Title: Quantum in Gravity?

Date: Jul 20, 2006 03:45 PM

URL: http://pirsa.org/06070054

Abstract: I will report my efforts to describe elementary Quantum behaviours, specifically single-particle interference and two-particle entanglement, in an accelerating frame.

Quantum Mechanics in a Non-Inertial Frame

About Abner

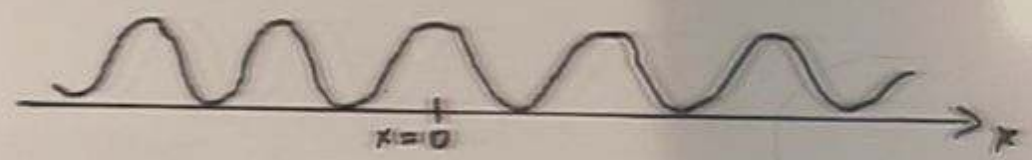
Why this topic?

What quantum mechanics?

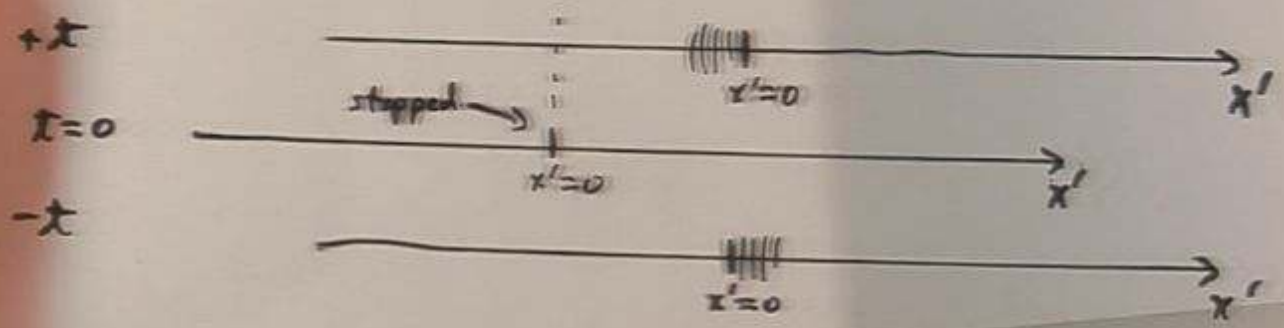
What non-inertial frame?

The Fringes and the Frames

Rest frame of Young's Fringes (Inertial, unprimed)



Non-inertial frame (primed)



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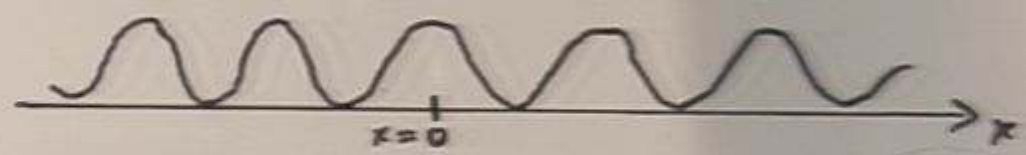

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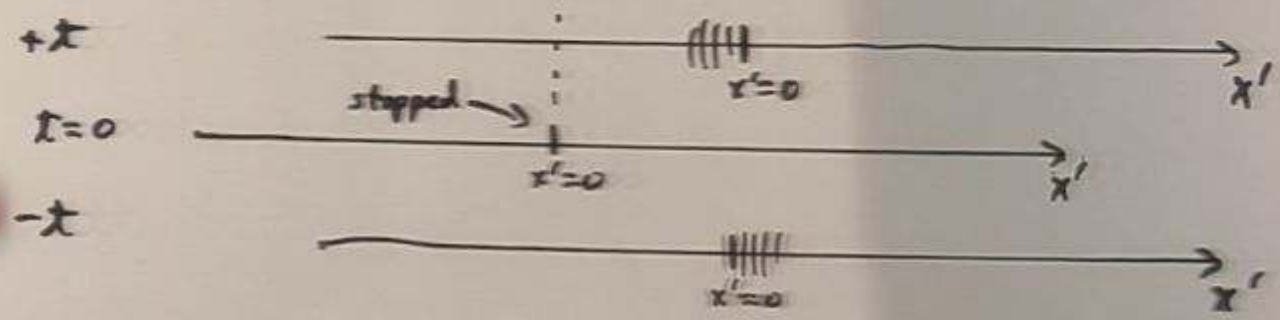

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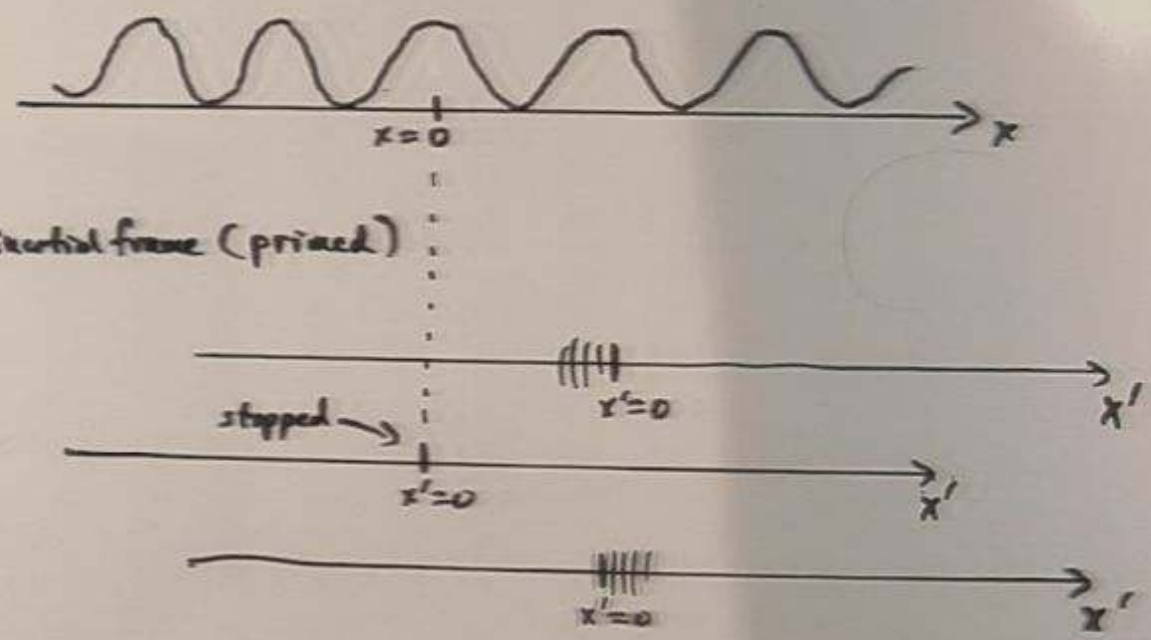

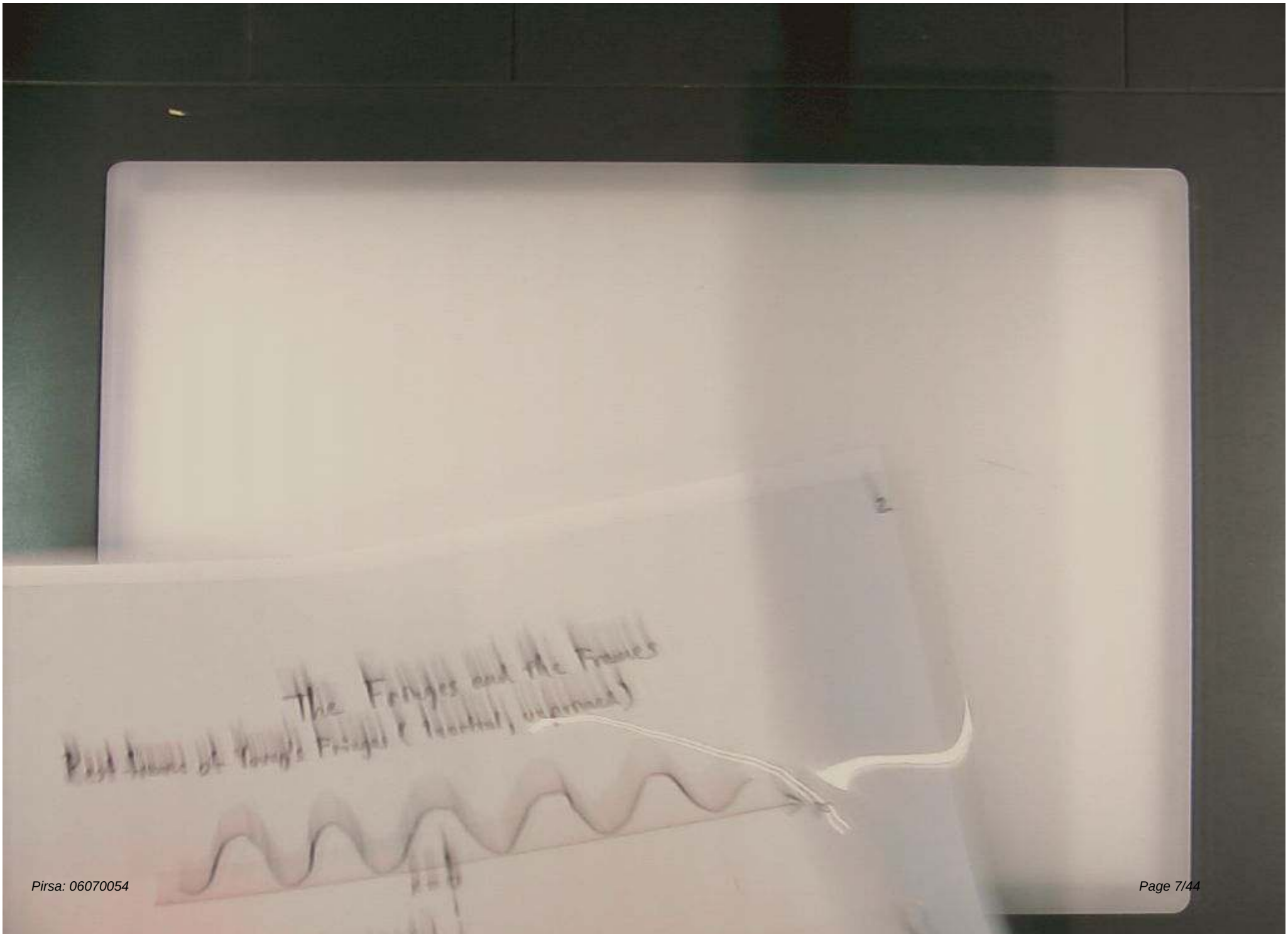
The Fringes and the Frames

Van Name Motion

Starting from rest, particle of mass M_0 experiences a constant force

(1) $x(t) = (1+t^2)^{1/2} - 1$

(2) $\beta(t) = \dfrac{t}{(1+t^2)^{1/2}}$

(3) $a(t) = \dfrac{a_0}{(1+t^2)^{3/2}}$

(4) $\beta(x) = \dfrac{(2x+x^2)^{1/2}}{1+x}$

where
$$\begin{cases} x \equiv \dfrac{F x_{conv}}{M_0 c^2} \\ t \equiv \dfrac{F t_{conv}}{M_0 c} \\ \beta \equiv \dfrac{v_{conv}}{c} \\ a_0 \equiv \dfrac{F}{M_0} \end{cases}$$

derivation

$$(F\,dt = dp)_{conv} \rightarrow dt = d(\gamma\beta) \Rightarrow t = \gamma\beta = \frac{\beta}{(1-\beta^2)^{\frac{1}{2}}} \Rightarrow (2) \Rightarrow \begin{cases} (1) \\ (3) \\ (4) \end{cases}$$

alternatively for (4)

$$(F\,dx = dE)_{conv} \rightarrow dx = d\gamma \Rightarrow x = \gamma - 1 = \frac{1}{(1-\beta^2)^{\frac{1}{2}}} - 1 \Rightarrow (4)$$

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Hyperbolic Motion

Time t as function of proper time τ of the moving point (clock)

$$t = \sinh \tau$$

⌈derivation

$$d\tau = \left(1-\beta^2(t)\right)^{\frac{1}{2}} dt \Rightarrow \tau = \int_0^t \left(1-\beta^2(t)\right)^{\frac{1}{2}} dt = \sinh^{-1} t \Rightarrow t = \sinh$$

⌋

Position x as function of τ

$$x = \cosh \tau - 1$$

⌈derivation

$$x = (1+t^2)^{\frac{1}{2}} - 1 = (1+\sinh^2\tau)^{\frac{1}{2}} - 1 = \cosh\tau - 1$$

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Hanging on x'-axis and Spreading the Time τ

$$x = (1+x')\cosh\tau - 1$$
$$t = (1+x')\sinh\tau$$

$$x' = \left[(1+x)^2 - t^2\right]^{1/2} - 1$$
$$\tau = \tanh^{-1}\left(\frac{t}{1+x}\right)$$

⌜derivation?⌟

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⌈derivation?⌋

Three Disappointments

6

1. Unequal relative speeds of the frames

$$\left.\frac{dx}{dt}\right|_{x'=\text{constant}} = \frac{t}{1+x} \qquad \left.\frac{dx'}{d\tau}\right|_{x=\text{constant}} = -(1+x')\left(\frac{t}{1+x}\right)$$

ugh!

2. Breakdown of speed addition:

$$\frac{\left(\frac{u'}{c}\right)}{(1+x')} = \frac{\left(\frac{u}{v}\right) - \left(\frac{t}{1+x}\right)}{1 - \left(\frac{u}{v}\right)\left(\frac{t}{1+x}\right)} \qquad \text{where} \begin{cases} u' \equiv \frac{dx'}{d\tau} \\ u \equiv \frac{dx}{dt} \end{cases}$$

ugh! → $(1+x')$

3. Alternative calculations of momentum (and energy) inconsistent

momentum $\quad \gamma\left(\frac{u'}{c}\right)\left(\frac{u'}{c}\right) \neq \gamma\left(\frac{u}{c}\right)\gamma\left(\frac{t}{1+x}\right)\left[\frac{u}{c} - \left(\frac{t}{1+x}\right)\right]$

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Three Disappointments

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$$\left.\frac{dx}{dt}\right|_{x'=\text{constant}} = \frac{\dot{x}}{1+x} \qquad \left.\frac{dx'}{d\tau}\right|_{x=\text{constant}} = -(1+x')\left(\frac{\dot{x}}{1+x}\right)$$

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Three Disappointments

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$$\left.\frac{dx}{dt}\right|_{x'=\text{constant}} = \frac{\alpha}{1+x} \qquad \left.\frac{dx'}{d\tau}\right|_{x=\text{constant}} = -(1+x')\left(\frac{\alpha}{1+x}\right)$$

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The Quick Fix

Prime-frame clock at x' keeps t' time, not τ time, where

$$dt' \equiv (1+x')\,d\tau$$

then

1. $$\left.\frac{dx'}{d\tau}\right|_{x=\text{const}} \longrightarrow \left.\frac{dx'}{dt'}\right|_{x=\text{const}} = -\left(\frac{\tau}{1+x}\right) = -\left.\frac{dx}{dt}\right|_{x'=\text{const}}$$

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3. The identities involving γ's and β's are restored

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On to the Quantum Mechanics

Single-Particle Fringes

$$\psi = e^{i(p_R x - E_R t)/\hbar} + e^{i(p_L x - E_L t)/\hbar}$$

$$\Rightarrow \psi^* \psi = 2 + 2\cos\underbrace{\left\{(p_R - p_L)x/\hbar - (E_R - E_L)t/\hbar\right\}}_{\equiv\ \Phi}$$

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Φ in Three Frames

1. Rest frame of the fringes (particle mass m_0, speed $\approx u$), inertial

$$\Phi = 2\gamma\left(\frac{u}{c}\right) m_0 u x/\hbar \Rightarrow \text{fringe width} \equiv W_0 = \left[1-\left(\frac{u}{c}\right)^2\right]^{\frac{1}{2}} \frac{2\pi\hbar}{m_0 u}$$

2. Another inertial frame, moving right at speed v

$$\Phi' = \gamma\left(\frac{v}{c}\right)\gamma\left(\frac{u}{c}\right)\left[2m_0 u x'/\hbar + 2\left(\frac{u}{c}\right)\left(\frac{v}{c}\right) m_0 c^2 t'/\hbar\right]$$

$$\Rightarrow \begin{cases} \text{fringe width} \equiv W = \left[1-\left(\frac{u}{c}\right)^2\right]^{\frac{1}{2}} W_0 \\ \text{fringe speed} = \left.\frac{dx}{dt}\right|_{\Phi=\text{const}} = \frac{E_R - E_L}{p_R - p_L} = -\frac{2\left(\frac{u}{c}\right)\left(\frac{v}{c}\right) m_0 c^2}{2m_0 u} = -v \end{cases}$$

3. The non-inertial frame

Φ = same as 2., with $\frac{v}{c} \to \left(\frac{x}{1+x}\right)$

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Two-Particle Fringes (Conditional)

$$\psi = e^{i\phi_R^1} e^{i\phi_L^2} + e^{i\phi_L^1} e^{i\phi_R^2}$$

$$\Rightarrow \psi^*\psi = 2 + 2\,[\underbrace{(\phi_R^1 - \phi_L^1) - (\phi_R^2 - \phi_L^2)}_{\equiv \Phi}]$$

For mass m_0 and speed $\pm u$ in the inertial frame, Φ in the non-inertial frame is

$$\Phi' = \Big\{ \gamma\left(\frac{t_1}{1+x_1}\right)\gamma\left(\frac{u}{c}\right)\left[2m_0 u\, x_1'/\hbar + 2\left(\frac{t_1}{1+x_1}\right)\left(\frac{u}{c}\right) m_0 c^2 t_1'/\hbar\right]$$
$$- \gamma\left(\frac{t_2}{1+x_2}\right)\gamma\left(\frac{u}{c}\right)\left[2m_0 u\, x_2'/\hbar + 2\left(\frac{t_2}{1+x_2}\right)\left(\frac{u}{c}\right) m_0 c^2 t_2'/\hbar\right]\Big\}$$

Since

$$\rho(x_2, t_2) \sim \int_{\text{one spatial cycle}} \psi^*\psi\, dx_1' = \int_{\text{one temp cycle}} \psi^*\psi\, dt_1' = \underline{\text{constant}},$$

$$\rho(x_1, t_1 \mid x_2, t_2) = \rho(x_1, t_1; x_2, t_2)/\rho(x_2, t_2) \sim \psi^*\psi \text{ itself!}$$

$\Rightarrow$ conditional fringes have correct <u>width</u>, <u>speed</u>, and <u>space-time location</u>.

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