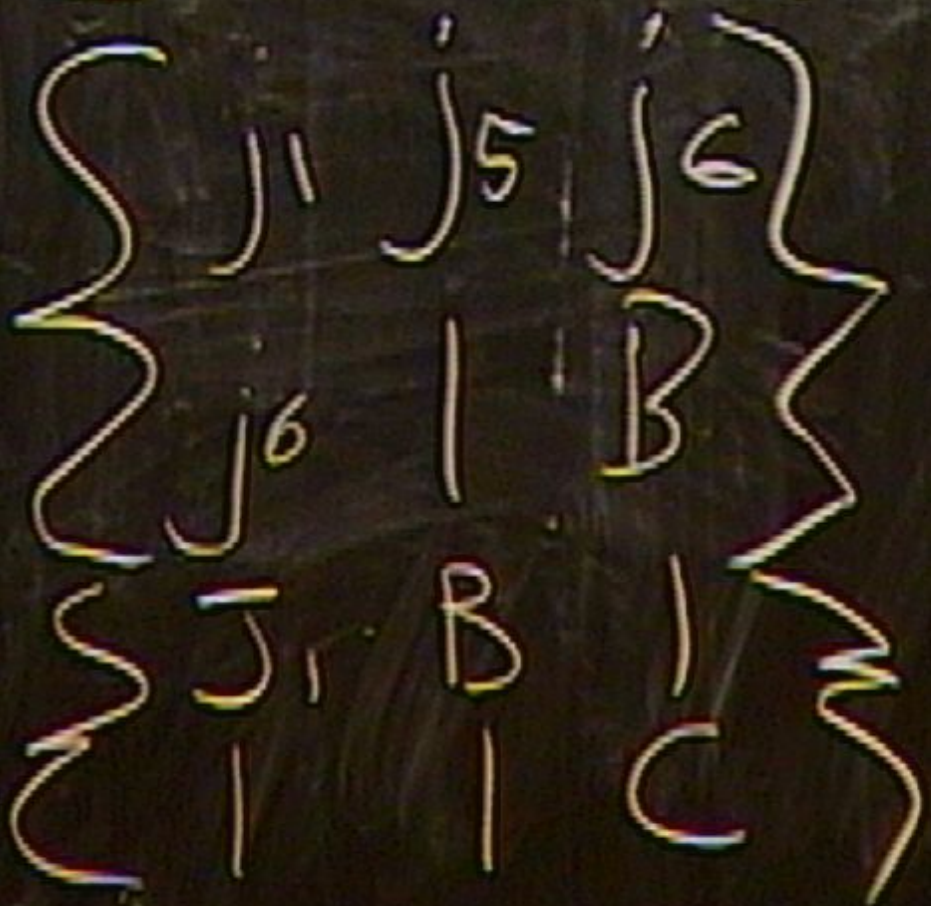
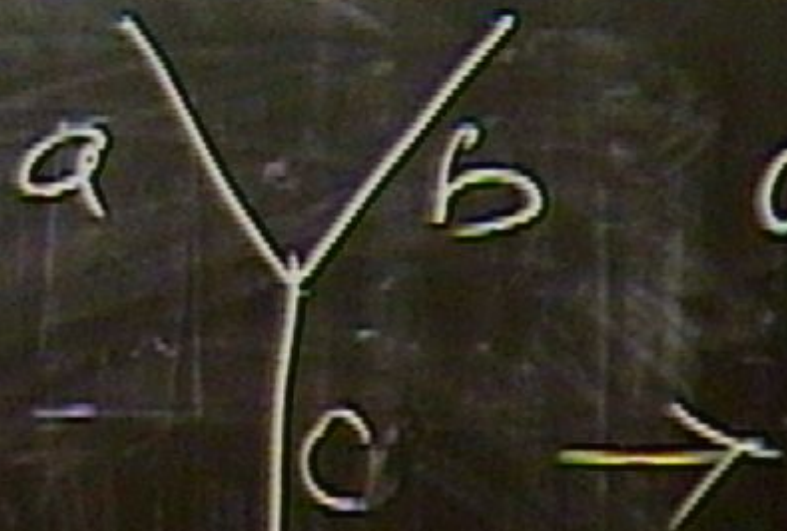
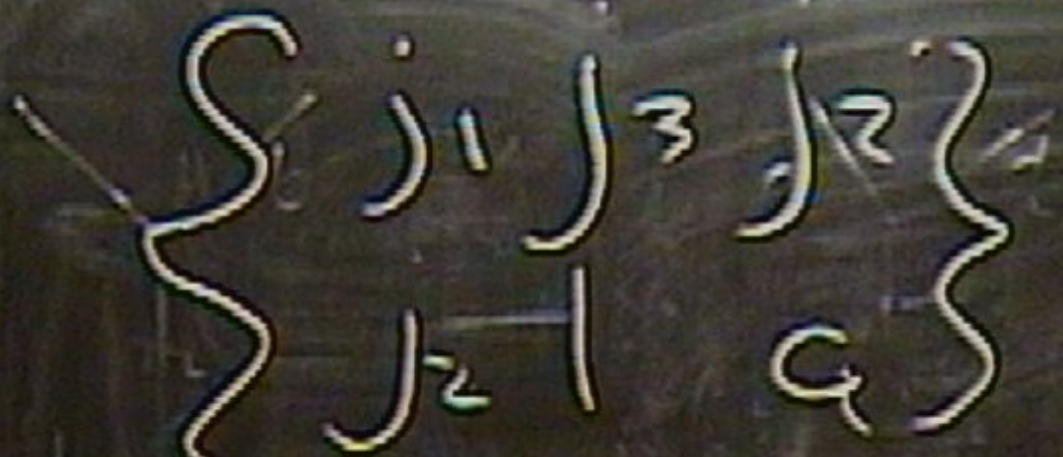


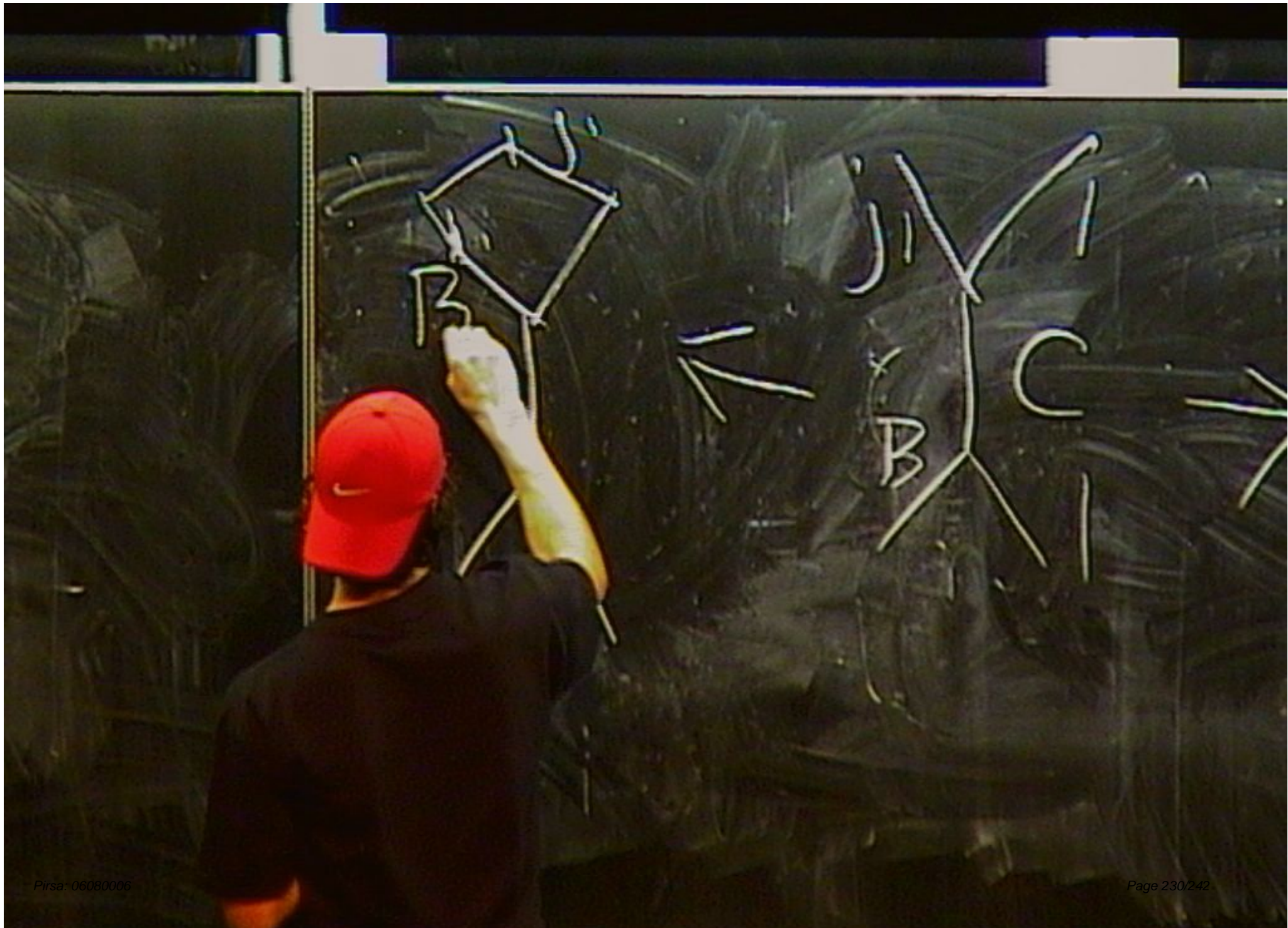


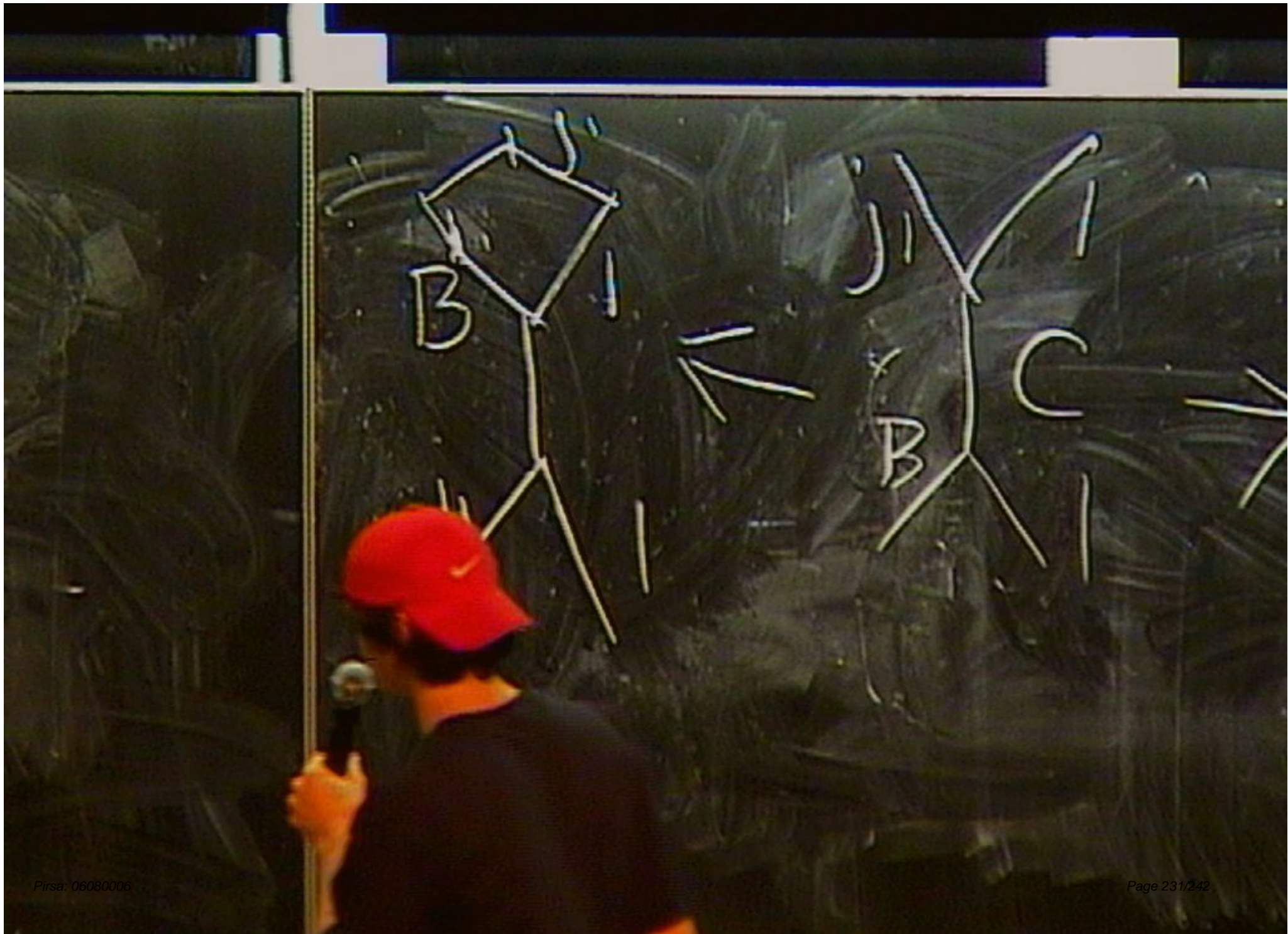


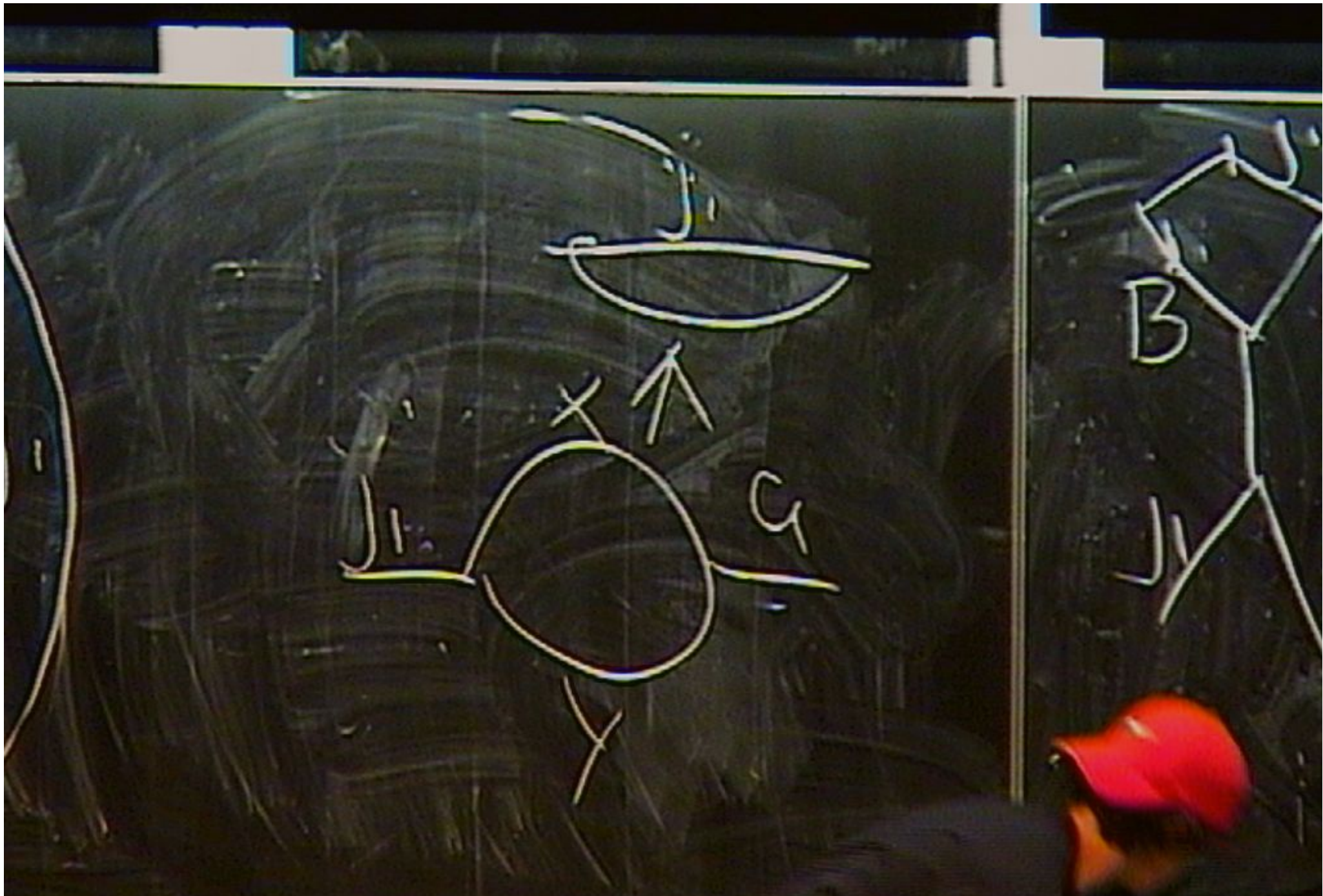




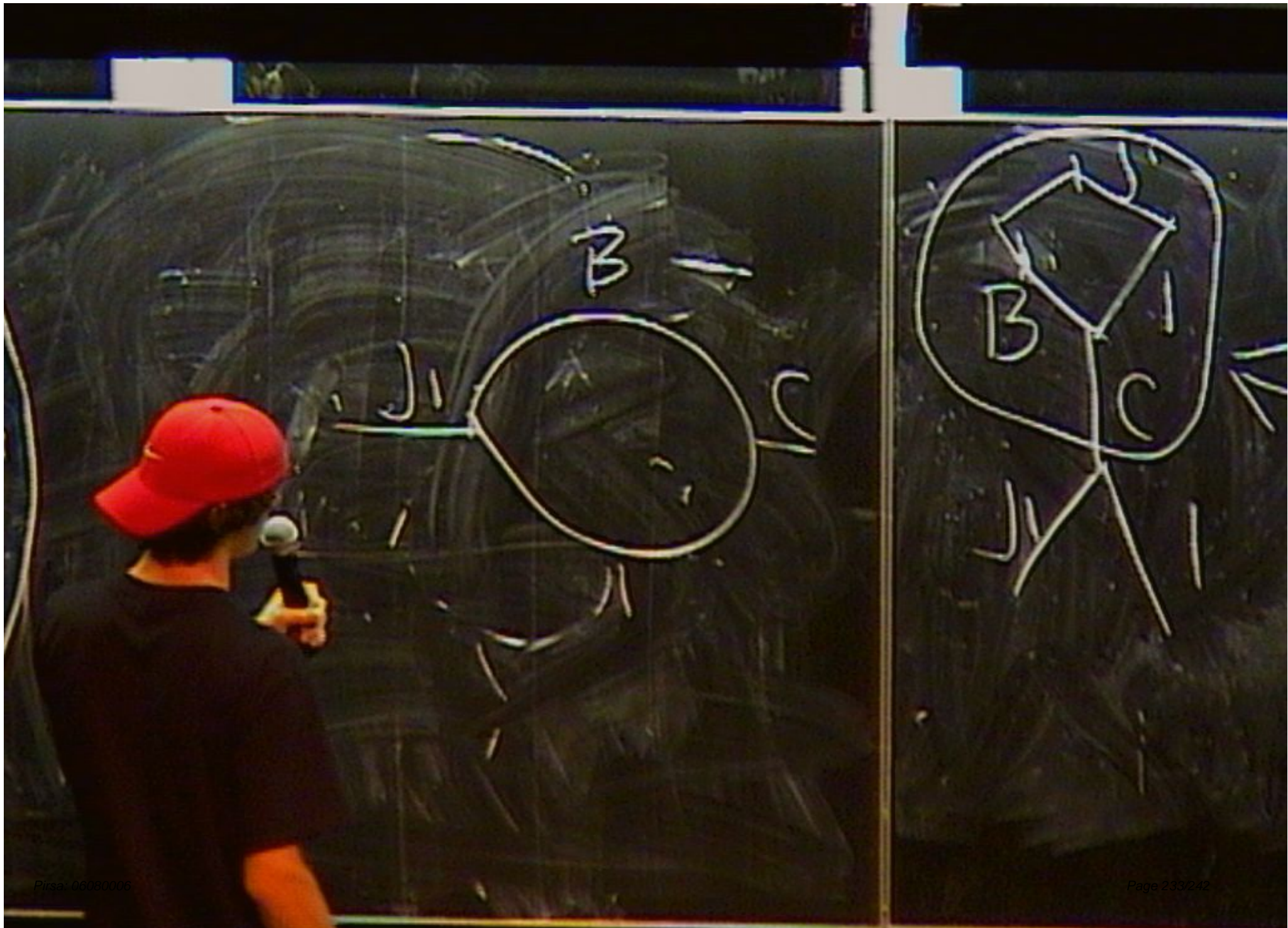


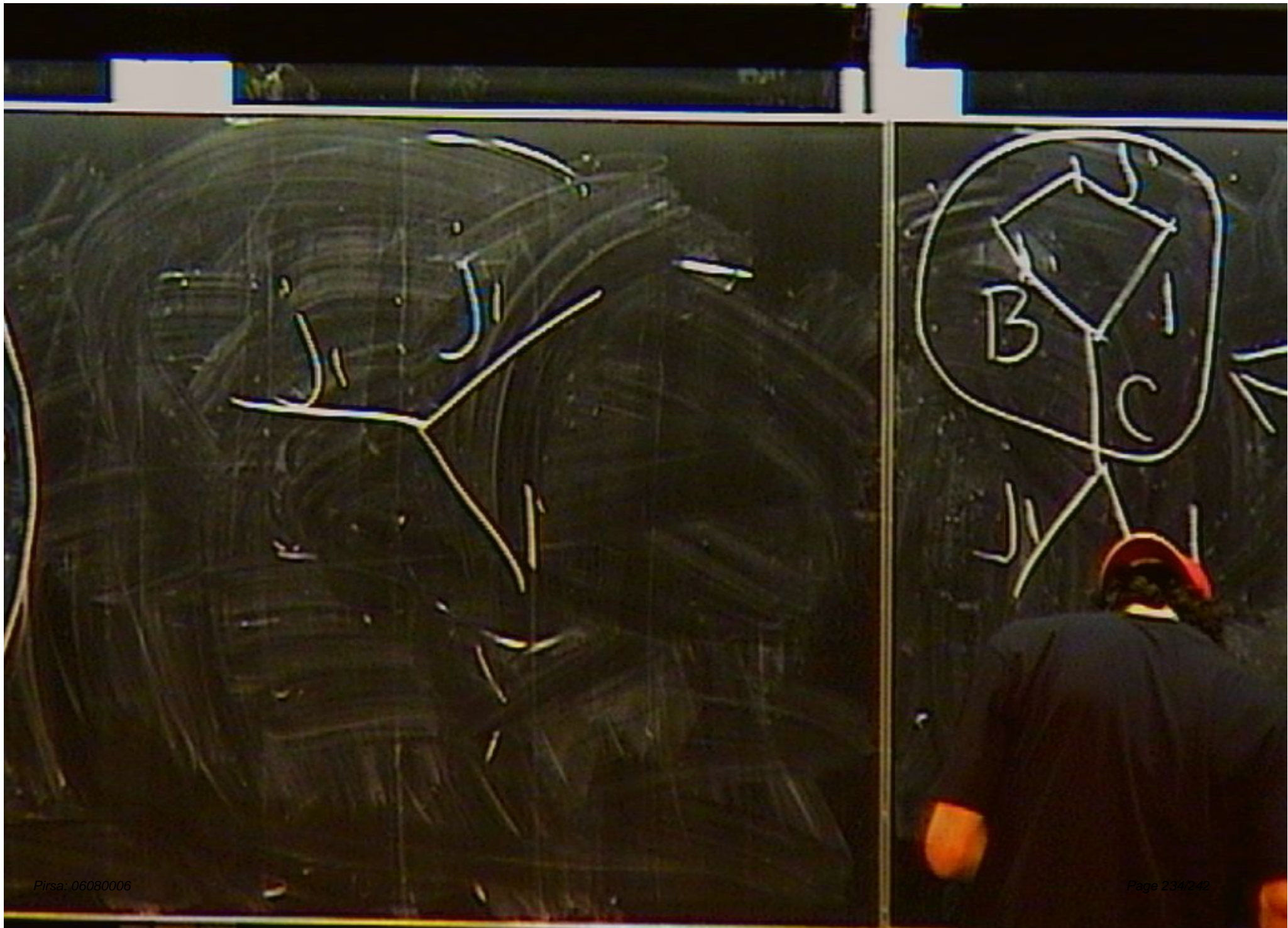


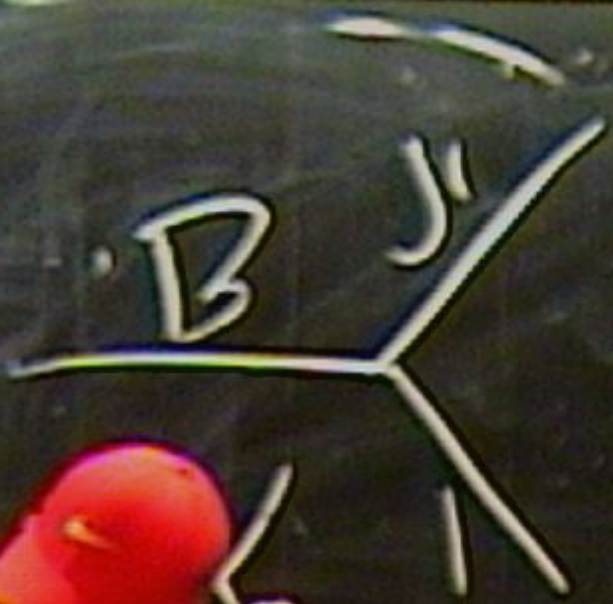
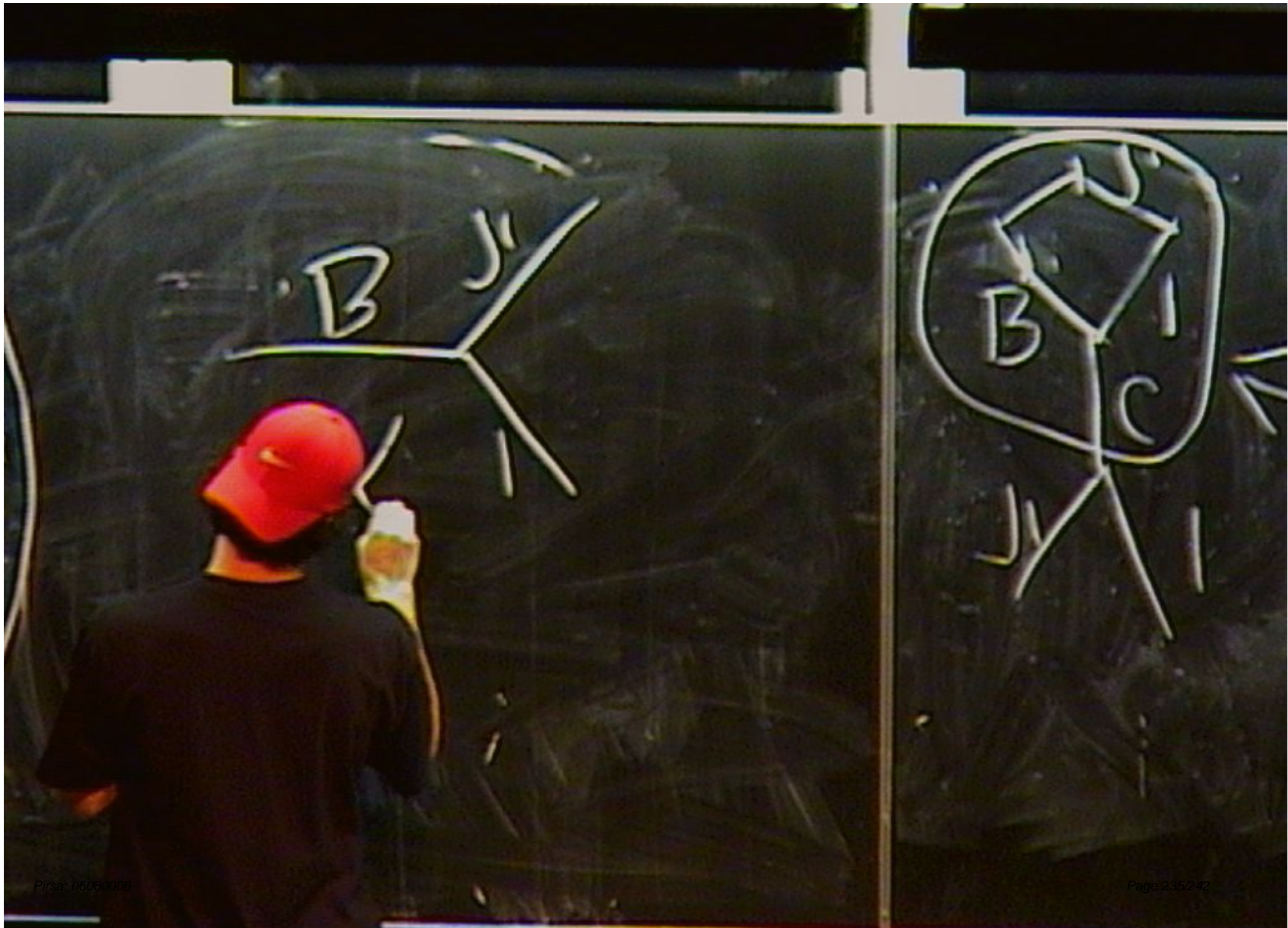


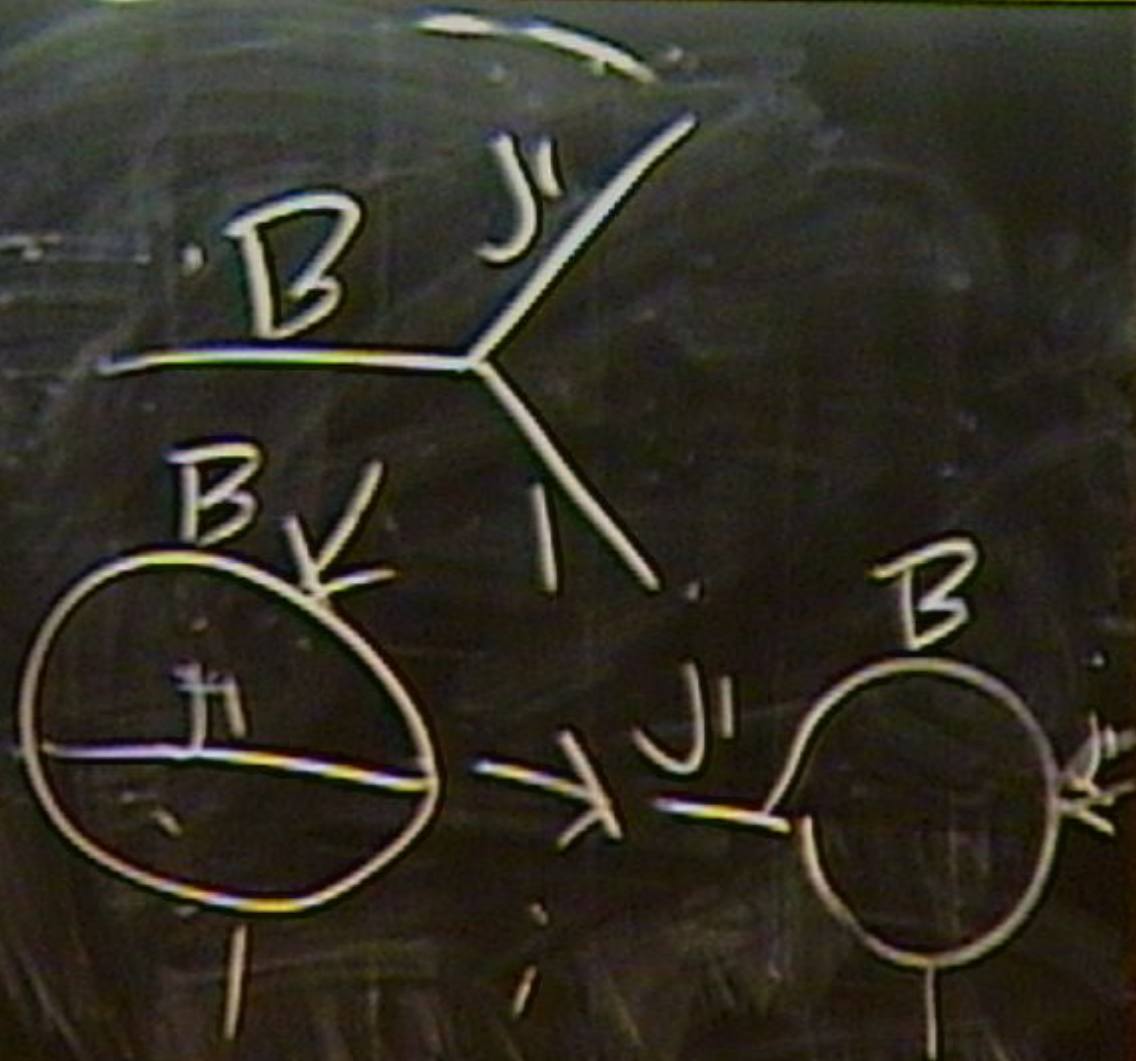


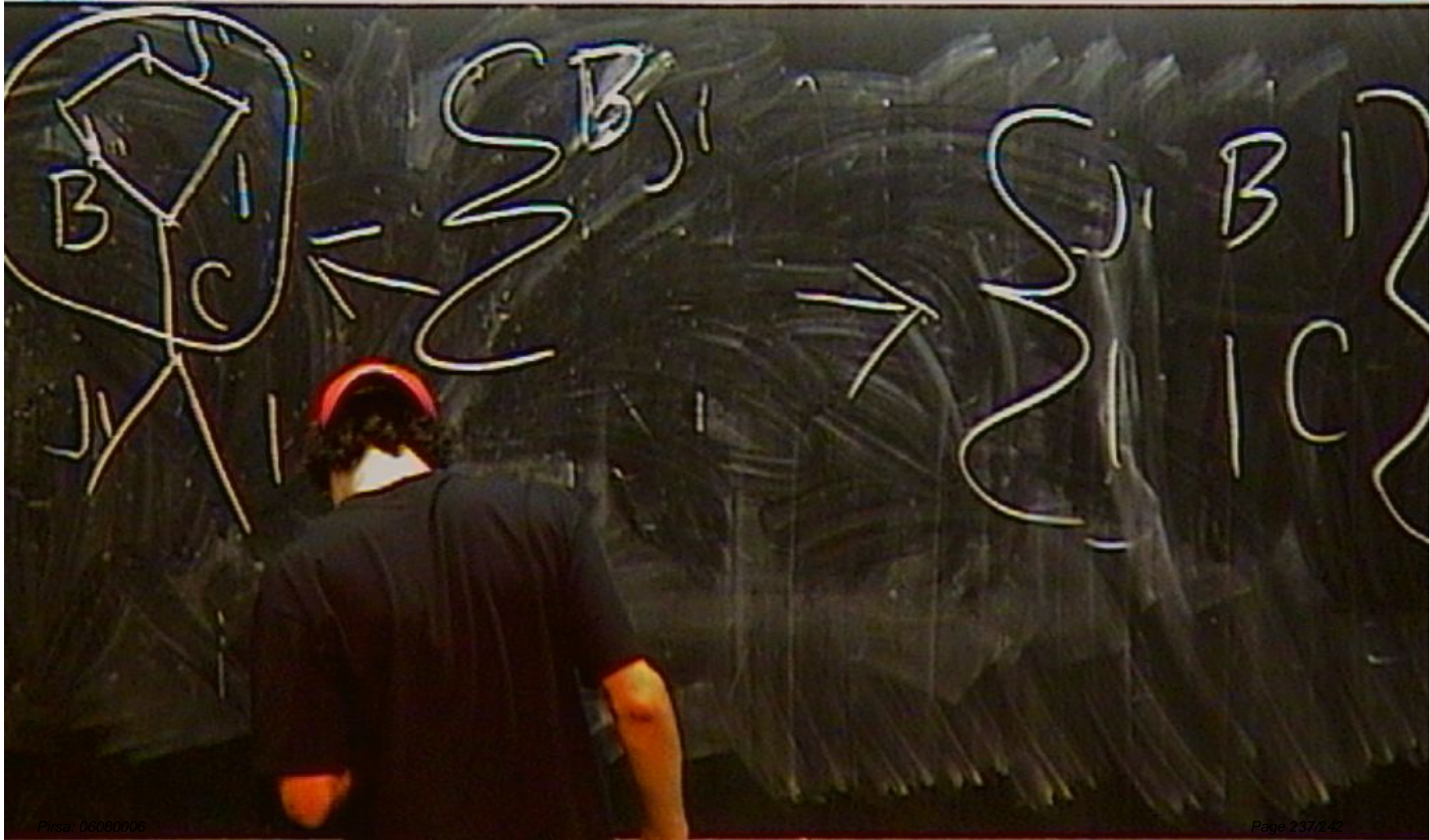


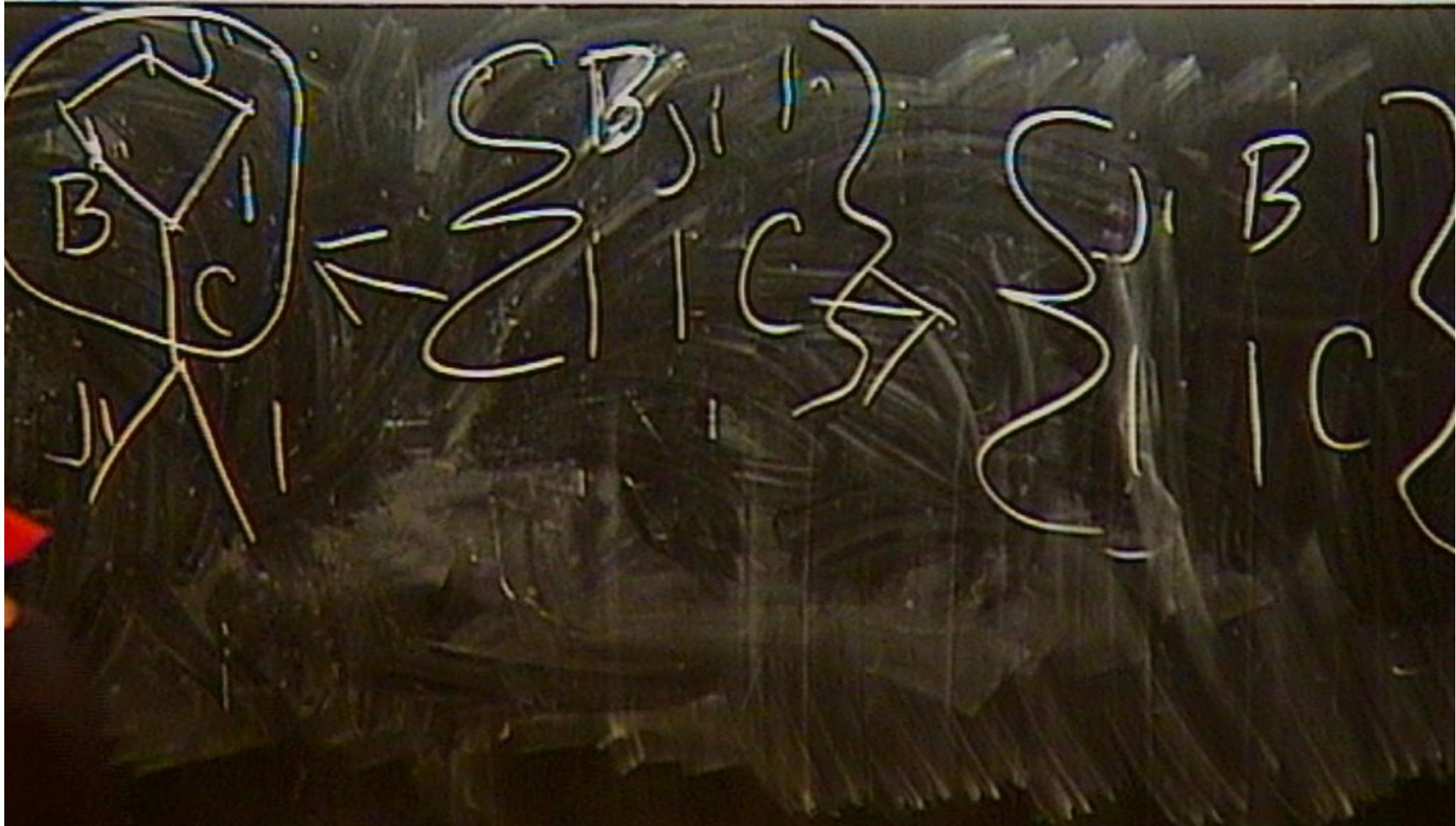


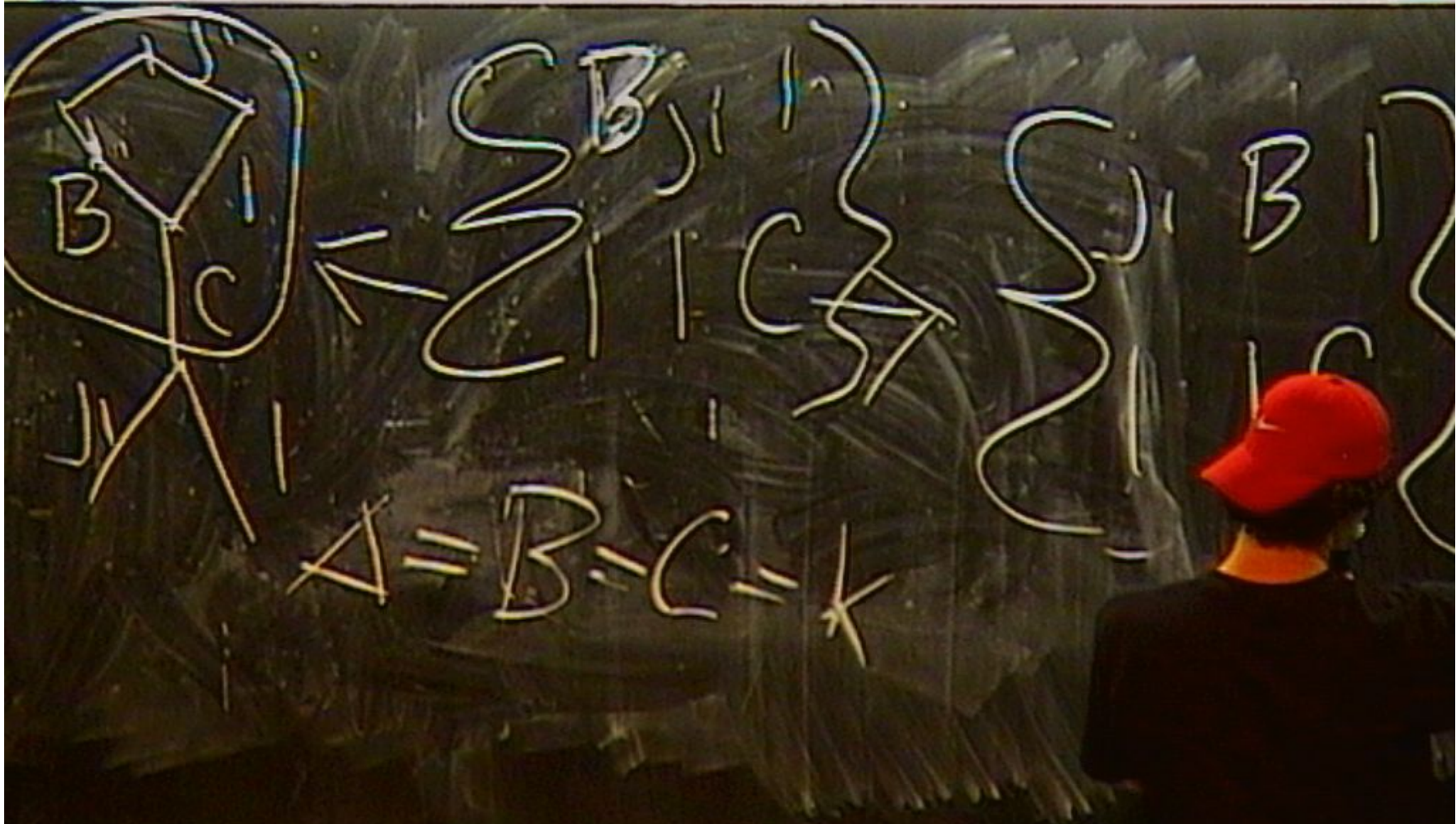














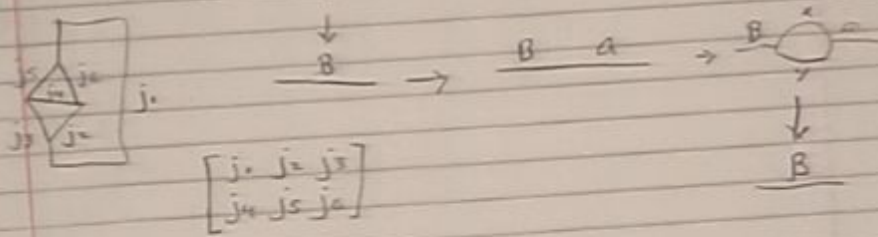
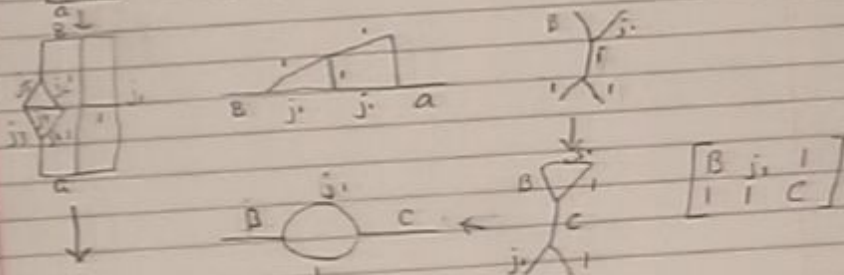
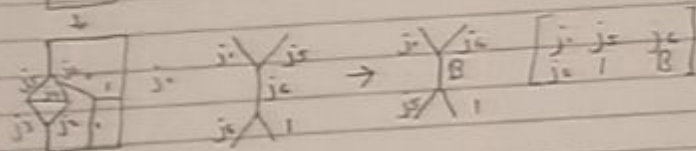
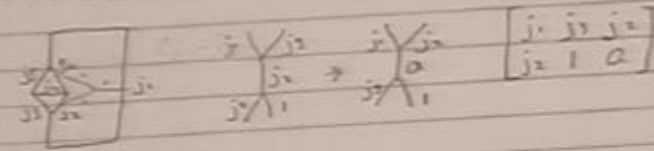
Handwritten mathematical expressions on the chalkboard, enclosed in large curly braces. The top row contains  $k$ ,  $j_2$ , and  $j_3$ . The bottom row contains  $j_1$ ,  $j_5$ , and  $j_6$ .





# Pongoro - Raggio Model

Presentation



constant  $\sqrt{6 \cdot d_{j1} \cdot C^2(j_{j2}) \cdot d_{j2} \cdot C^2(j_{j3}) \cdot d_{j3} \cdot C^2(j_{j1})}$

Heron's formula

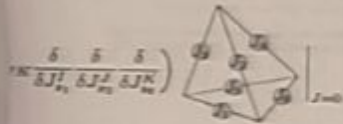
$$A_{\text{area}} = \frac{1}{4} \sqrt{(a+b+c)(a-b+c)(a+b-c)(a-b-c)}$$

Fongaro - Racine Model

21

$$\left\{ \begin{matrix} j_1 - 1 & j_1 & 1 \\ 1 & 1 & j_1 \end{matrix} \right\} = \frac{(-1)^{-2j_1+1}(j_1+1)}{\sqrt{6C^2(j_1)\Delta_{j_1}}} \quad (E11)$$

$$\left\{ \begin{matrix} j_1 + 1 & j_1 & 1 \\ 1 & 1 & j_1 \end{matrix} \right\} = \frac{(-1)^{-2j_1}j_1}{\sqrt{6C^2(j_1)\Delta_{j_1}}} \quad (E12)$$



$$= \sqrt{6 d_{j_1} C^2(j_1) d_{j_2} C^2(j_2) d_{j_3} C^2(j_3)} \sum_k \dim k \left\{ \begin{matrix} 1 & 1 & 1 \\ j_1 & k & j_1 \end{matrix} \right\} \left\{ \begin{matrix} 1 & j_2 & j_2 \\ j_2 & j_1 & k \end{matrix} \right\} \left\{ \begin{matrix} j_1 & j_2 & j_2 \\ j_2 & 1 & k \end{matrix} \right\} \left\{ \begin{matrix} k & j_2 & j_2 \\ j_4 & j_2 & j_2 \end{matrix} \right\} \quad (E13)$$

the simplified evaluations of the 6-j symbols containing values of 1 we get a result of

$$\begin{aligned} \delta_1 \delta_2 \delta_3 \left\{ \begin{matrix} j_1 & j_2 & j_2 \\ j_4 & j_2 & j_2 \end{matrix} \right\} &= \frac{-(j_1+1)}{4j_1(2j_1+1)} \sqrt{(j_1+j_2-j_2)(j_1-j_2+j_2)(1-j_1+j_2+j_2)(1+j_1+j_2+j_2)} \\ &\quad \sqrt{(j_1+j_2-j_2)(j_1-j_2+j_2)(1-j_1+j_2+j_2)(1+j_1+j_2+j_2)} \left\{ \begin{matrix} j_1 - 1 & j_2 & j_2 \\ j_4 & j_2 & j_2 \end{matrix} \right\} \\ &\quad + \frac{(C^2(j_1) + C^2(j_2) - C^2(j_2))(C^2(j_1) + C^2(j_2) - C^2(j_2))}{4j_1(j_1+1)} \left\{ \begin{matrix} j_1 & j_2 & j_2 \\ j_4 & j_2 & j_2 \end{matrix} \right\} \\ &\quad + \frac{j_1}{4(j_1+1)(2j_1+1)} \sqrt{(1+j_1+j_2-j_2)(1+j_1-j_2+j_2)(-j_1+j_2+j_2)(2+j_1+j_2+j_2)} \\ &\quad \sqrt{(1+j_1+j_2-j_2)(1+j_1-j_2+j_2)(-j_1+j_2+j_2)(2+j_1+j_2+j_2)} \left\{ \begin{matrix} j_1 + 1 & j_2 & j_2 \\ j_4 & j_2 & j_2 \end{matrix} \right\} \quad (E14) \end{aligned}$$