

Title: Quantum Information Theory 5

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Abstract:



Decoherence-free subspaces and spontaneous emission cancellation: necessity of Dicke limit

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Outline



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- Why decoherence-free subspaces (DFS)?
- Dark states and spontaneous emission cancellation
- Decoherence-free subspaces
- A theorem about limitations of DFS
- Conclusions

Why DFS?



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In quantum information processing qubits should be well isolated from their environment

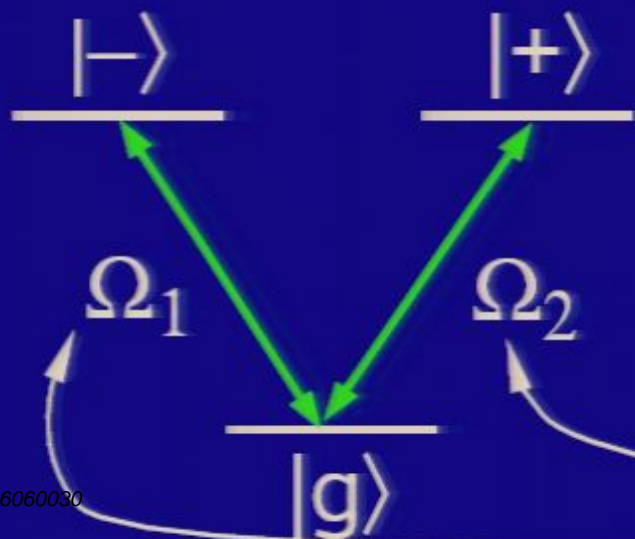
Decoherence-free subspaces consist of states that are not coupled to the environment

DFS were first suggested by Duan & Guo '97, Zanardi & Rasetti '97, Lidar, Chuang & Whaley '98

They could be used as an alternative/addition to quantum error correction

Decoherence-Free Subspaces appear if a particular superposition of states is not coupled to the environment

Consider a V-system with two excited states



The state

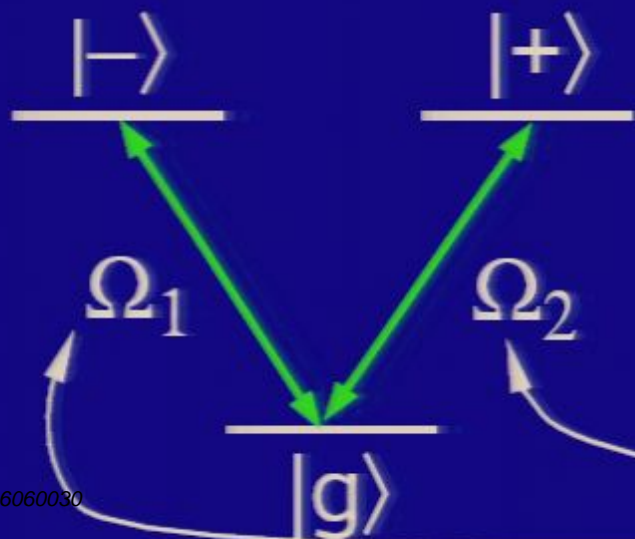
$$|D\rangle = \Omega_2 |-\rangle - \Omega_1 |+\rangle$$

is not coupled to the ground state by *this* particular superposition of laser light

Hamiltonian: $H_{\text{int}} = |e\rangle (\Omega_1 \langle -| + \Omega_2 \langle +|) + \text{H.c.}$

Ω_i = light field amplitude

Obviously $H_{\text{int}} |D\rangle = H_{\text{int}} (\Omega_2 |-\rangle - \Omega_1 |+\rangle) = 0$
(very much like dark states)



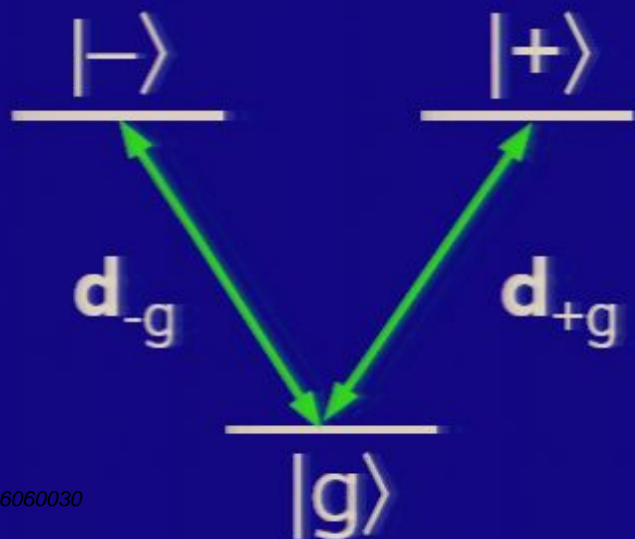
The state

$$|D\rangle = \Omega_2 |-\rangle - \Omega_1 |+\rangle$$

is not coupled to the ground state by *this* superposition of laser light

To be immune against decoherence $|D\rangle$ must not couple to *any* mode configuration

This is possible if the transition matrix elements are the same: $\mathbf{d}_{-g} = \mathbf{d}_{+g} \Leftrightarrow$
spontaneous emission cancellation [Zhu & Scully '96]

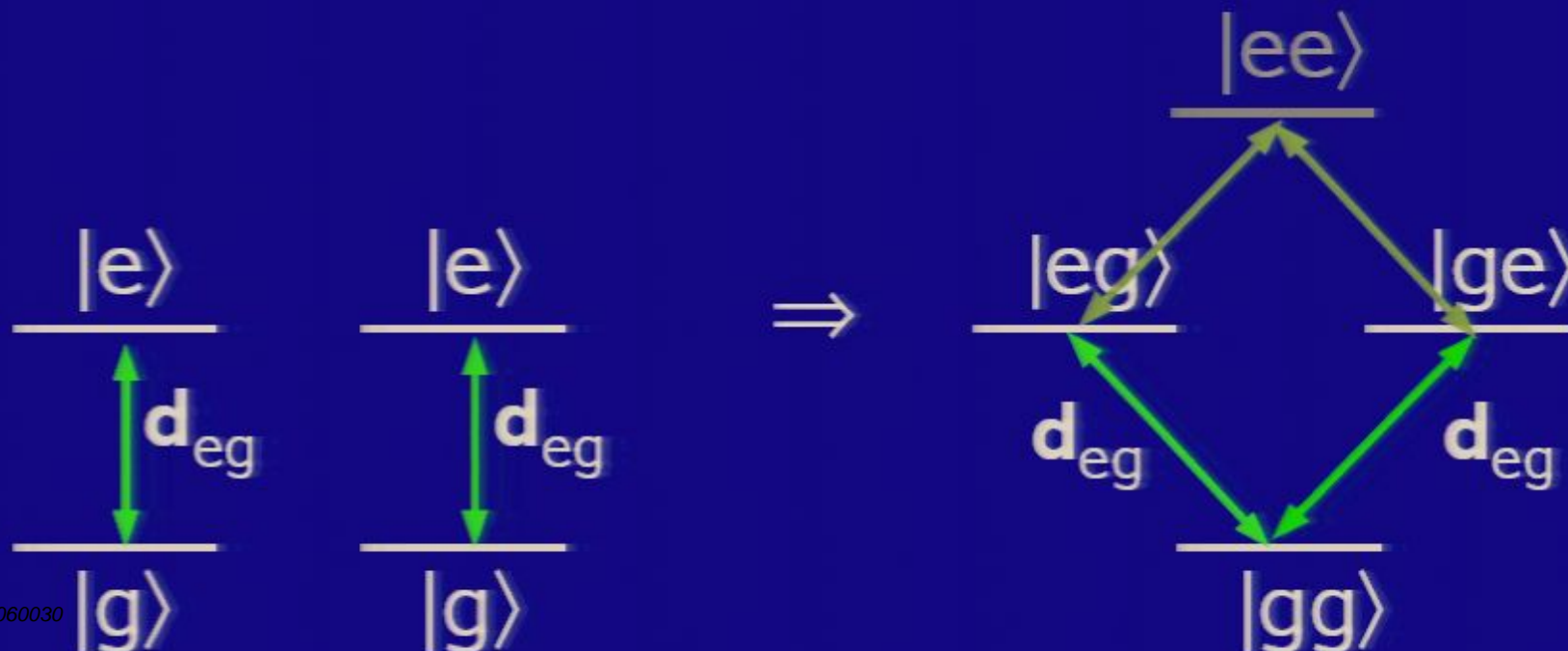


However, selection rules forbid this for atoms

Molecules seem to be necessary, but

... another trick to achieve such a cancellation is to use more than one atom. Consider two 2-level atoms

The energy eigenstates form a V-system with equal dipole matrix elements. $|eg\rangle - |ge\rangle$ forms a DFS

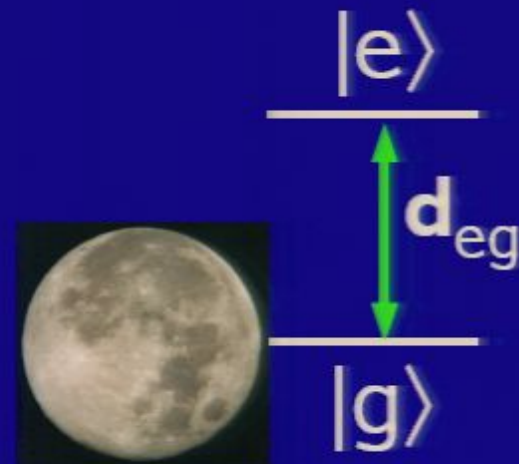
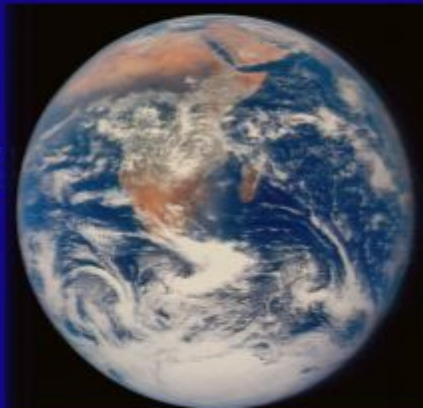


This procedure can be generalized to many N-level atoms [Duan & Guo '97, Zanardi & Rasetti '97, Lidar, Chuang & Whaley '98]

Theories usually employ master equations and Lie groups or Dicke states

However, there's a catch:

How far can the atoms be apart? Surely there's a limit



We have proven that, under very general conditions, a real DFS can only be obtained in the Dicke limit

Dicke limit: particles are co-located

Main assumptions of the theorem:

- System is composed of particles with a finite-dimensional internal Hilbert space located at a fixed position
- Markovian reservoir (no memory)
- Reservoir invariant under translations
- Energy is exchanged in system-reservoir interaction





Some details: $H_{\text{int}} = \sum_n \sum_i E^i(x_n) d_{i,n}$
with reservoir operators E^i and system operators $d_{i,n}$

Translational invariance is used to employ the Lehmann representation for reservoir operators

$$\hat{E}^i(\mathbf{x}_0 + \mathbf{x}) = e^{i\hat{\mathbf{P}} \cdot \mathbf{x} / \hbar} \hat{E}^i(\mathbf{x}_0) e^{-i\hat{\mathbf{P}} \cdot \mathbf{x} / \hbar}$$

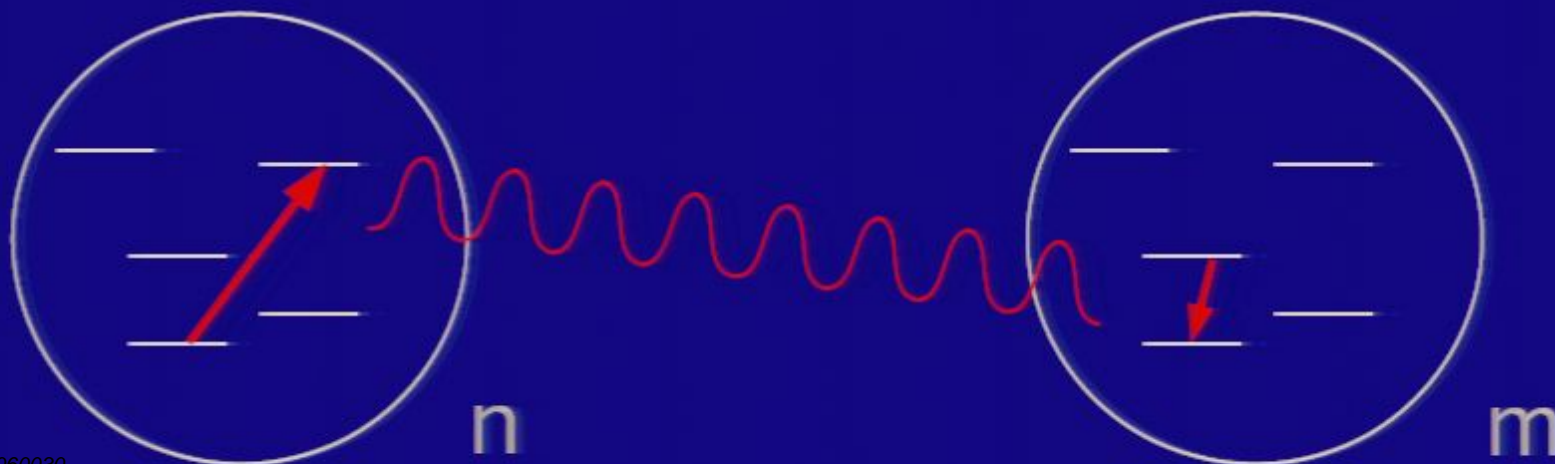
The Markovian master equation can be written as

$$\dot{\rho} = -\hat{\Gamma}\rho - \rho\hat{\Gamma} - i[H, \rho] + \text{jump terms}$$

where the decoherence matrix $\hat{\Gamma}$ describes the (de-) excitation of any particle in the system

Energy conservation (or time averaging) is needed to keep the master equation consistent [Dumcke & Spohn 1979]

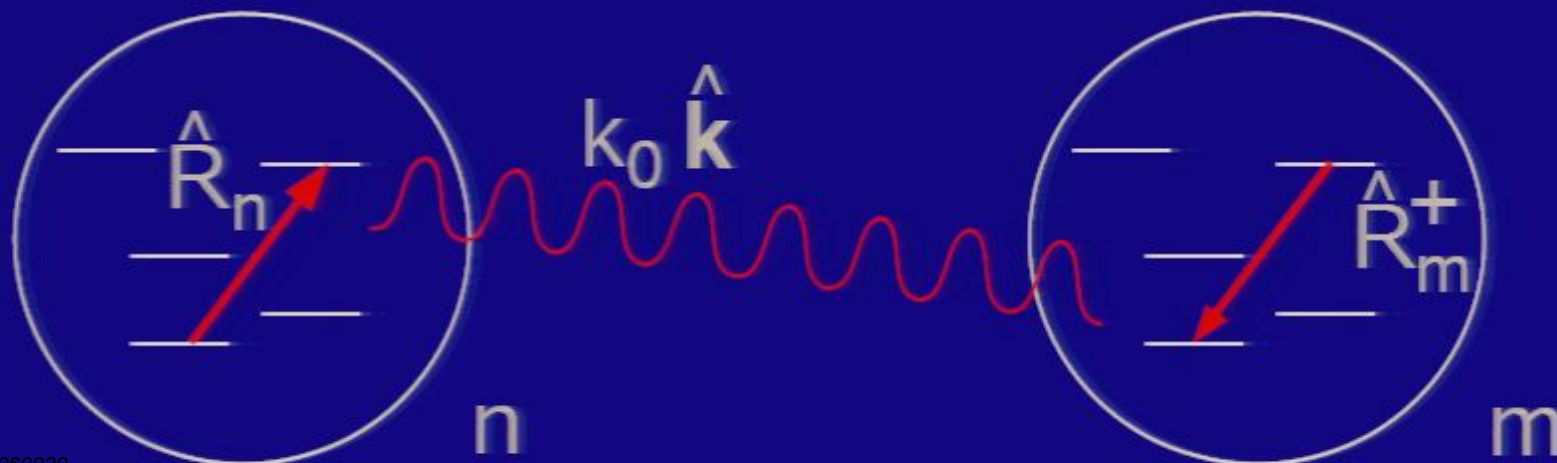
Without it the system would gain or lose energy in processes like photon-exchange between two atoms.



After lengthy calculations $\hat{\Gamma}^{\Lambda}$ can be written as

$$\hat{\Gamma} = \sum_{\Delta E} \int d^2 \hat{k} \sum_{n,m} e^{ik_0 \hat{k} \cdot (\mathbf{x}_n - \mathbf{x}_m)} \hat{R}_n(\Delta E) \hat{R}_m^{\dagger}(\Delta E)$$

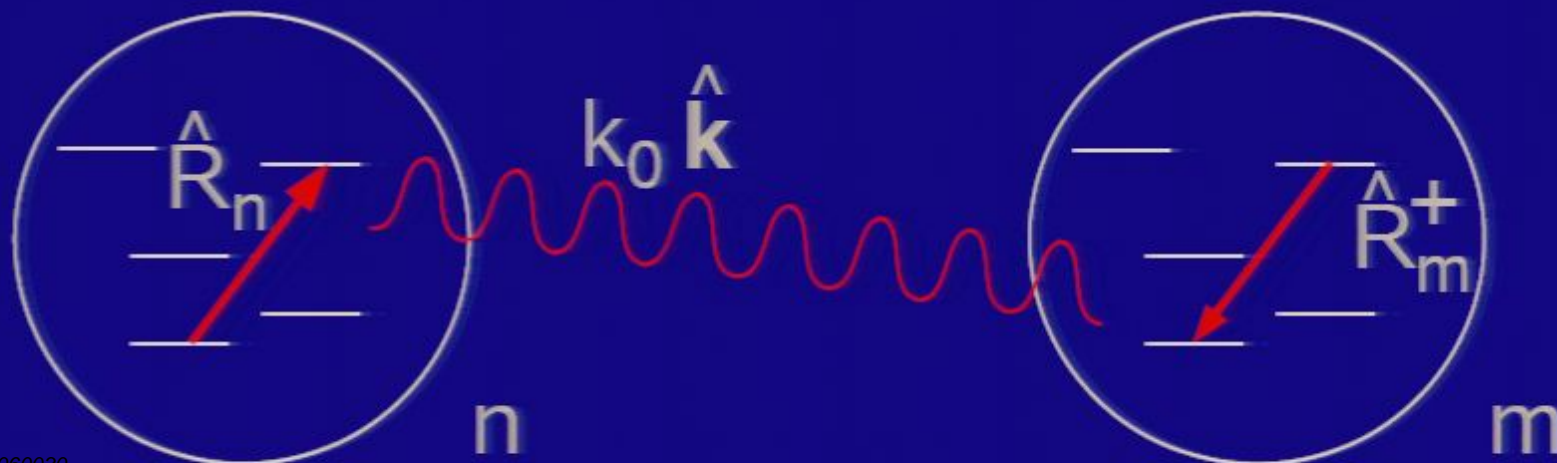
- \hat{k} is a unit vector (direction of the reservoir momentum)
- $\hat{R}_n(\Delta E)$ changes the energy of particle n by ΔE



$$\hat{\Gamma} = \sum_{\Delta E} \int d^2 \hat{k} \sum_{n,m} e^{ik_0 \hat{\mathbf{k}} \cdot (\mathbf{x}_n - \mathbf{x}_m)} \hat{R}_n(\Delta E) \hat{R}_m^\dagger(\Delta E)$$

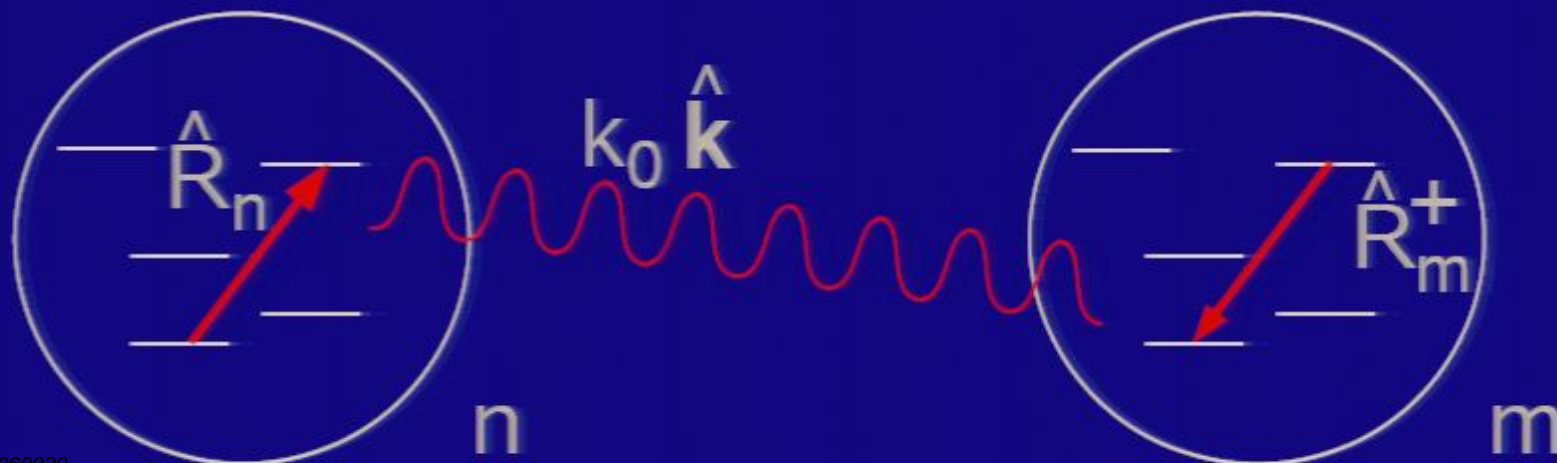
Our proof exploits that $\hat{\Gamma}$ can only have zero eigenvalues for all reservoir modes $\hat{\mathbf{k}}$ if $\mathbf{x}_n = \mathbf{x}_m$, i.e., in the Dicke limit

Otherwise the integral will always contain nonzero parts



This result points out where DFS may be realized:

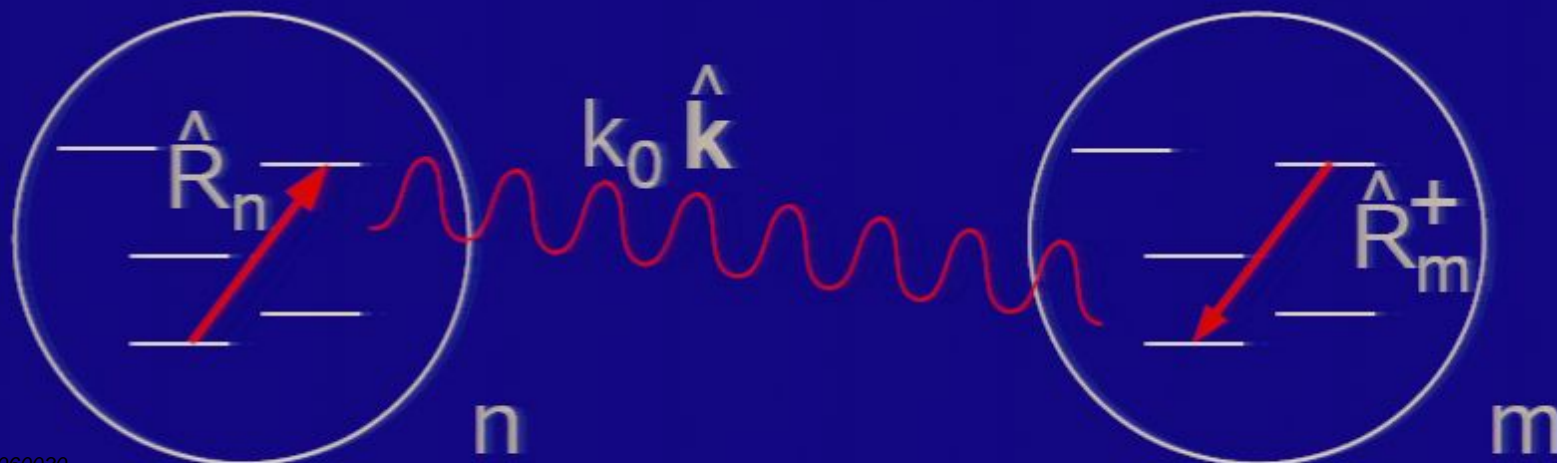
- Single-particle DFS (spontaneous emission cancellation)
- The integral disappears for a 1D reservoir.
Waveguides may therefore allow for ordinary DFS outside the Dicke limit



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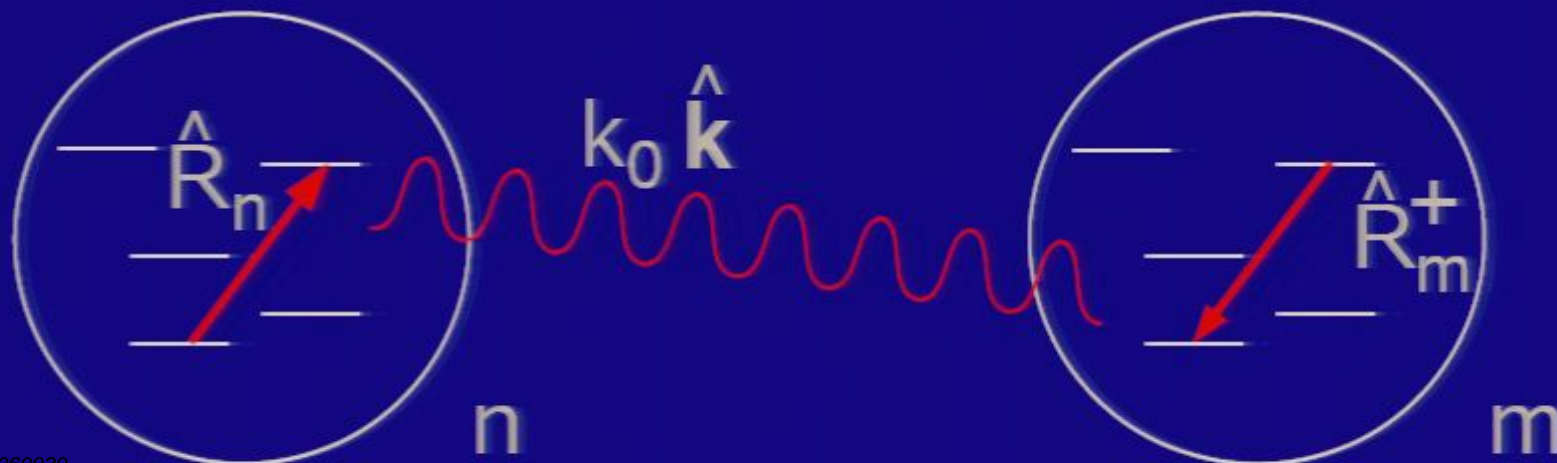
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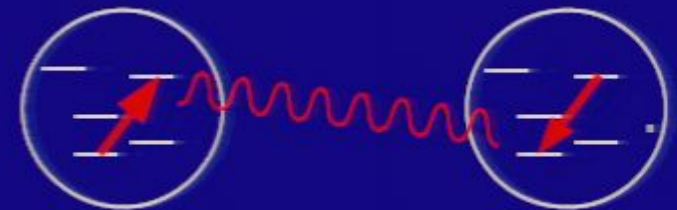
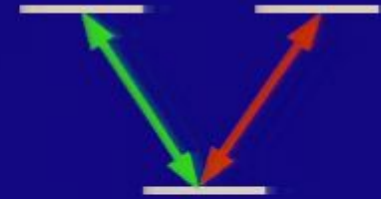


Conclusions



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- DFS are a tool to suppress decoherence
- DFS can only exist in Dicke limit,
- single-particle DFS, or 1D reservoirs may be a way around this problem



Thanks



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