

Title: Strings/Quantum Gravity 4

Date: Jun 08, 2006 03:30 PM

URL: <http://pirsa.org/06060015>

Abstract:

Hairy Branes and D-Branes
Behind the Horizon.

Moshe Rozali, UBC

with Kazumi Okuyama hep-th / 0602060

JHEP 0603: 071. (more refs
in paper).

Outline:

- * 2dim Strings and holographic (Matrix) description.
- * 2dim Black holes, Euclidean and Lorentzian.
(holographic description? look at D-branes)
- * D-Branes: trajectory, smearing, radiation.
- * D-Branes: analytic continuation, world-volume theory.
- * Conclusions: towards holographic description of 2dim Lorentzian black hole.

* Conclusions: towards holographic description of 2dim
Lorentzian black hole.

2 dim Strings, or Linear Dilaton Backgrounds

Main interest is in backgrounds which are asymptotical
of the form

metric = flat (in string frame).

dilaton = $\Phi = -\frac{Q}{2}x$ along single spacelike
direction

2 dim Strings, or Linear Dilaton Backgrounds

Main interest is in backgrounds which are asymptotical
of the form

metric = flat (in string frame).

$$\text{dilaton} = \Phi = -\frac{Q}{2} x$$

along single spacetime
direction

For example 2dim string theory

$$\left\{ ds^2 = dx^2 + dy^2 \right\}$$

string coupling

2 dim Strings, or Linear Dilaton Backgrounds

Main interest is in backgrounds which are asymptotical
of the form

metric = flat (in string frame).

$$\text{dilaton} = \Phi = -\frac{\alpha'}{2} x$$

along single spacetime
direction

For example 2dim string theory

$$\left\{ \begin{array}{l} ds^2 = dx^2 + dy^2 \\ \Phi = -x \end{array} \right.$$

string coupling
is e^Φ

[string coupling becomes large, need to add the Liouville wall].

Main interest is in backgrounds which are asymptotical
of the form

metric = flat (in string frame).

dilaton = $\Phi = -\frac{\sigma}{2} x$ along single spacelike
direction

For example 2dim string theory

$$\left\{ \begin{array}{l} ds^2 = dx^2 + dy^2 \\ \Phi = -x \end{array} \right\} \quad \begin{array}{l} \text{string coupling} \\ \text{is } e^{\Phi} \end{array}$$

[string coupling becomes large, need to add the Liouville wall].

Main interest is in backgrounds which are asymptotical
of the form

metric = flat (in string frame).

dilaton = $\Phi = -\frac{\sigma}{2} x$ along single spacelike
direction

For example 2dim string theory

$$\left\{ \begin{array}{l} ds^2 = dx^2 + dy^2 \\ \Phi = -x \end{array} \right\} \quad \begin{array}{l} \text{string coupling} \\ \text{is } e^{\Phi} \end{array}$$

[string coupling becomes large, need to add the Liouville wall].

metric = flat (in string frame).

$$\text{dilaton} = \Phi = -\frac{\alpha}{2} x$$

along single spacetime
direction

For example 2dim string theory

$$\left\{ \begin{array}{l} ds^2 = dx^2 + dy^2 \\ \Phi = -x \end{array} \right\} \quad \begin{array}{l} \text{string coupling} \\ \text{is } e^{\Phi} \end{array}$$

[string coupling becomes large, need to add the Liouville wall].

Which is the simplest example of an interesting set of theories arising in near horizon limit of NS5-branes. Their

For example 2dim string theory

$$\left\{ \begin{array}{l} ds^2 = dx^2 + dy^2 \\ \Phi = -x \end{array} \right\} \quad \begin{array}{l} \text{string coupling} \\ \text{is } e^\Phi \end{array}$$

[string coupling becomes large, need to add the Liouville wall].

Which is the simplest example of an interesting set of theories arising in near horizon limit of NS5-branes. Their holographic description nicely complements (and pre-dates)

AdS/CFT.

[string coupling becomes large, need to add the Liouville wall].

Which is the simplest example of an interesting set of theories arising in near horizon limit of NS5-branes. Their holographic description nicely complements (and pre-dates)

AdS/CFT.

Holographic Dual:

This is another example (e.g. AdS/CFT) of gravitational theories with holographic dual: non-gravitational "boundaries".

Which is the simplest example of an interesting set of theories arising in near horizon limit of NS5-branes. Their holographic description nicely complements (and pre-dates) AdS/CFT.

Holographic Dual:

This is another example (e.g. AdS/CFT) of gravitational theories with holographic dual: non-gravitational "boundary theory" encodes holographically the bulk gravitational

holographic description nicely complements (and pre-dates)

AdS/CFT.

Holographic Dual:

This is another example (e.g. AdS/CFT) of gravitational theories with holographic dual: non-gravitational "boundary" theory encodes holographically the bulk gravitational theory.

Most generally the duals are LST, for the simplest

This is another example (e.g. AdS/CFT) of gravitational theories with holographic dual: non-gravitational "boundary" theory encodes holographically the bulk gravitational theory.

Most generally the duals are LST, for the simplest case it is simply gauged matrix model.

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right]$$

in suitable (double-scaling limit)

This is another example (e.g. AdS/CFT) of gravitational theories with holographic dual: non-gravitational "boundary" theory encodes holographically the bulk gravitational theory.

Most generally the duals are LST, for the simplest case it is simply gauged matrix model.

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right]$$

in suitable (double-scaling
limit

theory encodes holographically the bulk gravitational theory.

Most generally the duals are LST, for the simplest case it is simply gauged matrix model.

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right] \quad \text{in suitable (double-scaling) limit}$$

This can be understood a la AdS/CFT as theory on

D-branes near maximally supersymmetric in the linear dilaton background

Most generally the duals are LST, for the simplest case it is simply gauged matrix model.

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right] \quad \text{in suitable (double-scaling) limit}$$

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides

case it is simply gauged matrix model.

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right] \quad \text{in suitable (double-scaling) limit}$$

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right] \quad \text{in suitable (double-scaling) limit}$$

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic

$$S = \int dt \left[\frac{1}{2} (\dot{x})^2 - V(x) \right]$$

in suitable (double-scaling
limit

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic dual].

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic dual].

limit

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic dual].

The theory has no propagating gravitons, but it does have

This can be understood a la AdS/CFT as theory on D-branes, now moving in the linear dilaton backgrounds.

[roughly, these correspond to the eigenvalues of the matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic dual].

The theory has no propagating gravitons, but it does have black holes (at least w/o the Liouville wall) ...

matrix X ; theory on world-volume of D-branes provides the holographic description of string background, though there's no precise algorithmic way to deduce the holographic dual].

The theory has no propagating gravitons, but it does have black holes (at least w/o the Liouville wall) ...

Black Holes:

Euclidean CFT:

$$\begin{cases} ds^2 = k (d\rho^2 + \tanh^2(\rho) d\theta^2) + \dots \\ e^{\Phi} = \frac{1}{\cosh(\rho)} \end{cases}$$



- k = curvature scale : $\frac{1}{2}$ correction under central (extremal)

Black Holes:

Euclidean CFT:

$$\begin{cases} ds^2 = k (d\rho^2 + \tanh^2(\rho) d\vartheta^2) + \dots \\ e^{\Phi} = \frac{1}{\cosh(\rho)} \end{cases}$$



- k = curvature scale; $\frac{1}{k}$ corrections under control (= strings corrections)
- Euclidean BH corresponds to known deformation of matrix model.
- holographic dual for Lorentzian BH is unknown.
- Some evidence BH don't form in 2dim case.
- Suggestion BH may be related to multi-particle-hole pairs.

Black Holes:

Euclidean CFT:

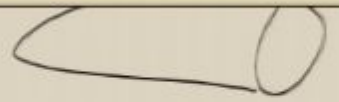
$$\begin{cases} ds^2 = k (d\rho^2 + \tanh^2(\rho) d\theta^2) + \dots \\ e^{\Phi} = \frac{1}{\cosh(\rho)} \end{cases}$$



- k = curvature scale; $\frac{1}{k}$ corrections under control (= strings corrections)
- Euclidean BH corresponds to known deformation of matrix model.
- holographic dual for Lorentzian BH is unknown.
- Some evidence BH don't form in 2dim case.
- Suggestion BH may be related to multi-particle-hole pairs.

Euclidean BH

$$e^{\Phi} = \frac{1}{\cosh(\rho)}$$



- u = curvature scale; $\frac{1}{u}$ corrections under control (= strings corrections)
- Euclidean BH corresponds to known deformation of matrix model.
- holographic dual for Lorentzian BH is unknown.
- Some evidence BH don't form in 2dim case.
- Suggestion BH may be related to multi-particle-hole pairs.

$\Rightarrow$ look at D-Branes in BH background.

These are also interesting as time-dependent (open)

- κ = curvature scale; $\frac{1}{\kappa}$ corrections under control (= strings corrections)
- Euclidean BH corresponds to known deformation of matrix model.
- holographic dual for Lorentzian BH is unknown.
- Some evidence BH don't form in 2dim case.
- Suggestion BH may be related to multi-particle-hole pairs.

$\Rightarrow$ look at D-Branes in BH background.

These are also interesting as time-dependent (open)

- $\ell =$ curvature scale; $\frac{1}{\ell}$ corrections under control (= strings corrections)
- Euclidean BH corresponds to known deformation of matrix model.
- holographic dual for Lorentzian BH is unknown.
- Some evidence BH don't form in 2dim case.
- Suggestion BH may be related to multi-particle-hole pairs.

$\Rightarrow$ look at D-Branes in BH background.

Those are also interesting as time-dependent (open) string background, complementing rolling tachyon discussion.

⇒ look at D-Branes in BH background.

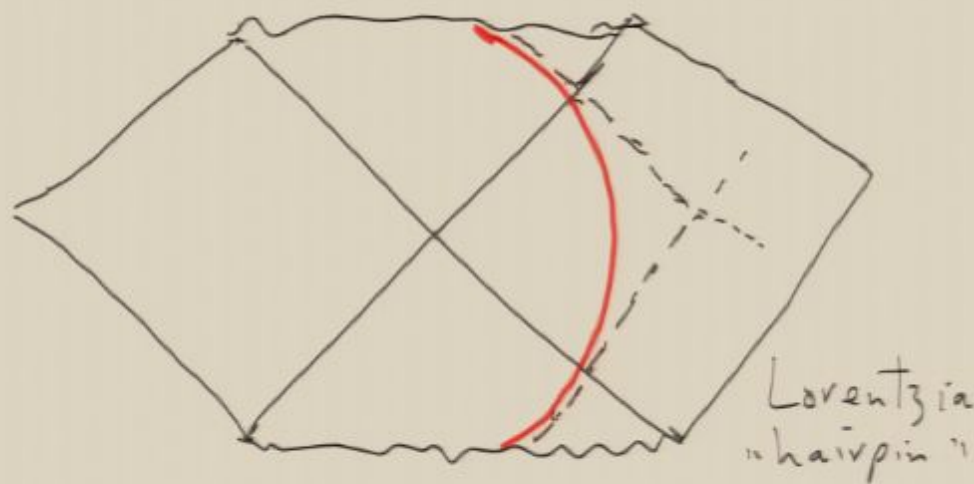
Those are also interesting as time-dependent (open) string background, complementing rolling tachyon discussion.

Hairpin Branes:

trajectory in:

$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due



String background, complementing rolling Tachyon discussion.

Hairpin Branes:

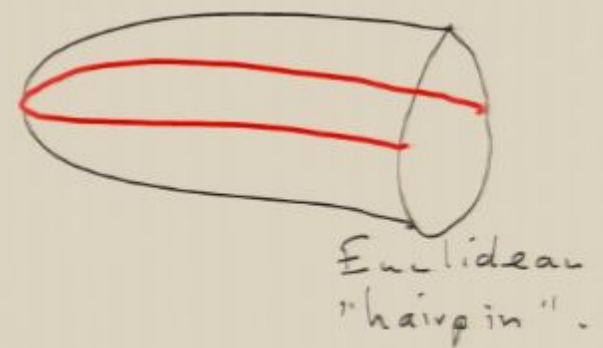
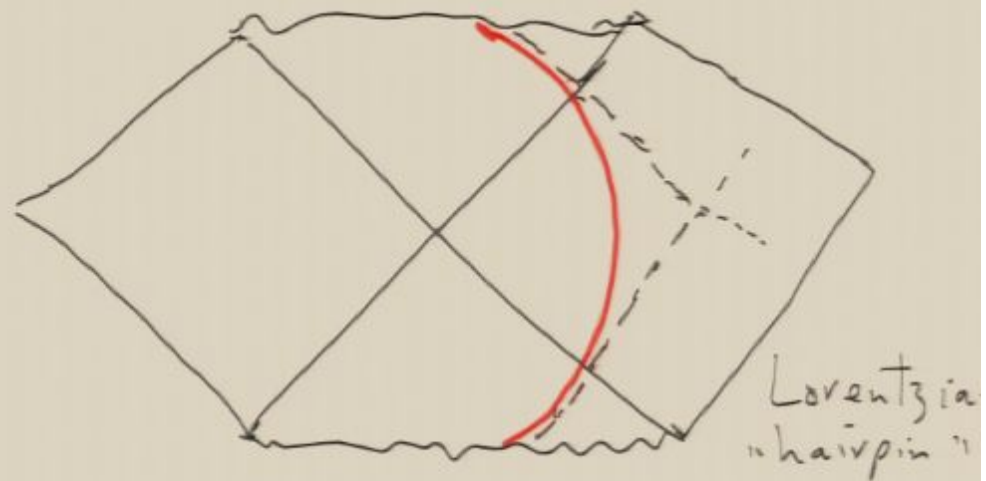
trajectory is:

$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due

to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)

The acceleration determined by parameter Q .



Hairpin Branes:

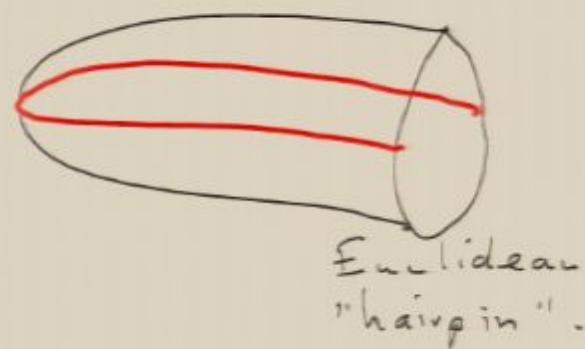
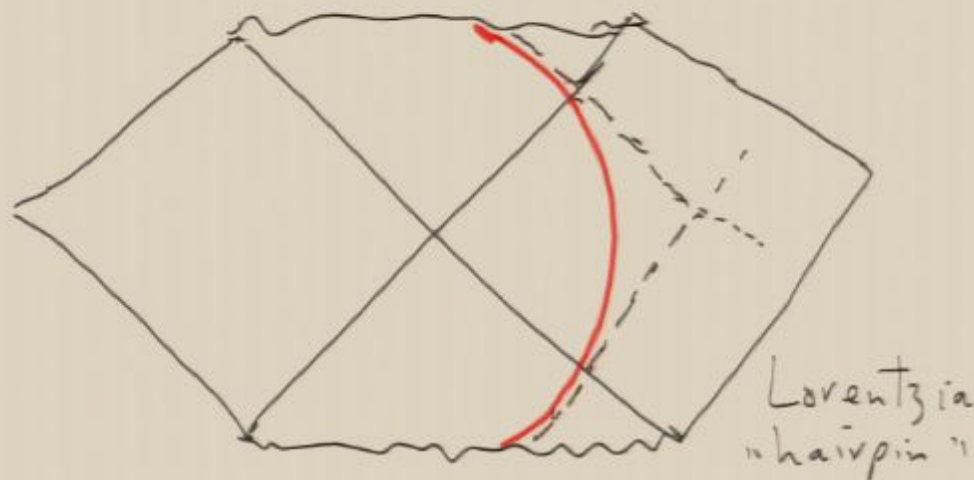
trajectory u:

$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due

to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)

The acceleration determined by parameter Q .

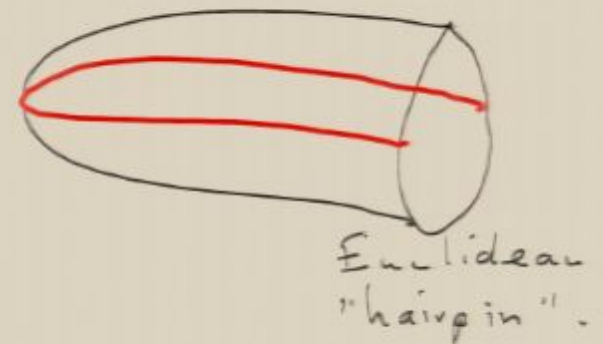


Trajectory u:



$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due
to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)



The acceleration determined by parameter Q .

Due to constant acceleration there is also Rindler
horizon (drawn).

Exact boundary state (Euclidean) is known, so strings

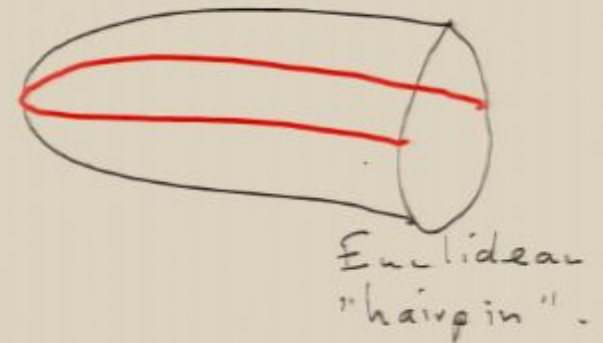
trajectory is:



$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due

to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)



The acceleration determined by parameter Q .

Due to constant acceleration there is also Rindler horizon (drawn).

Exact boundary at $t = 0$ (Euclidean) is known so string

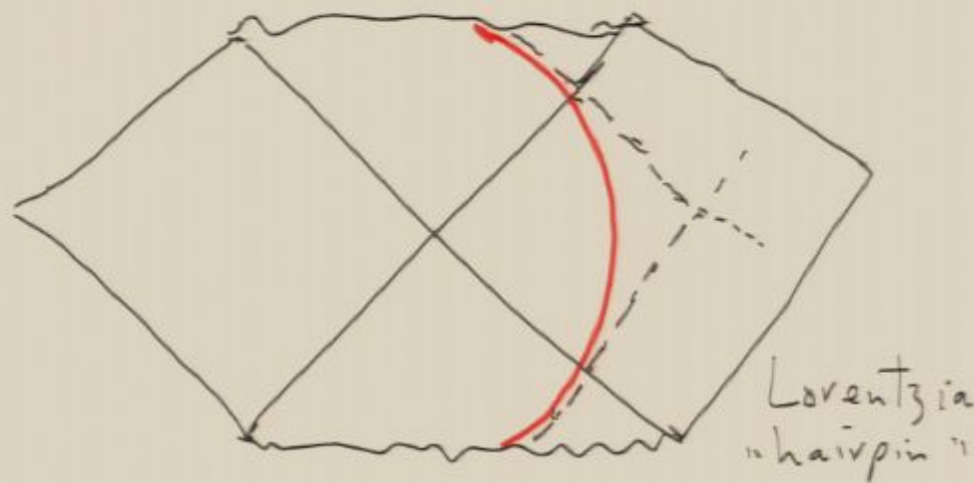
$\Rightarrow$ look at D-Branes in BH background.

Those are also interesting as time-dependent (open) string background, complementing rolling tachyon discussion.

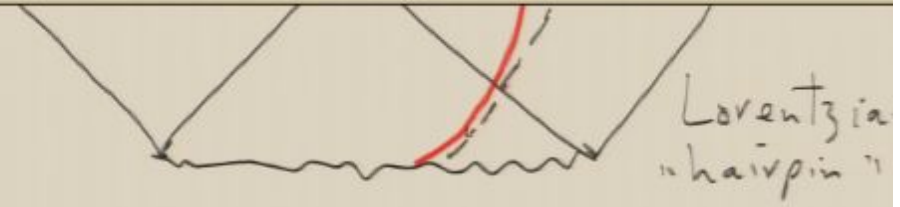
Hairpin Branes:

trajectory is:

$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

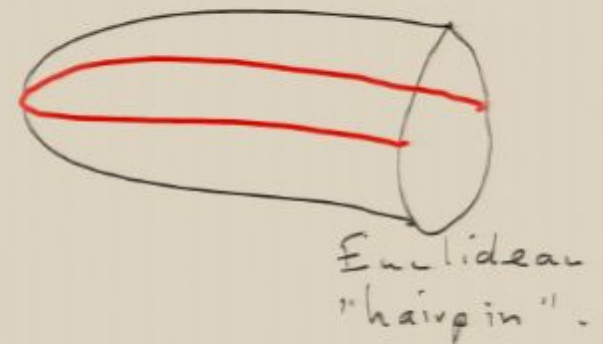


Trajectory u:



$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due
to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)



The acceleration determined by parameter Q .

Due to constant acceleration there is also Rindler
horizon (drawn).

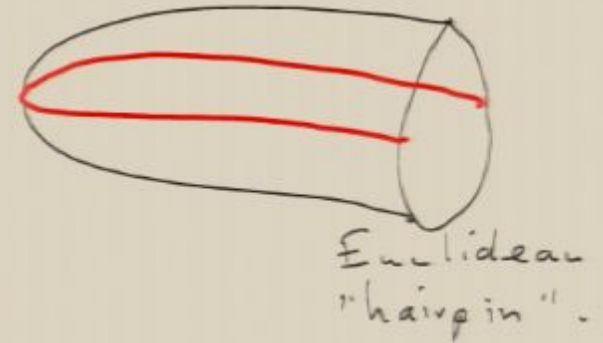
Exact boundary state (Euclidean) is known, so strings

$$\cosh(t) \sinh(\rho) = \text{Constant (modulus)}$$

Accelerating trajectory, force due

to variation of dilaton (since $\text{mass} \sim \frac{1}{g_{\text{str}}}$)

The acceleration determined by parameter Q .



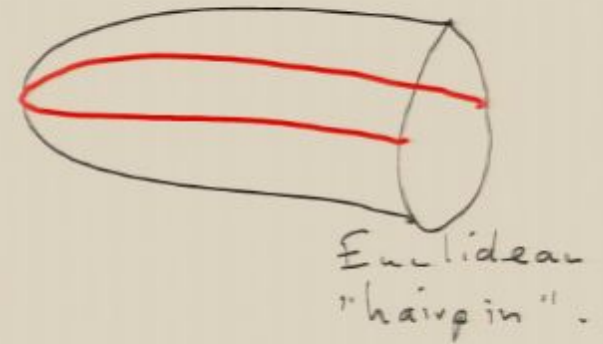
Due to constant acceleration there is also Rindler horizon (drawn).

Exact boundary state (Euclidean) is known, so strings corrections are under control, and are fairly simple.

Constant (modulus)

Accelerating trajectory, force due
to variation of dilaton (since $mass \sim \frac{1}{g_{str}}$)

The acceleration determined by parameter Q .



Due to constant acceleration there is also Rindler
horizon (drawn).

Exact boundary state (Euclidean) is known, so strings
corrections are under control, and are fairly simple.

Boundary State: all details about Euclidean Q -brane

to variation of dilaton (since $mass \sim \frac{1}{g_{str}}$)

Euclidean
"hairpin".

The acceleration determined by parameter Q .

Due to constant acceleration there is also Rindler horizon (drawn).

Exact boundary state (Euclidean) is known, so string corrections are under control, and are fairly simple.

Boundary State: all details about Euclidean D-brane

are summarized by boundary state: $|B\rangle = \int d^2p B(p) |p\rangle$

horizon (drawn).

Exact boundary state (Euclidean) is known, so string corrections are under control, and are fairly simple.

Boundary State: all details about Euclidean D-brane

are summarized by boundary state: $|B\rangle\rangle = \int d^2p B(p) |p\rangle\rangle$

$$\psi(p) \sim \frac{\Gamma(-z + \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)}$$

$$\alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by

corrections are under control, and are fairly simple.

Boundary State: all details about Euclidean D-brane

are summarized by boundary state: $|B\rangle\rangle = \int d^2p \, B(p) |p\rangle\rangle$

$$\psi(p) \sim \frac{\Gamma(-z: \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string field. For example it summarizes the shape of the D-brane as probed by gravitons.

corrections are under control, and are fairly simple.

Boundary State: all details about Euclidean D-brane are summarized by boundary state: $|B\rangle\rangle = \int d^2p \, B(p) |p\rangle\rangle$

$$\psi(p) \propto \frac{\Gamma(-z + \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string field. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s. includes information about the total

Boundary State: all details about Euclidean D-brane

are summarized by boundary state: $|B\rangle = \int d^2p \, B(p) |p\rangle\rangle$

$$\psi(p) \sim \frac{\Gamma(-z - \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total

Boundary State: all details about Euclidean D-brane

are summarized by boundary state: $|B\rangle = \int d^2p \, B(p) |p\rangle\rangle$

$$\psi(p) \sim \frac{\Gamma(-z: \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string field. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total

mass of D-brane, as such it has some positivity properties

are summarized by boundary state: $|B\rangle = \int d^2 p B(p) |p\rangle\rangle$

$$\psi(p) \sim \frac{\Gamma(-z - \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total mass of D-brane, as such it has some positivity properties.

$$\psi(p) \sim \frac{\Gamma(-z: \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total mass of D-brane, as such it has some positivity properties. We'll see that upon analytic continuation those are problematic.

$$\psi(p) \sim \frac{\Gamma(-z: \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total mass of D-brane, as such it has some positivity properties. We'll see that upon analytic continuation those are problematic.

$$\psi(p) \sim \frac{\Gamma(-z - \frac{p_x}{Q}) \Gamma(1 - i p_x Q)}{\Gamma(\frac{1}{2} - i \frac{p_x}{Q} - \alpha p_y) \Gamma(\frac{1}{2} - i \frac{p_x}{Q} + \alpha p_y)} \quad \alpha = \sqrt{1 + \frac{1}{Q^2}}$$

The coefficients $\psi(p)$ correspond to the source supplied by the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total mass of D-brane, as such it has some positivity properties. We'll see that upon analytic continuation those are problematic.

the D-brane to all closed string fields. For example it summarizes the shape of the D-brane as probed by gravitons.

In particular the b.s includes information about the total mass of D-brane, as such it has some positivity properties. We'll see that upon analytic continuation these are problematic.

Features: (Euc. or Lorentzian Branes).

Stringy smearing: exact boundary state is obtained

by smearing the classical trajectory over its position

mass of D-brane, as such it has some positive properties.

We'll see that upon analytic continuation those are problematic.

Features: { Euclidean Branes }
Features: { Lorentzian Branes }.

Stringy smearing: exact boundary state is obtained

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/\alpha}\right]$

This comes about due to open string halo around the

Features: (Eu. or Lorentzian Branes).

Stringy smearing: exact boundary state is obtained by smearing the classical trajectory over its position

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/Q}\right]$

This comes about due to open string halo around the trajectory. That halo becomes of infinite range at finite Q .

Closed String Radiation:

Features: (Eucl. or Lorentzian Branes).

Stringy smearing: exact boundary state is obtained by smearing the classical trajectory over its position

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/\alpha}\right]$

This comes about due to open string halo around the trajectory. That halo becomes of infinite range at finite α .

Closed String Radiation:

Stringy smearing: exact boundary state is obtained by smearing the classical trajectory over its position

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/\alpha}\right]$

This comes about due to open string halo around the trajectory. That halo becomes of infinite range at finite α .

Closed String Radiation:

Due to time dependence, the brane loses energy to

Stringy Smearing: exact boundary state is obtained by smearing the classical trajectory over its position

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/Q}\right]$

This comes about due to open string halo around the trajectory. That halo becomes of infinite range at finite Q.

Closed String Radiation:

Due to time dependence, the brane loses energy to closed string radiation; the total energy is finite.

modulus with weight: $\exp\left[-\frac{s}{\alpha} - e^{-s/\alpha}\right]$

This comes about due to open string halo around the trajectory. That halo becomes of infinite range at finite α .

Closed String Radiation:

Due to time dependence, the brane loses energy to closed string radiation; the total energy is finite.

Due to smearing the radiation extends beyond the would-be Rindler horizon, all the way to infinity.

Due to time dependence, the brane loses energy to closed string radiation; the total energy is finite.

Due to smearing the radiation extends beyond the would-be Rindler horizon, all the way to infinity.

Analytic Continuation:

Lorentzian quantities are obtained by analytic continuation, with some ambiguity reflecting vacuum choice. The boundary state presents a new such ambiguity.

Due to smearing the radiation extends beyond the would-be Rindler horizon, all the way to infinity.

Analytic Continuation:

Lorentzian quantities are obtained by analytic continuation, with some ambiguity reflecting vacuum choice. The boundary state presents a new such ambiguity.

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent"

Analytic Continuation:

Lorentzian quantities are obtained by analytic continuation, with some ambiguities reflecting vacuum choice. The boundary state presents a new such ambiguity.

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

Analytic Continuation:

Lorentzian quantities are obtained by analytic continuation, with some ambiguities reflecting vacuum choice. The boundary state presents a new such ambiguity.

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

$$\Psi_{\text{Lor.}}(p) \propto \frac{\Gamma(\tau + ipQ)}{\Gamma(1 + i\frac{2p}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

with some ambiguity reflecting vacuum choice. The boundary state presents a new such ambiguity.

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(\tau - ipQ)}{\Gamma(1 + i\frac{2p}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

We find that positivity properties of b.s. are not maintained by this prescription (conceivably most other prescriptions).

state presents a new such ambiguity.

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(r - ipQ)}{\Gamma(1 + i\frac{r}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

We find that positivity properties of b.s. are not maintained by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon

In this case the exact boundary state is almost entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(r - ipQ)}{\Gamma(1 + i\frac{rP}{Q})} \Gamma(\frac{1}{2} + i(\frac{p}{Q} + a\omega)) \Gamma(\frac{1}{2} + i(\frac{p}{Q} - a\omega)).$$

We find that positivity properties of b.s. are not maintained by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip.

entirely geometrical, easy to formulate "transparent horizon" boundary state. The b.s. coefficients are then:

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(r - ipQ)}{\Gamma(1 + i\frac{r}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

We find that positivity properties of b.s. are not maintained by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip.

In that sense it is "ghost-like" - its excitation will

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(r - ipQ)}{\Gamma(1 + i\frac{2p}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

We find that positivity properties of b.s are not maintained by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip. In that sense it is "ghost-like", its excitations will have opposite sign in their kinetic energy (or the wrong spin-statistics relation). This is more or less inevitable.

$$\Psi_{\text{hor.}}(p) \propto \frac{\Gamma(\tau - i p Q)}{\Gamma(1 + i \frac{2p}{Q})} \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} + a\omega\right)\right) \Gamma\left(\frac{1}{2} + i\left(\frac{p}{Q} - a\omega\right)\right).$$

We find that positivity properties of b.s are not maintained by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip. In that sense it is "ghost-like", its excitations will have opposite sign in their kinetic energy (or the wrong spin-statistics relation). This is more or less inevitable.

by this prescription (conceivably most other prescriptions).

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip. In that sense it is "ghost-like", its excitations will have opposite sign in their kinetic energy (or the wrong spin-statistics relation). This is more or less inevitable.

Interpretation and Conclusions:

Eternal black holes, with their two asymptotic regions,

In fact, the part of boundary state behind the horizon is related simply to the portion in front by sign flip. In that sense it is "ghost-like", its excitations will have opposite sign in their kinetic energy (or the wrong spin-statistics relation). This is more or less inevitable.

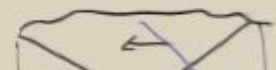
Interpretation and Conclusions:

Eternal black holes, with their two asymptotic regions, are dual to the thermal double of the QFT

is related simply to the portion in front by sign flip.
In that sense it is "ghost-like", its excitations will
have opposite sign in their kinetic energy (or the wrong
spin-statistics relation). This is more or less inevitable.

Interpretation and Conclusions:

Eternal black holes, with their two asymptotic regions,
are dual to the thermal double of the QFT.

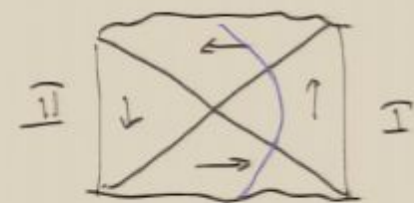


spin-statistics relation). This is more or less inevitable.

Interpretation and Conclusions:

Eternal black holes, with their two asymptotic regions, are dual to the thermal double of the QFT.

excitations of region II look ghost-like (as they are in real-time thermal QFT).

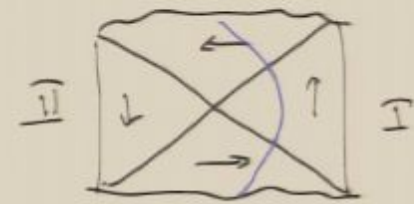


The part of D-brane behind horizon is then excitation of the "ghosts" of region II, or the thermal doublers. Absence

Interpretation and Conclusions:

Eternal black holes, with their two asymptotic regions, are dual to the thermal double of the QFT.

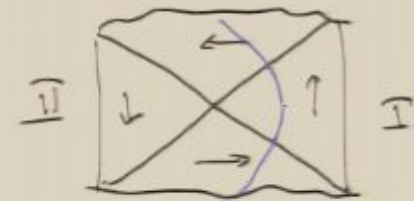
excitations of region II look ghost-like (as they are in real-time thermal QFT).



The part of D-brane behind horizon is then excitation of the "ghosts" of region II, or the thermal doublers. Absence

Eternal black holes, with their two asymptotic regions, are dual to the thermal double of the QFT.

excitations of region II look ghost-like (as they are in real-time thermal QFT).



The part of D-brane behind horizon is then excitation of the "ghosts" of region II, or the thermal doublers. Absence of positivity properties is natural in full real time field Theory.

of positivity properties is natural in full real time field theory.

Dual to Lorentzian Black Holes:

This suggests the dual to black hole ought to include such ghost-like excitations, this vibrates with other suggestions:

- Sen: large number of particle-hole excitations (holes = ghost branes...)
- Drizzling in the eternal BH Penrose diagram.

Dual to Lorentzian Black Holes:

This suggests the dual to black hole ought to include such ghost-like excitations, this agrees with other suggestions:

- Sen: large number of particle-hole excitations (holes = ghost branes...)
- Doubling in the eternal BH Penrose diagram.
- The world-volume theory of D-branes includes such excitations and suggests the correct structure is based on $SU(N|N)$

Dual to Lorentzian Black Holes:

This suggests the dual to black hole ought to include such ghost-like excitations, this vibes with other suggestions:

- Sen: large number of particle-hole excitations (holes = ghost branes...)
- Doubling in the eternal BH Penrose diagram.
- The world-volume theory of D-branes includes such excitations and suggests the correct structure is based on $SU(N|N)$ matrix model in the large N limit.

Dual to Lorentzian Black Holes:

This suggests the dual to black hole ought to include such ghost-like excitations, this agrees with other suggestions:

- Sen: large number of particle-hole excitations (holes = ghost branes...)
- Doubling in the eternal BH Penrose diagram.
- The world-volume theory of D-branes includes such excitations and suggests the correct structure is based on $SU(N|N)$ matrix model in the large N limit.

Dual to Lorentzian Black Holes:

This suggests the dual to black hole ought to include such ghost-like excitations, this vibrates with other suggestions:

- Sen: large number of particle-hole excitations (holes = ghost branes...)
- Doubling in the eternal BH Penrose diagram.
- The world-volume theory of D-branes includes such excitations and suggests the correct structure is based on $SU(N|N)$ matrix model in the large N limit.

- This suggests the dual to black hole ought to include such ghost-like excitations, this vibrates with other suggestions:
- Sen: large number of particle-hole excitations (holes = ghost branes...)
 - Doubling in the eternal BH Penrose diagram.
 - The world-volume theory of D-branes includes such excitations and suggests the correct structure is based on $SU(N|N)$ matrix model in the large N limit.