Title: Constructions and distributions of string vacua

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Abstract:

# Constructions and distributions of string vacua

Frederik Denef

PI, October 2005

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#### Motivation

The Landscape

The good news

The bad news

What can we do?

#### Construction of vacua

IIB KKLT vacua

IIB nonsusy AdS vacua with exponentially large volume

I/IIB with gauge fluxes

M-theory and IIA flux vacua

More models: heterotic, non-geometric, ...

de Sitter vacua

#### Statistics of vacua

Susy IIB

Nonsusy IIB

M-theory

Vacua with enhanced (R-)symmetries

Pirsa: 05100054 Intersecting brane and Gepner models

# Motivation

Not everything that can be counted counts, and not everything that counts can be counted.

- Albert Einstein

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String theory has an infinite number of supersymmetric vacua: e.g.  $[AdS_5 \times S^5]_N$ ;  $N = 1, ..., \infty$ .

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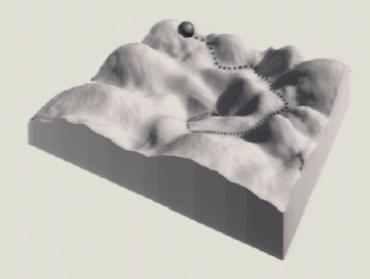
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Picture: String theory Landscape [Susskind]



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In any case, important to know e.g. how small  $\Lambda$  can actually get, or more generally, what a priori number distribution of vacua over parameter space is.

If sufficiently finely scanned, landscape picture offers at least Pirsa: 051@0@sssibility for a consistent explanation for a number of absurd finatunings of parameters in our universe

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 $\infty = a$  lot of vacua!

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 or 0?

Type of problem: given parameter range, find discrete quanta (e.g. fluxes) such that vacuum ends up in this range.

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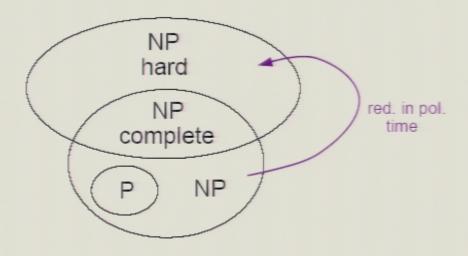
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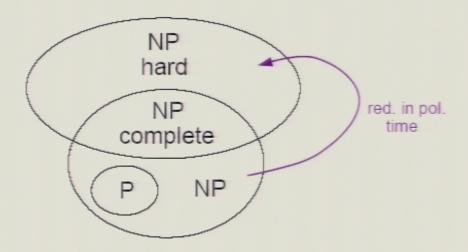
Type of problem: given parameter range, find discrete quanta (e.g. fluxes) such that vacuum ends up in this range.  $\rightsquigarrow$  Hard!

How hard? --> quantified in computational complexity theory.

Even in simple Bousso-Polchinski toy model, the problem to find the flux vectors  $N^{\alpha}$  such that  $0 < \Lambda(N) < \epsilon$  is NP-hard.

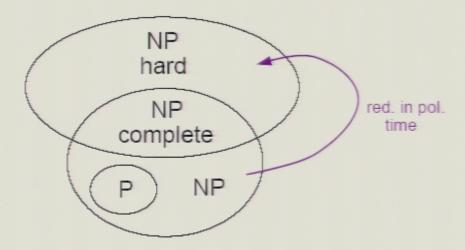


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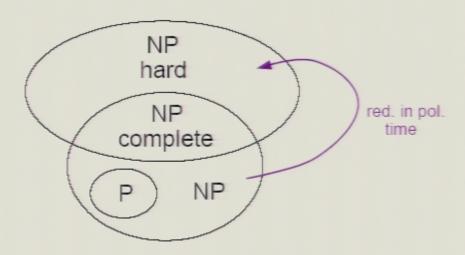
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- ightharpoonup NP-complete = NP  $\cap$  NP-hard (e.g. subset-sum)

Presumably: NP  $\neq$  P, but no proof to date (Clay prize problem).

⇒ if you find a polynomial time algorithm to identify string vacua from parameter data, you're rich...

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  - Distributions also more robust than individual solutions under corrections.
- Try to compute dynamical probabilities on parameter space

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# Construction of vacua

Any intelligent fool can make things bigger and more complex...

It takes a touch of genius, and a lot of courage, to move in the opposite direction.

— Albert Einstein

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# The first part of the program includes

- Intersecting brane models, including Kähler potentials, Yukawa couplings, susy breaking soft terms [Aldazabal, Angelantonj, Antoniadis, Blumenhagen, Camara, Cremades, Cvetic, Dudas, Franco, Görlich, Graña, Grimm, Ibañez, Jockers, Körs, Langacker, Liu, Louis, Lüst, Mayr, Marchesano, Rabadan, Reffert, Richter, Sagnotti, Shiu, Stieberger, Taylor, Uranga, Wang]
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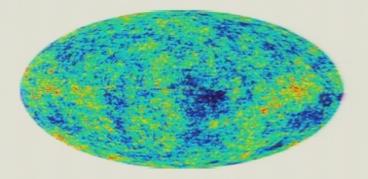
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- ⇒ want mechanism to stabilize moduli at sufficiently high scale.

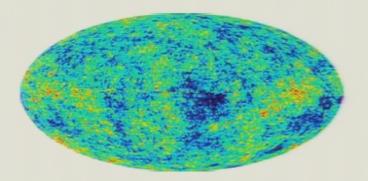
(Quintessence appears very hard to realize in string theory.)

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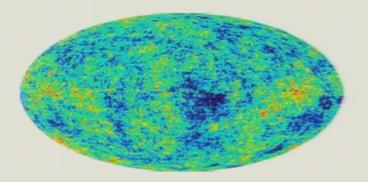


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Cosmological observations are in excellent agreement with inflation and a present time nonzero cosmological constant.

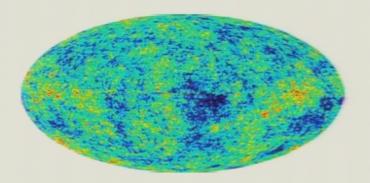
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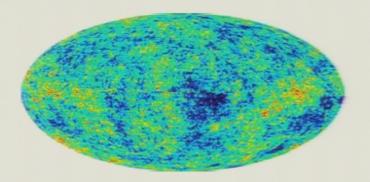


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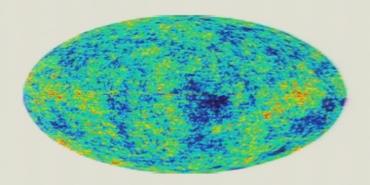
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Problem: find 4d string compactification with large volume, no massless moduli,  $\mathcal{N}=1$  unbroken susy and  $R_{KK}\ll R_{AdS}$ 

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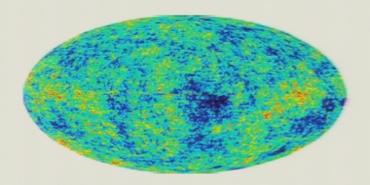
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 $\stackrel{\text{\tiny{00054}}}{\Rightarrow}$  Need higher order curvature corrections or instanton effects.

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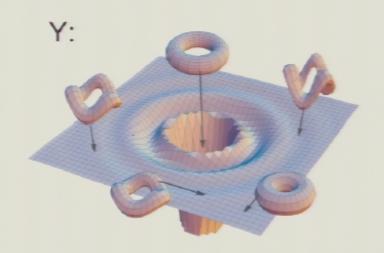
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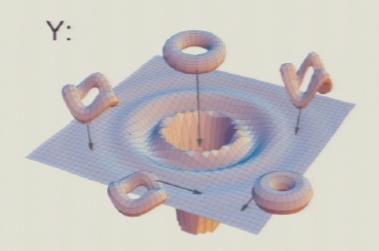
KKLT [Kachru-Kallosh-Linde-Trivedi]: IIB on warped CY<sub>3</sub> orientifold  $Y/\mathbb{Z}_2+\mathsf{RR}\ \mathsf{flux}\ F_3+\mathsf{NS}\ \mathsf{flux}\ H_3$ 



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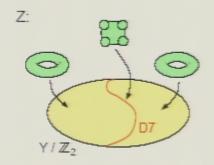
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Moduli: complex (shape) and Kähler (size)

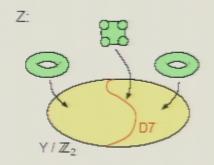
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 $\Leftrightarrow$  M/F-theory on elliptically fibered CY<sub>4</sub> Z + flux  $G_4$ .



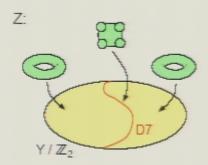
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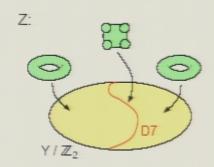


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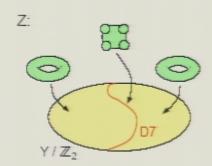
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Pirsa: 0510054 controlled regime? Yes, if  $W_{flux}\ll 1$ ; possible in IIB because  $_{PQ}f_{62/135}$ . Bousso-Polchinski like tuning by flux discretuum.

Explicit examples?

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Explicit examples? Took some time, mainly because of technical complications to prove sufficient number of contributions to W from D3-instantons and gaugino condensates (or from M5-instantons in M-theory dual).

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# By now, several examples known:

- ▶ [Denef-Douglas-Florea]: various constructions of models with a sufficient number of D3 instanton divisors with exactly 2 fermion zeromodes  $(h^{0,i}(M5) = 0$  [Witten]).
- ▶ [Denef-Douglas-Florea-Grassi-Kachru]: completely explicit, simple model:  $T^6/\mathbb{Z}_2 \times \mathbb{Z}_2$ ; all moduli (open, closed, untwisted and twisted) fixed. Gen. by [Lüst-Reffert-Schulgin-Stieberger].
- [Aspinwall-Kallosh]: Stabilize M-theory on K3 × K3, making use of prev. work of [Saulina, Kallosh Kashani-Poor Tomasiello] showing topological conditions on divisors to contribute to W substantially relaxed in the presence of flux. Refined recently by [Lüst-Reffert-Schulgin-Tripathy].

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# IIB nonsusy AdS vacua with exponentially large volume

In [Balasubramanian-Berglund, Balasubramanian-Berglund-Conlon-Quevedo, Conlon-Quevedo-Suruliz] it was shown that, when taking into account  $\alpha'$  corrections to the Kähler potential, a new branch of vacua can appear as nonsusy AdS minima of the potential.

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Rough idea: keep some divisor volumes  $\rho_i \sim O(1)$  while sending overall vol to infinity, and balance nonperturbative  $e^{-\rho_i}$  off against perturbative  $\alpha'^3/Vol$  corrections.

⇒ Volume stabilized at exponentially large value:

$$Vol \sim W_{flux} e^{c/g_s}$$

where  $W_{flux}$  and  $g_s$  are fixed by the fluxes.

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Unlike KKLT, apparently also in well-controlled regime for O(1) values of  $W_{flux}$ .

# I/IIB with gauge fluxes

[Antoniadis-Kumar-Maillard] considered IIB on  $T^6/\mathbb{Z}_2$  orientifold with closed string fluxes and magnetized D9/D7-branes.

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 $\leadsto$  Gauge flux + B-field  $\mathcal{F} \equiv F - B$  gives D-terms, constraining Kähler moduli

$$\operatorname{Im}[e^{-i\theta}e^{\mathcal{F}+iJ}]=0$$

( = mirror to slag cond.), and F-terms, constraining complex structure and open string moduli:

$$\mathcal{F}^{(0,2)} = 0$$

[Brunner-Douglas-Fiol-Römelsberger, Mariño-Minasian-Moore-Strominger, Jockers-Louis, Gomis-Marchesano-Mateos].

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Original AM/AKM models have some problems with tadpole cancellation (also only closed string moduli stabilization). Scenario Pirsa: 0516662 uld still be realizable, though perhaps only with broken susy? Page 71/135

# M-theory flux vacua

M-theory on  $G_2$  hol. manifold X: turning on  $G_4$ -flux in X gives

$$W=\int G_4\wedge (\frac{1}{2}C_3+\Phi_3)$$

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but leads to runaway potential and no susy vacua [Beasley-Witten].

[Acharya] realized that replacing  $W \to W + c$ ,  $\operatorname{Im} c \neq 0$ , does produce susy vacua, and proposed such a c can be produced by holomorphic Chern-Simons contribution living on singularity fibered over certain 3-manifold.

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Rather different compactification is obtained when taking  $G_4$  along 4d spacetime  $\rightsquigarrow$  Freund-Rubin;  $X = \text{weak } G_2$  (Einstein), e.g.  $X = AdS_4 \times S^7$ . Often moduli-free, and can support chiral fermions [Acharya-Denef-Hofman-Lambert], but typically  $R_{KK} \sim R_{AdS}$ .

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More general M-theory compactifications on weak  $G_2$  + fluxes  $G_2$  + fluxes  $G_3$  been discussed e.g. by [Lambert] and in more detail recently [Dall'Agata-Prezas].

## Type IIA flux vacua

Studied in [Derendinger-Kounnas-Petropoulos-Zwirner], in particular for  $T^6/\mathbb{Z}_2 \times \mathbb{Z}_2$  with RR, NS-NS, and *metric* fluxes (torsion), by relating it to 4d gauged sugra. Find all untwisted geometrical moduli can be stabilized.

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Analysis of toroidal case with metric fluxes refined and generalized in [Camara-Ibañez-Font], including tadpole cancellation conditions involving metric fluxes, and inclusion of intersecting brane models.

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Heterotic moduli stabilzation has been studied by [Gukov-Kachru-Liu-McAllister, Buchbinder-Ovrut, Cardoso-Curio-Dall'Agata-Lüst, Becker-Becker-Dasgupta-Green-Sharpe, Curio-Krause-Lüst, Gurrieri-Lukas-Micu]. Complications due to lack of tuning with only H-flux and difficulty describing metric flux (non-Kähler).

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Much harder to establish, because of control issues and many more possible decay channels (perturbative and nonperturbative).

No fully established concrete examples known, although there is at this point no fundamental reason to doubt their existence, and there are various plausible proposals for construction, including:

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- [Saltman-Silverstein]: Susy broken at KK scale (flux compactifications on product of Riemann surfaces).
- [Maloney-Silverstein-Strominger, Silverstein]: Susy broken at string scale (noncritical string theories).

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Metastability of nonsusy vacua? Some progress in [Hong-Mitra]:  $AdS_4 \times X_3 \times Y_3$  orientifold + flux + (anti-)D3 with Einstein  $X_3$  and  $Y_3$  is in general perturbatively unstable.

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- Particle physics constraints.
- Cosmological constraints. Strong! E.g. [Kofman-Yi]: reheating after brane-antibrane inflation in Giddings-Kachru-Polchinski type warped compactifications produces angular KK modes severely overclosing the universe.

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## Statistics of vacua

We can't solve problems by using the same kind of thinking we used when we created them

- Albert Einstein

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[Douglas, Ashok-Douglas, Denef-Douglas, Douglas-Shiffman-Zelditch]

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Vacuum characterized by discrete (compactification) data  $\vec{N}$  and critical point of effective potential  $V_N(z)$ :

$$(\vec{N},z):V_N'(z)=0, V_N''(z)>0$$

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$$\mathcal{N}_{vac}(z \in \mathcal{S}) = \sum_{\vec{N}} \int_{\mathcal{S}} d^n z \, \delta^n(V_N'(z)) |\det V_N''(z)|$$

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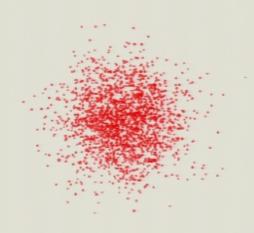
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Pirsa: 05100054 Page 102/135

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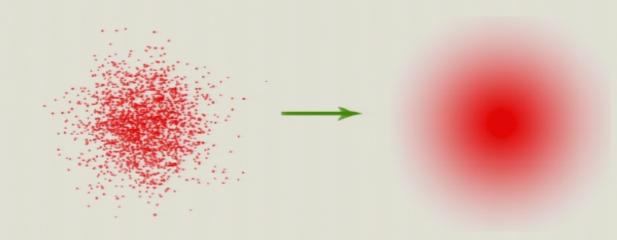


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→ not very practical: need some more structure + approx. -

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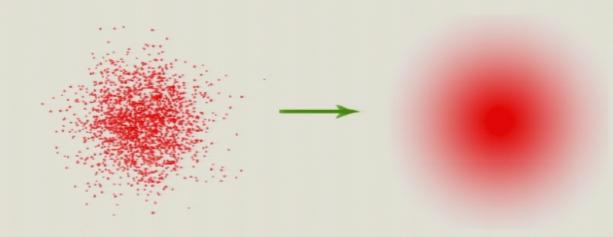
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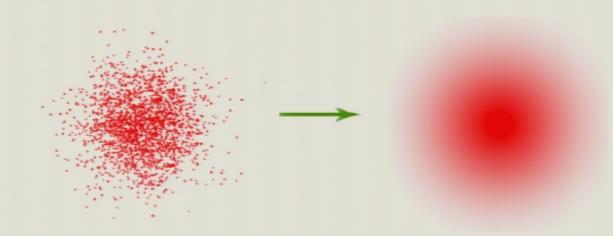


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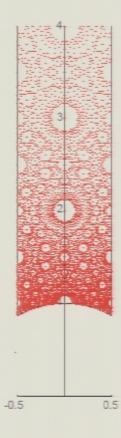
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 $\rho(z) = \int d\mu_0 [W, F, M, Y]_z f(W, F, M, Y)_z \rightarrow \text{finite dim. irRage}^{108/135}$ 

# Toy example

Type IIB on rigid CY ( $\Rightarrow$  only dilaton-axion  $\tau$ ).



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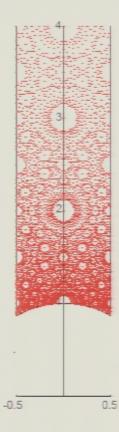
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$$\mathcal{N}_{vac}(\mathcal{R}) = \int_{\mathcal{R}} \frac{d^{2}\tau}{(\text{Im}\,\tau)^{2}} \int_{|W|^{2}-|F|^{2} \leq L_{*}} d^{2}W \, d^{2}F \, \delta^{2}(F)|W|^{2} \\
= \frac{\pi L_{*}^{2}}{2} \int_{\mathcal{R}} \frac{d^{2}\tau}{(\text{Im}\,\tau)^{2}} e^{-2\pi i \pi T} e^{-2\pi i T$$

▶ Number of flux vacua in region S of moduli space

$$\mathcal{N}_{\mathcal{S}}(L \leq L_*) \approx \frac{(2\pi L_*)^{b_3}}{b_3!} \int_{\mathcal{S}} \frac{1}{\pi^m} \det(R + \omega \mathbf{1})$$

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Example [Giryavets-Kachru-Tripathy-Trivedi]:  $X_3 = CY$  hypersurface in WP[1,1,1,1,4],  $X_4 = CY$  hypersurface in WP[1,1,1,1,8,12]. Has  $\chi/24 = 972$ ,  $b_3 = 300$ , so

$$\mathcal{N}_{vac} \sim 10^{500}$$

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▶ Number of flux vacua in region S of moduli space

$$\mathcal{N}_{\mathcal{S}}(L \leq L_*) \approx \frac{(2\pi L_*)^{b_3}}{b_3!} \int_{\mathcal{S}} \frac{1}{\pi^m} \det(R + \omega \mathbf{1})$$

where  $L_* = \chi(X_4)/24$ .

Example [Giryavets-Kachru-Tripathy-Trivedi]:  $X_3 = \text{CY}$  hypersurface in WP[1,1,1,1,4],  $X_4 = \text{CY}$  hypersurface in WP[1,1,1,1,8,12]. Has  $\chi/24 = 972$ ,  $b_3 = 300$ , so

$$\mathcal{N}_{vac} \sim 10^{500}$$

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 $\Rightarrow$  smallest c.c.  $\sim M_s^4/N_{vac}$ .

String coupling  $g_s$ : again uniformly distributed.

Vacua cluster near conifold degenerations:

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Pirsa: 05100054 of above tested by Monte Carlo simulations

# Nonsusy IIB

F-breaking vacua, 
$$F=:M_{susy}^2\ll M_p^2$$
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 $\rightarrow$  low breaking scale disfavored (but much less than naive guess  $dF^{2n}$ )

Above worked out in detail in [Denef-Douglas 2], but there is more intuitive argument in 1-field model [Dine-O'Neil-Sun]:

$$W = W_0 + Fz + Mz^2 + Yz^3$$

Solving V'(0) = 0 requires  $M = 2W_0$ , c.c. near zero requires  $|F|^2 \approx 3|W_0|^2$  and stability in this case requires |Y| < |F|.

 $\Rightarrow$  in addition to tuning c.c. three complex parameters need to be tuned small, of order |F|. In generic ensmemble: independent

Pirsa: 05100054  $dN \sim d\Lambda d|F|^6$ .

# M-theory on $G_2$

#### [Acharya-Denef-Valandro]

ightharpoonup Number of susy flux vacua in region S of moduli space

$$\mathcal{N}_{\mathcal{S}} \sim c_2^{b_3} \int_{\mathcal{S}} \det g$$

Large volumes strongly suppressed:

$$d\mathcal{N}[V] \sim (kc_2)^{b_3} dV^{-3b_3/7} \Rightarrow V < (kc_2)^{7/3}$$

Small cc's strongly suppressed because

$$\Lambda \sim 1/V^3$$

 Small F-breaking susy breaking scales strongly suppressed because

$$M_{susv}^2 \sim 1/V^{3/2}$$

# Vacua with enhanced (R-)symmetries

[DeWolfe-Giryavets-Kachru-Taylor, DeWolfe, Dine-Sun]: study explicit constructions and statistics of IIB flux vacua with discrete symmetries (mostly R-symmetries).

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Open string sector statistics:

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#### Open string sector statistics:

▶ [Gmeiner-Blumenhagen-Honecker-Lüst-Weigand]: statistics of intersecting brane models on  $T^6/\mathbb{Z}_2 \times \mathbb{Z}_2$  (counting number of solutions to tadpole condition [NP-hard!]).

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#### Fraction of models with

- ► SU(3): 0.08
- ► SU(2): 0.99
- ▶ U(1): 0.42
- ▶ No symm. rep: 0.84
- ▶ 3 gen. quarks: 3 × 10<sup>-5</sup>
- ▶ 3 gen. leptons: 2 × 10<sup>-3</sup>

Total:  $10^{-9}$ .

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