

Title: Cosmological Model Selection and the Inflationary Cosmology

Date: Oct 24, 2005 03:00 PM

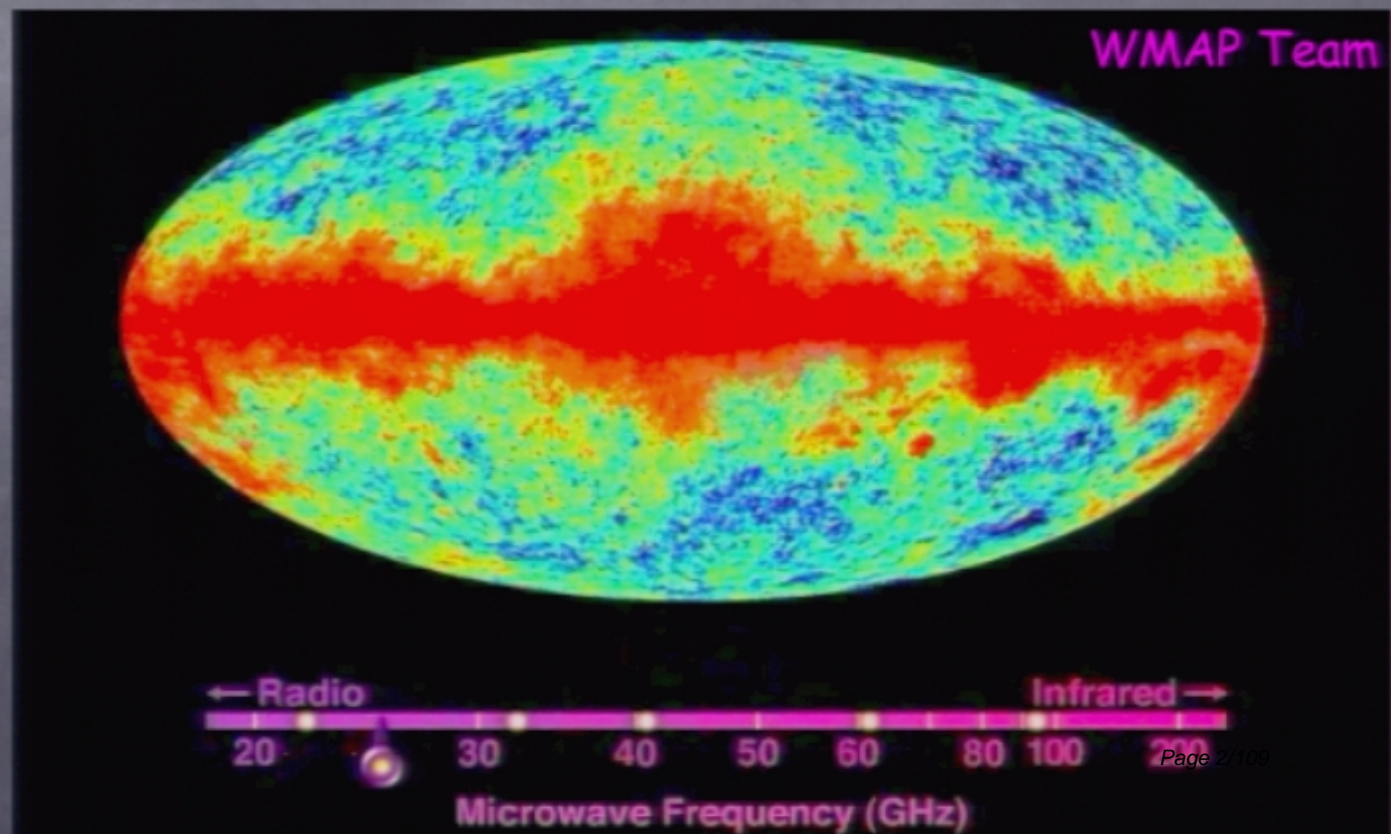
URL: <http://pirsa.org/05100041>

Abstract:

Cosmological model selection and inflation

US

University of Sussex



Microsoft-free
presentation



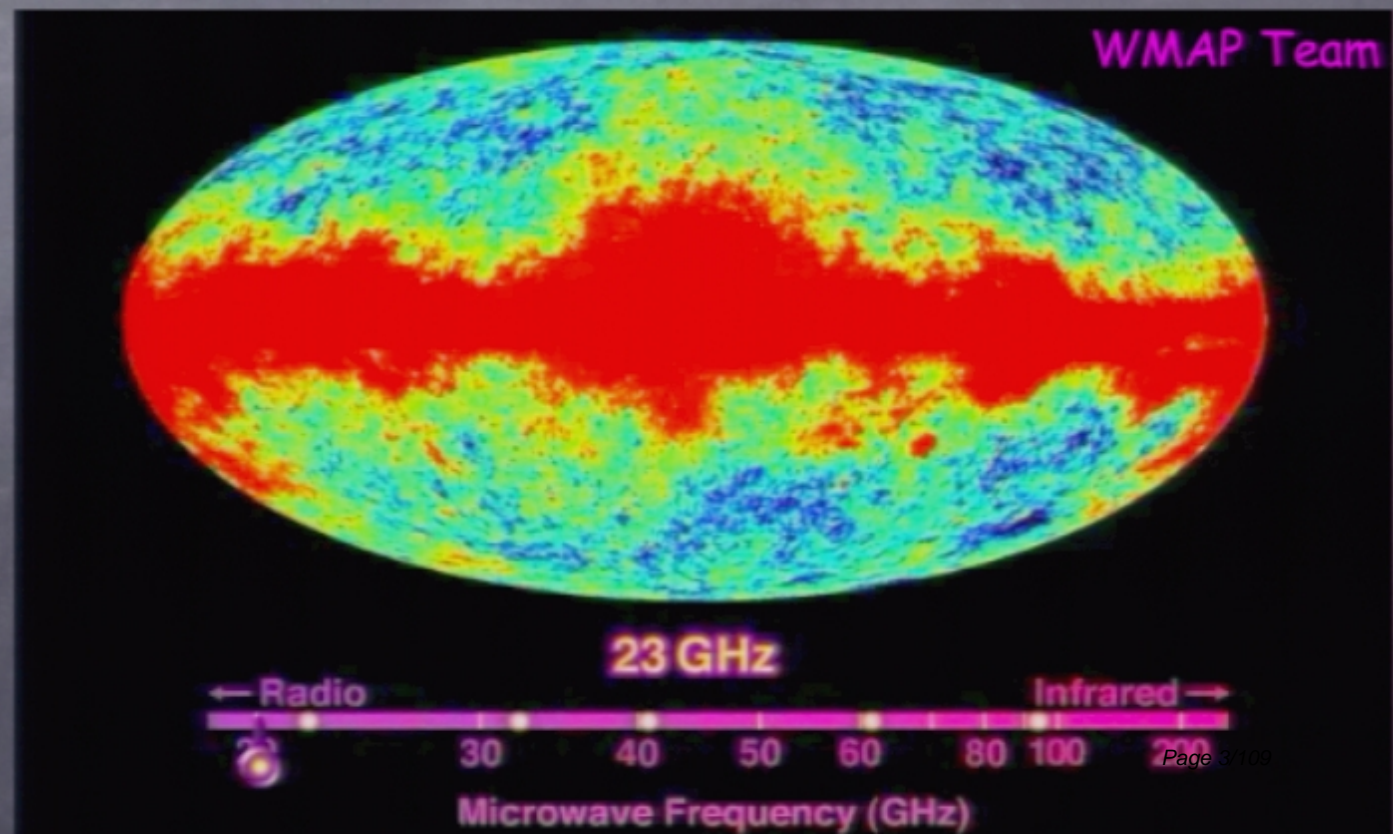
Cosmological model selection and inflation

US

University of Sussex



Microsoft-free
presentation



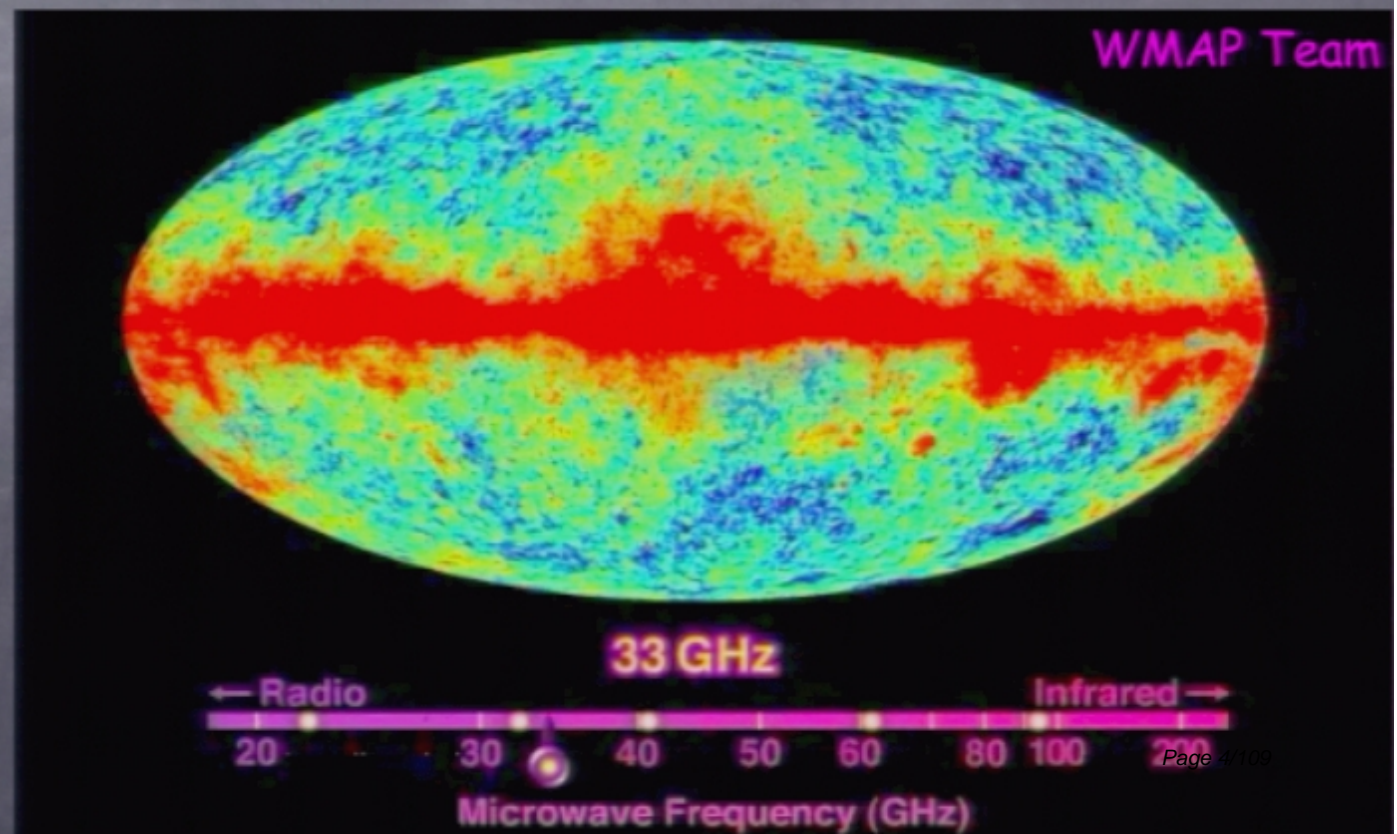
Cosmological model selection and inflation

US

University of Sussex



Microsoft-free
presentation



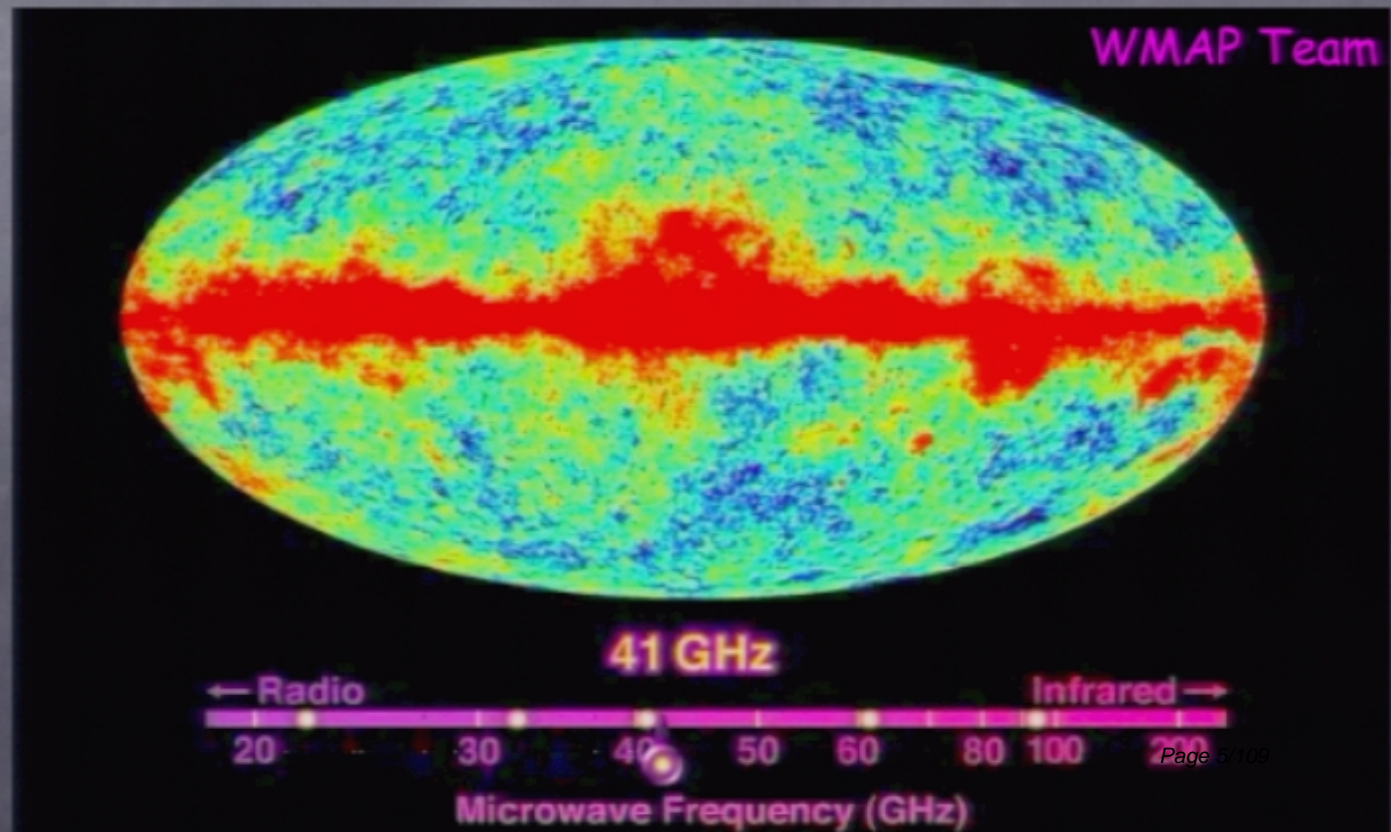
Cosmological model selection and inflation

US

University of Sussex



Microsoft-free
presentation



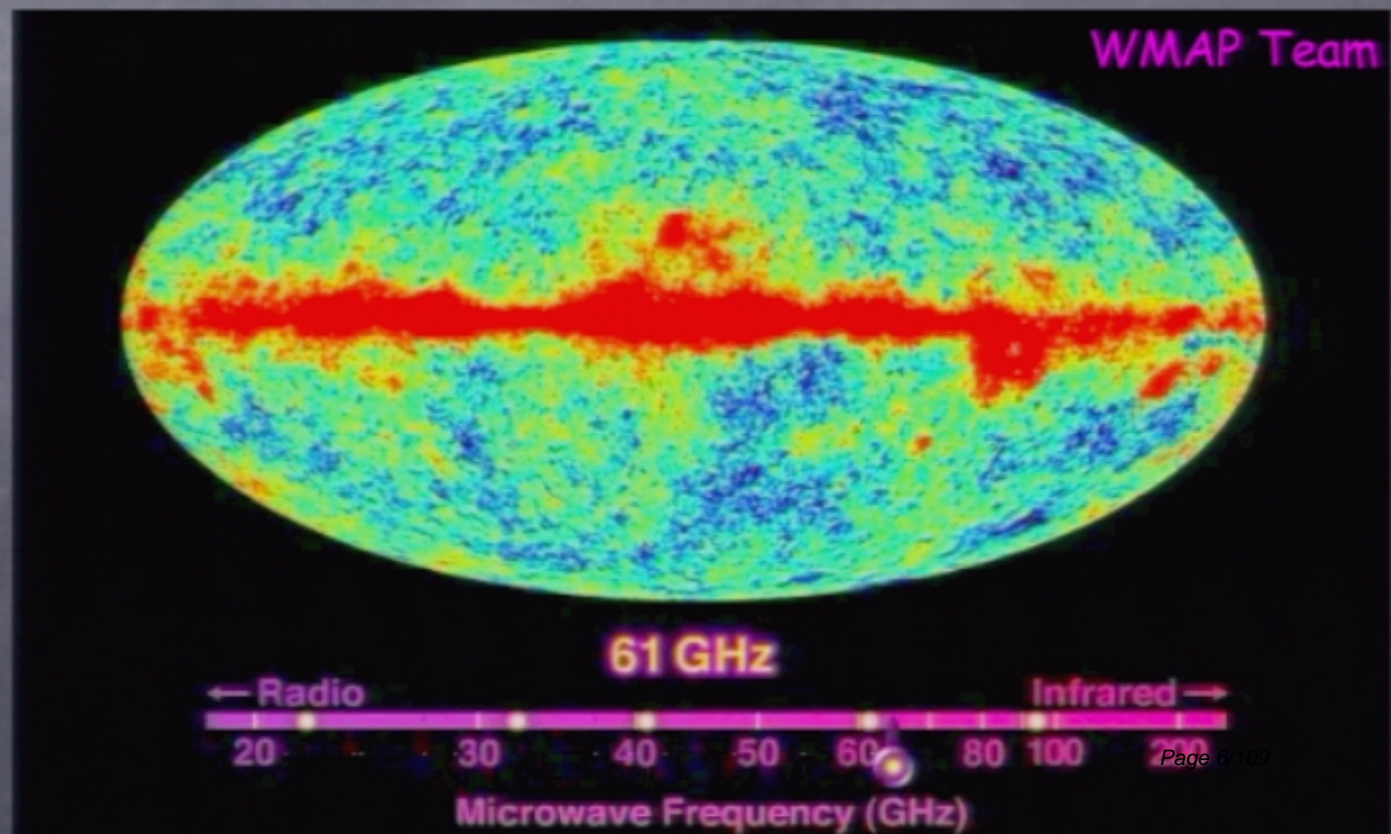
Cosmological model selection and inflation

US

University of Sussex



Microsoft-free
presentation



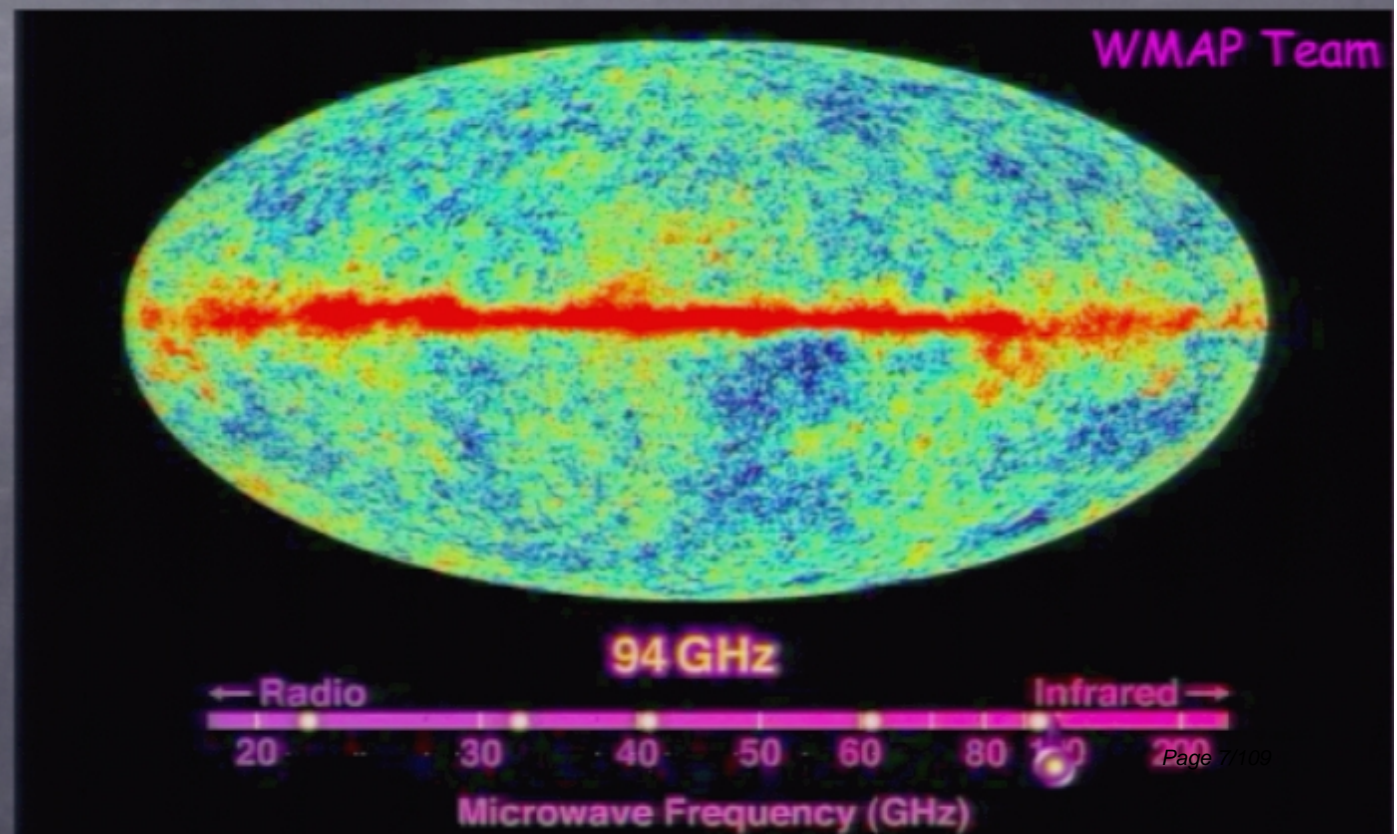
Cosmological model selection and inflation

US

University of Sussex

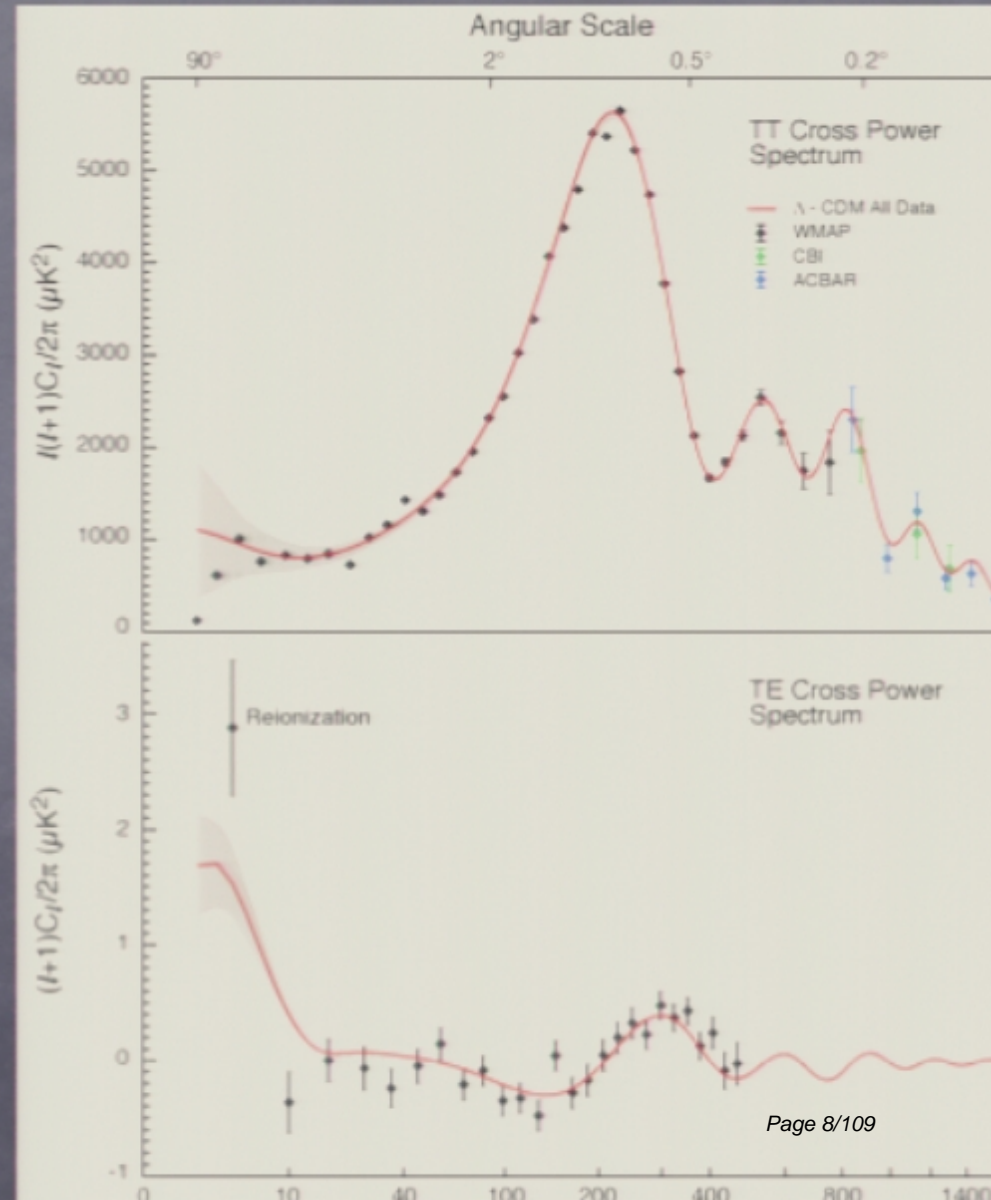


Microsoft-free
presentation



Conclusions from WMAP:

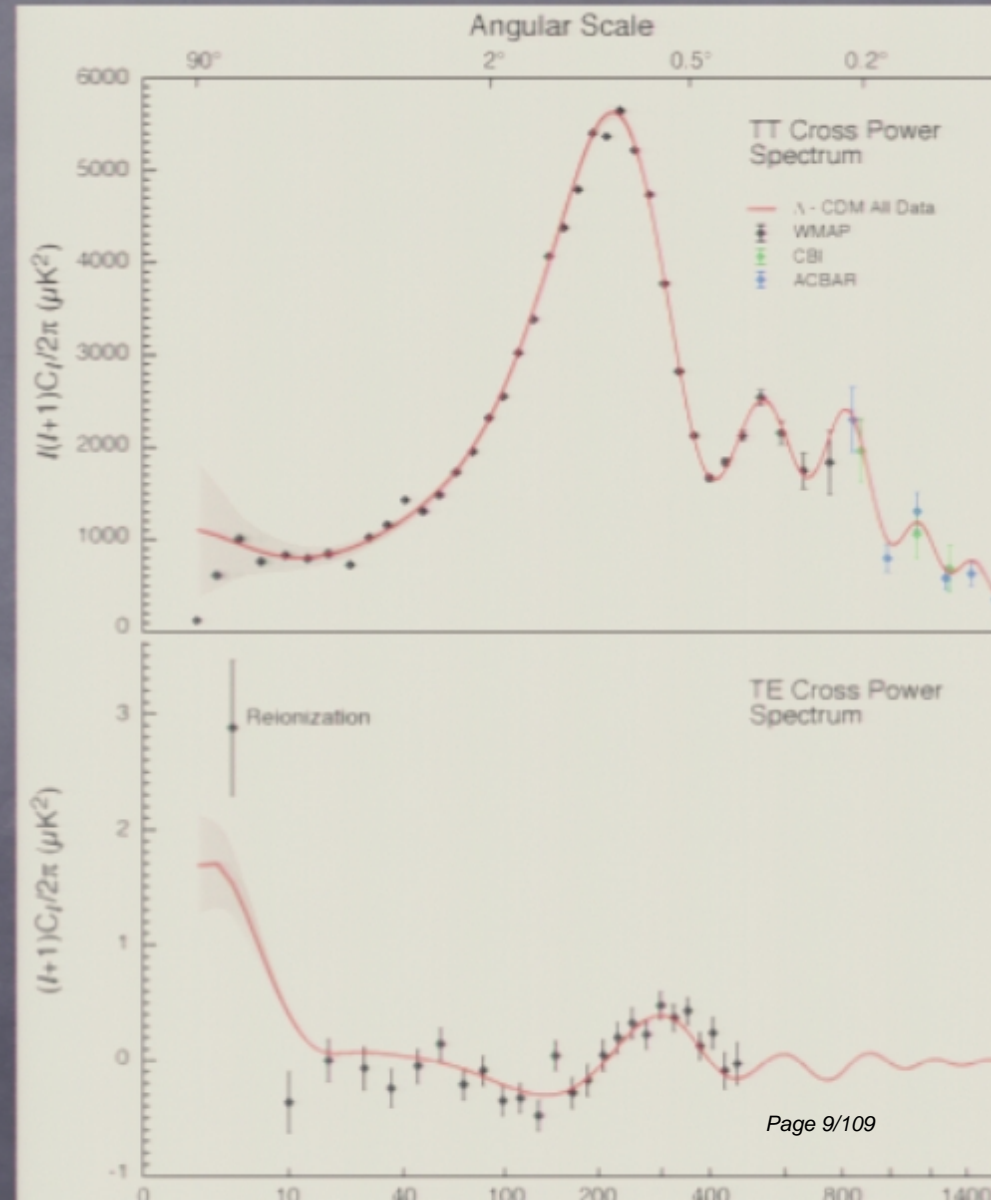
If you want to explain this data, the simplest way is ...



Conclusions from WMAP:

If you want to explain this data, the simplest way is ...

- A spatially-flat Universe
- Dark matter and dark energy
- Initial perturbations which are gaussian, adiabatic and nearly scale-invariant, e.g. as given by inflation.

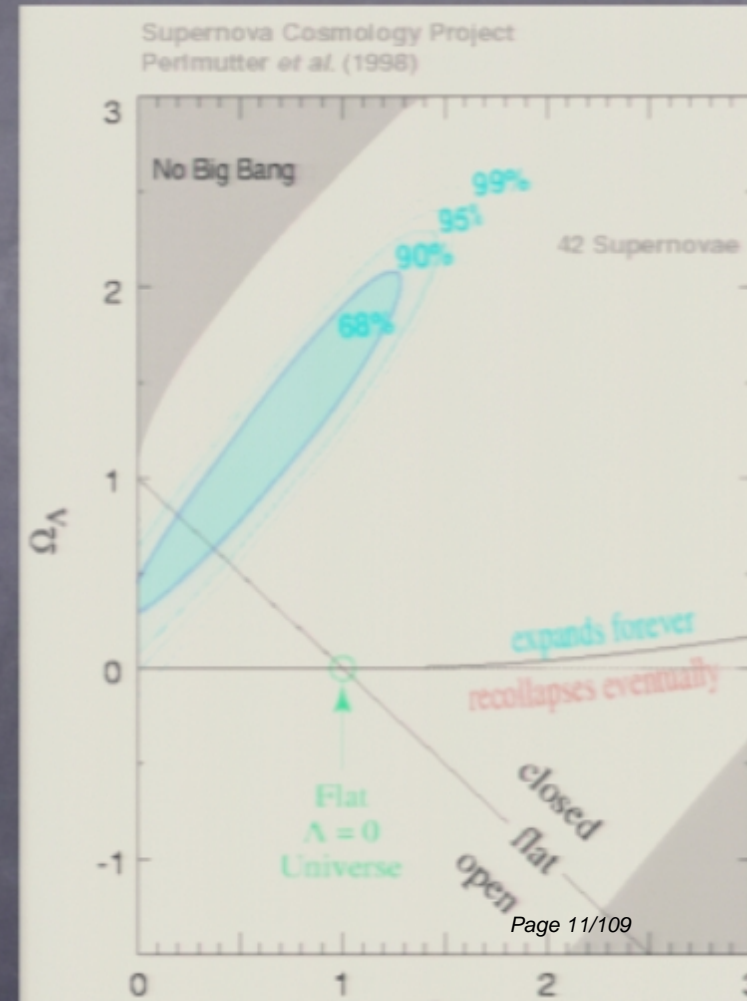


What is inflation?

Inflation is any period of the Universe's evolution during which the Universe is accelerating $\ddot{a} > 0$

What is inflation?

Inflation is any period of the Universe's evolution during which the Universe is accelerating $\ddot{a} > 0$

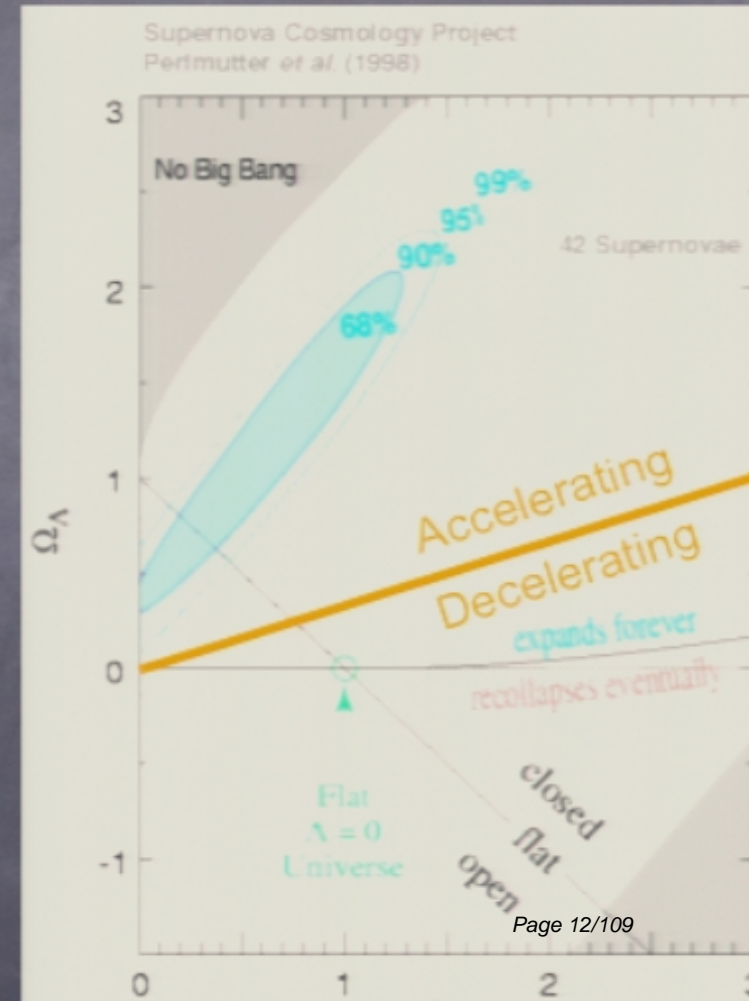


What is inflation?

Inflation is any period of the Universe's evolution during which the Universe is accelerating $\ddot{a} > 0$

This can also be written in terms of the comoving Hubble length as

$$\frac{d(H^{-1}/a)}{dt} < 0$$



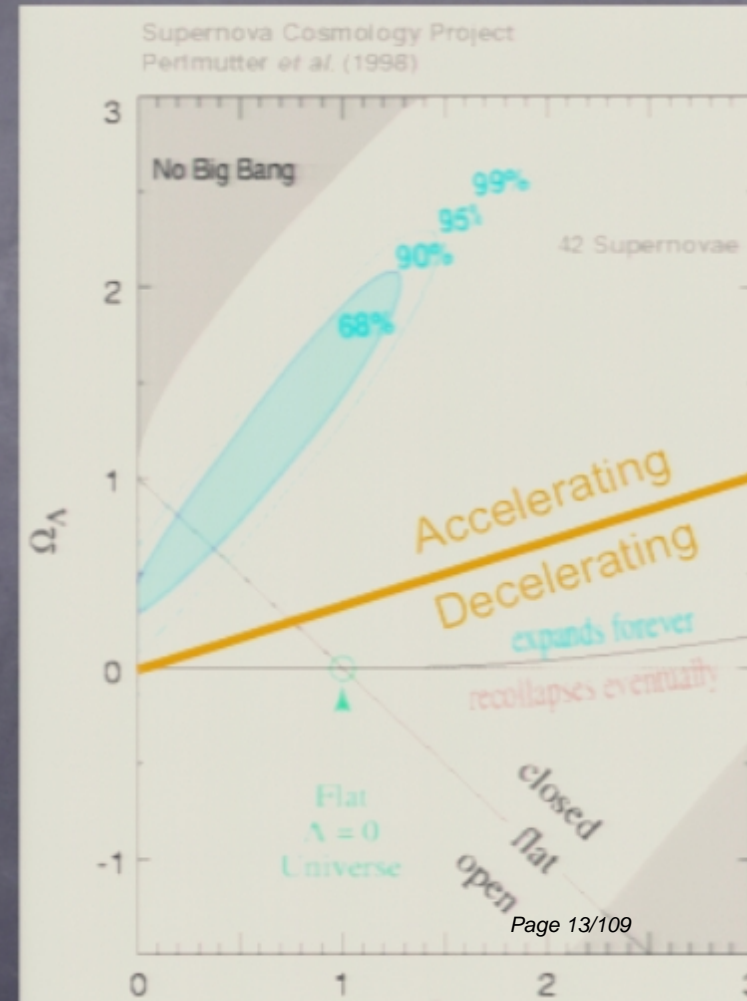
What is inflation?

Inflation is any period of the Universe's evolution during which the Universe is accelerating $\ddot{a} > 0$

This can also be written in terms of the comoving Hubble length as

$$\frac{d(H^{-1}/a)}{dt} < 0$$

Early Universe inflation is the most plausible explanation we have for the origin of structure.



Predictions of the simplest models

Predictions of the simplest models

The simplest models of inflation predict nearly scale-invariant spectra of adiabatic gaussian density perturbations and gravitational waves, in their growing mode, in a spatially-flat Universe.

Predictions of the simplest models

The simplest models of inflation predict nearly scale-invariant spectra of adiabatic gaussian density perturbations and gravitational waves, in their growing mode, in a spatially-flat Universe.

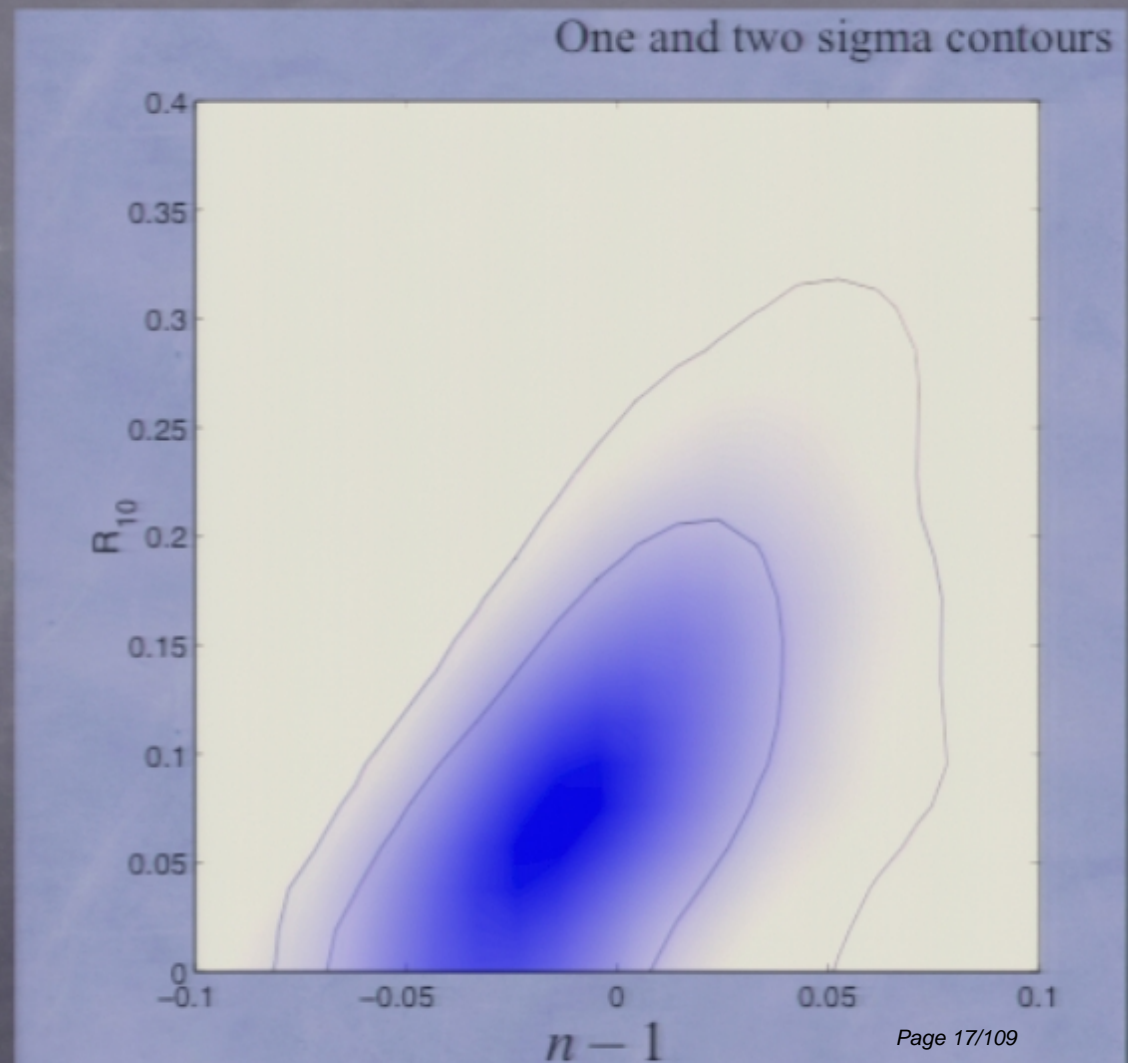
WMAP does not provide any evidence against any of these, and gives support to all but the gravitational waves. As such, it gives strong general support to the inflationary paradigm (but not uniquely to inflation).

Current constraints

Leach & Liddle, PRD, astro-ph/0306305

Comparison with observations:

- Fit to data compilation of WMAP, other CMB experiments (VSA, CBI and ACBAR), and 2dF galaxy survey.
- Use CAMB plus CosmoMC plus WMAP likelihood code plus slow-roll inflation module.

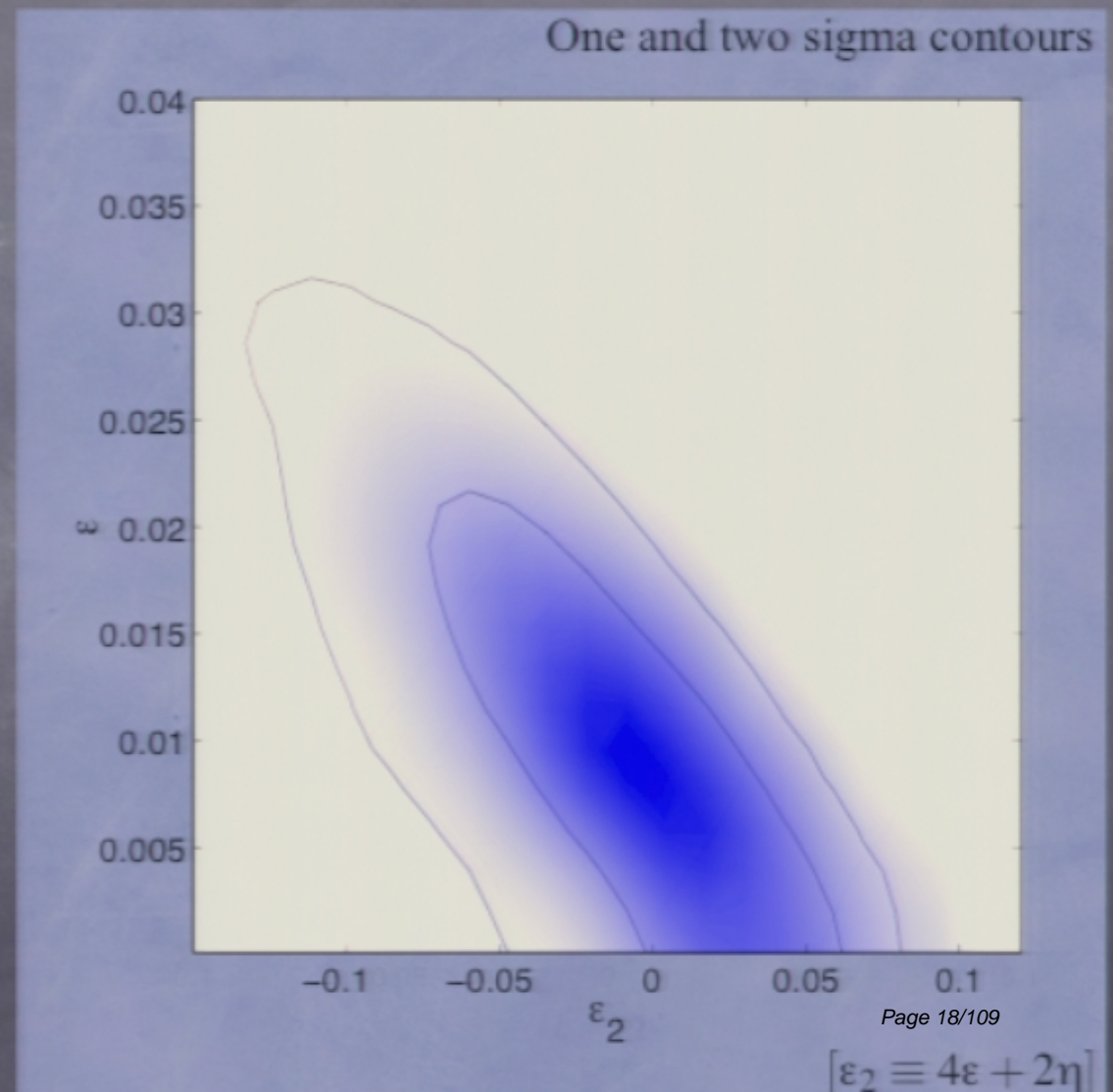



Current status of single-field inflation models

Leach & Liddle, PRD, astro-ph/0306305

Comparison with observations:

- Fit to data compilation of WMAP, other CMB experiments (VSA, CBI and ACBAR), and 2dF galaxy survey.
- Use CAMB plus CosmoMC plus WMAP likelihood code plus slow-roll inflation module.





Nearly scale-invariant spectrum $n = 0.99 \pm 0.04$
(for a power-law fit to the data)

Nearly scale-invariant spectrum $n = 0.99 \pm 0.04$
(for a power-law fit to the data)

Adiabatic perturbations Good fit to data assuming these dominate.

Nearly scale-invariant spectrum $n = 0.99 \pm 0.04$
(for a power-law fit to the data)

Adiabatic perturbations Good fit to data assuming these dominate.

Gaussian perturbations No unambiguous evidence of primordial non-gaussianity.

Gravitational waves

Nearly scale-invariant spectrum $n = 0.99 \pm 0.04$
(for a power-law fit to the data)

Adiabatic perturbations Good fit to data assuming these dominate.

Gaussian perturbations No unambiguous evidence of primordial non-gaussianity.

Gravitational waves Not detected: $r < 0.43$
(for a power-law fit to the data)

Nearly scale-invariant spectrum $n = 0.99 \pm 0.04$
(for a power-law fit to the data)

Adiabatic perturbations Good fit to data assuming these dominate.

Gaussian perturbations No unambiguous evidence of primordial non-gaussianity.

Gravitational waves Not detected: $r < 0.43$
(for a power-law fit to the data)

Growing mode Good fit to data assuming no decaying mode.
Temperature-polarization anti-correlation.

Spatial flatness $\Omega_{\text{tot}} = 1.02 \pm 0.02$

Conclusion: the simplest inflation models are doing very well!!

More complicated models

There is presently no observational need to consider more complicated models, but here's some possibilities:

More complicated models

There is presently no observational need to consider more complicated models, but here's some possibilities:

- **Multi-field models.** These allow isocurvature perturbations, which may be correlated with the usual adiabatic ones.

More complicated models

There is presently no observational need to consider more complicated models, but here's some possibilities:

- **Multi-field models.** These allow isocurvature perturbations, which may be correlated with the usual adiabatic ones.
- **The Curvaton.** In this model, negligible adiabatic perturbations are produced during inflation, being later generated from isocurvature perturbations generated during inflation. This is a natural way to introduce some non-gaussianity.

More complicated models

There is presently no observational need to consider more complicated models, but here's some possibilities:

- **Multi-field models.** These allow isocurvature perturbations, which may be correlated with the usual adiabatic ones.
- **The Curvaton.** In this model, negligible adiabatic perturbations are produced during inflation, being later generated from isocurvature perturbations generated during inflation. This is a natural way to introduce some non-gaussianity.
- **Braneworld Inflation.** In standard braneworld inflation, the form of perturbations generated from a given potential changes, though the general predictions are unharmed.

More complicated models

There is presently no observational need to consider more complicated models, but here's some possibilities:

- **Multi-field models.** These allow isocurvature perturbations, which may be correlated with the usual adiabatic ones.
- **The Curvaton.** In this model, negligible adiabatic perturbations are produced during inflation, being later generated from isocurvature perturbations generated during inflation. This is a natural way to introduce some non-gaussianity.
- **Braneworld Inflation.** In standard braneworld inflation, the form of perturbations generated from a given potential changes, though the general predictions are unharmed.

Even if effects from these more complex models are never seen, they introduce degeneracies in interpreting observations.

What is the Standard Cosmological Model?

What is the Standard Cosmological Model?

While there is broad consensus that the standard cosmological model gives an excellent description of the observed data, there isn't actually agreement on what the standard cosmological model is!

What is the Standard Cosmological Model?

While there is broad consensus that the standard cosmological model gives an excellent description of the observed data, there isn't actually agreement on what the standard cosmological model is!

The precise constraints obtained depend on

What is the Standard Cosmological Model?

While there is broad consensus that the standard cosmological model gives an excellent description of the observed data, there isn't actually agreement on what the standard cosmological model is!

The precise constraints obtained depend on

- The observational datasets used.

What is the Standard Cosmological Model?

While there is broad consensus that the standard cosmological model gives an excellent description of the observed data, there isn't actually agreement on what the standard cosmological model is!

The precise constraints obtained depend on

- The observational datasets used.
- The set of cosmological parameters used to define the cosmological model.

What is the Standard Cosmological Model?

While there is broad consensus that the standard cosmological model gives an excellent description of the observed data, there isn't actually agreement on what the standard cosmological model is!

The precise constraints obtained depend on

- The observational datasets used.
- The set of cosmological parameters used to define the cosmological model.

There have been a variety of choices made for both of these.

Table 7. Best Fit Parameters: Power Law Λ CDM

WMAP: Spergel et al

| | WMAP | WMAPext ^{16a} | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|---------------------------|---------------------------|---------------------------|---------------------------------|
| A | 0.9 ± 0.1 | 0.8 ± 0.1 | 0.8 ± 0.1 | $0.75^{+0.08}_{-0.07}$ |
| n_s | 0.99 ± 0.04 | 0.97 ± 0.03 | 0.97 ± 0.03 | 0.96 ± 0.02 |
| τ | $0.166^{+0.076}_{-0.071}$ | $0.143^{+0.071}_{-0.062}$ | $0.148^{+0.073}_{-0.071}$ | $0.117^{+0.057}_{-0.053}$ |
| h | 0.72 ± 0.05 | 0.73 ± 0.05 | 0.73 ± 0.03 | 0.72 ± 0.03 |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.13 ± 0.01 | 0.134 ± 0.006 | 0.133 ± 0.006 |
| $\Omega_b h^2$ | 0.024 ± 0.001 | 0.023 ± 0.001 | 0.023 ± 0.001 | 0.0226 ± 0.0008 |
| χ^2_{eff}/ν | 1429/1341 | 1440/1352 | 1468/1381 | ... ^b |

Table 7. Best Fit Parameters: Power Law Λ CDM

WMAP: Spergel et al

| | WMAP | WMAPext ^{16a} | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|---------------------------|---------------------------|---------------------------|---------------------------------|
| A | 0.9 ± 0.1 | 0.8 ± 0.1 | 0.8 ± 0.1 | $0.75^{+0.08}_{-0.07}$ |
| n_s | 0.99 ± 0.04 | 0.97 ± 0.03 | 0.97 ± 0.03 | 0.96 ± 0.02 |
| τ | $0.166^{+0.076}_{-0.071}$ | $0.143^{+0.071}_{-0.062}$ | $0.148^{+0.073}_{-0.071}$ | $0.117^{+0.057}_{-0.053}$ |
| h | 0.72 ± 0.05 | 0.73 ± 0.05 | 0.73 ± 0.03 | 0.72 ± 0.03 |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.13 ± 0.01 | 0.134 ± 0.006 | 0.133 ± 0.006 |
| $\Omega_b h^2$ | 0.024 ± 0.001 | 0.023 ± 0.001 | 0.023 ± 0.001 | 0.0226 ± 0.0008 |
| χ^2_{eff}/ν | 1429/1341 | 1440/1352 | 1468/1381 | ... ^b |

Table 7. Best Fit Parameters: Power Law Λ CDM

| | WMAP | WMAPext ^{16a} | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|---------------------------|---------------------------|---------------------------|---------------------------------|
| A | 0.9 ± 0.1 | 0.8 ± 0.1 | 0.8 ± 0.1 | $0.75^{+0.08}_{-0.07}$ |
| n_s | 0.99 ± 0.04 | 0.97 ± 0.03 | 0.97 ± 0.03 | 0.96 ± 0.02 |
| τ | $0.166^{+0.076}_{-0.071}$ | $0.143^{+0.071}_{-0.062}$ | $0.148^{+0.073}_{-0.071}$ | $0.117^{+0.057}_{-0.053}$ |
| h | 0.72 ± 0.05 | 0.73 ± 0.05 | 0.73 ± 0.03 | 0.72 ± 0.03 |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.13 ± 0.01 | 0.134 ± 0.006 | 0.133 ± 0.006 |
| $\Omega_b h^2$ | 0.024 ± 0.001 | 0.023 ± 0.001 | 0.023 ± 0.001 | 0.0226 ± 0.0008 |
| χ^2_{eff}/ν | 1429/1341 | 1440/1352 | 1468/1381 | ... ^b |

Table 8. Best Fit Parameters for the Running Spectral Index Λ CDM Model

| | WMAP | WMAPext | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|------------------------|--------------------|----------------------------|---------------------------------|
| A | 0.92 ± 0.12 | 0.9 ± 0.1 | 0.84 ± 0.09 | $0.83^{+0.09}_{-0.08}$ |
| n_s | $0.93^{+0.07}_{-0.07}$ | 0.91 ± 0.06 | $0.93^{+0.04}_{-0.05}$ | 0.93 ± 0.03 |
| $dn_s/d \ln k$ | -0.047 ± 0.04 | -0.055 ± 0.038 | $-0.031^{+0.023}_{-0.025}$ | $-0.031^{+0.016}_{-0.017}$ |
| τ | 0.20 ± 0.07 | 0.20 ± 0.07 | 0.17 ± 0.06 | 0.17 ± 0.06 |
| h | 0.70 ± 0.05 | 0.71 ± 0.06 | 0.71 ± 0.04 | $0.71^{+0.04}_{-0.03}$ |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.14 ± 0.01 | 0.136 ± 0.009 | $0.135^{+0.008}_{-0.009}$ |
| $\Omega_b h^2$ | 0.023 ± 0.002 | 0.022 ± 0.001 | 0.022 ± 0.001 | 0.0224 ± 0.0009 |
| χ^2_{eff}/ν | 1431/1342 | 1437/1350 | 1465/1380 | ... |

Table 7. Best Fit Parameters: Power Law Λ CDM

| | WMAP | WMAPext ^{16a} | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|---------------------------|---------------------------|---------------------------|---------------------------------|
| A | 0.9 ± 0.1 | 0.8 ± 0.1 | 0.8 ± 0.1 | $0.75^{+0.08}_{-0.07}$ |
| n_s | 0.99 ± 0.04 | 0.97 ± 0.03 | 0.97 ± 0.03 | 0.96 ± 0.02 |
| τ | $0.166^{+0.076}_{-0.071}$ | $0.143^{+0.071}_{-0.062}$ | $0.148^{+0.073}_{-0.071}$ | $0.117^{+0.057}_{-0.053}$ |
| h | 0.72 ± 0.05 | 0.73 ± 0.05 | 0.73 ± 0.03 | 0.72 ± 0.03 |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.13 ± 0.01 | 0.134 ± 0.006 | 0.133 ± 0.006 |
| $\Omega_b h^2$ | 0.024 ± 0.001 | 0.023 ± 0.001 | 0.023 ± 0.001 | 0.0226 ± 0.0008 |
| χ^2_{eff}/ν | 1429/1341 | 1440/1352 | 1468/1381 | ... ^b |

Table 8. Best Fit Parameters for the Running Spectral Index Λ CDM Model

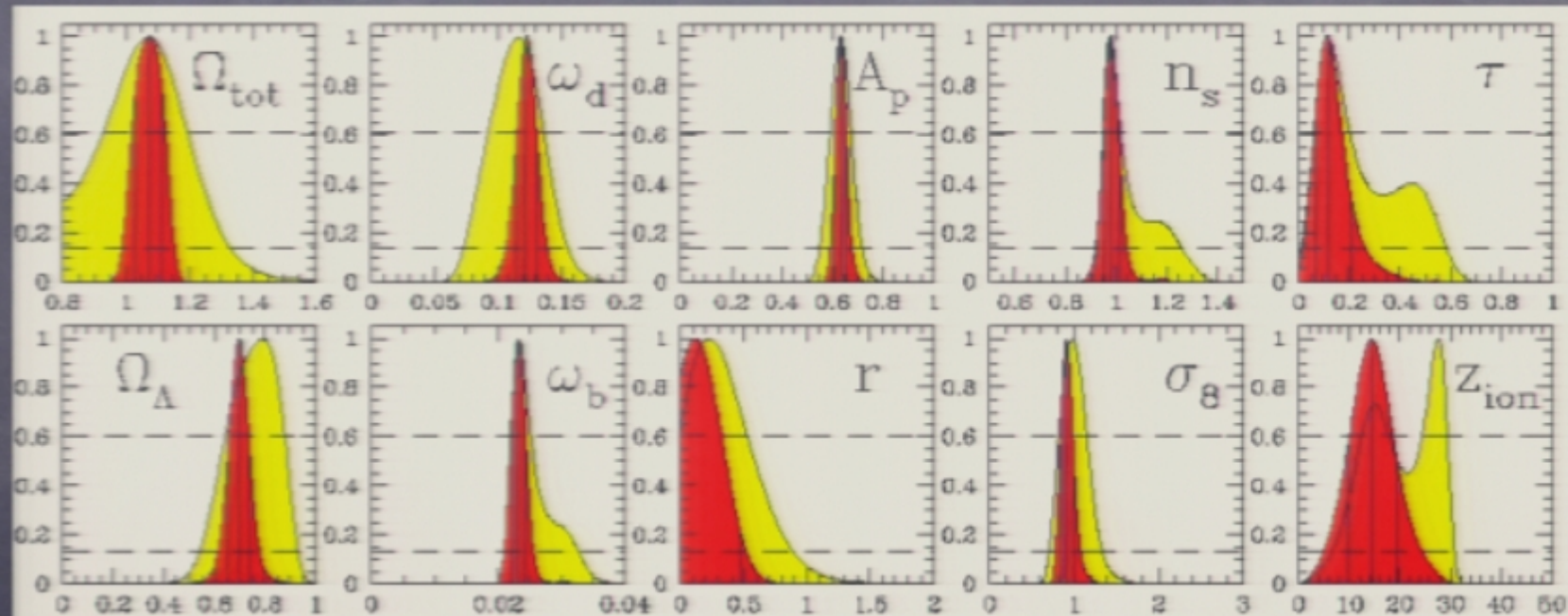
| | WMAP | WMAPext | WMAPext+2dFGRS | WMAPext+ 2dFGRS+ Lyman α |
|--------------------|------------------------|--------------------|----------------------------|---------------------------------|
| A | 0.92 ± 0.12 | 0.9 ± 0.1 | 0.84 ± 0.09 | $0.83^{+0.09}_{-0.08}$ |
| n_s | $0.93^{+0.07}_{-0.07}$ | 0.91 ± 0.06 | $0.93^{+0.04}_{-0.05}$ | 0.93 ± 0.03 |
| $dn_s/d \ln k$ | -0.047 ± 0.04 | -0.055 ± 0.038 | $-0.031^{+0.023}_{-0.025}$ | $-0.031^{+0.016}_{-0.017}$ |
| τ | 0.20 ± 0.07 | 0.20 ± 0.07 | 0.17 ± 0.06 | 0.17 ± 0.06 |
| h | 0.70 ± 0.05 | 0.71 ± 0.06 | 0.71 ± 0.04 | $0.71^{+0.04}_{-0.03}$ |
| $\Omega_m h^2$ | 0.14 ± 0.02 | 0.14 ± 0.01 | 0.136 ± 0.009 | $0.135^{+0.008}_{-0.009}$ |
| $\Omega_b h^2$ | 0.023 ± 0.002 | 0.022 ± 0.001 | 0.022 ± 0.001 | 0.0224 ± 0.0009 |
| χ^2_{eff}/ν | 1431/1342 | 1437/1350 | 1465/1380 | ... |

Parameter Estimation

In parameter estimation, the choice of parameters has already been made and we aim to constrain their values, for example by a likelihood analysis.

Parameter Estimation

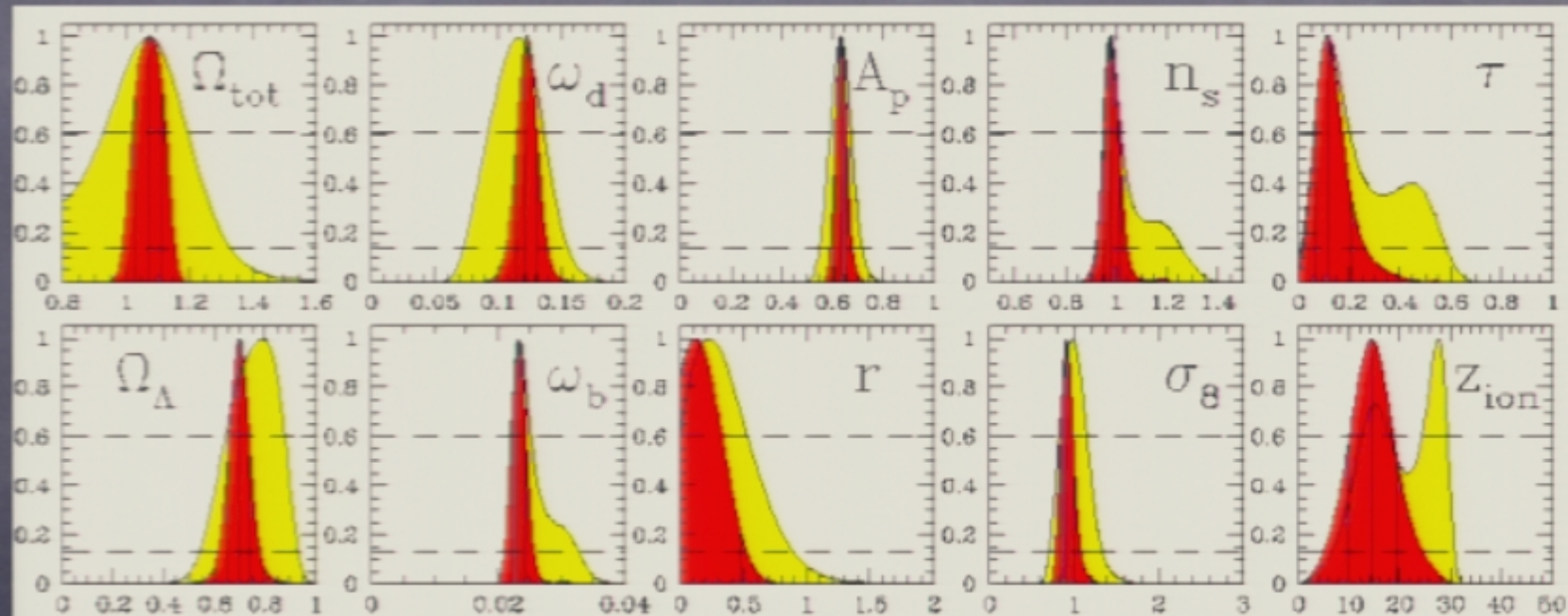
In parameter estimation, the choice of parameters has already been made and we aim to constrain their values, for example by a likelihood analysis.



Tegmark et al. (2003)

Parameter Estimation

In parameter estimation, the choice of parameters has already been made and we aim to constrain their values, for example by a likelihood analysis.



Tegmark et al. (2003)

The maximum likelihood gives the best values for the parameters, and the neighbouring behaviour gives the confidence limits.

Model Selection

In model selection, we aim to distinguish different cosmological models, meaning different choices of the parameters to be varied. In particular we need to allow for model dimensionality: that different models may have different numbers of parameters.

Model Selection

In model selection, we aim to distinguish different cosmological models, meaning different choices of the parameters to be varied. In particular we need to allow for model dimensionality: that different models may have different numbers of parameters.

A suitable baseline cosmological model to consider is the simplest one giving an adequate fit to current data. It is a spatially-flat adiabatic Λ CDM model with five fundamental parameters and two phenomenological ones.

Model Selection

In model selection, we aim to distinguish different cosmological models, meaning different choices of the parameters to be varied. In particular we need to allow for model dimensionality: that different models may have different numbers of parameters.

A suitable baseline cosmological model to consider is the simplest one giving an adequate fit to current data. It is a spatially-flat adiabatic Λ CDM model with five fundamental parameters and two phenomenological ones.

| | |
|------------|--|
| Ω_m | matter density |
| Ω_b | baryon density |
| Ω_r | radiation density |
| h | hubble parameter |
| A | adiabatic density perturbation amplitude |
| <hr/> | |
| τ | reionization optical depth |
| b | bias parameter (or parameters) |

There are many, many ways in which this base cosmological model can be extended.

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations. Of the latter set, the first six refer to adiabatic perturbations, the next three to tensor perturbations, and the remainder to isocurvature perturbations.

| | |
|------------------------------|---|
| Ω_k | spatial curvature |
| $N_\nu - 3.04$ | effective number of neutrino species (CMBFAST definition) |
| m_{ν_i} | neutrino mass for species 'i' [or more complex neutrino properties] |
| m_{dm} | (warm) dark matter mass |
| $w + 1$ | dark energy equation of state |
| dw/dz | redshift dependence of w [or more complex parametrization of dark energy evolution] |
| $c_s^2 - 1$ | effects of dark energy sound speed |
| $1/r_{\text{top}}$ | topological identification scale [or more complex parametrization of non-trivial topology] |
| $d\alpha/dz$ | redshift dependence of the fine structure constant |
| dG/dz | redshift dependence of the gravitational constant |
| <hr/> | |
| $n - 1$ | scalar spectral index |
| $dn/d \ln k$ | running of the scalar spectral index |
| k_{cut} | large-scale cut-off in the spectrum |
| A_{feature} | amplitude of spectral feature (peak, dip or step) ... |
| k_{feature} | ... and its scale [or adiabatic power spectrum amplitude parametrized in N bins] |
| f_{NL} | quadratic contribution to primordial non-gaussianity [or more complex parametrization of non-gaussianity] |
| r | tensor-to-scalar ratio |
| $r + 8n_T$ | violation of the inflationary consistency equation |
| $dn_T/d \ln k$ | running of the tensor spectral index |
| \mathcal{P}_S | CDM isocurvature perturbation ... |
| n_S | ... and its spectral index ... |
| $\mathcal{P}_{S\mathcal{R}}$ | ... and its correlation with adiabatic perturbations ... |
| $n_{S\mathcal{R}} - n_S$ | ... and the spectral index of that correlation [or more complicated multi-component isocurvature perturbation] |
| $G\mu$ | cosmic string component of perturbations |

How do we compare different cosmological models
(i.e. different choices of fundamental parameters)?
Can we say which model is best?

How do we compare different cosmological models
(i.e. different choices of fundamental parameters)?
Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

How do we compare different cosmological models (i.e. different choices of fundamental parameters)?
Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

Problem 2: as we add extra parameters, the uncertainties on existing parameters increase, and eventually we learn nothing useful about anything.

How do we compare different cosmological models (i.e. different choices of fundamental parameters)?
Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

Problem 2: as we add extra parameters, the uncertainties on existing parameters increase, and eventually we learn nothing useful about anything.

We need a way of penalizing use of extra parameters - an implementation of Ockham's razor.

Model Selection Statistics

Liddle, MNRAS, astro-ph/0401198

How do we compare different cosmological models
(i.e. different choices of fundamental parameters)?
Can we say which model is best?

Model Selection Statistics

Liddle, MNRAS, astro-ph/0401198

Model Selection Statistics

Liddle, MNRAS, astro-ph/0401198

- Akaike information criterion (Akaike 1974)
- Bayesian information criterion (Schwarz 1978)
- Bayesian evidence (Jeffreys 1961 etc)

Model Selection Statistics

Liddle, MNRAS, astro-ph/0401198

- Akaike information criterion (Akaike 1974)

$$\text{AIC} = -2\ln \mathcal{L}_{\text{max}} + 2k$$

k = number of parameters

- Bayesian information criterion (Schwarz 1978)

$$\text{BIC} = -2\ln \mathcal{L}_{\text{max}} + k\ln N$$

N = number of datapoints

- Bayesian evidence (Jeffreys 1961 etc)

Model Selection Statistics

Liddle, MNRAS, astro-ph/0401198

- Akaike information criterion (Akaike 1974)

$$\text{AIC} = -2 \ln \mathcal{L}_{\text{max}} + 2k$$

k = number of parameters

- Bayesian information criterion (Schwarz 1978)

$$\text{BIC} = -2 \ln \mathcal{L}_{\text{max}} + k \ln N$$

N = number of datapoints

- Bayesian evidence (Jeffreys 1961 etc)

$$E = \int d\theta \mathcal{L}(\theta) \text{pr}(\theta)$$

θ = parameter vector, pr = prior

The preferred model is the one which minimizes the information criterion, or maximizes the evidence.

The **Bayesian evidence** is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The **Bayesian evidence** is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The **Bayesian Information Criterion** was derived using Bayesian statistics. It gives a crude approximation to the **Bayesian evidence**. While it can give guidance, the assumptions of its validity are questionable in cosmological applications (eg parameter degeneracies).

The **Bayesian evidence** is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The **Bayesian Information Criterion** was derived using Bayesian statistics. It gives a crude approximation to the **Bayesian evidence**. While it can give guidance, the assumptions of its validity are questionable in cosmological applications (eg parameter degeneracies).

The **Akaike Information Criterion** was derived using information theory techniques. It gives an approximate minimization of the so-called **Kullback-Leibler information entropy**, which is a measure of the difference between two probability distributions

Model selection techniques are essential when considering whether or not new data requires the addition of new parameters to describe it.

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

Assuming that the density is the only parameter, with a uniform prior from 0.1 to 2, and likelihood

$$\mathcal{L} = \mathcal{L}_0 \exp\left(-\frac{(\Omega - 1.02)^2}{2 \times 0.02^2}\right)$$

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

Assuming that the density is the only parameter, with a uniform prior from 0.1 to 2, and likelihood

$$\mathcal{L} = \mathcal{L}_0 \exp\left(-\frac{(\Omega - 1.02)^2}{2 \times 0.02^2}\right)$$

■ **Flat:** Evidence = $\mathcal{L}(\Omega = 1) = 0.6\mathcal{L}_0$

■ **Curved:** Evidence = $\frac{1}{1.9} \int \mathcal{L}(\Omega) d\Omega \simeq 0.03\mathcal{L}_0$

According to the evidence, the flat model is a better description of the data, with odds of about 20:1 **against** the curved model.

Note that this assumes flat and curved were thought equally likely before the data came along.

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

Assuming that the density is the only parameter, with a uniform prior from 0.1 to 2, and likelihood

$$\mathcal{L} = \mathcal{L}_0 \exp\left(-\frac{(\Omega - 1.02)^2}{2 \times 0.02^2}\right)$$

■ **Flat:** Evidence = $\mathcal{L}(\Omega = 1) = 0.6\mathcal{L}_0$

■ **Curved:** Evidence = $\frac{1}{1.9} \int \mathcal{L}(\Omega) d\Omega \simeq 0.03\mathcal{L}_0$

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

Assuming that the density is the only parameter, with a uniform prior from 0.1 to 2, and likelihood

$$\mathcal{L} = \mathcal{L}_0 \exp\left(-\frac{(\Omega - 1.02)^2}{2 \times 0.02^2}\right)$$

■ Flat: Evidence = $\mathcal{L}(\Omega = 1) = 0.6\mathcal{L}_0$

■ Curved: Evidence = $\frac{1}{1.9} \int \mathcal{L}(\Omega) d\Omega \simeq 0.03\mathcal{L}_0$

Notes:

1) Even if parameter estimation had given $\Omega_{\text{tot}} = 1.05 \pm 0.02$ the flat case would still have been preferred.

2) Someone adamantly insisting before WMAP that the total density was 1.02, to the exclusion of all other values, could claim WMAP supported them better than flat.

A simple example: spatial curvature

WMAP says $\Omega_{\text{tot}} = 1.02 \pm 0.02$

This has been widely interpreted as supporting the idea of a flat Universe, but actually favouring a slightly closed Universe.

Assuming that the density is the only parameter, with a uniform prior from 0.1 to 2, and likelihood

$$\mathcal{L} = \mathcal{L}_0 \exp\left(-\frac{(\Omega - 1.02)^2}{2 \times 0.02^2}\right)$$

■ Flat: Evidence = $\mathcal{L}(\Omega = 1) = 0.6\mathcal{L}_0$

■ Curved: Evidence = $\frac{1}{1.9} \int \mathcal{L}(\Omega) d\Omega \simeq 0.03\mathcal{L}_0$

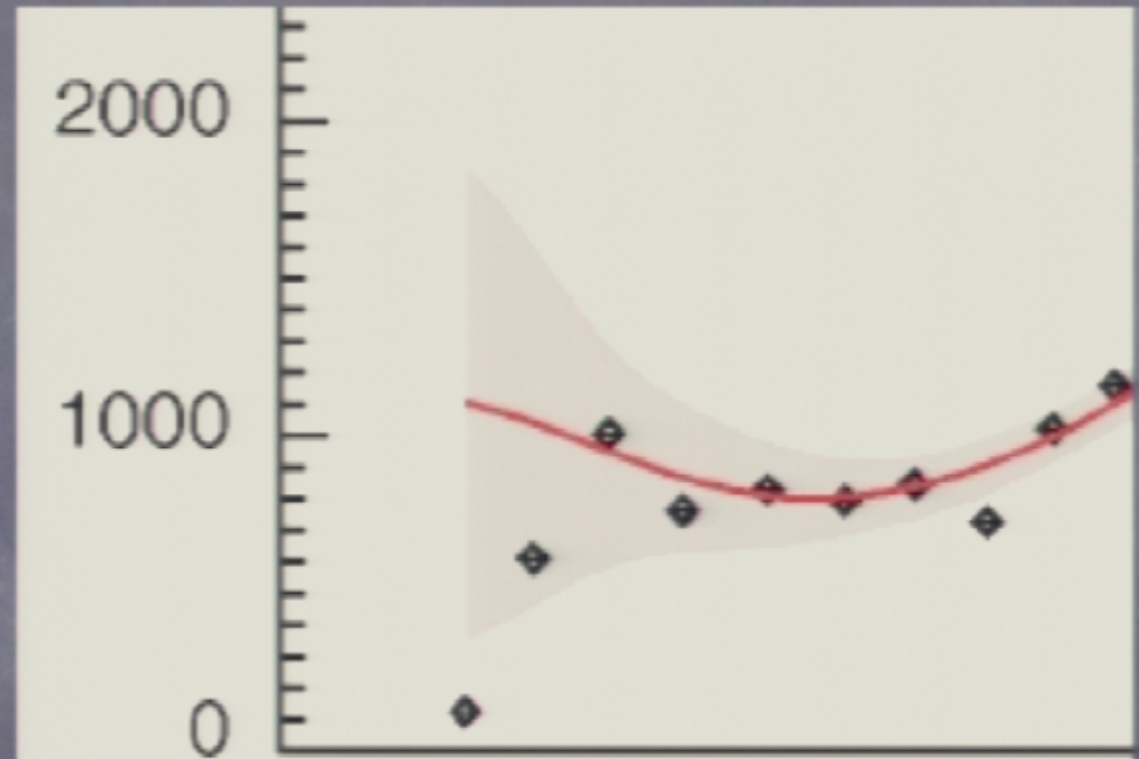
Notes:

1) Even if parameter estimation had given $\Omega_{\text{tot}} = 1.05 \pm 0.02$ the flat case would still have been preferred.

2) Someone adamantly insisting before WMAP that the total density was 1.02, to the exclusion of all other values, could claim WMAP supported them better than flat.

New physics from low quadrupole??

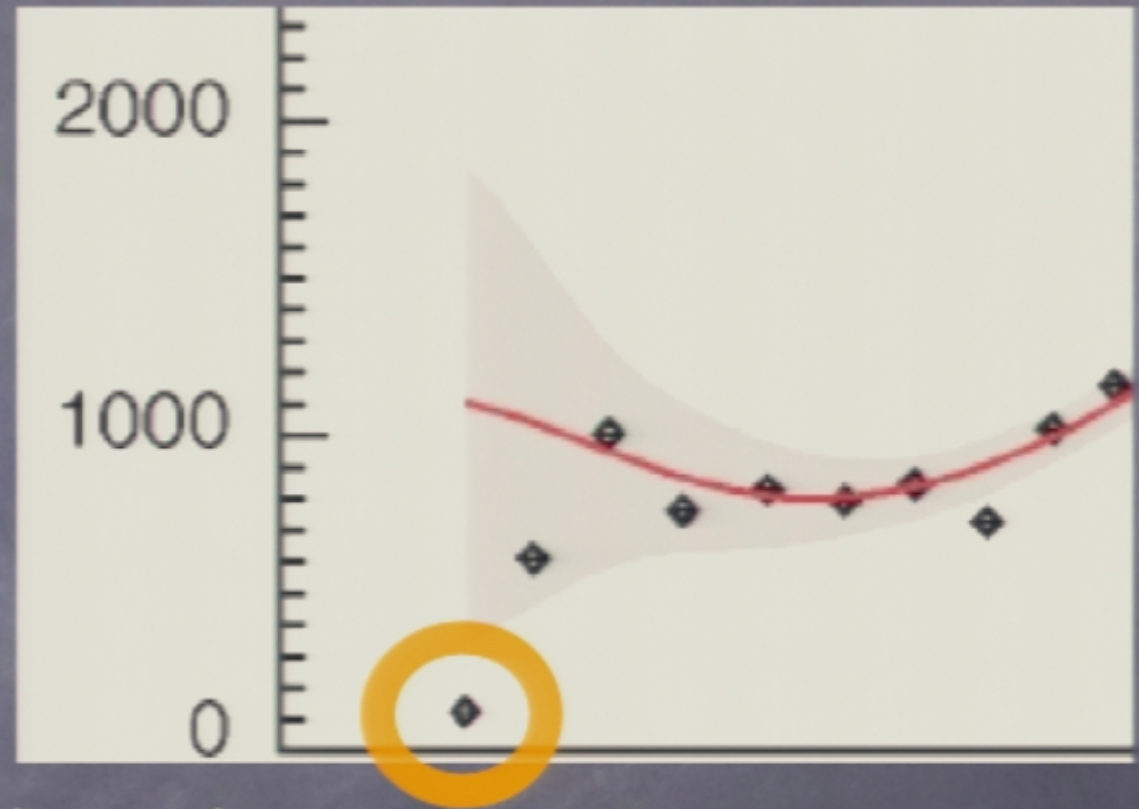
[Argument roughly following Niarchou et al]



New physics from low quadrupole??

[Argument roughly following Niarchou et al]

If you want to explain this with new physics, you have to introduce new parameters, for which you will be penalized. As the discrepancy is only at the 95% level, the gain in fit will never compensate for this penalty.



How do we compare different cosmological models
(i.e. different choices of fundamental parameters)?
Can we say which model is best?

Something like 95% of all 95% confidence "detections" turn out to be wrong. Why?

Something like 95% of all 95% confidence "detections" turn out to be wrong. Why?

- **Statistical fluke:** By definition important only if people do their error analysis wrongly

Something like 95% of all 95% confidence "detections" turn out to be wrong. Why?

- **Statistical fluke:** By definition important only if people do their error analysis wrongly.
- **Publication bias:** Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).

Something like 95% of all 95% confidence "detections" turn out to be wrong. Why?

- **Statistical fluke:** By definition important only if people do their error analysis wrongly.
- **Publication bias:** Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).
- **Inappropriate "a posteriori" reasoning:** choosing "interesting" features from the data and assessing their significance via Monte Carlo analyses.

Something like 95% of all 95% confidence "detections" turn out to be wrong. Why?

- **Statistical fluke:** By definition important only if people do their error analysis wrongly.
- **Publication bias:** Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).
- **Inappropriate "a posteriori" reasoning:** choosing "interesting" features from the data and assessing their significance via Monte Carlo analyses.
- **Neglect of model dimensionality:** using parameter estimation rather than model selection.

Model Selection and Isocurvature Modes

Beltran, Garcia-Bellido, Lesgourgues, Liddle, Slosar, PRD, astro-ph/0501477

Even if the real perturbations are adiabatic, some level of isocurvature perturbations will always be allowed.

While parameter estimation techniques can only place upper limits on the isocurvature modes, model selection can give positive support to simpler models.

Model Selection and Isocurvature Modes

Beltran, Garcia-Bellido, Lesgourgues, Liddle, Slosar, PRD, astro-ph/0501477

Even if the real perturbations are adiabatic, some level of isocurvature perturbations will always be allowed.

While parameter estimation techniques can only place upper limits on the isocurvature modes, model selection can give positive support to simpler models.

We consider the three observationally-distinct classes of isocurvature mode, **CDI**, **NID** and **NIV**. Only one type of mode is permitted per model, but with arbitrary spectral index and correlation to adiabatic: **4 extra parameters**. We compare with two adiabatic models, one with $n=1$ and one with n varying.

Model Selection and Isocurvature Modes

The Bayesian Evidence was computed using a technique called **thermodynamic integration**. This is an MCMC method where the chains are heated in order to fully explore the prior space (parameter estimation chains sample the posterior which is usually localized to a small fraction of the prior).

We tested several variants on this scheme. Accurate determination of the evidence required approximately 10^7 likelihood evaluations per model, making it a supercomputer class problem.

Model Selection and Isocurvature Modes

The Bayesian Evidence was computed using a technique called **thermodynamic integration**. This is an MCMC method where the chains are heated in order to fully explore the prior space (parameter estimation chains sample the posterior which is usually localized to a small fraction of the prior).

We tested several variants on this scheme. Accurate determination of the evidence required approximately 10^7 likelihood evaluations per model, making it a supercomputer class problem.

Jeffreys Scale:

| | |
|--------------------------|------------------------------------|
| $\Delta \ln E < 1$ | Not worth more than a bare mention |
| $1 < \Delta \ln E < 2.5$ | Substantial evidence |
| $2.5 < \Delta \ln E < 5$ | Strong to very strong evidence |
| $5 < \Delta \ln E$ | Decisive evidence |

Model Selection and Isocurvature Modes

Note that the results depend on the priors chosen. Our prior range covers the complete range from all adiabatic to all isocurvature using the relative fraction. We use two different parametrizations to test robustness.

Model Selection and Isocurvature Modes

Note that the results depend on the priors chosen. Our prior range covers the complete range from all adiabatic to all isocurvature using the relative fraction. We use two different parametrizations to test robustness.

| Model | ln(Evidence)* | |
|--------|-------------------|-------------------|
| | Parametrization 1 | Parametrization 2 |
| AD-HZ | 0.0 ± 0.1 | |
| AD-n | 0.0 ± 0.1 | |
| AD-iso | -1.0 ± 0.2 | -1.0 ± 0.2 |
| AD-iso | -1.0 ± 0.2 | -2.0 ± 0.2 |
| AD-iso | -1.0 ± 0.3 | -2.3 ± 0.2 |

Model Selection and Isocurvature Modes

Note that the results depend on the priors chosen. Our prior range covers the complete range from all adiabatic to all isocurvature using the relative fraction. We use two different parametrizations to test robustness.

| Model | ln(Evidence)* | |
|-------|-------------------|-------------------|
| | Parametrization 1 | Parametrization 2 |
| AD-HZ | 0.0 ± 0.1 | |
| AD-n | 0.0 ± 0.1 | |
| CDI | -1.0 ± 0.2 | -1.0 ± 0.2 |
| NID | -1.0 ± 0.2 | -2.0 ± 0.2 |
| NIV | -1.0 ± 0.3 | -2.3 ± 0.2 |

Nested Sampling

Mukherjee, Parkinson and Liddle, [astro-ph/0508461](#)

The main lesson from that work is that a more efficient algorithm is needed to make computations feasible. We have recently implemented Skilling's Nested Sampling algorithm for cosmology.

Nested Sampling

Mukherjee, Parkinson and Liddle, astro-ph/0508461

The main lesson from that work is that a more efficient algorithm is needed to make computations feasible. We have recently implemented Skilling's Nested Sampling algorithm for cosmology.

Skilling (2004) rewrote the evidence as

$$E = \int \mathcal{L}(\theta) \text{pr}(\theta) d\theta = \int_0^1 \mathcal{L}(X) dX$$

where X is the fractional prior mass.

Nested Sampling

Mukherjee, Parkinson and Liddle, astro-ph/0508461

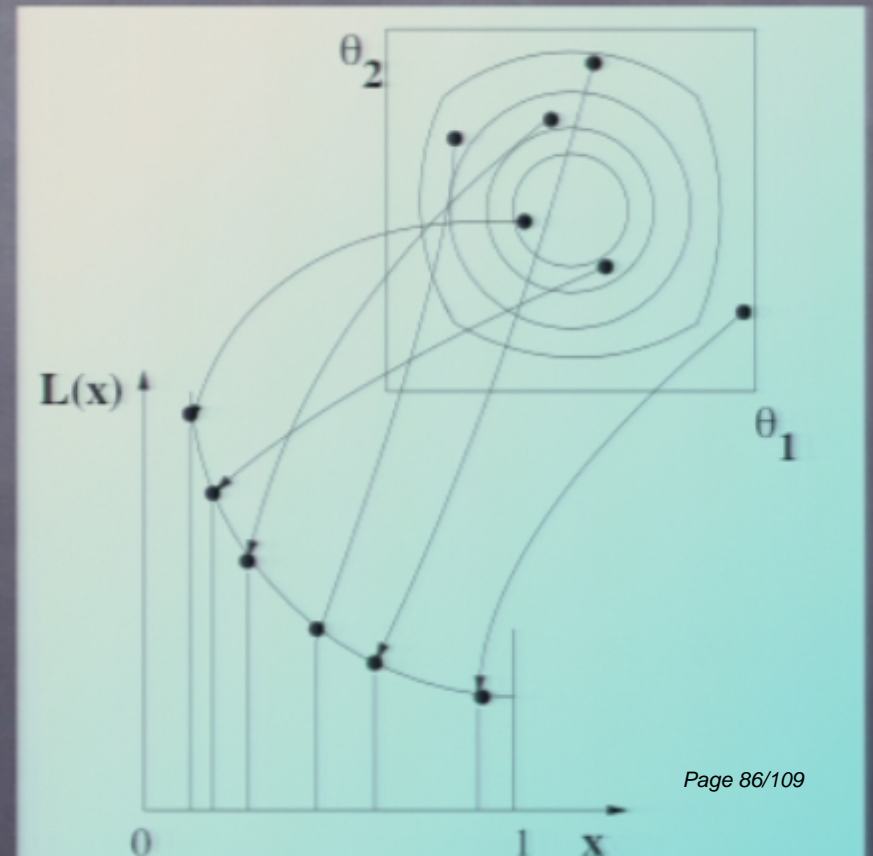
The main lesson from that work is that a more efficient algorithm is needed to make computations feasible. We have recently implemented Skilling's Nested Sampling algorithm for cosmology.

Skilling (2004) rewrote the evidence as

$$E = \int \mathcal{L}(\theta) \text{pr}(\theta) d\theta = \int_0^1 \mathcal{L}(X) dX$$

where X is the fractional prior mass.

This can then be evaluated using Monte Carlo samples to trace the variation of likelihood with prior mass, peeling away thin nested isosurfaces of equal likelihood.

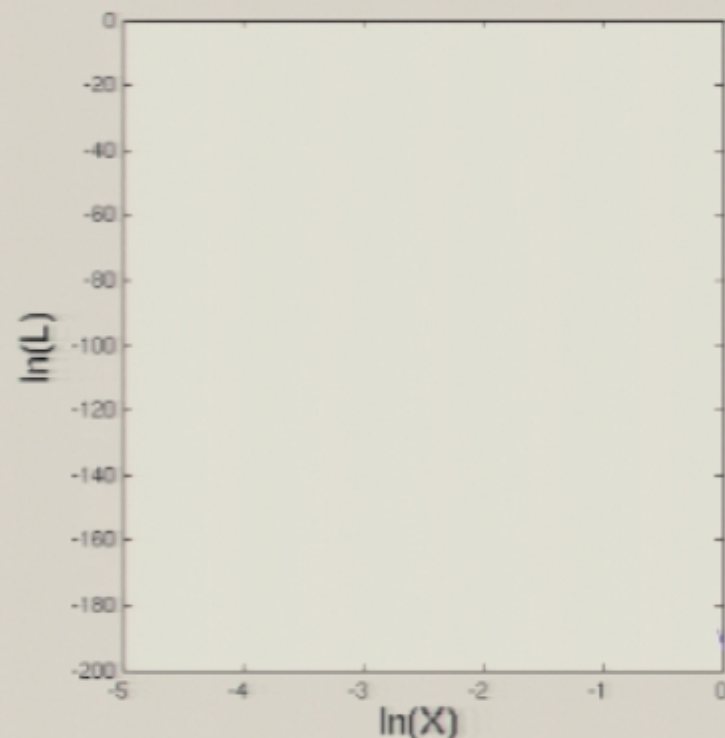
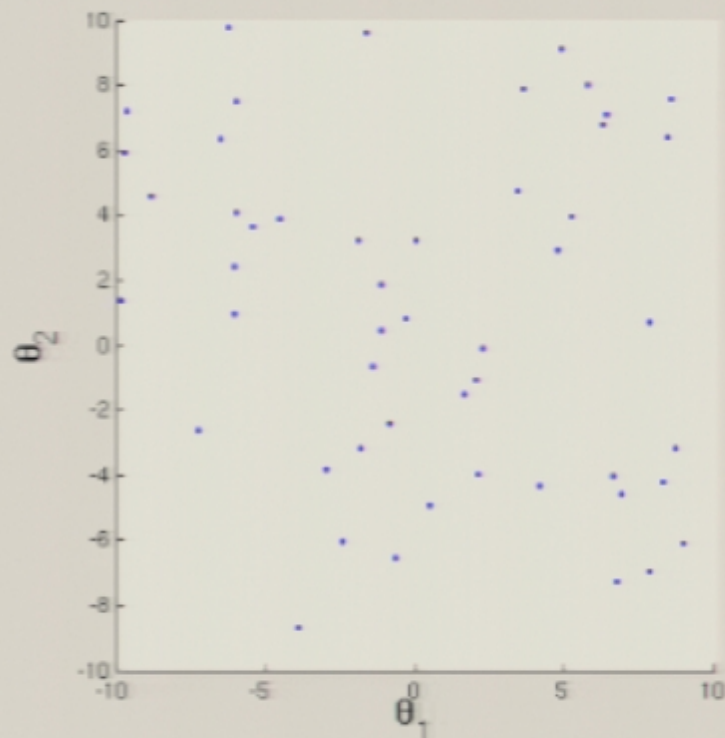


Nested Sampling

The method 'walks' a set of points (eg 300) into the high-likelihood region using replacement. The main difficulty in implementing the algorithm successfully is in efficiently generating replacement points which are uniformly sampled from the remaining prior volume.

Nested Sampling

The method 'walks' a set of points (eg 300) into the high-likelihood region using replacement. The main difficulty in implementing the algorithm successfully is in efficiently generating replacement points which are uniformly sampled from the remaining prior volume.



A model selection example

We used the Bayesian evidence to compare various cosmological models with the simplest one.

A model selection example

We used the Bayesian evidence to compare various cosmological models with the simplest one.

| Model | Λ CDM+HZ | Λ CDM+ n_s | Λ CDM+ n_s (wide prior) | HZ+ w | $w + n_s$ |
|---------------------------------|------------------|----------------------|--------------------------------------|---------------------|---------------------|
| n_s | 1 | 0.8 – 1.2 | 0.6 – 1.4 | 1 | 0.8 – 1.2 |
| w | -1 | -1 | -1 | $-\frac{1}{3} - -1$ | $-\frac{1}{3} - -1$ |
| e.f | 1.5 | 1.7 | 1.7 | 1.7 | 1.8 |
| $N_{\text{like}} (\times 10^4)$ | 8.4 | 17.4 | 16.7 | 10.6 | 18.0 |
| $\ln E$ | 0.00 ± 0.08 | -0.58 ± 0.09 | -1.16 ± 0.08 | -0.45 ± 0.08 | -1.52 ± 0.08 |

A model selection example

We used the Bayesian evidence to compare various cosmological models with the simplest one.

| Model | Λ CDM+HZ | Λ CDM+ n_s | Λ CDM+ n_s (wide prior) | HZ+ w | $w + n_s$ |
|---------------------------------|------------------|----------------------|--------------------------------------|---------------------|---------------------|
| n_s | 1 | 0.8 – 1.2 | 0.6 – 1.4 | 1 | 0.8 – 1.2 |
| w | -1 | -1 | -1 | $-\frac{1}{3} - -1$ | $-\frac{1}{3} - -1$ |
| e.f | 1.5 | 1.7 | 1.7 | 1.7 | 1.8 |
| $N_{\text{like}} (\times 10^4)$ | 8.4 | 17.4 | 16.7 | 10.6 | 18.0 |
| $\ln E$ | 0.00 ± 0.08 | -0.58 ± 0.09 | -1.16 ± 0.08 | -0.45 ± 0.08 | -1.52 ± 0.08 |

At the moment the more complex models are not excluded, but they are mildly disfavoured against the simplest model.

Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

- **Fisher matrix approach:**

- simulate data for a fiducial model (eg Λ CDM);
 - estimate expected parameter uncertainties about that model;
 - interpret as excluding models outside the contours

Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

- **Fisher matrix approach:**

simulate data for a fiducial model (eg Λ CDM);
estimate expected parameter uncertainties about that model;
interpret as excluding models outside the contours

- **Bayes factor approach:**

simulate data at each point in parameter plane;
compute Bayes factor (ie evidence ratio) of full model versus
eg Λ CDM at each point.

Fisher information matrix drawbacks

Fisher information matrix drawbacks

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (eg dark energy evolution).

Fisher information matrix drawbacks

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (eg dark energy evolution).
- Usually interpreted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however **not** that simulated.

Fisher information matrix drawbacks

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (eg dark energy evolution).
- Usually interpreted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however **not** that simulated.
- The criterion for ruling out LambdaCDM is exactly the same as that used to rule out any other value in the plane, eg $w = -0.99$. Special status of LambdaCDM is not recognised.

Fisher information matrix drawbacks

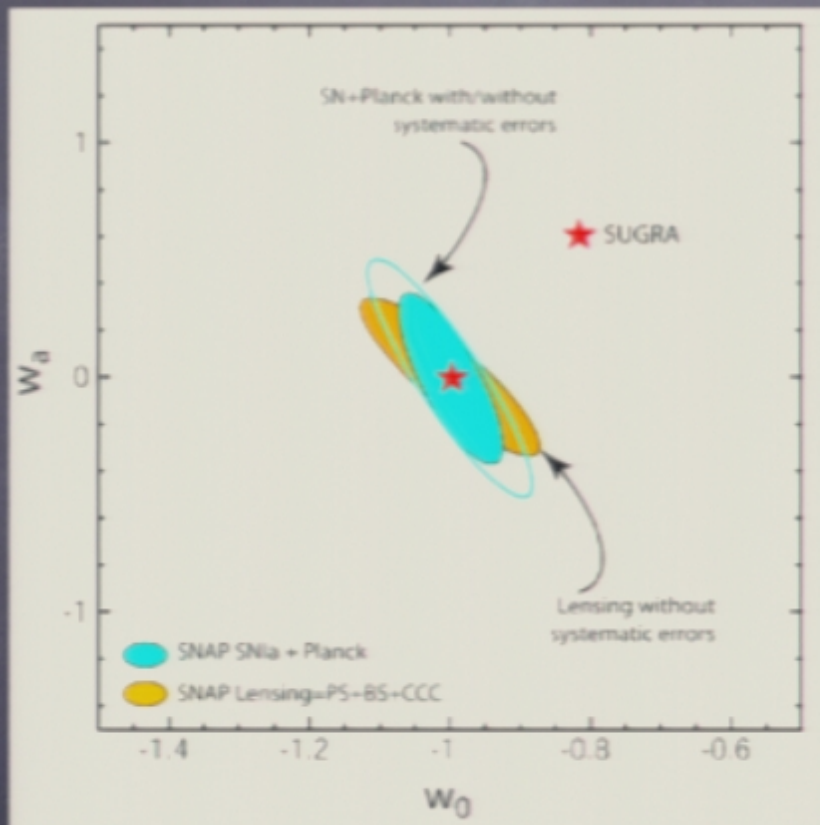
- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (eg dark energy evolution).
- Usually interpreted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however **not** that simulated.
- The criterion for ruling out LambdaCDM is exactly the same as that used to rule out any other value in the plane, eg $w = -0.99$. Special status of LambdaCDM is not recognised.
- Fisher matrix approach assumes a gaussian likelihood.

An unfair provisional comparison

$$w = w_0 + (1 - a)w_a$$

An unfair provisional comparison

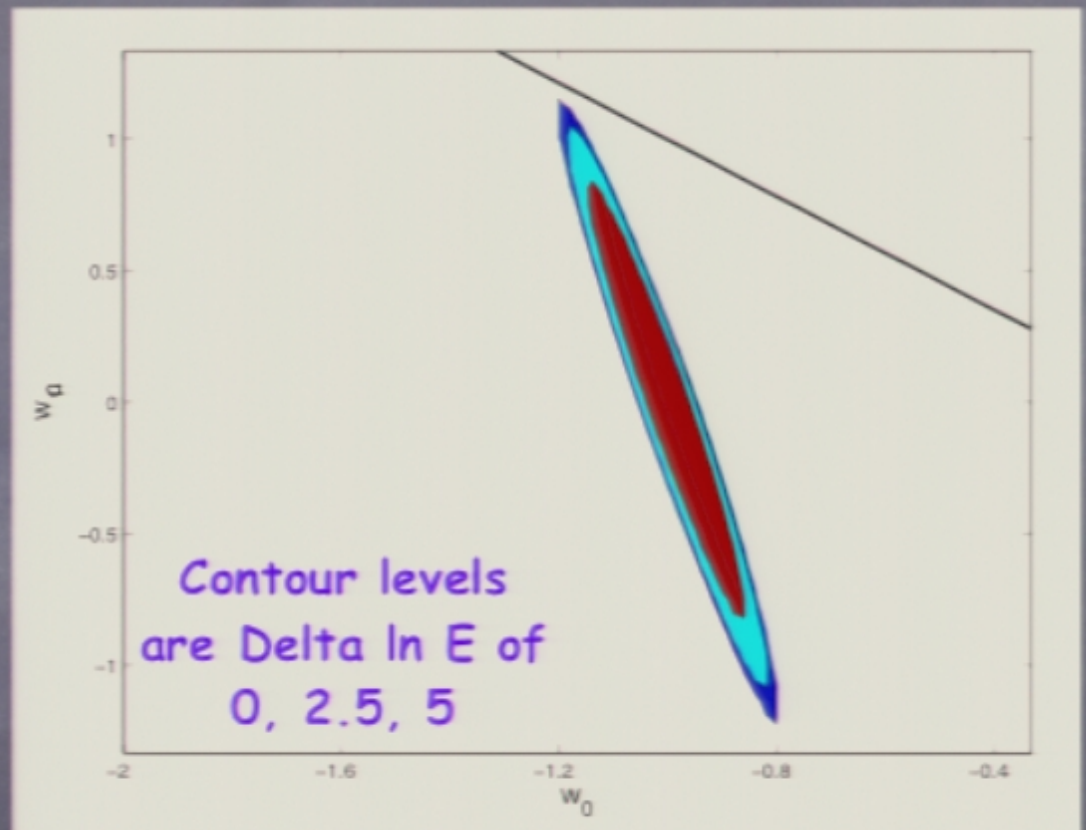
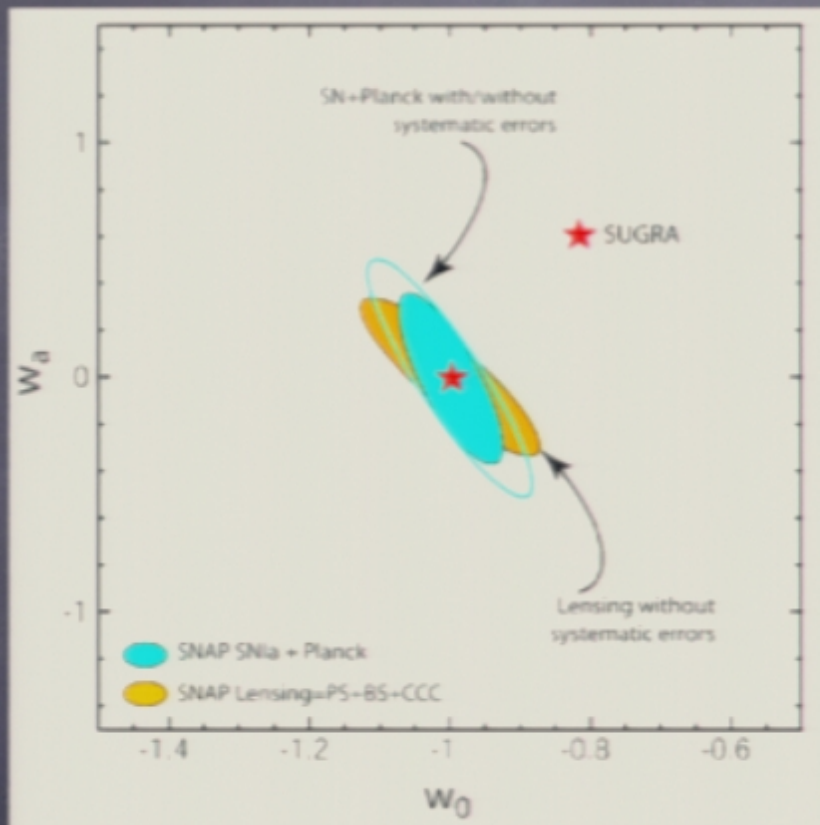
$$w = w_0 + (1 - a)w_a$$



Projected Fisher matrix
uncertainties about LambdaCDM
(SNAP collaboration)

An unfair provisional comparison

$$w = w_0 + (1 - a)w_a$$

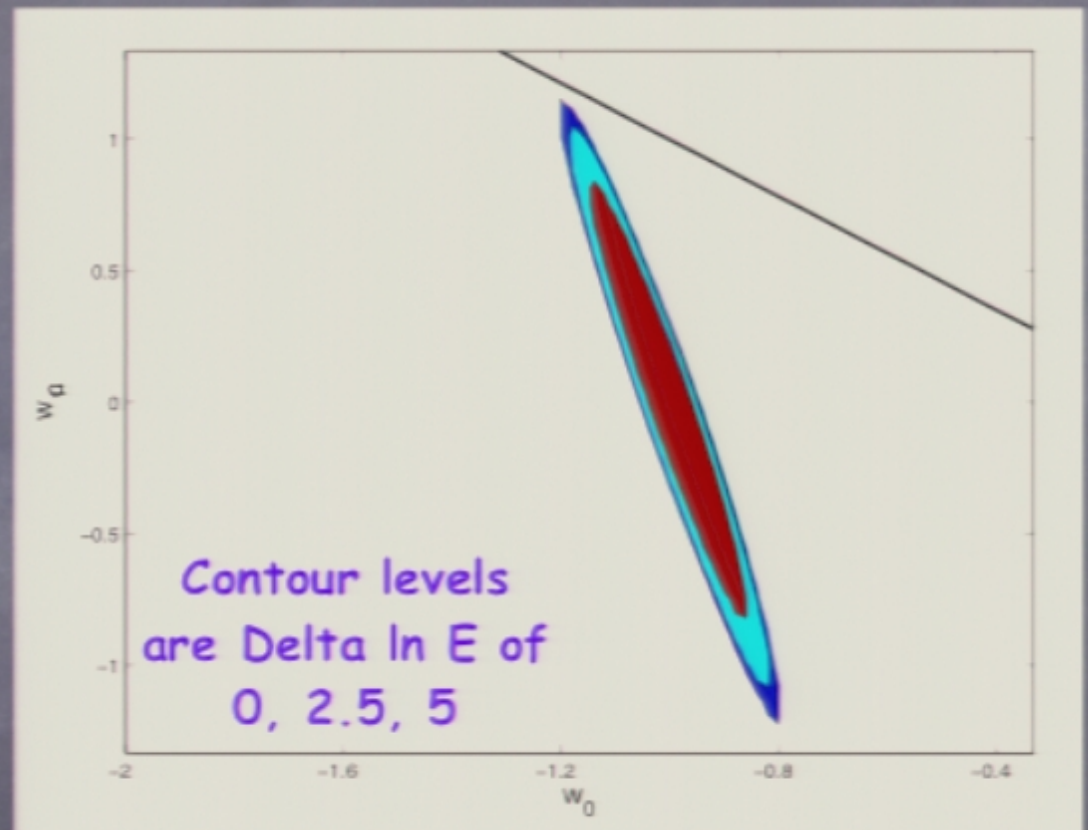
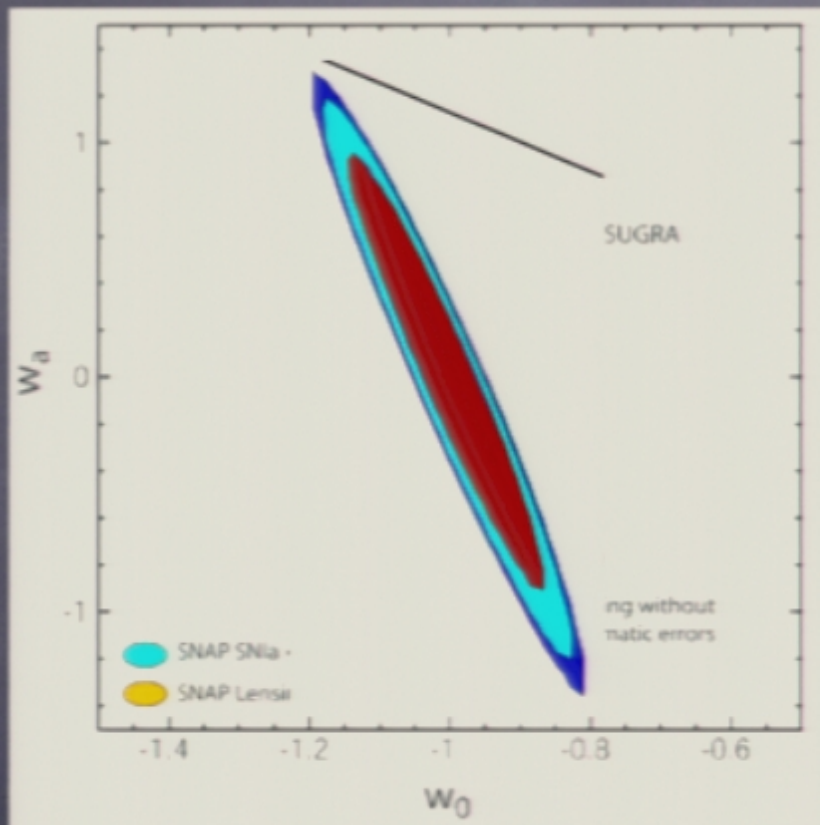


Projected Fisher matrix
uncertainties about LambdaCDM
(SNAP collaboration)

Projected Bayes factor plot against LambdaCDM,
SNAP supernovae only with Omega_matter prior
(myself and collaborators, in prep)

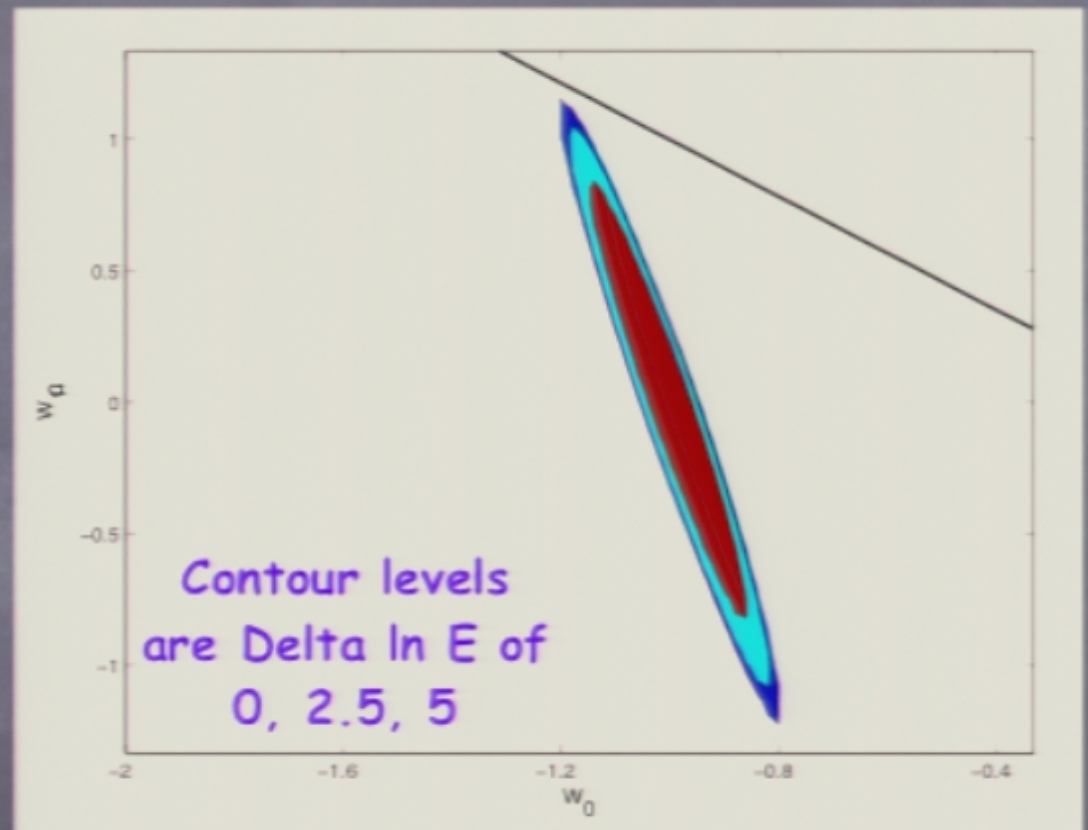
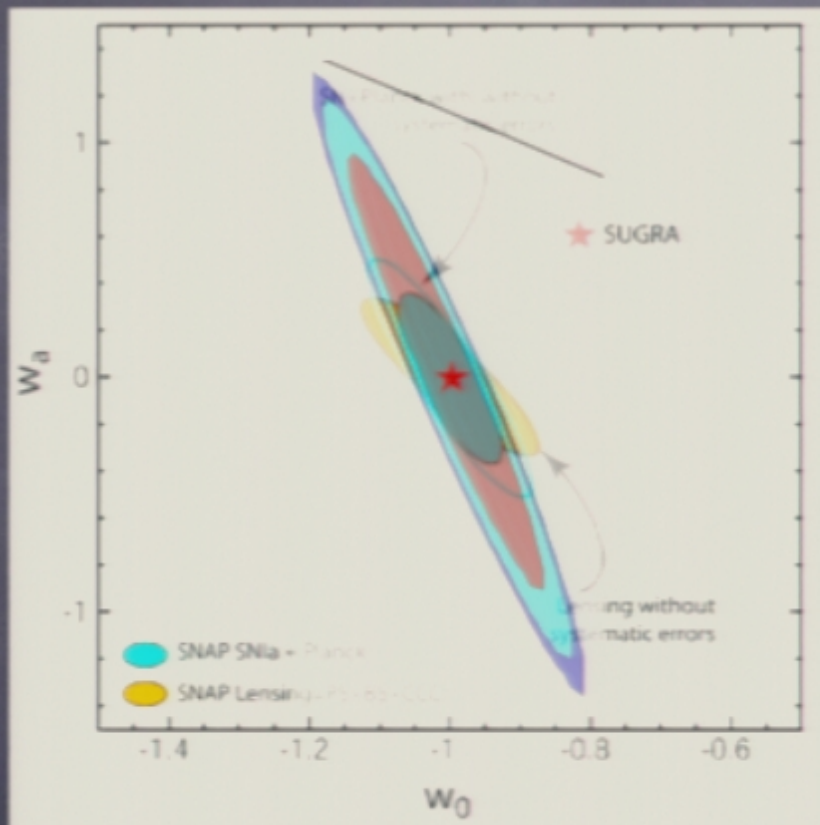
An unfair provisional comparison

$$w = w_0 + (1 - a)w_a$$



An unfair provisional comparison

$$w = w_0 + (1 - a)w_a$$



Projected Fisher matrix
uncertainties about LambdaCDM
(SNAP collaboration)

Projected Bayes factor plot against LambdaCDM,
SNAP supernovae only with Ω_{matter} prior
(myself and collaborators, in prep)

Conclusions

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The **Bayesian evidence** is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model (prior $0.8 < n < 1.2$).

Conclusions

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The **Bayesian evidence** is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model (prior $0.8 < n < 1.2$).

Conclusions

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The **Bayesian evidence** is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model (prior $0.8 < n < 1.2$).
- Isocurvature models are disfavoured, but not at a conclusive level.

Conclusions

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The **Bayesian evidence** is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model (prior $0.8 < n < 1.2$).
- Isocurvature models are disfavoured, but not at a conclusive level.

