Title: Fractional branes and dynamical SUSY breaking

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Abstract: Fractional branes and dynamical SUSY breaking

Fractional branes and dynamical SUSY breaking

Sebastián Franco

Princeton University

September 2005

Based on: hep-th/0505040: Franco, Hanany, Saad and Uranga

See also: hep-th/0502113: Franco, Hanany and Uranga

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Outline

- AdS/CFT correspondence and extensions
- Fractional branes and RG flows
- The conifold cascade
- Complex deformations and toric geometry
- Classifying fractional branes
- DSB from fractional branes
 - Complex cone over dP₁
 - Combining fractional branes
 - Cones over Yp.q manifolds

AdS/CFT correspondence and extensions

AdS/CFT Correspondence

$$\mathcal{N} = 4 \; SU(N) \; \text{SYM} \; \longleftrightarrow \; \text{Type IIB on } AdS_5 \times S^5$$



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Extend the duality by looking for:

Less SUSY

- Place D3-branes at a singularity
- SUGRA: $AdS_5 \times S^5 \rightarrow AdS_5 \times X^5$



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Break conformal invariance

- Add fractional D3-branes (wrapped D5-branes)
- SUGRA: turn on 3-form flux



Quiver theories

Gauge theories that typically arise in our constructions:

- Product gauge group
- $\prod_{i=1}^k SU(d^i)$
- Bifundamental matter

SU(N₁)	SU(N ₂)
\overline{N}_1	N_2

Quiver theories

Gauge theories that typically arise in our constructions:

Product gauge group

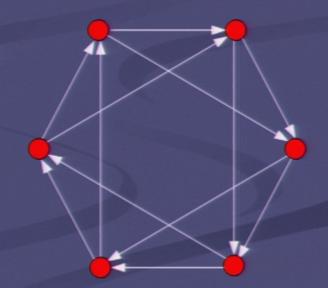
$$\prod_{i=1}^k SU(d^i)$$

Bifundamental matter

$$\begin{array}{c|c} SU(N_1) & SU(N_2) \\ \hline \hline N_1 & N_2 \\ \end{array}$$

Quiver diagram

simple diagrammatic way of encoding the matter content of a gauge theory



Regular and fractional branes

← Anomaly free rank assignments

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Quiver theories

Gauge theories that typically arise in our constructions:

Product gauge group

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Bifundamental matter

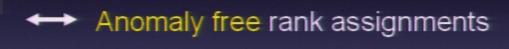
$$\frac{SU(N_1)}{N_1} \qquad \frac{SU(N_2)}{N_2}$$

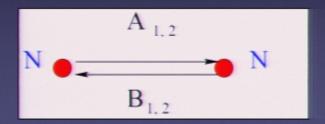
Quiver simple diagrammatic way of encoding diagram the matter content of a gauge theory

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Superpotential terms: oriented polygons in the quiver

Regular and fractional branes

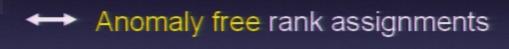


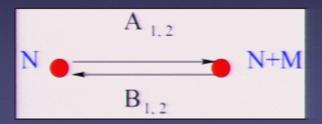


	Gauge group	
N D3-branes	SU(N) × SU(N)	conformal

$$W = \frac{\lambda}{2} \epsilon^{ij} \epsilon^{kl} \operatorname{Tr} A_i B_k A_j B_l$$

Regular and fractional branes



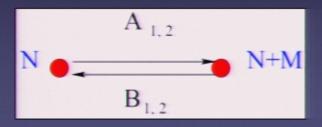


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Regular and fractional branes





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	Gauge group	
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N D3-branes M D5-branes	SU(N) × SU(N+M)	not conformal

$$\beta_i = \frac{d(8\pi^2/g_i^2)}{d\ln\mu}$$

$$\beta_1 = 3 [N + (R_A - 1)(N + M) + (R_B - 1)(N + M)]$$

$$\beta_2 = 3 [(N + M) + (R_A - 1)N + (R_B - 1)N]$$

$$R(A) = 1/2 + \mathcal{O}(M/N)^2 \qquad \beta_1 = -3M$$

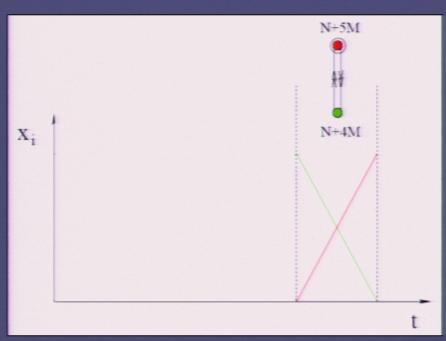
$$\Rightarrow$$

$$R(B) = 1/2 + \mathcal{O}(M/N)^2 \qquad \beta_2 = 3M$$

Duality Cascade

go beyond infinite coupling by switching to an alternative Seiberg dual description Klebanov and Strassler

The conifold example

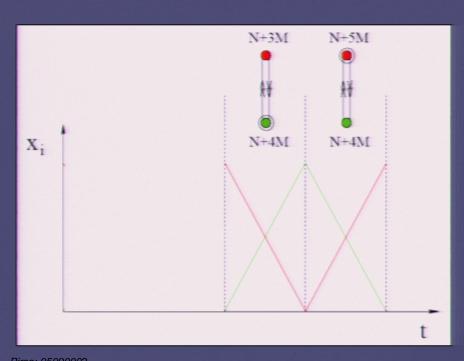


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Duality Cascade

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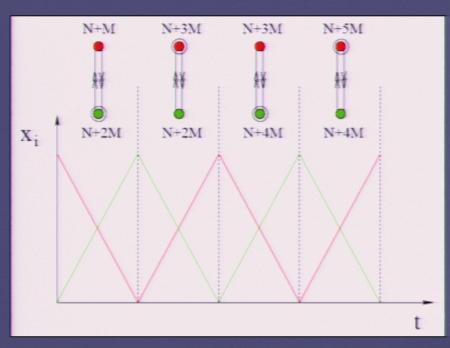
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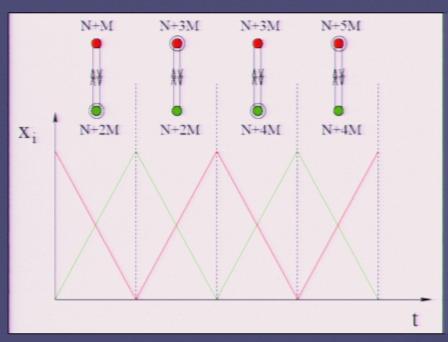
The conifold example



Duality Cascade

go beyond infinite coupling by switching to an alternative Seiberg dual description Klebanov and Strassler

The conifold example



- Periodic flow
- # D5-branes:

M = const

D3-branes:

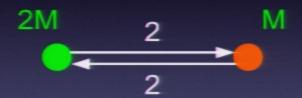
 $N \sim t = \log \mu$

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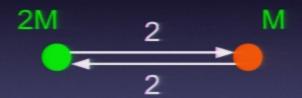
• : 2M colors with 2M flavors

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 : 2M colors with 2M flavors —— quantum corrected moduli space

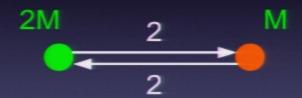
$$\det M - B\tilde{B} = \Lambda^{4M}$$



 : 2M colors with 2M flavors —— quantum corrected moduli space

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$$\downarrow$$
Deformed conifold



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Deformed conifold

- Logarithmic RG flow
- In the IR: confinement and chiral symmetry breaking



 : 2M colors with 2M flavors —— quantum corrected moduli space

$$\det M - B\tilde{B} = \Lambda^{4M}$$

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Deformed conifold

- Logarithmic RG flow
- In the IR: confinement and chiral symmetry breaking

This behavior is nicely captured by a SUGRA dual

Klebanov and Strassle

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Is the conifold behavior generic?

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Is the conifold behavior generic?

Can it be exploited to generate cascades with more structure? Phenomenological models?

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- Is the conifold behavior generic?
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- Infinite families of explicit Sasaki-Einstein geometries and their gauge theory duals are now known

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- Infinite families of explicit Sasaki-Einstein geometries and their gauge theory duals are now known
- New supergravity solutions dual to the UV cascades are known

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Related developments

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Related developments

Z-minimization: geometric dual of a-maximization

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Related developments

- Z-minimization: geometric dual of a-maximization
- Brane tilings: towards the gauge theory dual of an arbitrary toric singularity

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Toric varieties

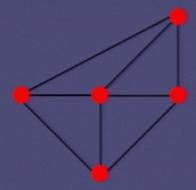
- Admit an U(1)^d action. T^d fibrations
- Described by specifying shrinking cycles and relations

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Toric varieties

- Admit an U(1)^d action. T^d fibrations
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Toric diagrams

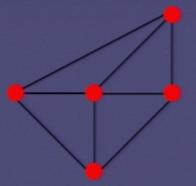


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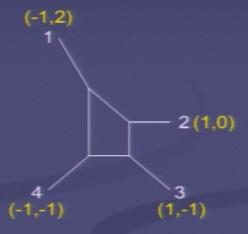
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Toric diagrams



(p,q) webs

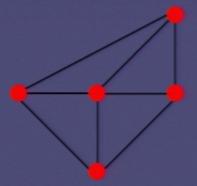


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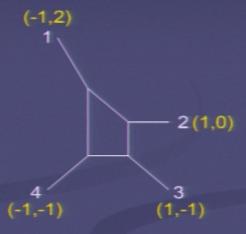
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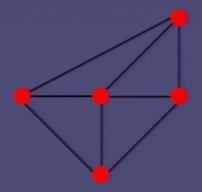
(p,q) webs



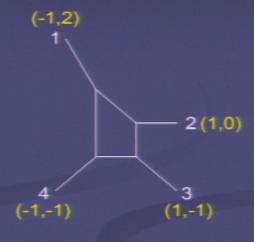
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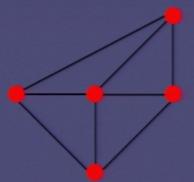
Complex deformations

 Decomposition of (p,q) web into subwebs in equilibrium

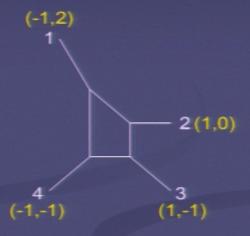
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(p,q) webs



Complex deformations

- Decomposition of (p,q) web into subwebs in equilibrium
- Decomposition of toric polytope into Minkowsky sum Altmann

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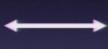
Unobstructed complex deformations



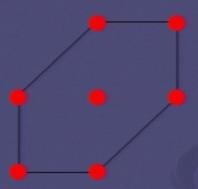
Decomposition of toric polytope into Minkowsky sum

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Unobstructed complex deformations



Decomposition of toric polytope into Minkowsky sum



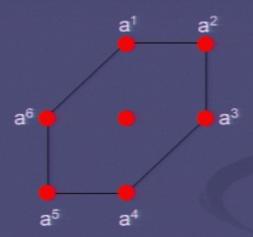
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Unobstructed complex deformations



Decomposition of toric polytope into Minkowsky sum

ai: vertices of toric diagram i = 1 to N



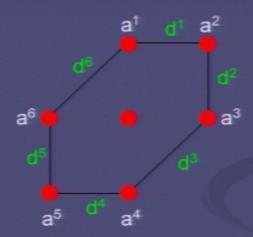
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Unobstructed complex deformations



Decomposition of toric polytope into Minkowsky sum

- ai: vertices of toric diagram i = 1 to N
- di: edges of toric diagram di = ai+1 ai



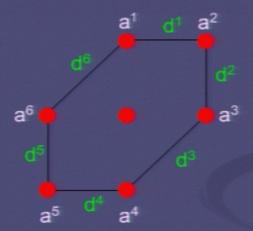
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Unobstructed complex deformations



Decomposition of toric polytope into Minkowsky sum

- ai: vertices of toric diagram i = 1 to N
- d': edges of toric diagram di = ai+1 ai



$$\sum_{i=1}^{N} (t_i)^k d^i = 0$$

Deformation unobstructed upto order K

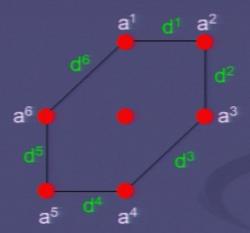
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Unobstructed complex deformations



Decomposition of toric polytope into Minkowsky sum

- ai: vertices of toric diagram i = 1 to N
- dⁱ: edges of toric diagram dⁱ = aⁱ⁺¹ aⁱ



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Deformation unobstructed upto order K

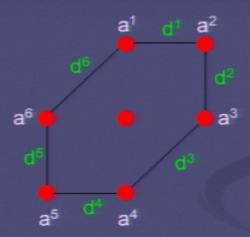
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Unobstructed complex deformations



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Overall scaling plus first order equation: N-3 deformations at first order

Unobstructed complex deformations



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Unobstructed complex deformations



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Deformation unobstructed upto order K

Overall scaling plus first order equation: N-3 deformations at first order

Fractional branes can be classified according to the IR dynamics they trigger

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Fractional branes can be classified according to the IR dynamics they trigger

Deformation fractional branes

 Quiver in the absence of D3-branes corresponds to either isolated nodes without bifundamentals or a gauge invariant in the superpotential

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Fractional branes can be classified according to the IR dynamics they trigger

Deformation fractional branes

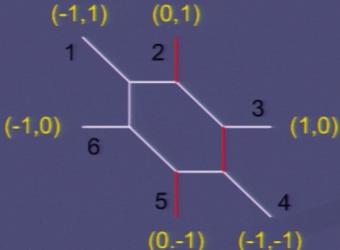
- Quiver in the absence of D3-branes corresponds to either isolated nodes without bifundamentals or a gauge invariant in the superpotential
- SUGRA: complex deformation that (partially) smoothes the singularity

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Fractional branes can be classified according to the IR dynamics they trigger

Deformation fractional branes

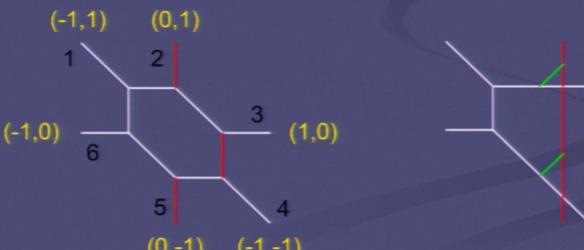
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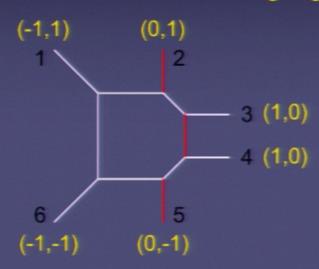


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- flat directions along which the dynamics generically reduces to an N=2 theory
- Quiver in the absence of D3-branes corresponds to a gauge invariant not present in the superpotential

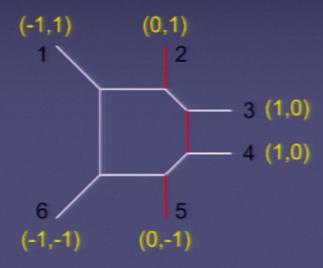
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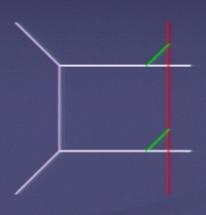
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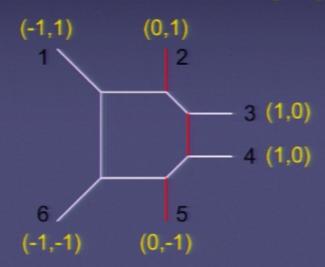
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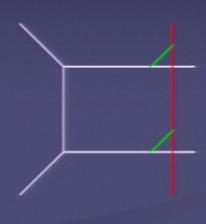


But sometimes the complex deformation associated to a given set of fractional branes is obstructed!

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- flat directions along which the dynamics generically reduces to an N=2 theory
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But sometimes the complex deformation associated to a given set of fractional branes is obstructed!

DSB fractional branes

 All other rank assignments. Typically a gauge group generates an ADS superpotential

Franco, Hanany, Saad and Uranga Berenstein, Herzog, Ouyang and Pinansky

Complex cone over dP₁

5d base is Y^{2,1}, member of the Y^{p,q} family

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Complex cone over dP₁

5d base is Y^{2,1}, member of the Y^{p,q} family

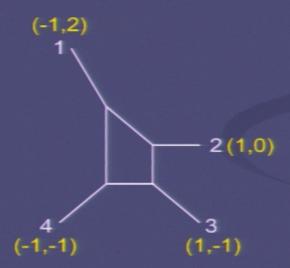
- Y^{2,1}: S² × S³ topology → one single type of fractional branes
- Logarithmic cascade in the UV Franco, Hanany and Uranga
 Ejaz, Herzog and Klebanov

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Complex cone over dP₁

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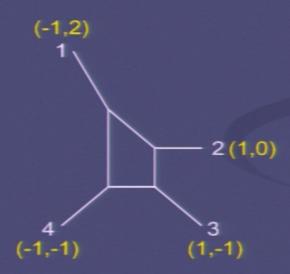


No decomposition of the (p,q) web

Complex cone over dP₁

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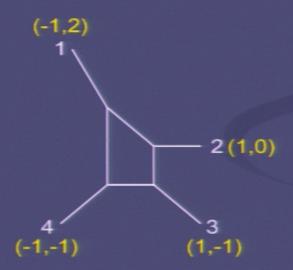


No decomposition of the (p,q) web Complex deformation is obstructed

Complex cone over dP₁

5d base is Y^{2,1}, member of the Y^{p,q} family

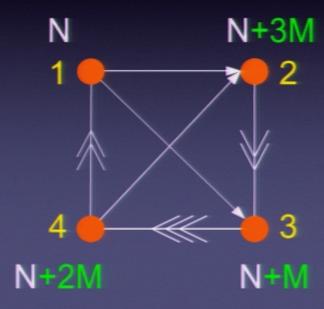
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No decomposition of the (p,q) web

Complex deformation is obstructed

What happens at the → IR bottom of the cascade? Page 61/114



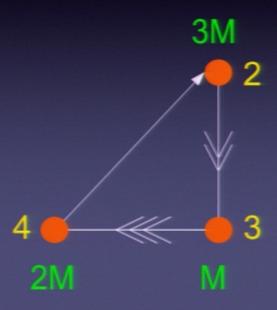
Global symmetry:

$$SU(2) \times U(1)_F \times U(1)_R \times U(1)_B$$

$$W = \epsilon_{\alpha\beta} X_{23}^{\alpha} X_{34}^{\beta} X_{42} + \epsilon_{\alpha\beta} X_{34}^{\alpha} X_{41}^{\beta} X_{13} - \epsilon_{\alpha\beta} X_{12} X_{23}^{\alpha} X_{34}^{\beta} X_{41}^{\beta}$$

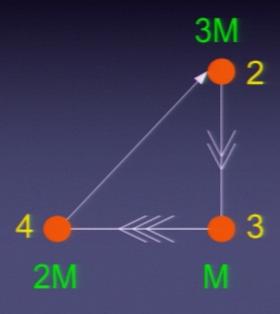
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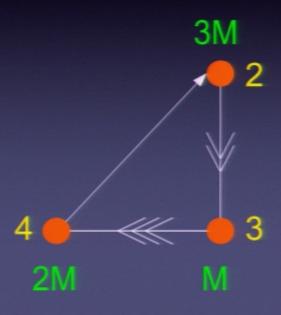
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Fractional branes

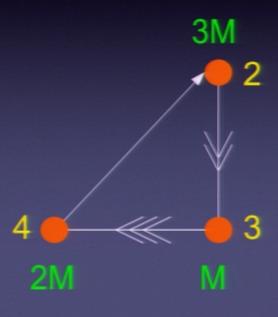
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Fractional branes

| the continuous continuo

$$W = X_{42}X_{23}Y_{34} - X_{42}Y_{23}X_{34}$$



Fractional branes

| pinched to the singularity

$$W = X_{42}X_{23}Y_{34} - X_{42}Y_{23}X_{34}$$

D-flat directions:

$$X_{42} X_{23} X_{34}$$

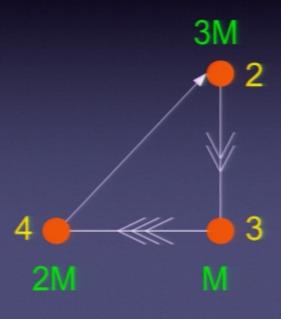
 $X_{42} Y_{23} X_{34}$

$$X_{42} X_{23} Y_{34}$$

 $X_{42} Y_{23} Y_{34}$

$$X_{42} X_{23} Z_{34}$$

 $X_{42} Y_{23} Z_{34}$



Fractional branes

| the continuous continuo

$$W = X_{42}X_{23}Y_{34} - X_{42}Y_{23}X_{34}$$

D-flat directions:

$$X_{42} X_{23} X_{34}$$

 $X_{42} Y_{23} X_{34}$

$$X_{42} X_{23} Y_{34}$$

$$X_{42} X_{23} Z_{34}$$

$$X_{42}\,Y_{23}\,Y_{34}$$

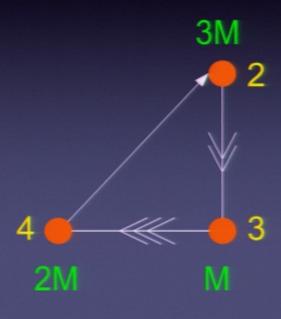
$$X_{42} Y_{23} Z_{34}$$

F-flatness:

$$\frac{\partial W}{\partial Y_{34}} = X_{42} X_{23} = 0$$

$$\frac{\partial W}{\partial X_{24}} = X_{42}Y_{23} = 0$$





Fractional branes

pinched to the singularity

$$W = X_{42}X_{23}Y_{34} - X_{42}Y_{23}X_{34}$$

D-flat directions:

$$X_{42} X_{23} X_{34}$$

 $X_{42} Y_{23} X_{34}$

$$X_{42} X_{23} Y_{34}$$

$${\rm X}_{42}\,{\rm X}_{23}\,{\rm Z}_{34}$$

$$X_{42} Y_{23} Y_{34}$$

$$X_{42} Y_{23} Z_{34}$$

F-flatness:

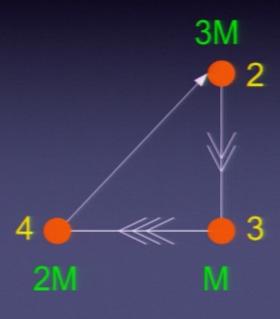
$$\frac{\partial W}{\partial Y_{34}} = X_{42} X_{23} = 0$$

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Classicaly: no mesonic flat directions

the only SUSY point is the origin

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Fractional branes

| pinched to the singularity

$$W = X_{42}X_{23}Y_{34} - X_{42}Y_{23}X_{34}$$

D-flat directions:

$$X_{42} X_{23} X_{34}$$

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Classicaly: no mesonic flat directions

the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors

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- Classicaly: no mesonic flat directions
 - the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

SUSY is broken dinamically

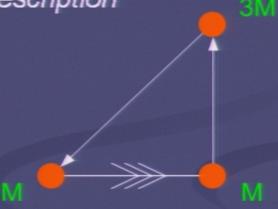
Pirsa: 05090002 Page 73/114

the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

SUSY is broken dinamically

$$W = X_{24}Z_{43}P_{32}$$



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- Classicaly: no mesonic flat directions
 - the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

SUSY is broken dinamically

Pirsa: 05090002 Page 75/114

the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

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$$W = X_{24}Z_{43}P_{32}$$



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SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

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$$W = X_{24}Z_{43}P_{32}$$



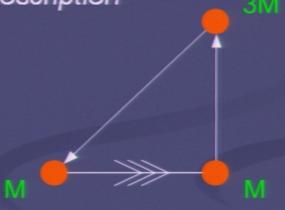
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the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

SUSY is broken dinamically

$$W = M_{34}Z_{43} + 2\left(\frac{\Lambda^8}{M_{34}}\right)^{1/2}$$



$$V_F = |M_{34}|^2 + |Z_{43} + \Lambda^4 M_{34}^{-3/2}|^2$$
 \longrightarrow $M_{34} \to 0$ $Z_{43} \to \infty$

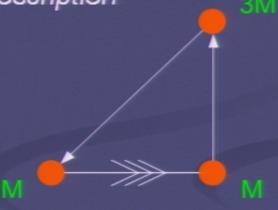
Pirsa: 05090002

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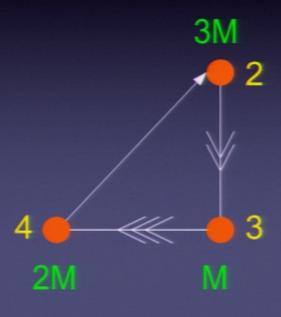
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$$V_D = (|Z_{43}|^2 - \xi)^2$$

Field theory analysis



Fractional branes

| the continuous continuo

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D-flat directions:

$$X_{42} X_{23} X_{34}$$

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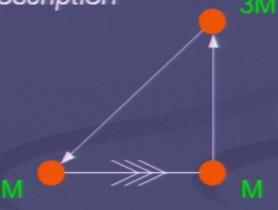
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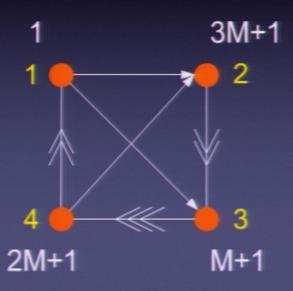
$$W = M_{34}Z_{43} + 2\left(\frac{\Lambda^8}{M_{34}}\right)^{1/2}$$



$$V_F = |M_{34}|^2 + |Z_{43} + \Lambda^4 M_{34}^{-3/2}|^2 \longrightarrow \begin{array}{c} M_{34} \to 0 \\ Z_{43} \to \infty \end{array}$$

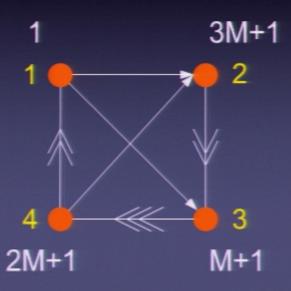
$$V_D = (|Z_{43}|^2 - \xi)^2$$

Non-susy minimum for fixed FI parameters



 $SU(3M+1) \times SU(M+1) \times SU(2M+1)$

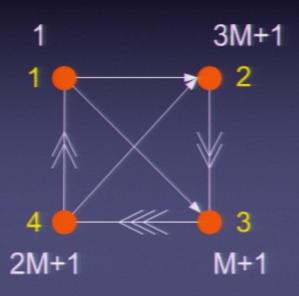
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Flat direction: X₄₂ Y₂₃ Z₃₄

 $SU(3M+1) \times SU(M+1) \times SU(2M+1)$

Pirsa: 05090002



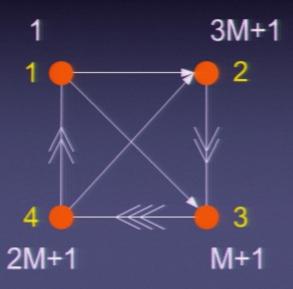
$$SU(3M+1) \times SU(M+1) \times SU(2M+1)$$

Flat direction: $X_{42} Y_{23} Z_{34}$

SU(3M+1) 2M+2 flavors

ADS superpotential for mesons

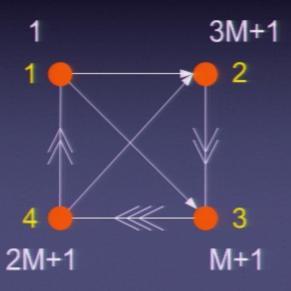
Pirsa: 05090002 Page 85/114



$$SU(3M+1) \times SU(M+1) \times SU(2M+1)$$

ADS superpotential for mesons

$$M = \begin{pmatrix} X_{12} X_{23} & X_{12} Y_{23} \\ X_{42} X_{23} & X_{42} Y_{23} \end{pmatrix}$$



$$SU(3M+1) \times SU(M+1) \times SU(2M+1)$$

ADS superpotential for mesons

$$M = \begin{pmatrix} X_{12} X_{23} & X_{12} Y_{23} \\ X_{42} X_{23} & X_{42} Y_{23} \end{pmatrix}$$

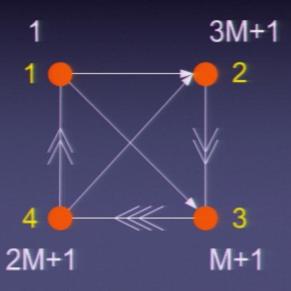
The D3-brane is repelled from the origin

Pirsa: 05090002

Not a complex deformation preserving CY with ISD fluxe^{89e 87/114}

different types of fractional branes that independently lead to SUSY RG flows

Pirsa: 05090002 Page 88/114



$$SU(3M+1) \times SU(M+1) \times SU(2M+1)$$

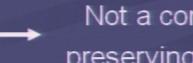
Flat direction:
$$X_{42} Y_{23} Z_{34}$$

ADS superpotential for mesons

$$M = \begin{pmatrix} X_{12} & X_{23} & X_{12} & Y_{23} \\ X_{42} & X_{23} & X_{42} & Y_{23} \end{pmatrix}$$

The D3-brane is repelled from the origin

Pirsa: 05090002



Not a complex deformation preserving CY with ISD fluxes 89/114

different types of fractional branes that independently lead to SUSY RG flows

Pirsa: 05090002 Page 90/114

the only SUSY point is the origin

SU(3M) dominates: SU(3M) has 2M flavors — ADS superpotential

SUSY is broken dinamically

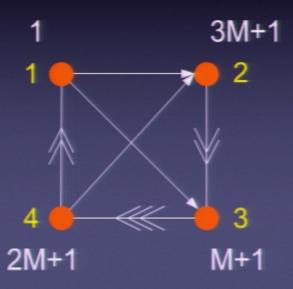
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Non-susy minimum for fixed FI parameters



 $SU(3M+1) \times SU(M+1) \times SU(2M+1)$

Pirsa: 05090002 Page 92/114

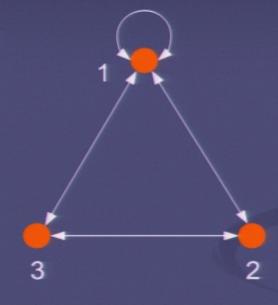
different types of fractional branes that independently lead to SUSY RG flows

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different types of fractional branes that independently lead to SUSY RG flows

can break SUSY when combined

Example: Suspended Pinch Point (SPP)



$$xy = zw^2$$

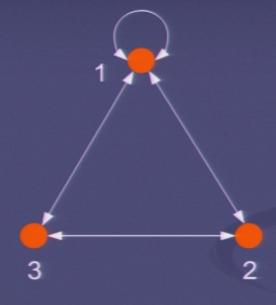
$$W = X_{21}X_{12}X_{23}X_{32} - X_{32}X_{23}X_{31}X_{13} + X_{13}X_{31}X_{11} - X_{12}X_{21}X_{11}$$

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different types of fractional branes that independently lead to SUSY RG flows

can break SUSY when combined

Example: Suspended Pinch Point (SPP)



$$xy = zw^2$$

$$W = X_{21}X_{12}X_{23}X_{32} - X_{32}X_{23}X_{31}X_{13} + X_{13}X_{31}X_{11} - X_{12}X_{21}X_{11}$$

(0,1,0) Complex deformation

(1,0,0)

N=2 brane

moves along line of A_1 singularities parametrized by z (X_{11})



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We expect runaway behavior, with the (1,0,0) brane escaping to infinity

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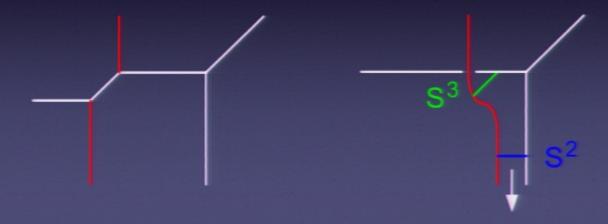


We expect runaway behavior, with the (1,0,0) brane escaping to infinity

E.g. consider:

of
$$(0,1,0)$$
 branes = P
of $(1,0,0)$ branes = M

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We expect runaway behavior, with the (1,0,0) brane escaping to infinity

E.g. consider:

SU(P) gauge group has M flavors and develops ADS superpotential

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We expect runaway behavior, with the (1,0,0) brane escaping to infinity

E.g. consider:

- SU(P) gauge group has M flavors and develops ADS superpotential
- There is no SUSY vacuum and X₁₁ → ••

Infinite family of explicit Sasaki-Einstein metrics

Gauntlett, Martelli, Sparks and Waldram

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Infinite family of explicit Sasaki-Einstein metrics

Gauntlett, Martelli, Sparks and Waldram

S² × S³ topology → one single type of fractional branes.

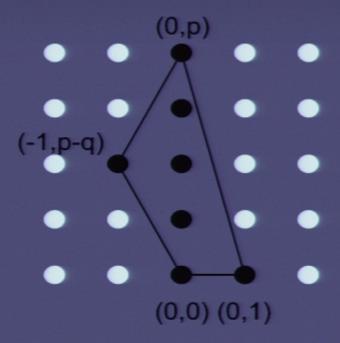
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Yp,q manifolds

Infinite family of explicit Sasaki-Einstein metrics

Gauntlett, Martelli, Sparks and Waldram

S² × S³ topology → one single type of fractional branes.



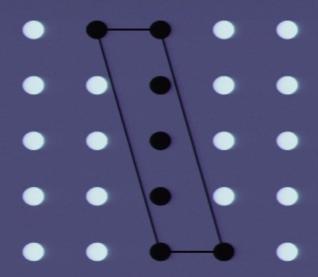
$$0 \le q \le p$$

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Infinite family of explicit Sasaki-Einstein metrics

Gauntlett, Martelli, Sparks and Waldram

S² × S³ topology → one single type of fractional branes.



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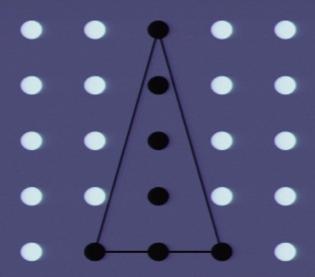
Special cases:

•
$$Y^{p,0} = T^{1,1}/Z_p$$

Infinite family of explicit Sasaki-Einstein metrics

Gauntlett, Martelli, Sparks and Waldram

 $S^2 \times S^3$ topology \rightarrow one single type of fractional branes.



$$0 \le q \le p$$

Y
$$^{p,0} = T^{1,1}/Z_p$$

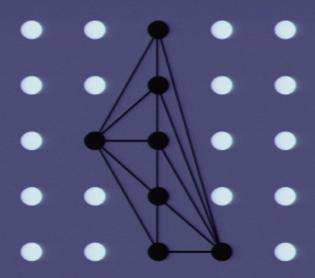
Y $^{p,p} = S^5/Z_{2p}$

•
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Infinite family of explicit Sasaki-Einstein metrics

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Y
$$^{p,p} = S^5 / Z_{2p}$$

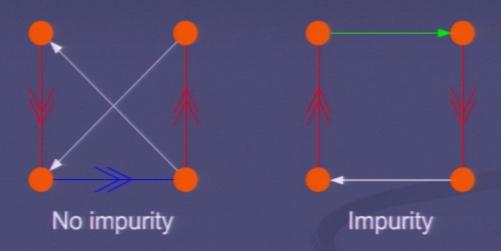
 Supergravity solutions dual to the UV region of the cascades have been constructed

Ejaz, Herzog and Klebanov

Not in contradiction with 1st order complex deformations

Burrington, Liu, Mahato and P. Zayas

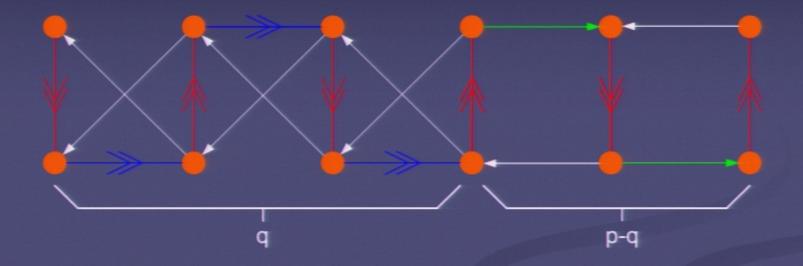
Start from Y p,p = S⁵ / Z_{2p} and add (p-q) impurities



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Not in contradiction with 1st order complex deformations

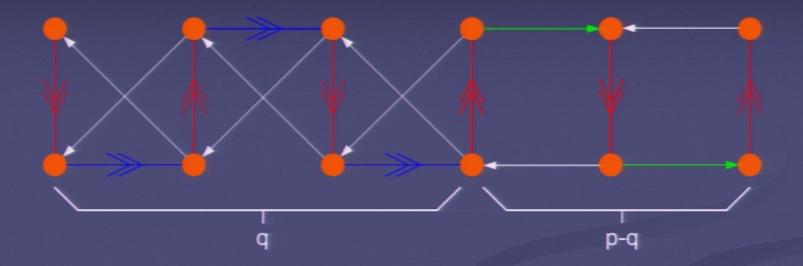
Burrington, Liu, Mahato and P. Zayas



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Not in contradiction with 1st order complex deformations

Burrington, Liu, Mahato and P. Zayas

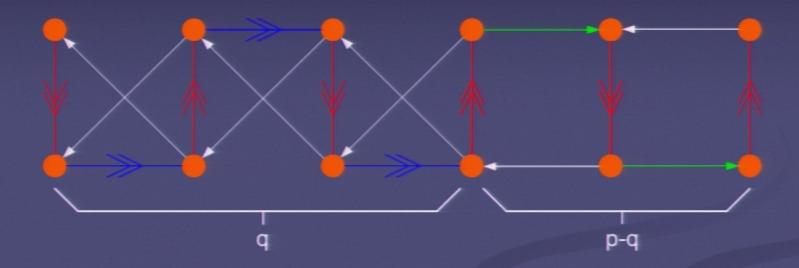


 Construct rank vector of fractional brane using baryonic charges of bifundamental fields

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Not in contradiction with 1st order complex deformations

Burrington, Liu, Mahato and P. Zayas



- Construct rank vector of fractional brane using baryonic charges of bifundamental fields
- For the last node:

$$N_F - N_C = -p + q < 0$$
 for $q < p$

An ADS superpotential is generated and DSB takes place

 We studied the gauge theory dynamics on fractional branes on toric CY singularities

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- We studied the gauge theory dynamics on fractional branes on toric CY singularities
- Three possible behaviors:
 - Complex deformation
 - N=2 dynamics
 - Supersymmetry breaking (associated complex deformation is obstructed)

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- We studied the gauge theory dynamics on fractional branes on toric CY singularities
- Three possible behaviors:
 - Complex deformation
 - N=2 dynamics
 - Supersymmetry breaking (associated complex deformation is obstructed)
- Generic fractional branes lead to SUSY breaking due to non-perturbative superpotentials
- When dynamics of FI terms (equivalently baryonic operators) is taken into account: runaway behavior

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Further directions

Turn these ideas into a simple mechanism to generate DSB in realistic models