

Title: Advanced AdS/CFT Topics

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URL: <http://pirsa.org/05060093>

Abstract:



Selected Topics in AdS/CFT

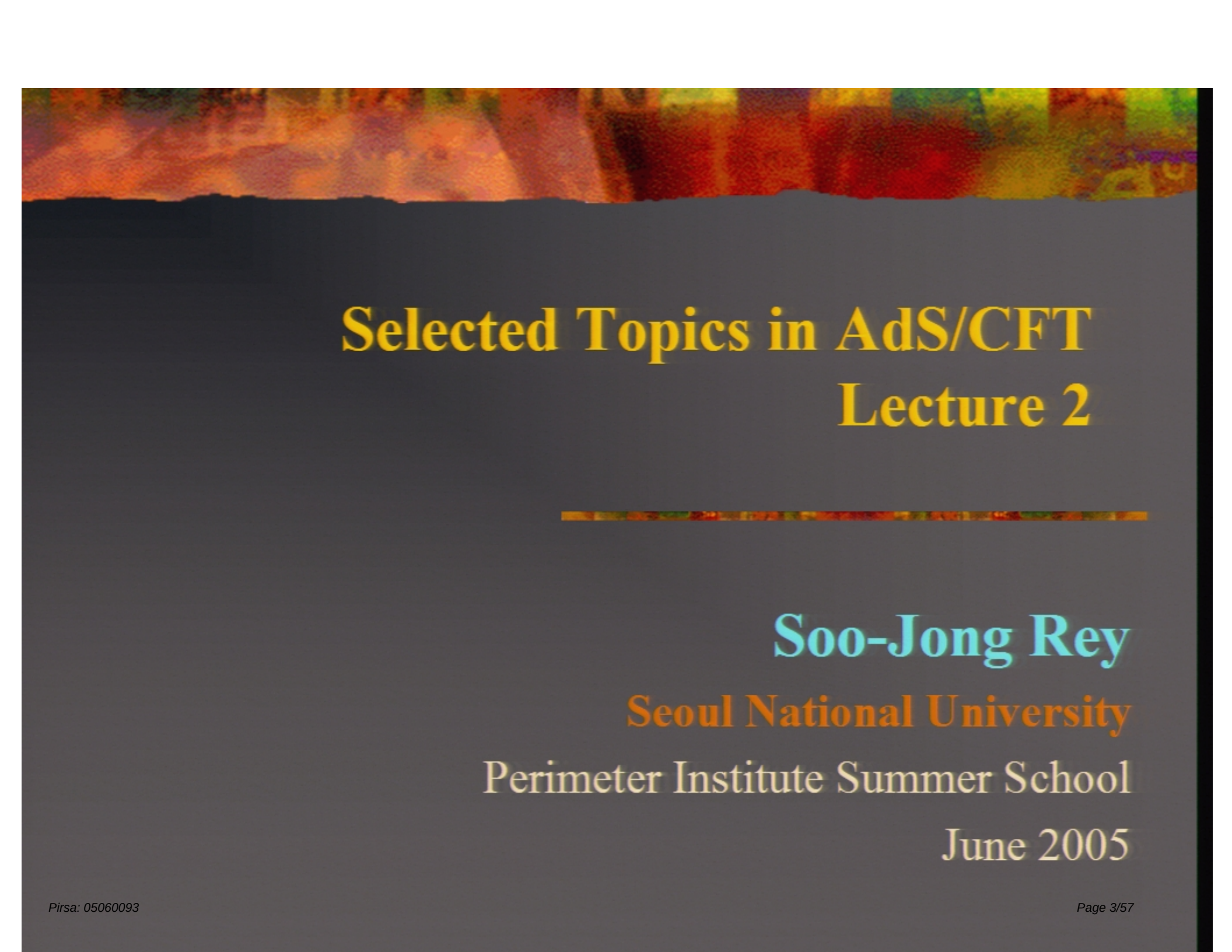
Lecture 2

Soo-Jong Rey

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Perimeter Institute Summer School

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$$\begin{aligned}
 & \underbrace{\text{AdS}_5 \times S_5} \\
 & -X_{-1}^2 - X_5^2 - \vec{X}^2 = R^2 \\
 & \text{SO}(4, 2) \times \text{SO}(6)
 \end{aligned}$$



$$\text{AdS}_5 \times S_5$$

$$-X_4^2 - X_5^2 - \vec{X}^2 = R^2$$

$$SO(4, 2) \times SO(6)$$

SYM_4



$N=4$ SUSY

$$SU(4) \simeq SO(6)$$

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SYM_4
 $\mathbb{R}_{3,1}$

Point



conf

$N=4$ SUSY

$(3,1) \quad SU(4) \simeq SO(6)$

symm

$SO(4,2)$

$\sim SU(2,2)$



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T=0 vs. T>0 SYM Theory

T=0 (zero-temperature 4d N=4 SYM):

$$ds^2 = r^2 (-dt^2 + dx^2) + dr^2/r^2 + (dS_5)^2$$

r = 5th dim // 4d energy scale: $0 < r < \infty$

T>0 (finite-temperature 4d N=4 SYM):

$$ds^2 = r^2 (-F dt^2 + dx^2) + F^{-1} dr^2/r^2 + (dS_5)^2$$

$$F = (1 - (kT)^4/r^4): \quad kT < r < \infty$$

5d AdS Schwarzschild BH = 4d heat bath

Static Quark Potential at Finite Temperature

$$V(r) \sim -(1.254\dots)\sqrt{g_{YM}^2 N} \left(\frac{1}{r} - \frac{1}{r_*} \right) \theta(r - r_*)$$

Notice:

- Nonanalyticity in λ persists
- exact $1/r$ persists
- potential vanishes beyond r_*

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- geometry produced by non-BPS D3-branes:

$$ds^2 = \frac{1}{\sqrt{Z}}(-H dt^2 + dx^2) + \sqrt{Z} \frac{1}{H} dr^2$$

$$H = 1 - \frac{M^4}{r^4} \quad M^4 = \frac{2^7 \pi^4}{3} (g^2 N)^2 \frac{\mathcal{F}}{N^2}$$

$$\text{N=4 SYM free energy} \quad \mathcal{F} = \frac{4\pi^2}{45} N^2 T^4$$

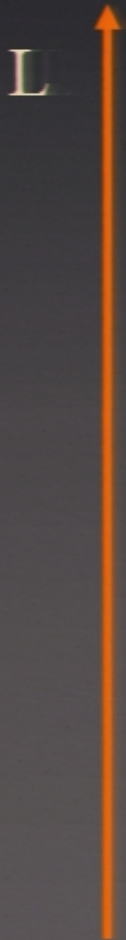
- first integral of motion:

$$\left(\frac{Z}{H}\right)^2 U'^2 + \frac{Z}{H} = \left(\frac{Z}{H}\right)_* = \frac{g^2 N}{M^4} \frac{1}{(U_*/M)^4 - 1}$$

- $T \neq 0$ geometric UV-IR relation:

$$L = \frac{\sqrt{g^2 N}}{U_*} f\left(\frac{U_*}{M}\right) \quad \text{with} \quad U_* \geq M$$

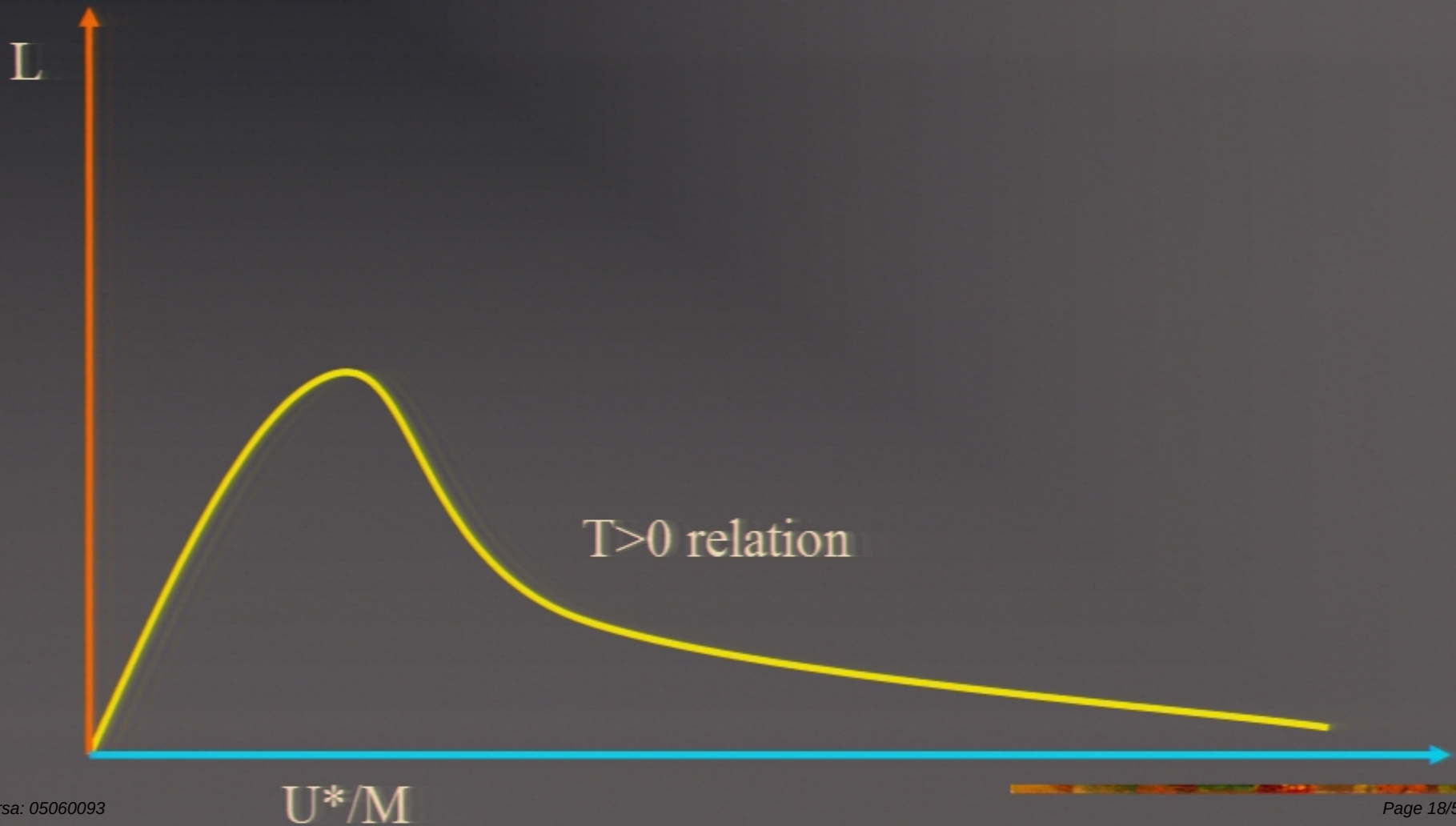
UV-IR relation for $T > 0$



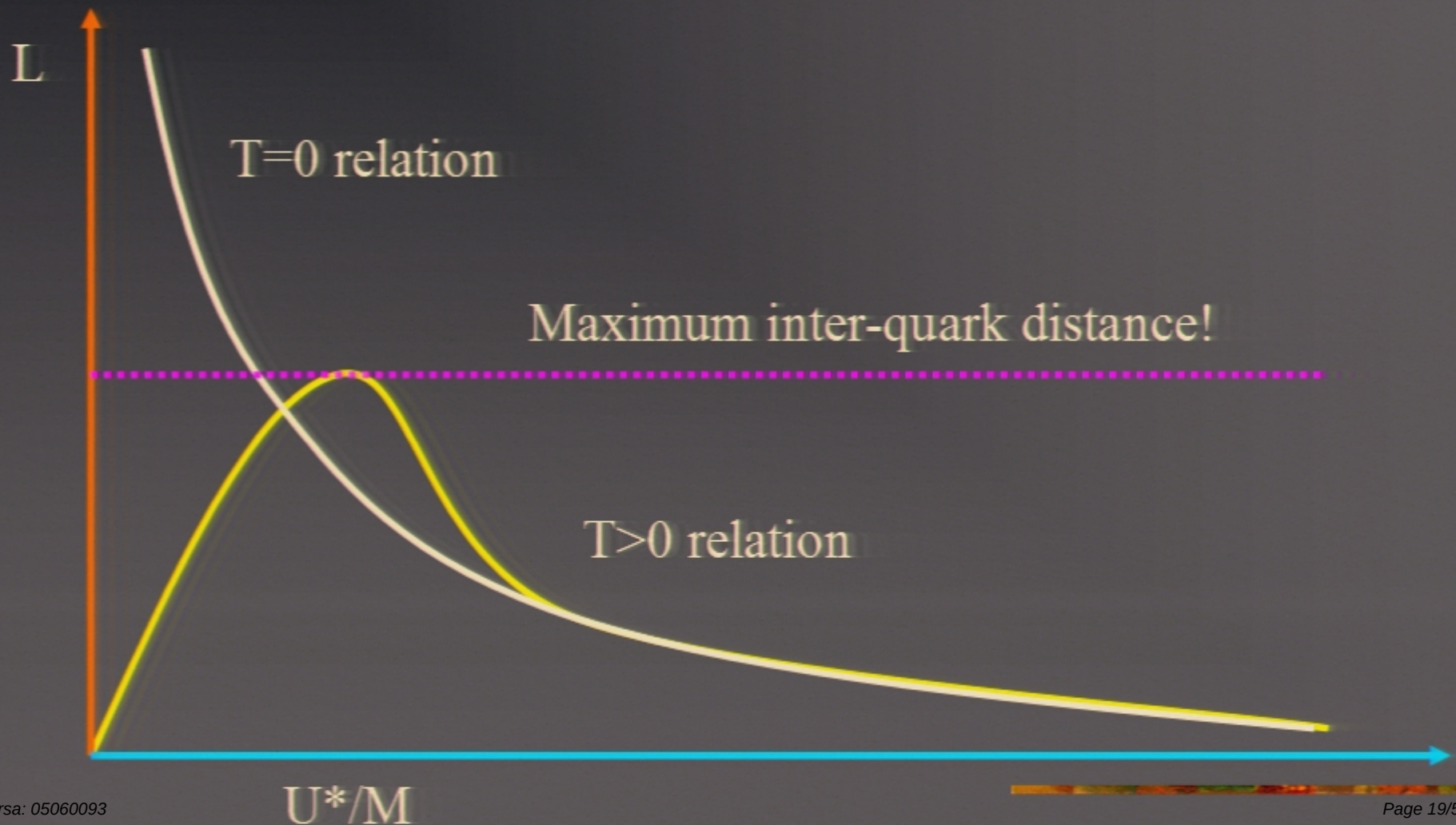
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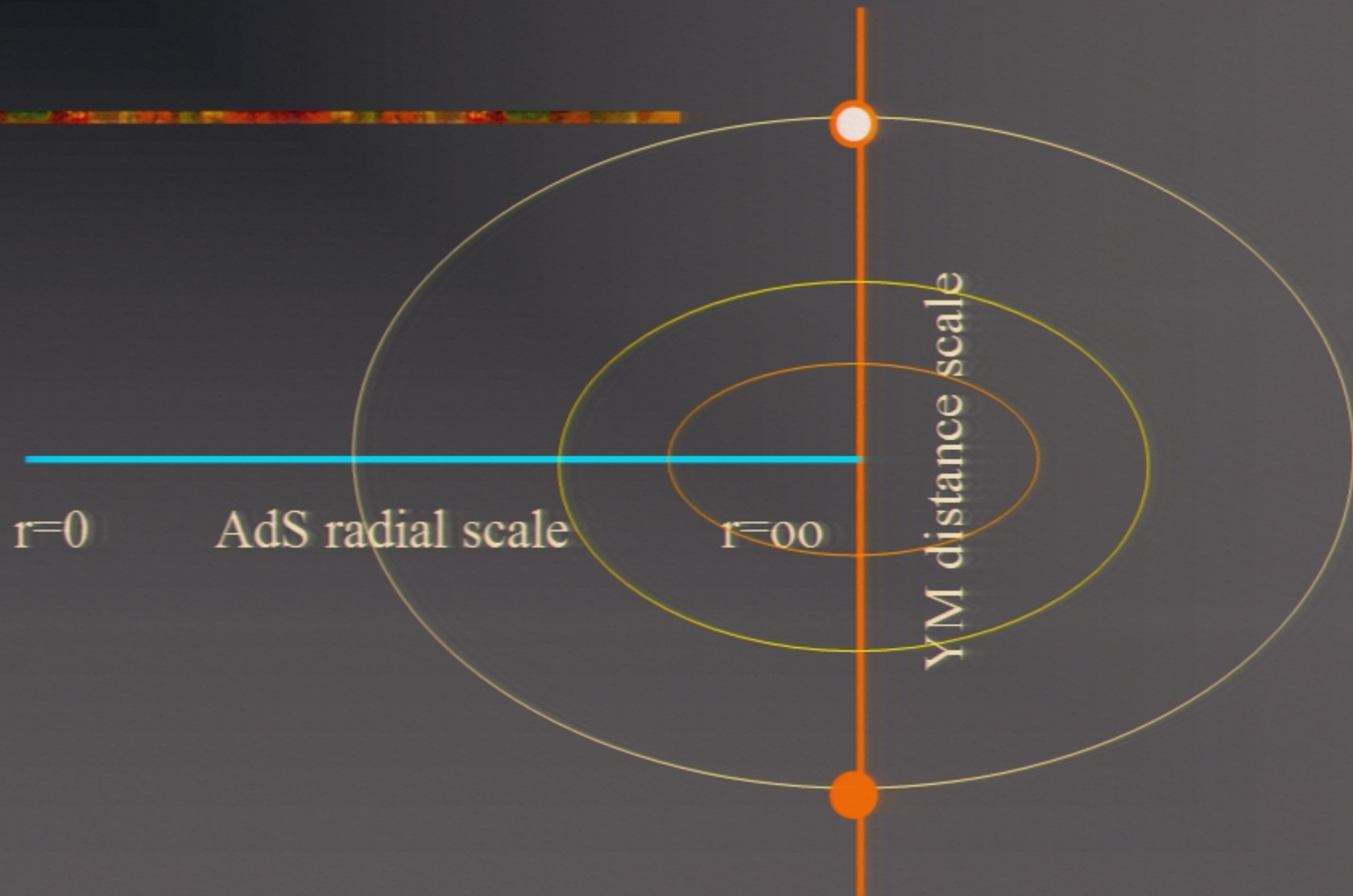
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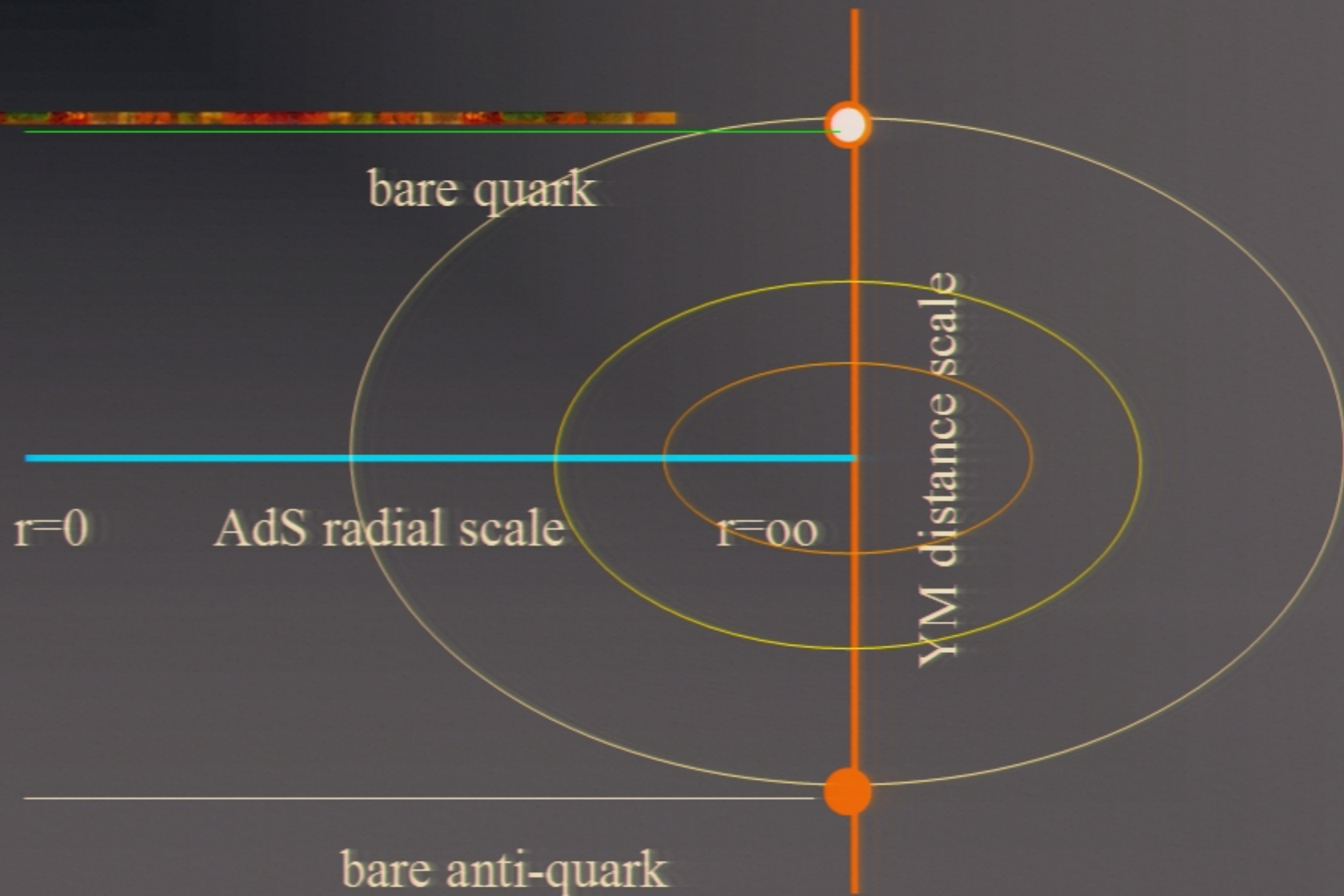
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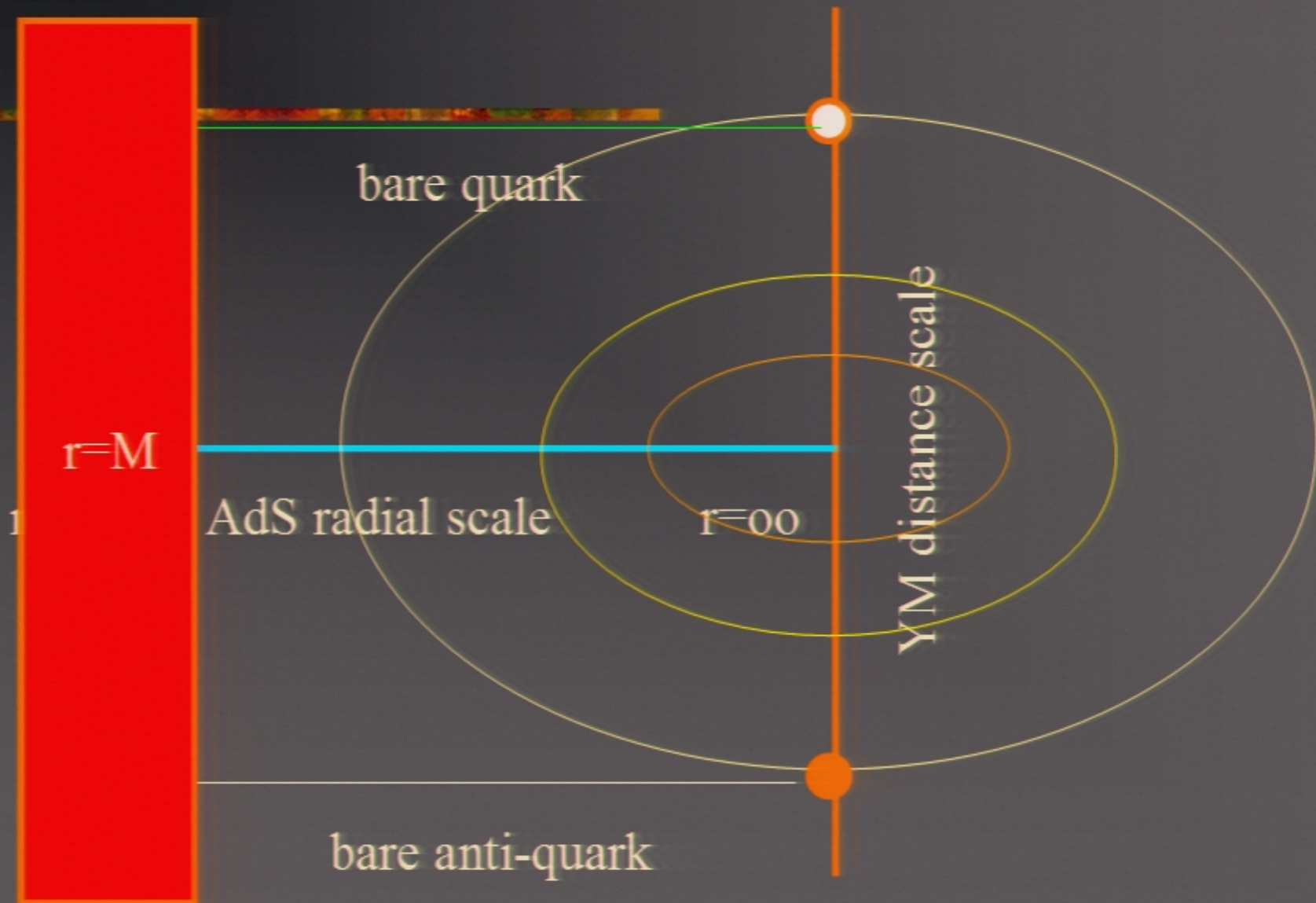
Heavy Meson Configuration ($T > 0$)



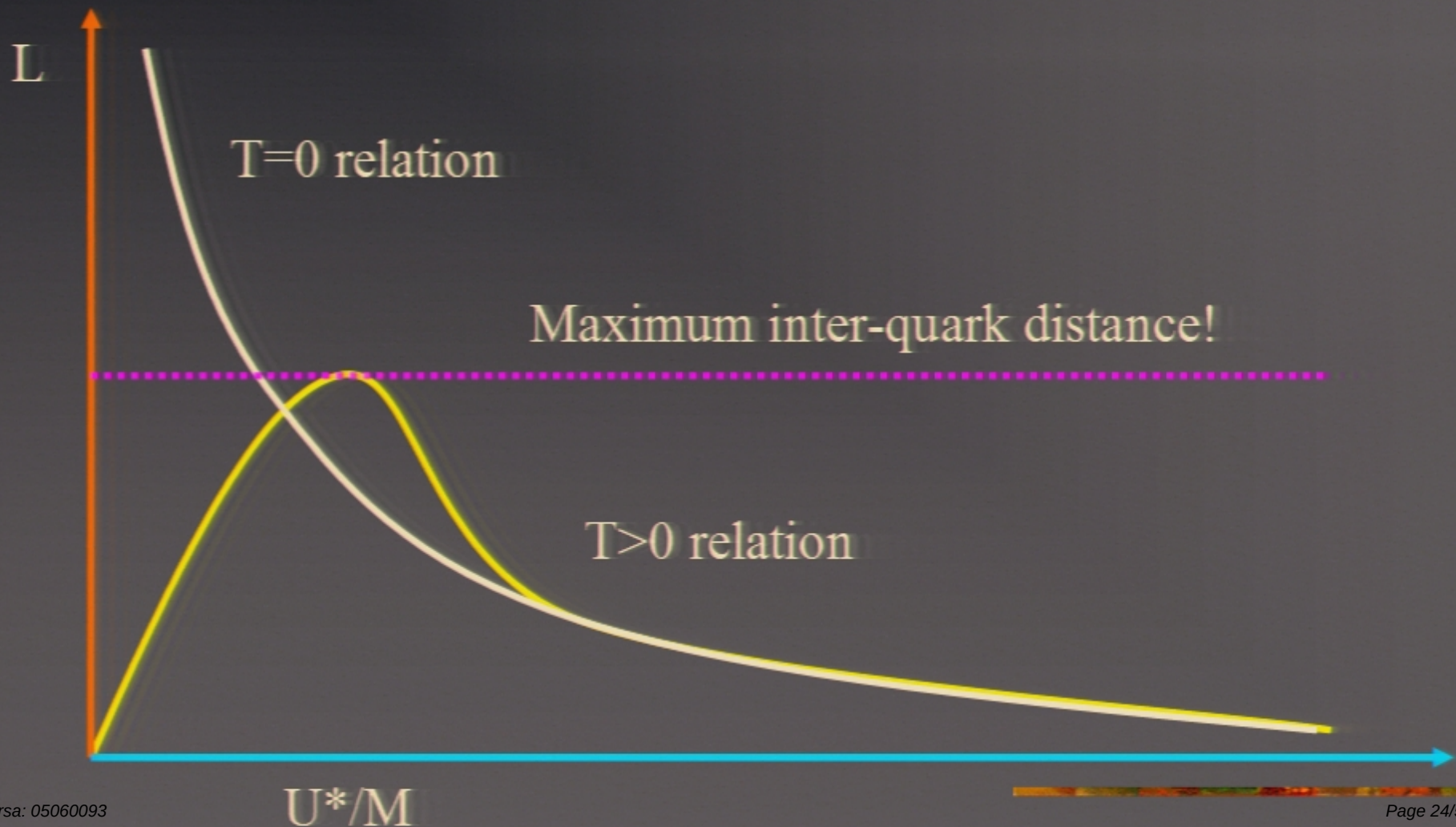
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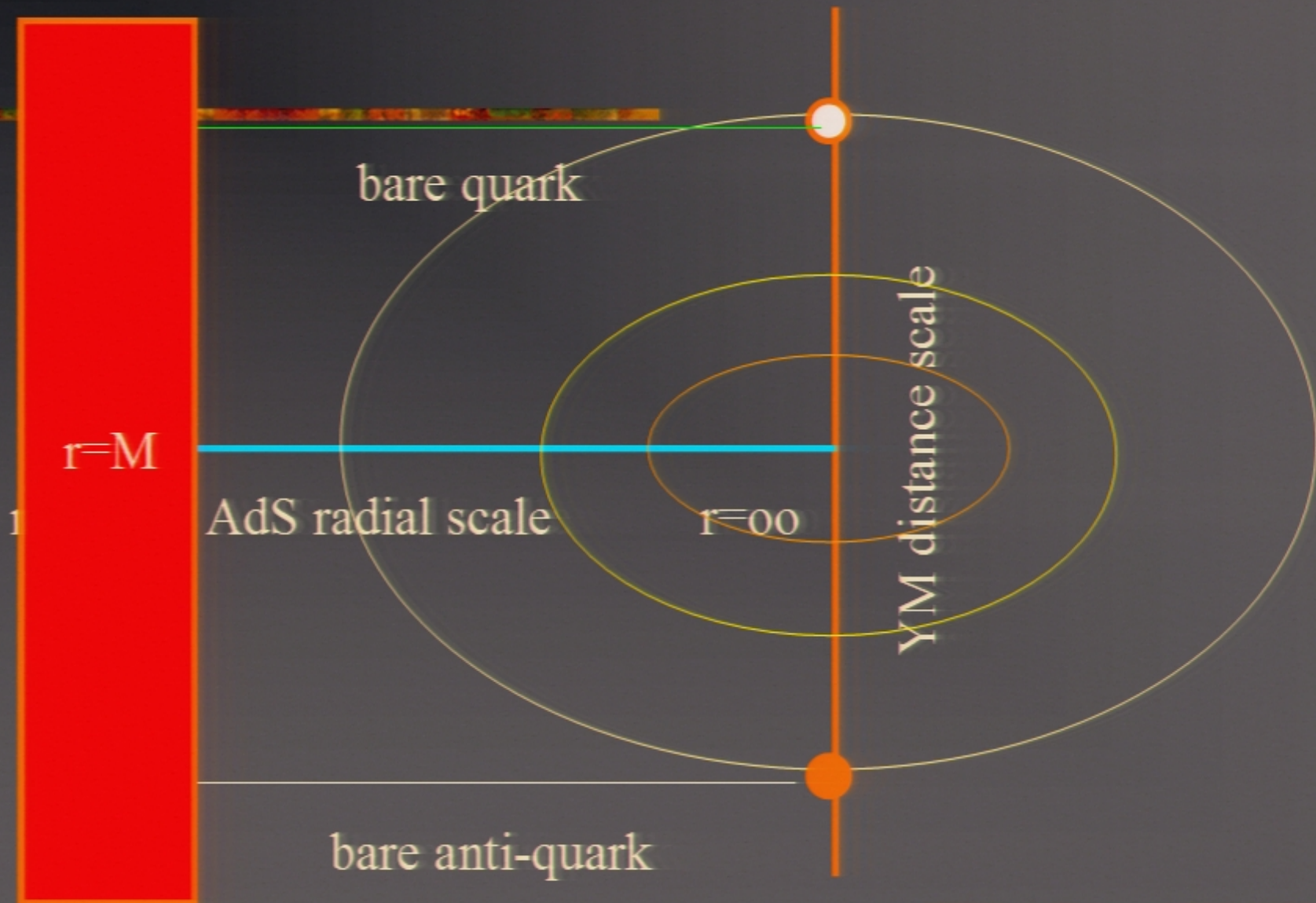
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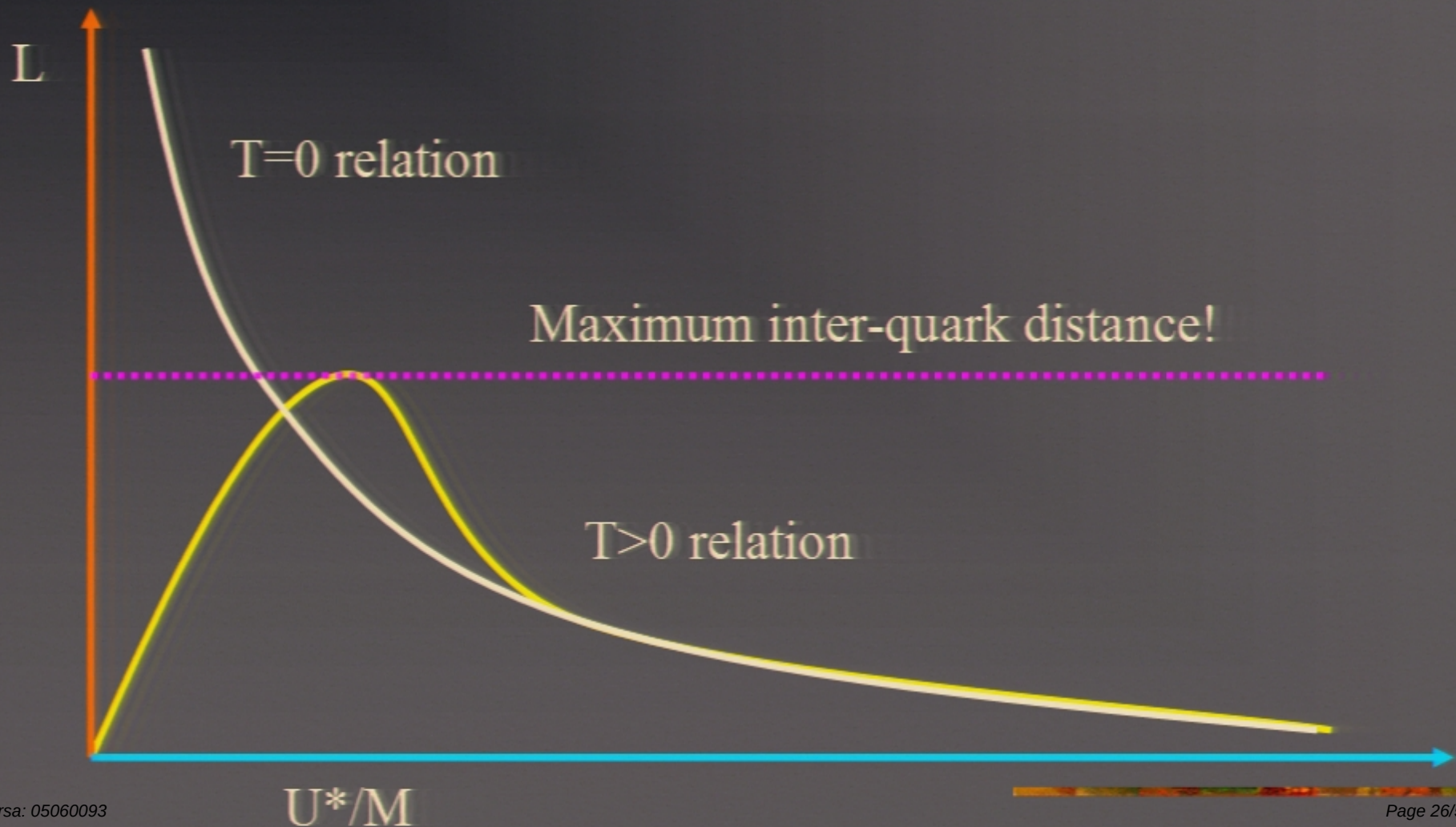
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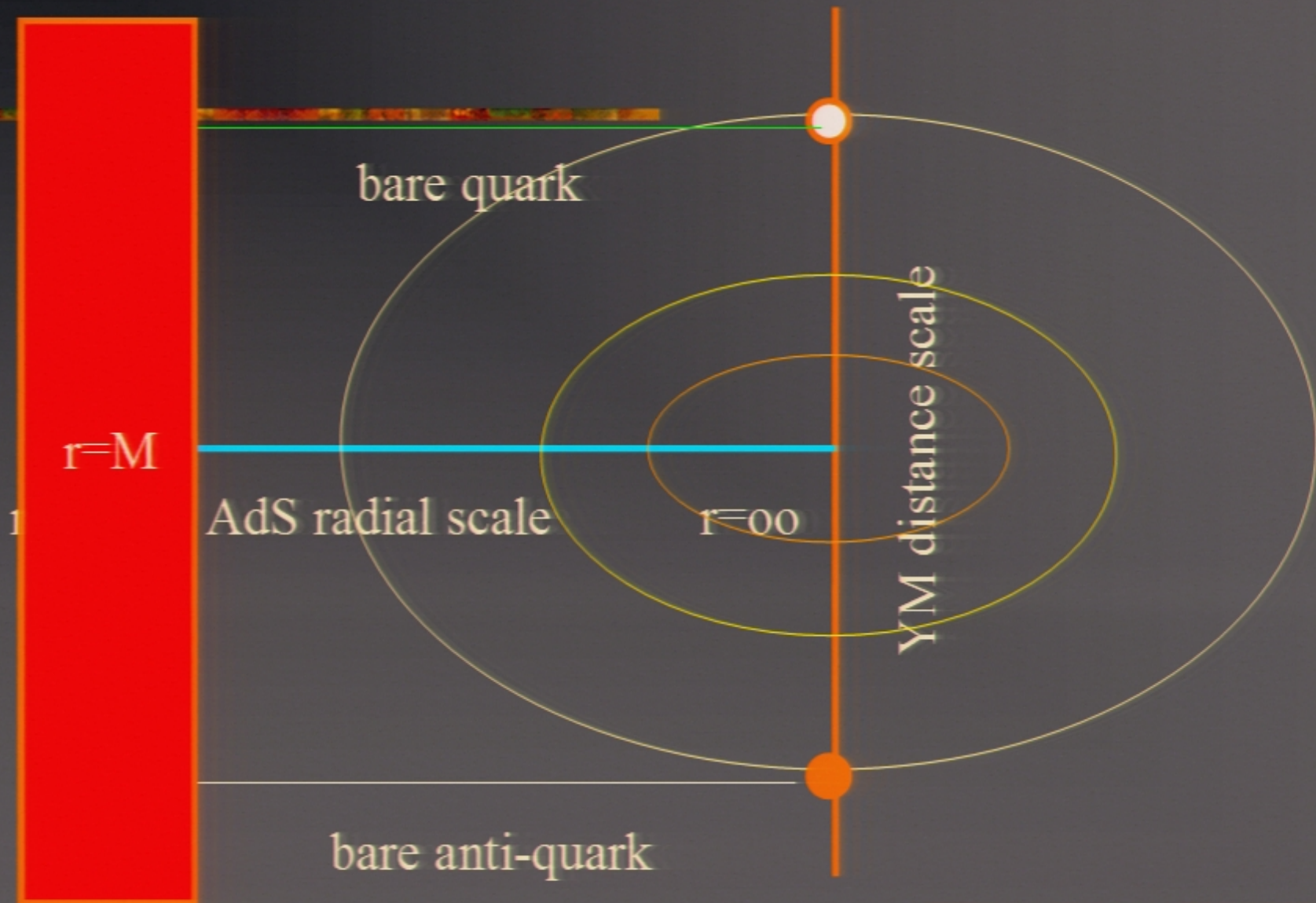
Heavy Meson Configuration ($T > 0$)



UV-IR relation for $T > 0$



Heavy Meson Configuration ($T > 0$)



- string energy:

$$\begin{aligned}
 E_{(Q\bar{Q})}(a = U_*/M) &= \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2 + H Z^{-1}} \\
 &= \lim_{u \rightarrow \infty} 2M \int_a^u dx \sqrt{\frac{x^4 - 1}{x^4 - a^4}}
 \end{aligned}$$

- heavy-quark potential:

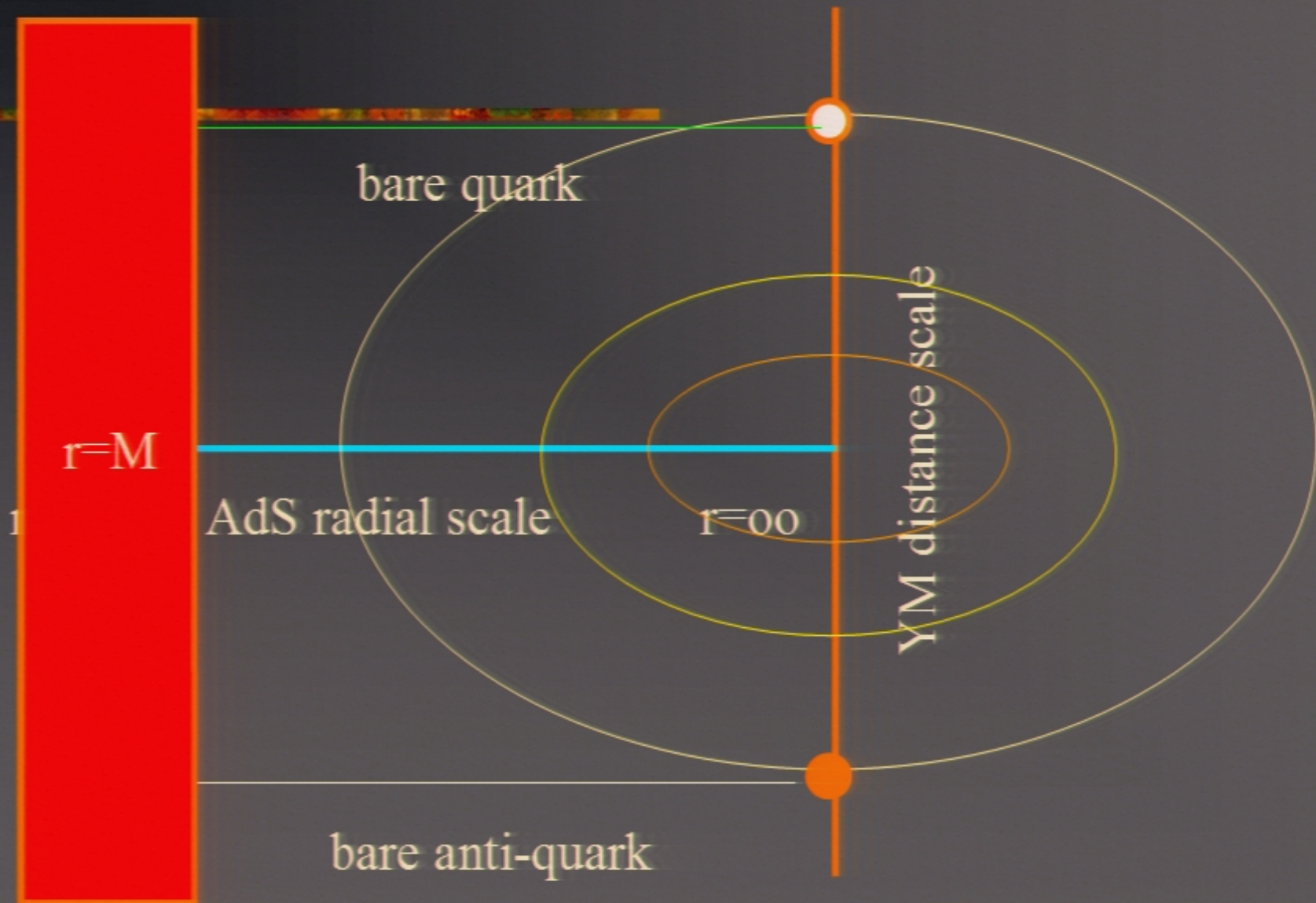
$$V_{(Q\bar{Q})}(U_*) \equiv E_{(Q\bar{Q})}(a) - (M_Q + M_{\bar{Q}})_{T>0}$$

- using "geometric UV-IR duality",

$$V_{Q\bar{Q}}(L) = -C \frac{\sqrt{g^2 N}}{L} h(LM) \quad C = 1.254\dots$$

$$h(LM) = 1 - \frac{1}{2C} \frac{LM}{\sqrt{g^2 N}} + \mathcal{O}\left(\frac{LM}{\sqrt{g^2 N}}\right)^4$$

Heavy Meson Configuration ($T > 0$)



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Compactify on “thermal” S1

- * Euclidean AdS5 Schwarzschild black hole:

$$ds^2 = r^2 (F dt^2 + dx^2) + F^{-1} dr^2/r^2 + (dS_5)^2$$

$$F = (1 - (kT)^4/r^4): \quad kT < r < \infty$$

- * antiperiodic boundary condition for fermions on τ
- * 3d x-space is now Lorentzian
- * Similar construction works for nonextreme D4-brane wrapped on “thermal” S1 $\rightarrow$ 4d Lorentzian x-space

Application:

D-branes on “thermal” S_1

- ❑ “thermal S_1 ” breaks $N=4$ susy completely
- ❑ At low-energy, 3d Yang-Mills + (junks)
- ❑ 5d AdS replaced by
5d Euclidean black hole (time $\leftrightarrow$ space)
- ❑ glueball spectrum is obtainable by studying
bound-state spectrum of gravity modes

Note:

4d space-time, topology: gravity // gauge

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YM₂₊₁ glueball spectrum

□ 0⁺⁺: solve dilaton eqn=2nd order linear ode

□ result:

	N=3 lattice	N=∞ lattice	AdS/CFT
.	4.329(41)	4.065(55)	4.07(input)
*	6.52 (9)	6.18 (13)	7.02
**	8.23 (17)	7.99 (22)	9.92
***	-	-	12.80

[M. Teper]

YM₃₊₁ glueball spectrum

- Use T>0 D4-brane instead
- 0⁺⁺: solve dilaton eqn=2nd-order linear ODE
- result:

	N=3 lattice	AdS/CFT
.	1.61(15)	1.61(input)
*	2.8	2.38
**	-	3.11
***	-	3.82

[M. Teper]

other glueballs fit reasonably well (why??)

The Story of Square-Root

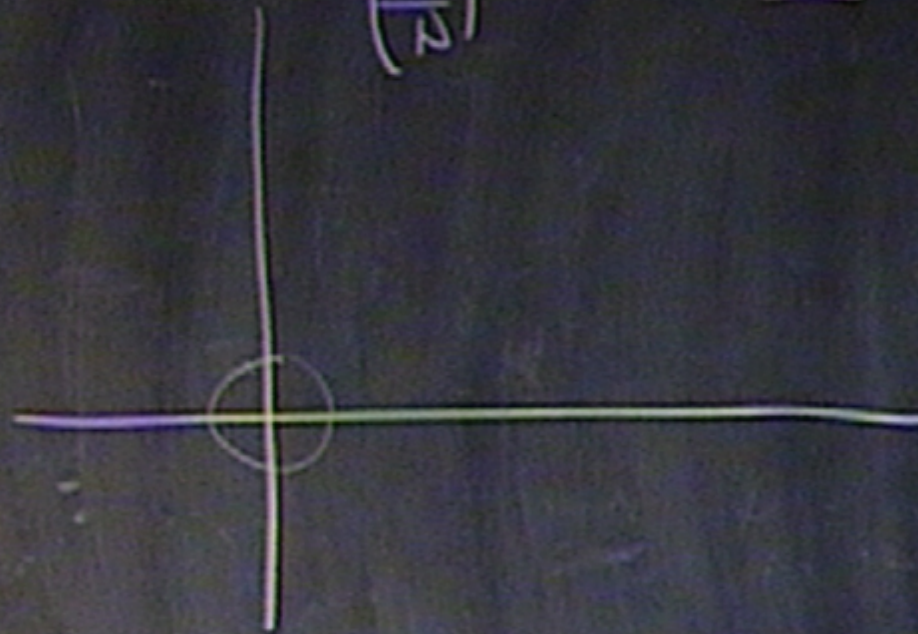
- branch cut from strong coupling?
- artifact of $N \rightarrow \infty$ limit
- heuristically, saddle-point of matrices

$$\langle \text{Tr } e^M \rangle = \int [dM] (\text{Tr } e^M) \exp(-\lambda^{-2} \text{Tr } M^2)$$

Recall modified Bessel function

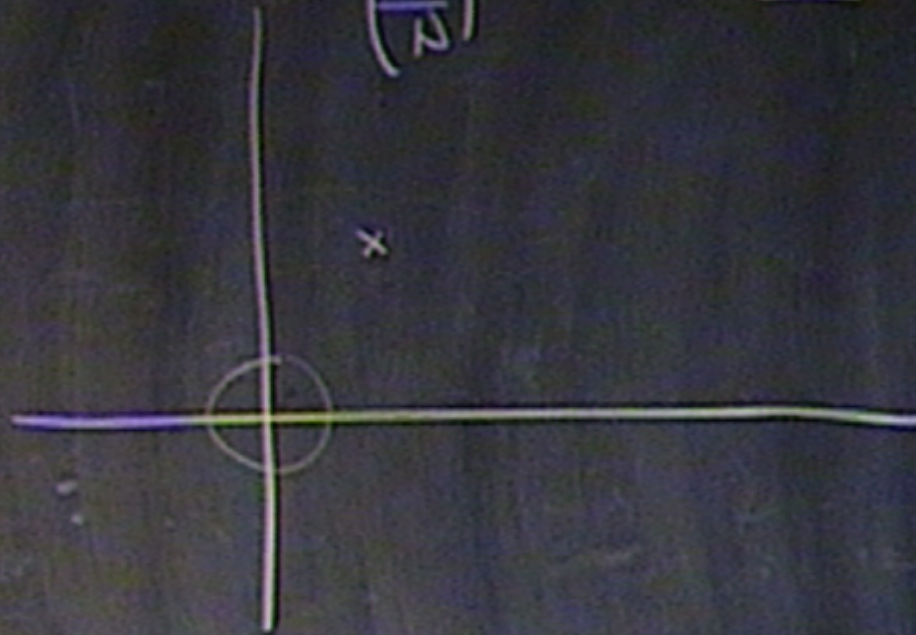
$O(6)$

$$\left(\frac{1}{N}\right)^{2h-2} \quad \underline{L\uparrow}$$



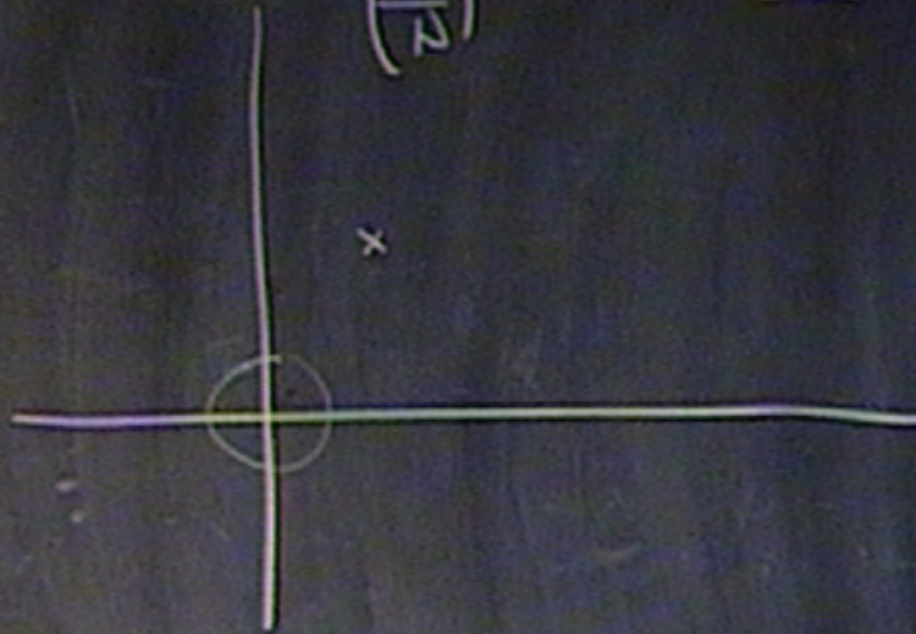
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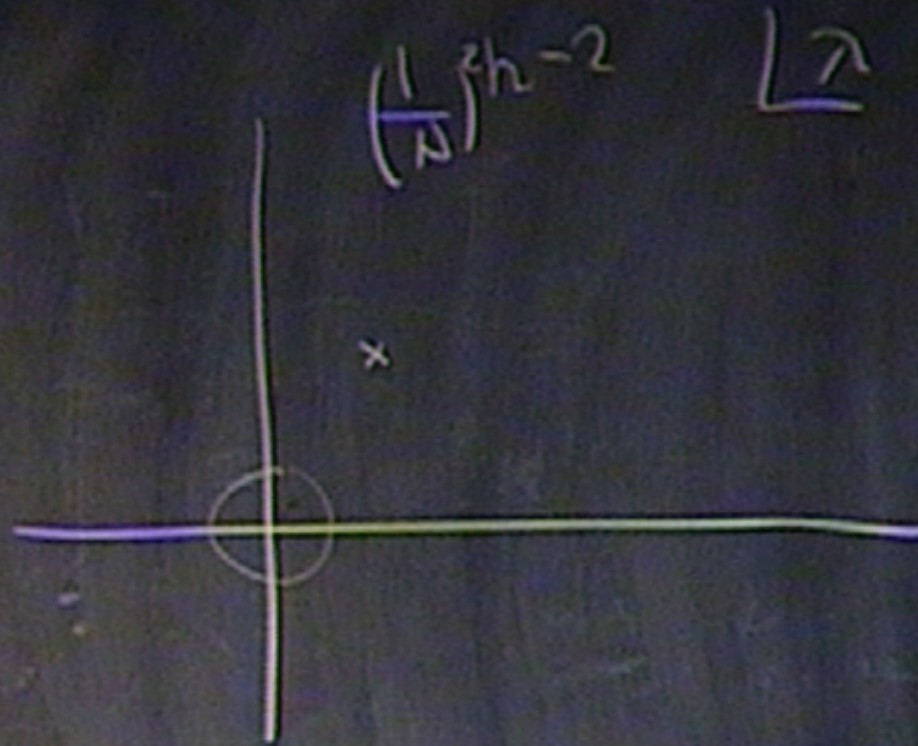


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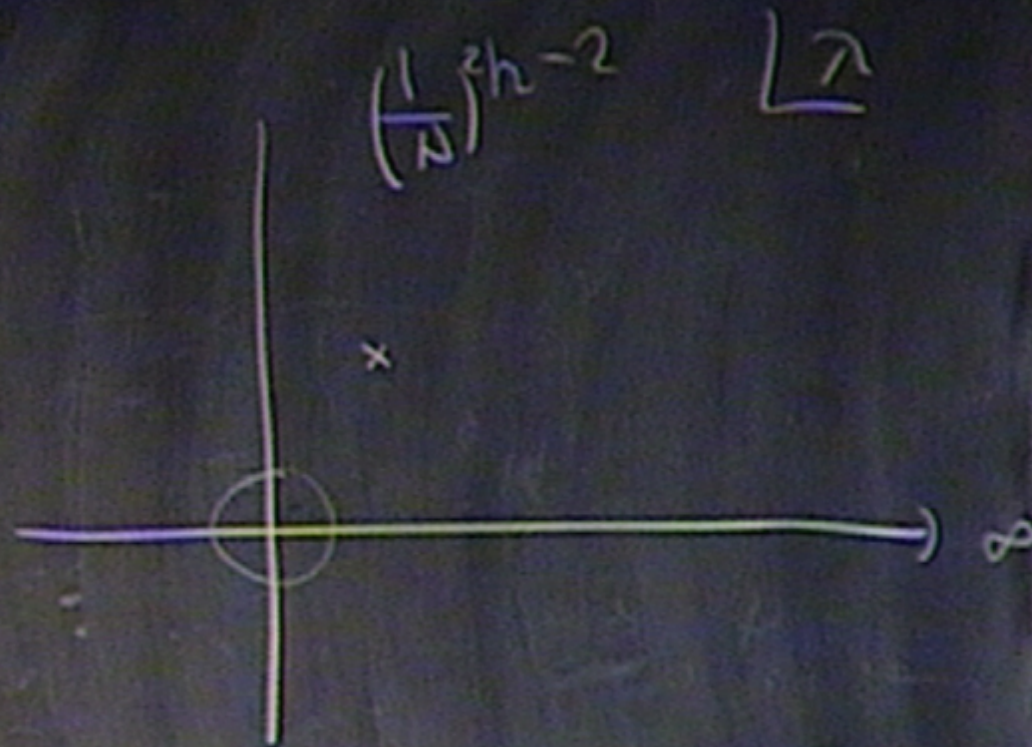
$$\left(\frac{1}{N}\right)^{2h-2} \quad \underline{L^7}$$



$O(6)$

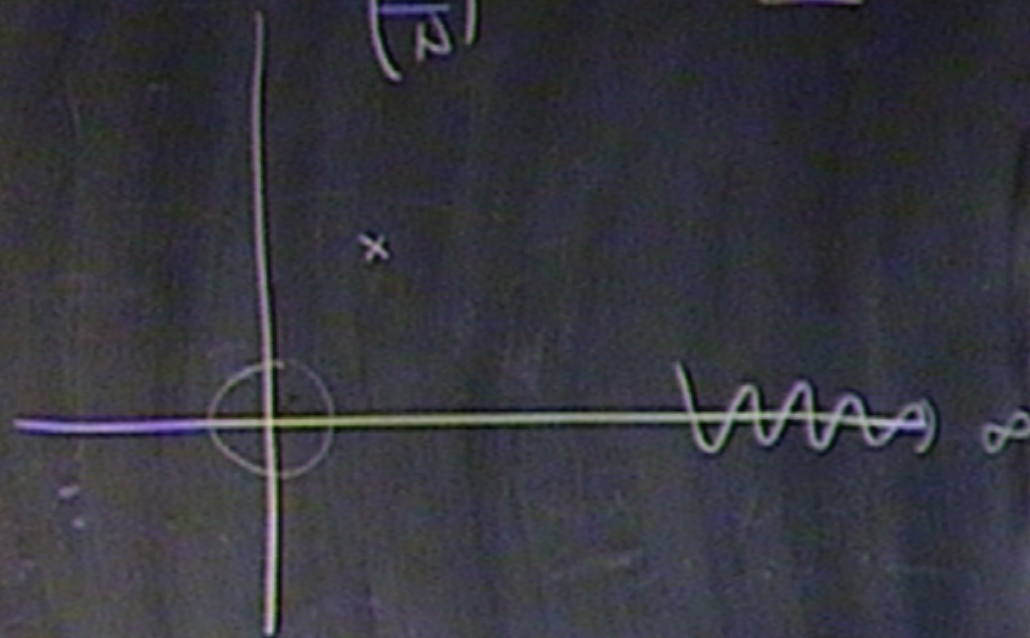


$O(6)$

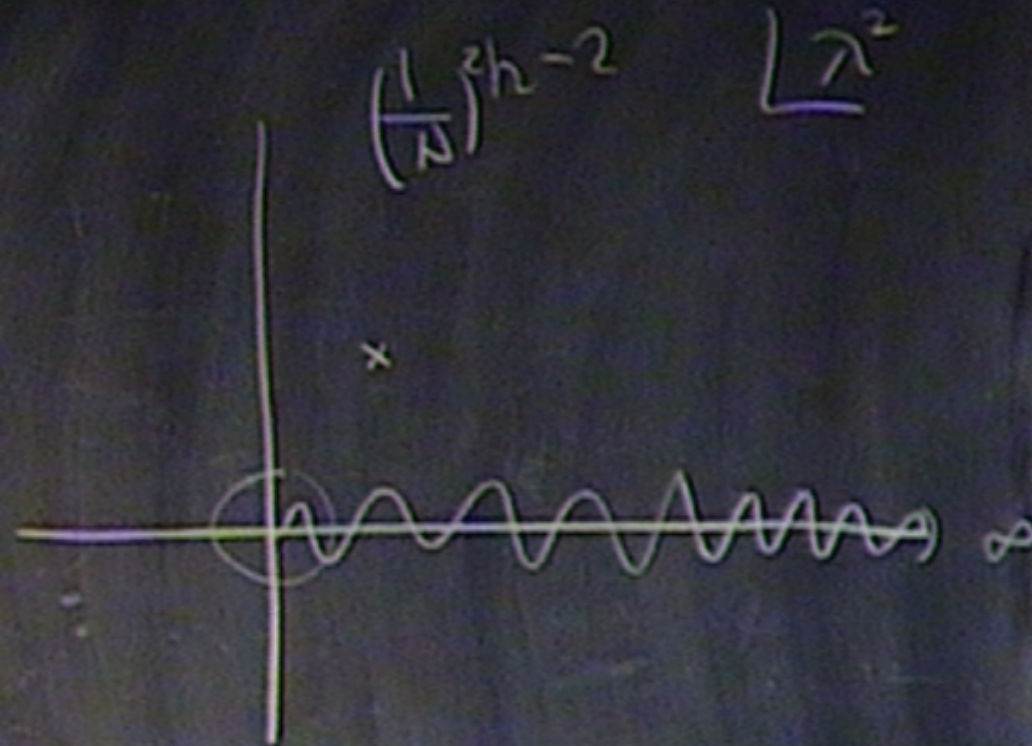


$O(6)$

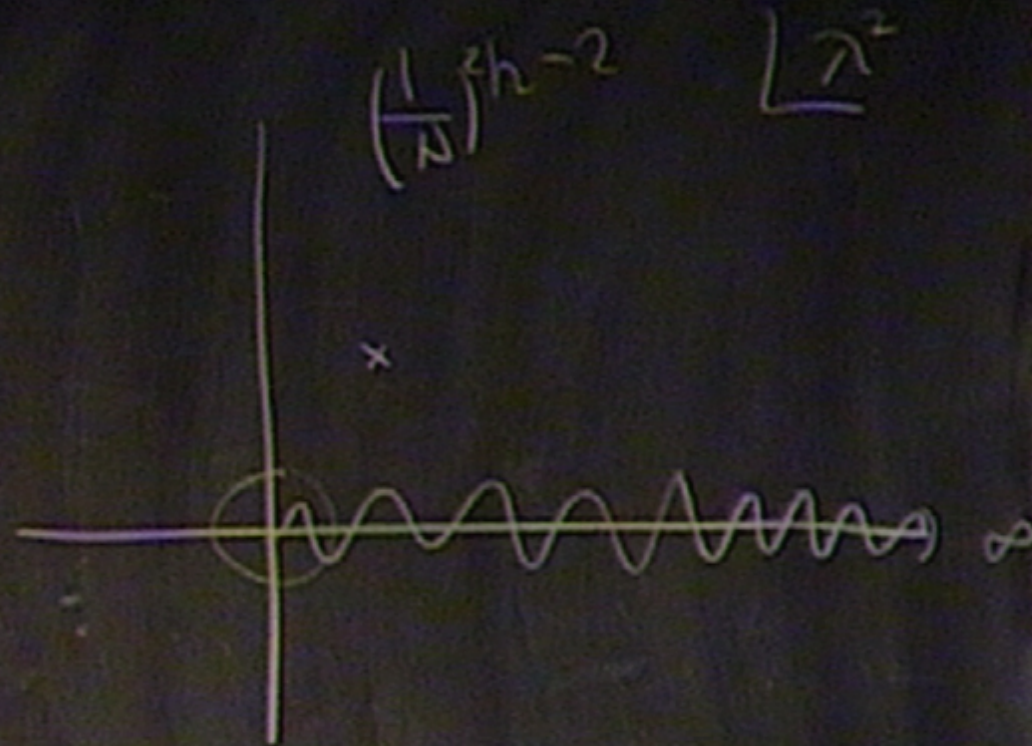
$$\left(\frac{1}{N}\right)^{2h-2} \quad \angle \lambda^2$$



$\phi(6)$



$O(6)$



The Story of Square-Root

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Baryons

- In SU(3) QCD, baryon is a color singlet state of 3 quarks $\varepsilon_{abc} Q^a Q^b Q^c$
- In SU(N) SYM, baryon is a color singlet state of N “static” quarks $\varepsilon_{ab\dots c} Q^a Q^b \dots Q^c$
- **N-ality symmetry:** $Q^a \rightarrow \omega Q^a$ ($\omega^N = 1$)
- **U(N) vs. SU(N):**
center of U(N) is U(1); center of SU(N) is Z_N
only SU(N) has baryon of N-ality

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Baryon vertex in AdS5



A diagram illustrating a baryon vertex in AdS5. It features a yellow circle with the text "D5 wrapped on S5" inside. A horizontal orange line passes through the bottom of the circle. A vertical cyan line is positioned to the right of the circle. A horizontal yellow line is located below the orange line.

D5 wrapped
on S5

N fundamental strings

On D5-brane worldvolume.....

- D5-brane couples to external B_2 and G_5 as

$$L_{D5} \sim - (dA + B_2)^2 + A \wedge G_5$$

- D5 is wrapped on S^5 , so Gauss' law for A_0 is

$$\int_{D5} \nabla \cdot E_{D5} = N \quad \text{where} \quad E_{D5} = d A_0$$

- bulk+D5-brane EOM for B_2 field

$$\nabla \cdot H_3 = \delta(D5) E_{D5} \quad \text{source for } B_2 \text{ field in the bulk}$$

- Combining the two equations, we conclude that N bundles of fundamental string ended on D5-brane

- Where are these strings from?

Baryon vertex in AdS5



D5 wrapped
on S^5

N fundamental strings

Closer look at AdS5 SUGRA

- 10d IIB supergravity action contains Chern-Simons term

$S_{CS} = \int (B_2 \wedge G_3 - C_2 \wedge H_3) \wedge G_5$ on $AdS_5 \times S^5$ yields

$$S_{5d} = \int_{AdS_5} [e^{-2\Phi} |H_3|^2 + |G_3|^2 + N (B_2 \wedge G_3 - C_2 \wedge H_3)]$$

- Inside AdS_5 , D5 in the bulk and D3-brane near infinity:

$$L_{D3} = - (dA' + B_2)^2 + A' \wedge G_3$$

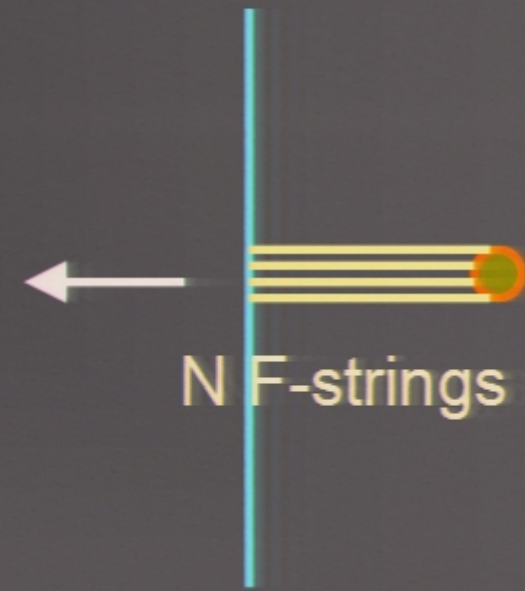
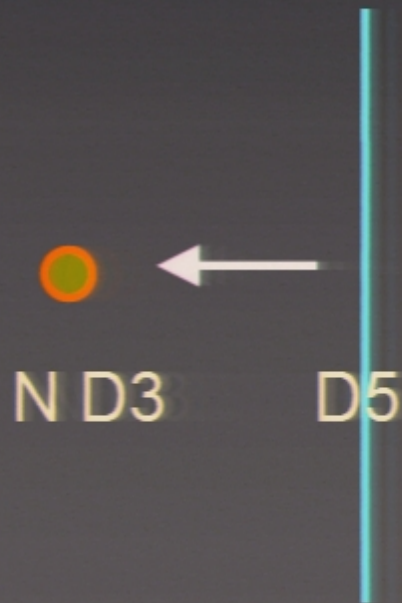
- Solve combined AdS SUGRA+D5+D3 EOM for C_2

$$\nabla \cdot (\wedge G_3 + N B_2 - E_{D3}) = 0 \quad (\text{but } B_2 \text{ sourced by D5})$$

- G_3 is trivial, so induces N unit of electric field on D3!!

8 is the magic number....

- Bring D5-brane from flat 10d infinity into N D3-branes
before crossing after crossing



Like quarks inside baryons, N F-strings obey Fermi statistics !

Intuitive picture

- Consider $SU(2) \rightarrow U(1) \rightarrow \emptyset$ gauge theory
- Add θ -term

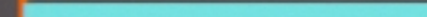
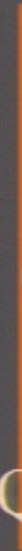
$$S_{U(1)} = \int [E^2 + B^2 + \theta E \cdot B]$$

- At nonzero θ , Gauss' law reads

$$\nabla \cdot (E + \theta B) = 0$$

- So, monopole acquires nonzero electric charge proportional to θ (Witten effect)

$$Q_E = -\theta Q_M / 2\pi$$



Monopole θ -jump

$-Q_E$

induced Q_E

- By electric charge conservation, $-Q_E$ must be deposited on θ -jump
- Between $+Q_E$ and $-Q_E$, a flux tube is formed

Baryon in AdS_5

- D5-brane wrapped on S^5 minimizes energy by falling into AdS_5 geometry
- Since S^5 size is constant inside AdS_5 , energy is constant (BPS state)
- Large worldvolume E field deforms the shape of D5-brane
- Test of finite N through N -ality property