

Title: Advanced AdS/CFT Topics

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Abstract:



Selected Topics in AdS/CFT

lecture 1

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June 2005

Topics:

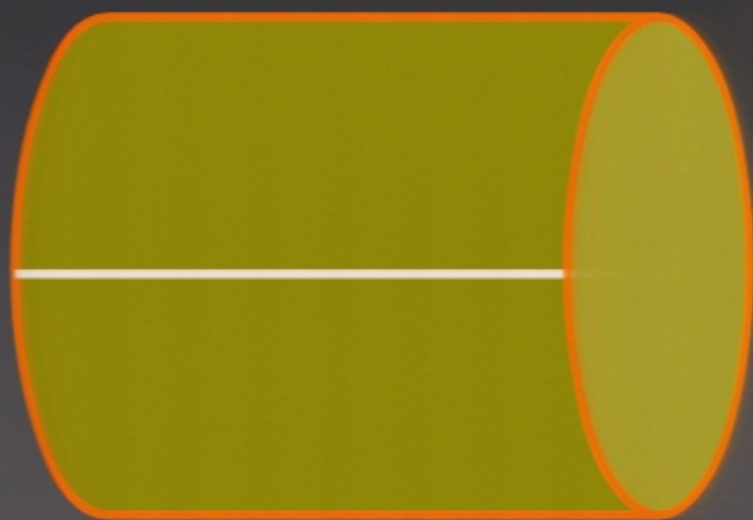
- Lightning Review of AdS5/SYM4
- Wilson loop (hep-th/9803001, 9803002)
- Polyakov loop (hep-th/9803135)
- Baryons
- Instantons
- Time-varying phenomena
- homeworks: to be distributed separately

Starting Point

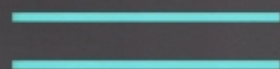
Starting Point

- string theory: open string + closed string
- low-energy limit of string theory:
 - open string $\rightarrow$ gauge theory (spin=1)
 - closed string $\rightarrow$ gravity theory (spin=2)
- channel duality: open & closed

Channel Duality



open string channel



time



closed string channel

Gauge – Gravity Duality

- keypoint:

- “gauge theory dynamics”

- (low energy limit of open string dynamics)

- is describable by

- “gravity dynamics”

- (low-energy limit of closed string dynamics)

- and vice versa

Tool Kits for AdS/CFT

- elementary excitations:
open string + closed string
- solitons:
D-branes (quantum) + NS-branes (classical)
- Mach's principle:
spacetime $\tilde{\Lambda}$ elementary + soliton sources
- “holography”:
gravity fluctuations = source fluctuations

p-Brane (gravity description)

- string effective action

$$\begin{aligned} \blacksquare S = s d^{10} \times [& e^{-2\Phi} (R_{(10)} + (d\Phi)^2 + |H_3|^2 + \dots) \\ & + \frac{1}{2} (\sum_p |G_{p+2}|^2 + \dots)] \\ & = (1/g_{st}^2)(\text{NS-NS sector}) + (\text{R-R sector}) \end{aligned}$$

- $H_3 = d B_2$, $G_{p+2} = d C_{p+1}$ etc.

- R-R sector is “quantum” of NS-NS sector

elementary and solitons

- fixed “magnetic” charge, energy-minimizing static configurations
(necessary condition for BPS states)

- NS 5-brane:

$$E = s g_{st}^{-2} [(r\Phi)^2 + H_3^2] > g_{st}^{-2} s_{S3} e^{-2\Phi} H_3$$

- F-string:

$$E = s [g_{st}^{-2} (r\Phi)^2 + g_{st}^2 K_7^2] > g_{st}^0 s_{S7} K_7$$

- Mass: $M(\text{string}) \sim 1$; $M(\text{NS5}) \sim (1/g_{st}^2)$

“elementary” ; “soliton”

p-brane: in between

- p-brane:

$$E = s g_{st}^{-2} (r \Phi)^2 + (G_{p+2})^2 > g_{st}^{-1} s e^{-\Phi} G_{p+2}$$

- $M(\text{p-brane}) \gg (1/g_{st})$

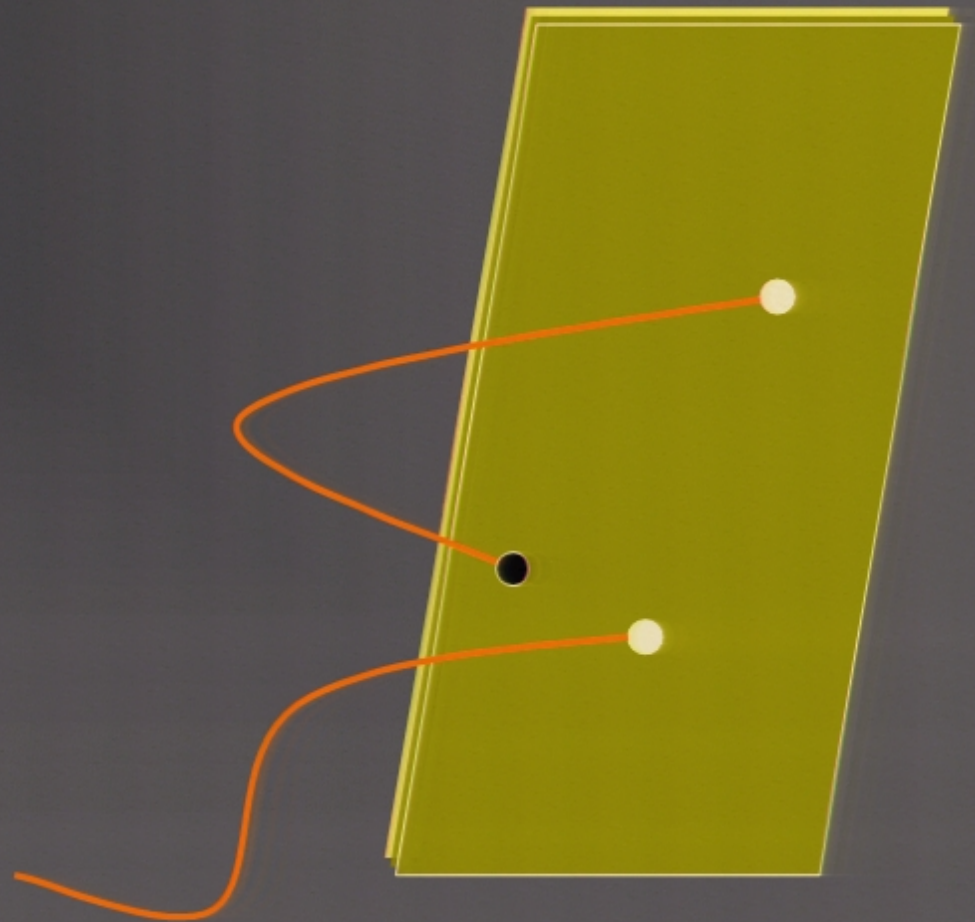
- p-brane = “quantum soliton”:

more solitonic than F-string but less solitonic than NS5-brane

- F-string $\$$ p-brane $\$$ NS5-brane

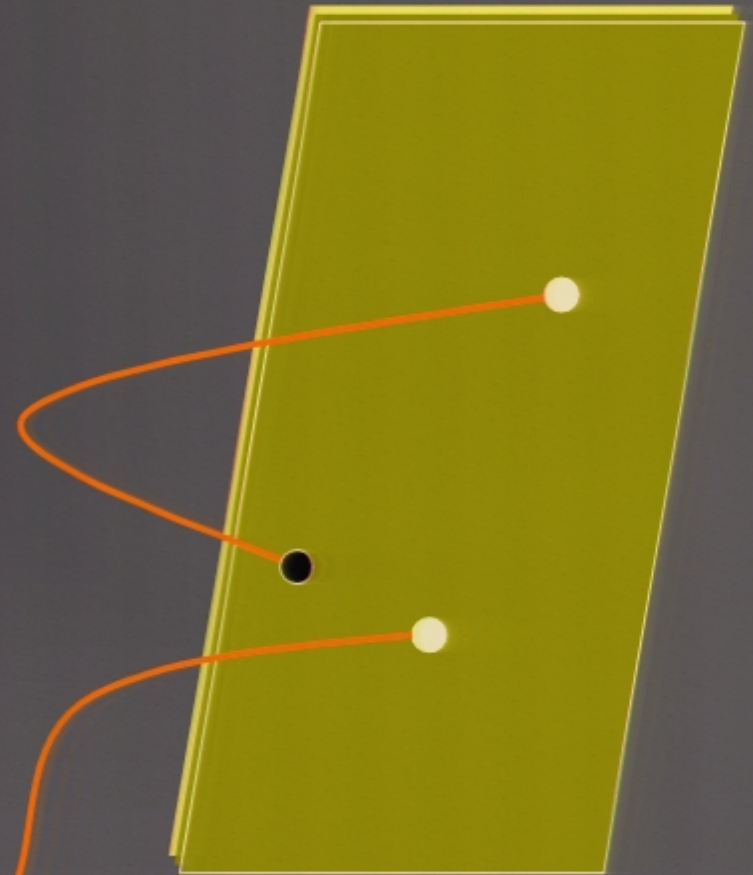
- quantum treatment of p-brane is imperative

Dp-brane (CFT description)



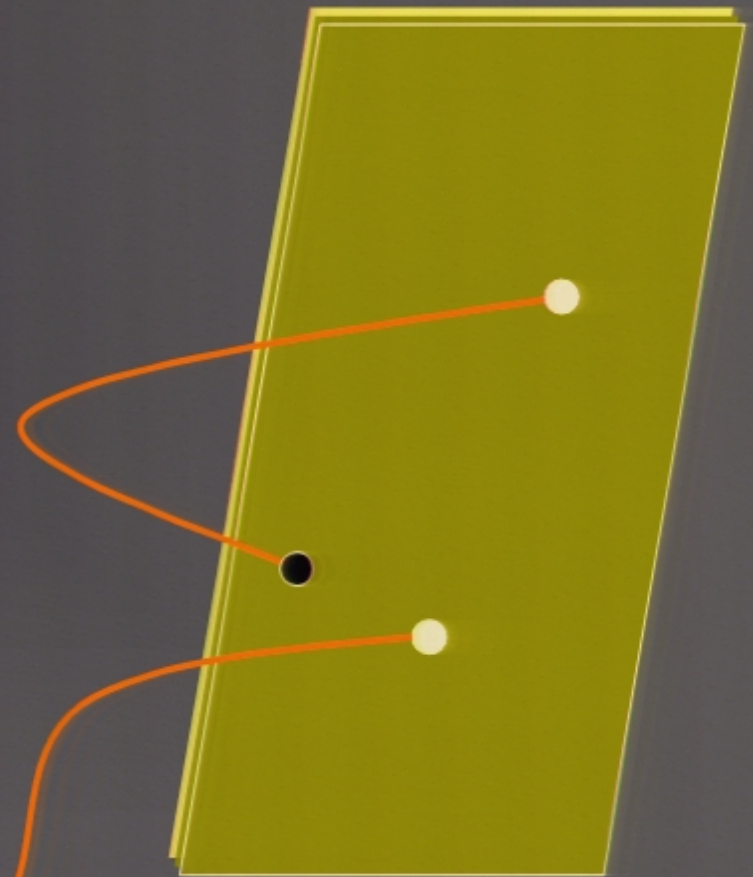
Dp-brane (CFT description)

- strings can end on it
 - open string
- string endpoints labelled by Chan-Paton factors
 - $(\quad) = (i,j)$
- mass set by disk diagram
 - $M \sim 1/g_{st}$
 - identifiable with p-brane
- $Q_L = Q_R$ (half BPS object)



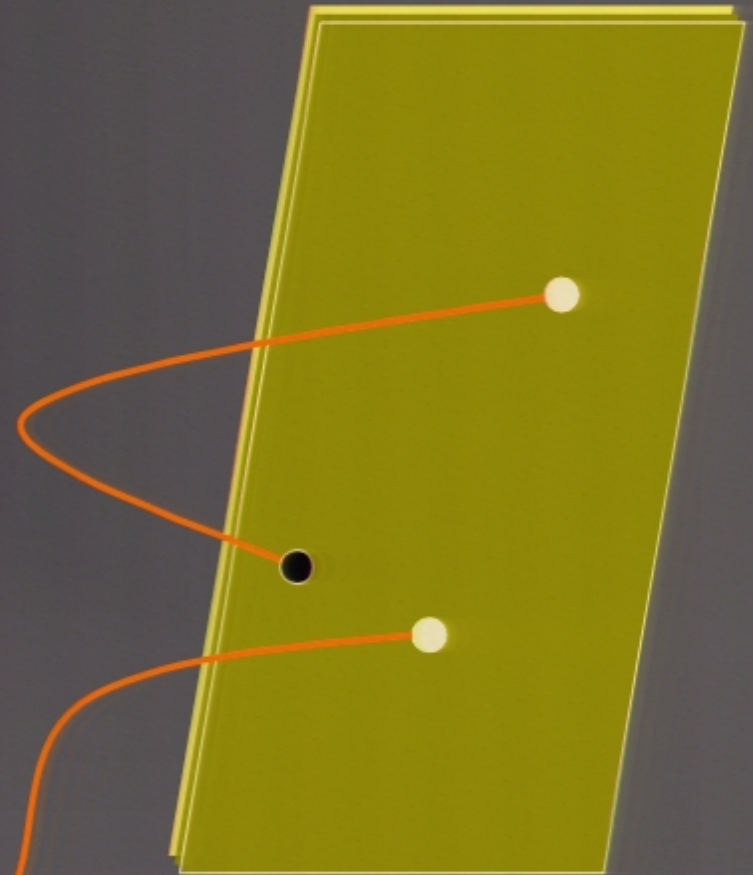
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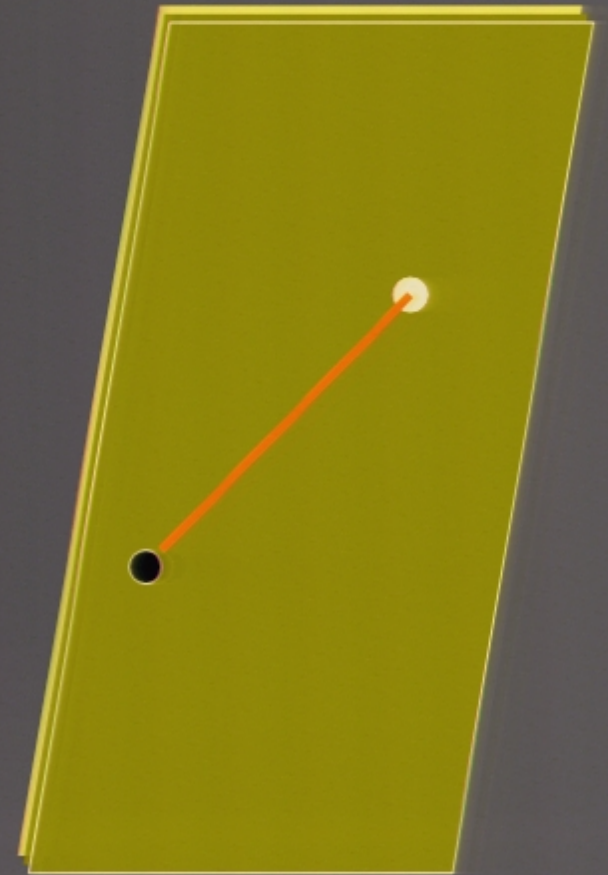


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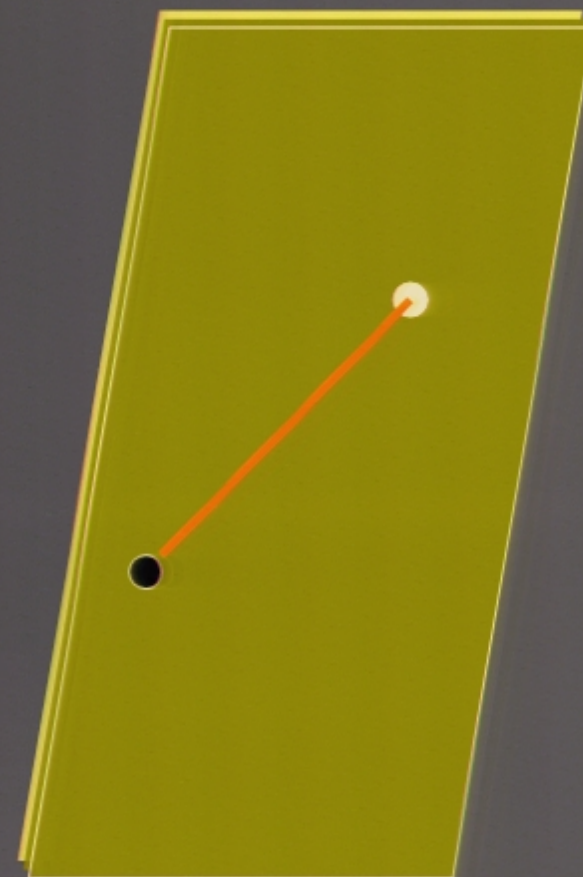
SYM_{p+1} at Low-Energy



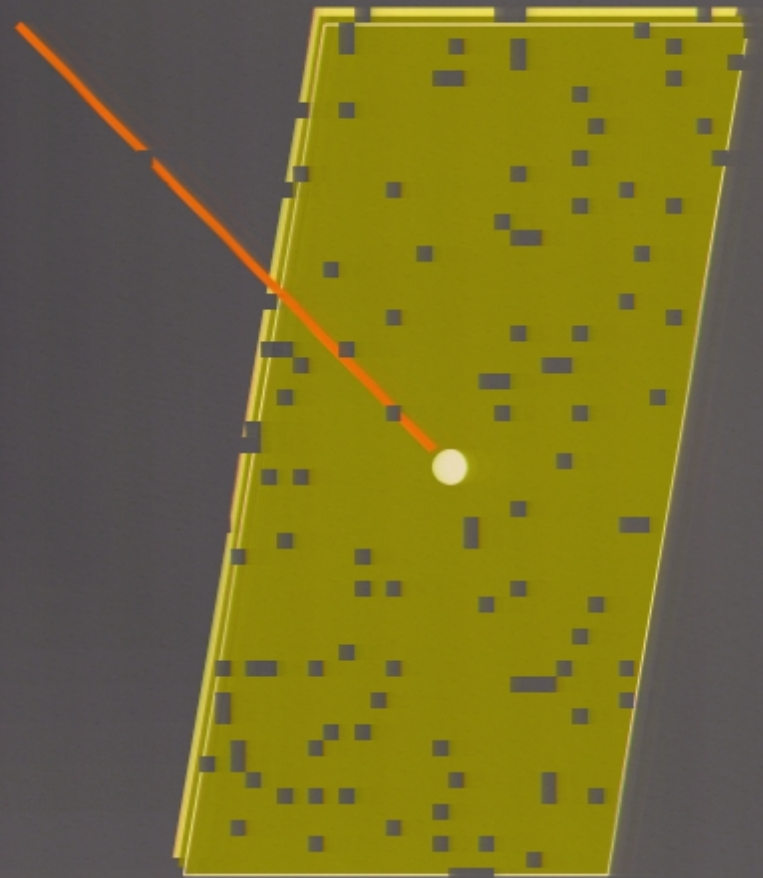
SYM_{p+1} at Low-Energy

- infinite tension limit
($\alpha' \rightarrow 0$)
- open strings $\rightarrow$ rigid rods
($Mw \sim \Delta r / \alpha'$)
- (N,N) string dynamics
 $U(N) \text{ SYM}(P+1)$
$$L = g_{\text{YM}}^{-2} \text{Tr} (F_{mn}^2 + (D_m \Phi^a)^2 + [\Phi^a, \Phi^b]^2 + \dots)$$

“||”
“?”

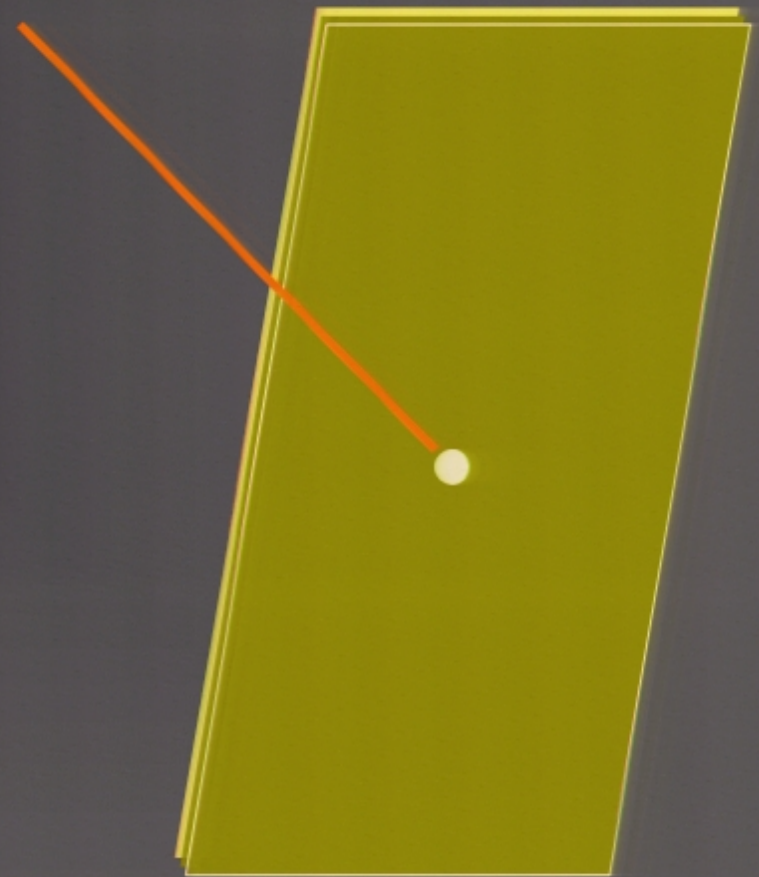


Static source: heavy quarks



Static source: heavy quarks

- semi-infinite string
($M \sim 1$)
- at rest or constant velocity
(static source)
- labelled by a single Chan-Paton factor
(i) = (i)
heavy (anti)quark
in (anti)fundamental rep.



Perturbation theory (1)

- large N conformal gauge theory :
double expansion in $1/N$ and $\lambda^2 = g_{YM}^2 N$
- $S_{YM} = (N/\lambda^2) \int d^4x \text{Tr} (F_{mn}^2 + (D_m \Phi)^2 + \dots)$
- planar expansion: $(N/\lambda^2)^{V-E} N^F = (1/N)^{2h-2} (\lambda^2)^{E-V}$
- λ^2 : nonlinear interactions
- $1/N$: quantum fluctuations (effective $\sim$)
- observable: $\sum_{h=0}^{\infty} \sum_{l=0}^{\infty} (1/N)^{2h-2} (\lambda^2)^l C_{l,h}$

Perturbation theory (2)

- closed string theory:

double expansion in g_{st} and α'

- α' : string worldsheet fluctuation

$$S_{\text{string}} = \frac{1}{2\pi\alpha'} \int d^2\sigma \partial X^\mu \bar{\partial} X_\mu$$

- g_{st} : string quantum loop fluctuation

$$S_{\text{SFT}} = \Phi * Q_B \Phi + g_{st} \Phi * \Phi * \Phi$$

$$\text{observables} = \sum_{h=0}^{\infty} \sum_{l=0}^{\infty} g_{st}^{2h-2} (2\pi\alpha' p^2)^l D_{l,h}$$

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Identifying the Two Sides....

- large- N YM and closed string theory have the same perturbation expansion structure
- $g_{\text{st}} \propto (1/N)$
- $(\alpha'/R^2) \propto 1/\lambda$ where R = characteristic scale
- Maldacena's AdS/CFT correspondence:
“near-horizon”(R) geometry of D3-brane
= large- N SYM_{3+1} at large but fixed λ
- not only perturbative level but also nonperturbatively (evidence?)

D3-Brane Geometry

- 10d SUGRA(closed string)+4d SYM (D3-brane):

$$S_{\text{total}} = S_{10\text{d}} + s \, d^4 x \, (-e^{-\Phi} \text{Tr} F_{mn}^2 + C_4 \dots)$$

- Solution

- $ds^2 = Z^{-1/2} dx_{3+1}^2 + Z^{1/2} dy_6^2$

- $G_5 = (1+\mathcal{A}) d\text{Vol}_4 \mathcal{A} dZ^{-1}$ where

$$Z = (1+R^4/r^4); \quad r = |y| \quad \text{and} \quad R^4 = 4\pi g_{\text{st}} N \alpha'^2$$

- $r \neq 1$: 10d flat spacetime

- $r \neq 0$: characteristic curvature scale = R

Identifications

- D-brane stress tensor grows with (N/g_{st})
- At large N , curvedness grows with $g_{\text{st}}^2 (N/g_{\text{st}}) = g_{\text{st}} N$
- Near D3-brane, spacetime = $\text{AdS}_5 \times S^5$
- 4d D3-brane fluct. \$ 10d spacetime fluct.
(from coupling of D3-brane to 10d fields)

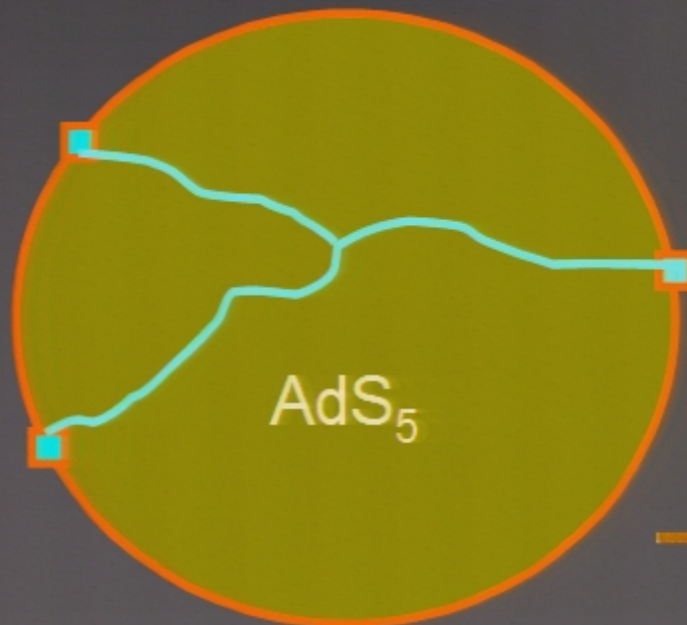
$$\text{Tr}(F_{mp}F_{np}) \quad \$ \text{ metric } g_{mn}$$

$$\text{Tr}(F_{mn}F_{mn}) \quad \$ \text{ dilaton } \Phi$$

$$\text{Tr}(F_{mp}F_{np}^*) \quad \$ \text{ Ramond-Ramond } C, C_{mn}$$

AdS/CFT correspondence

- Dirichlet problem in AdS_5 or $\text{EAdS}_5 = \text{H}_5$
- $Z_{\text{gravity}} \gg \exp(-S_{\text{AdS5}}(\phi_{\text{bulk},a}, \phi_a^1))$
 $= Z_{\text{SYM}} \gg \int [dA] \exp(-S_{\text{SYM}} - \int \sum_a \phi_a^1 O_a)$



AdS₅

- In flat 4+2 dimensional space

$$ds^2 = -dX_0^2 - dX_5^2 + \sum_{a=1}^4 dX_a^2$$

- embed hyperboloid

$$X_0^2 + X_5^2 - \sum_{a=1}^4 X_a^2 = R^2$$

- SO(4,2) inv, homogeneous, isotropic
- AdS₅ = induced geometry on hyperboloid
- <homework>

- Global coordinates:

$$ds^2 = R^2(-\cosh^2 \rho \, d\tau^2 + d\rho^2 + \sinh^2 \rho \, d\Omega_3^2)$$

$$\text{boundary} = R_t \times S_3$$

- Poincare coordinates:

$$ds^2 = R^2[r^2(-dt^2 + dx^2 + dy^2 + dz^2) + r^{-2}dr^2]$$

$$\text{boundary} = R_t \times R_3$$

Why AdS Throat = D3-Brane?

- D-brane absorption cross-section:
SUGRA computation = SYM computation



AdS₅ region in gravity description
= N D3-branes in gauge description

Another argument

- D-instantons probing (Euclidean) AdS5
- For U(N) gauge group, “homogeneous” instanton number $< N$ (otherwise inhomogeneous)
- Q D-instanton cluster in approx. flat region

$$S_{\text{Dinstanton}} = -(1/g_{\text{st}} \alpha'^2) \text{Tr}_Q [\Phi^1, \Phi^2]^2 + \dots$$



$$\text{■ } \langle \text{Tr}(\Phi^1)^2 \rangle \gg QL^2, \quad \langle \text{Tr}(\Phi^2)^2 \rangle \gg Q^2 g_{\text{st}} \alpha'^2 / L^2$$

$$\text{■ rotational symmetry ! } L^4 = Q g_{\text{st}} \alpha'^2 = N g_{\text{st}} \alpha'^2$$

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How can it be that $5d = 4d$?

- extensive quantities in 4d SYM theory scales as $[\text{length}]^4$
- Extensive quantities in 5d AdS gravity scales as $[\text{length}]^5$
- So, how can it be that quantities in 4d theory is describable by 5d theory??
- [Question] Show both area and volume of a ball of radius X in AdS_d scales as X^{d-1} !

- Answer:

Consider $(d+1)$ -dimensional AdS space. Take a finite comoving d -volume $(\text{vol})_d = \Delta t \Delta^{d-1} \mathbf{x}$.

"Volume" $V_{d+1}(r_*)$ of ball of radius $r = r_*$:

$$\begin{aligned} V_{d+1}(r_*) &= \int_0^{r_*} Z^{1/4}(r) dr \int_{\text{vol}_d} Z^{-d/4}(r) dt d^{d-1} \mathbf{x} \\ &= \frac{1}{4} r_*^4 R^{-3} \text{vol}_4 \quad \text{for } d = 4. \end{aligned}$$

"Area" $A_d(r)$ of ball of radius $r = r_*$:

$$\begin{aligned} A_d(r_*) &= \int_{r=r_*} Z^{-d/4}(r) dt d^{d-1} \mathbf{x} \\ &= r_*^4 R^{-4} \text{vol}_4 \quad \text{for } d = 4. \end{aligned}$$

Comparing the two, we find

$$\frac{1}{4} \frac{A_5(r_*)}{V_4(r_*)} = \frac{1}{R} = \text{AdS curvature radius.}$$

Notice that the ratio is independent of r_* , and holds for both $r_* \ll 1$ and $r_* \gg 1$!

Shall we test AdS/CFT?

- Recall that heavy quarks are represented by fundamental strings attached to D3-brane
- now strings are stretched and fluctuates inside AdS_5
- Let's compute interaction potential between quark and antiquark
- Do we obtain physically reasonable answers?

Static Quark Potential at Zero Temperature

$$V(r) = - (1.254\dots) \frac{\sqrt{g_{YM}^2 N}}{r}$$

Notice:

- Square-Root --- non-analyticity for λ
- exact $1/r$ --- conformal invariance

- geometry produced by D3-branes:

$$ds^2 = G_{mn}(x)dx^m dx^n = \frac{1}{\sqrt{Z}}(-dt^2 + d\mathbf{x}^2) + \sqrt{Z}dr^2$$

$$Z = 1 + \frac{R^4}{r^4} \rightarrow \frac{R^4}{r^4} \quad \text{near-horizon}$$

- string dynamics in the D3-brane background:

$$L_{\text{string}} = -\frac{1}{2\pi\alpha'} \int d\sigma \sqrt{-h} \quad (\text{Nambu-Goto})$$

$$h_{ab}(\sigma, \tau) = G_{mn}(X) \partial_a X^m \partial_b X^n.$$

- look for static string configuration:

$$\text{static gauge : } X^0 = \tau, \quad \mathbf{X} \equiv \mathbf{X}_{||} + U\hat{r}$$

- compute $h_{ab} \rightarrow$ plug back to Nambu-Goto

Notice:
 $R^4 = g_{\text{st}} N \alpha'^2$

α'^2 cancels
out!

- string Lagrangian:

$$L_{\text{string}} = -\frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2 + Z^{-1} \mathbf{X}'_{\parallel}{}^2}$$

- static-gauge constraint = energy-conservation
- eqns of motion:

$$\left(\frac{U'}{\sqrt{U'^2 + Z^{-1} \mathbf{X}'_{\parallel}{}^2}} \right)' = \mathbf{X}'_{\parallel}{}^2 (\partial_U Z^{-1})$$
$$\left(\frac{Z^{-1} \mathbf{X}'_{\parallel}}{\sqrt{U'^2 + Z^{-1} \mathbf{X}'_{\parallel}{}^2}} \right)' = 0.$$

- heavy quark

$$X^0 = \tau, \quad X_{\parallel}(\sigma) = 0, \quad \frac{1}{2\pi\alpha'} U(\sigma) = \sigma$$

- string energy = quark mass

$$E_{\text{string}} = \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2} = \Delta\sigma$$

- interpretation: radial distance = energy scale

$$r(\text{gravity}) \leftrightarrow \frac{1}{\lambda}(\text{gauge})$$

geometric UV-IR duality, holography

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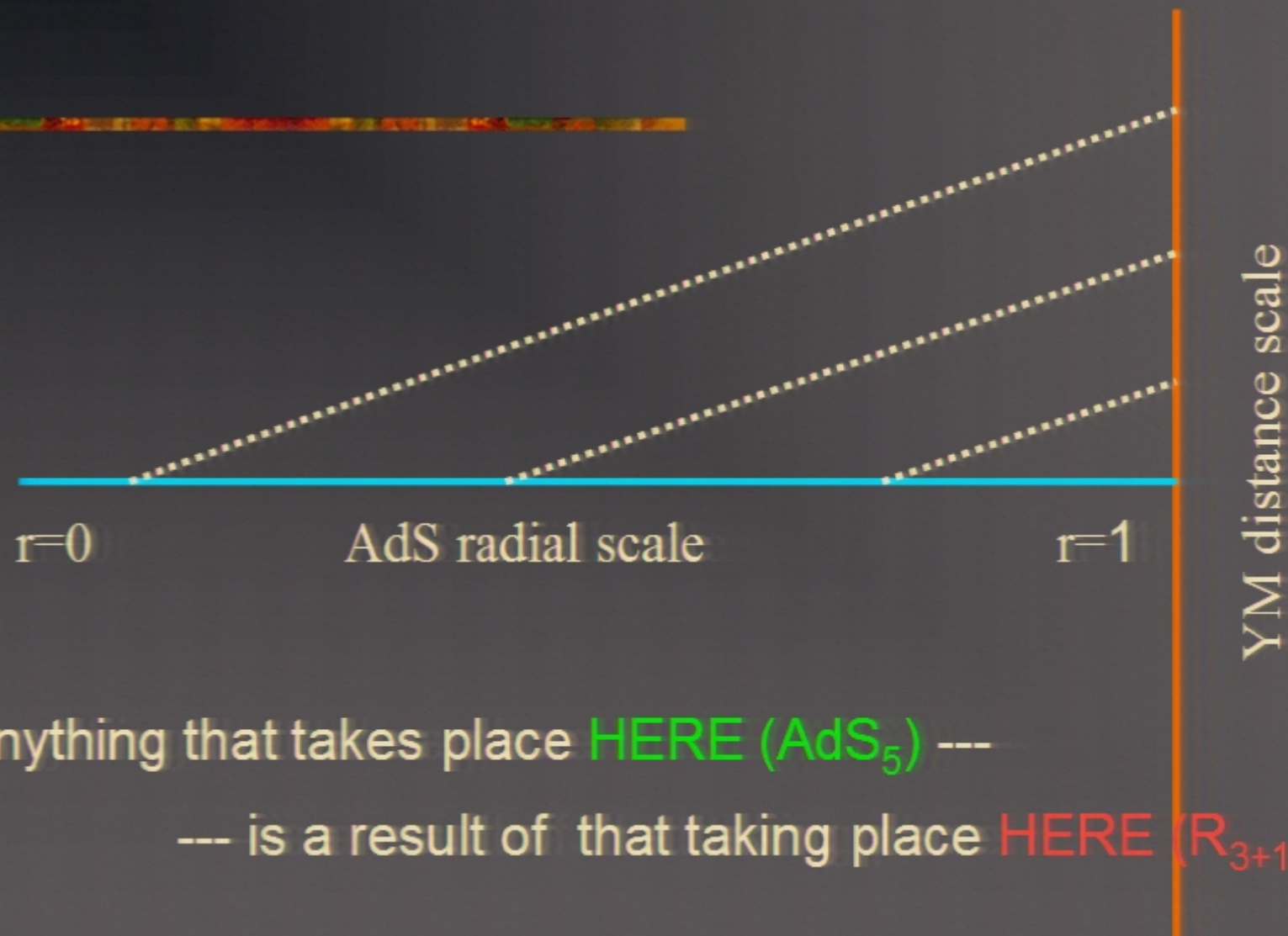
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Holography (boundary = bulk)



- heavy meson

$$X^0 = \tau, \quad \mathbf{X}_{\parallel}(\sigma) = \sigma \hat{\mathbf{r}}_3, \quad U = U(\sigma)$$

- U -eqn of motion

$$-2Z^{-1}U'' + (\partial_U Z^{-1})(2U'^2 + Z^{-1}) = 0$$

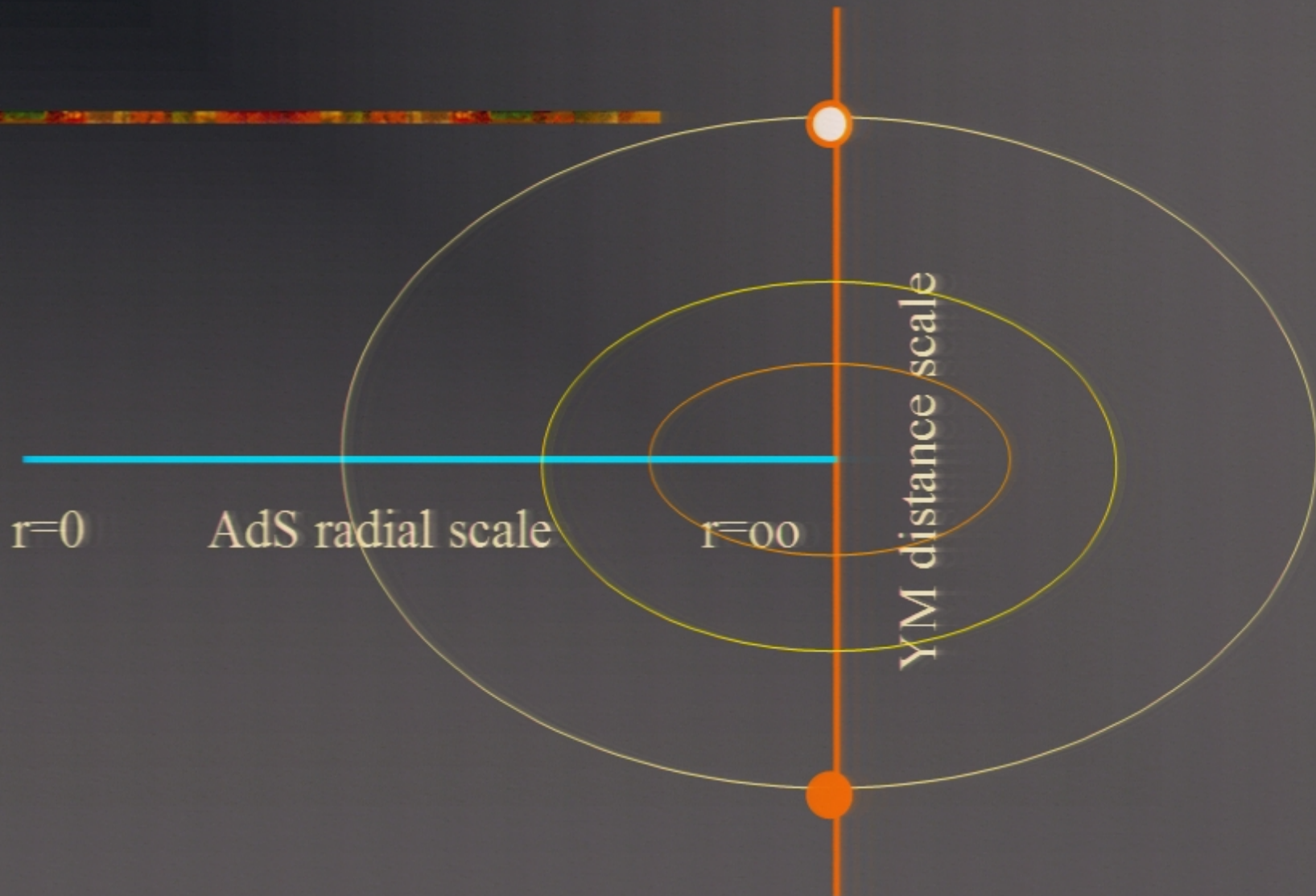
- first integral of motion:

$$Z^2 U'^2 + Z = \frac{g^2 N}{U_*^4} \quad \leftarrow \text{int. const.}$$

- solution $U(\sigma)$ via implicit function form:
easy! mechanical particle analog

$$q(t) := U_*/U(\sigma) \quad \rightarrow \text{bounded orbit}$$

Heavy Meson Configuration



- exact solution:

$$r - \frac{L}{2} = \pm \frac{\sqrt{g^2 N}}{U_*} \left[\sqrt{2} \mathbf{E} - \frac{1}{\sqrt{2}} \mathbf{F} \right] \left(\cos^{-1} \frac{U_*}{U}, \frac{1}{\sqrt{2}} \right)$$

- inter-quark distance = semi-minor axis of "closed orbit":

$$L = C \frac{\sqrt{g^2 N}}{U_*} \quad C = \frac{2\sqrt{\pi}\Gamma(3/4)}{\Gamma(1/4)} = 1.19814...$$

→ geometric UV-IR relation

- Note!

$$\begin{array}{ll} \text{non-analyticity :} & \sqrt{g^2 N} \\ \text{holography :} & L \propto 1/U_* \end{array}$$

Heavy Meson Configuration

bare quark

$r=0$

AdS radial scale

$r=\infty$

YM distance scale

bare anti-quark

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- string energy:

$$\begin{aligned} E_{(Q\bar{Q})}(U_*) &= \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2 + Z^{-1}} \\ &= \lim_{u \rightarrow \infty} 2U^* \int_0^u dx \frac{x^2}{\sqrt{x^4 - 1}} \end{aligned}$$

[used the first integral ("conserved energy")]

- heavy-quark potential:

$$\begin{aligned} V_{(Q\bar{Q})}(U_*) &\equiv E_{(Q\bar{Q})}(U_*) - (M_Q + M_{\bar{Q}}) \\ &= -2U_* \left[\frac{\sqrt{2}\pi^{3/2}}{\Gamma^2(1/4)} \right] \end{aligned}$$

- using "geometric UV-IR duality",

$$V_{Q\bar{Q}}(L) = -C \frac{\sqrt{g^2 N}}{L} \quad C = 1.254\dots$$

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bare anti-quark

- string energy:

$$\begin{aligned} E_{(Q\bar{Q})}(U_*) &= \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2 + Z^{-1}} \\ &= \lim_{u \rightarrow \infty} 2U^* \int_0^u dx \frac{x^2}{\sqrt{x^4 - 1}} \end{aligned}$$

[used the first integral ("conserved energy")]

- heavy-quark potential:

$$\begin{aligned} V_{(Q\bar{Q})}(U_*) &\equiv E_{(Q\bar{Q})}(U_*) - (M_Q + M_{\bar{Q}}) \\ &= -2U_* \left[\frac{\sqrt{2}\pi^{3/2}}{\Gamma^2(1/4)} \right] \end{aligned}$$

- using "geometric UV-IR duality",

$$V_{Q\bar{Q}}(L) = -C \frac{\sqrt{g^2 N}}{L} \quad C = 1.254\dots$$

- exact solution:

$$r - \frac{L}{2} = \pm \frac{\sqrt{g^2 N}}{U_*} \left[\sqrt{2} \mathbf{E} - \frac{1}{\sqrt{2}} \mathbf{F} \right] \left(\cos^{-1} \frac{U_*}{U}, \frac{1}{\sqrt{2}} \right)$$

- inter-quark distance = semi-minor axis of "closed orbit":

$$L = C \frac{\sqrt{g^2 N}}{U_*} \quad C = \frac{2\sqrt{\pi}\Gamma(3/4)}{\Gamma(1/4)} = 1.19814\dots$$

→ geometric UV-IR relation

- Note!

$$\begin{array}{ll} \text{non-analyticity :} & \sqrt{g^2 N} \\ \text{holography :} & L \propto 1/U_* \end{array}$$

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What have we evaluated?

- rectangular Wilson loop in N=4 SYM
- $W[C] = \text{Tr} P \exp \oint_C (i A_m dx^m + \Phi^a dy^a)$
- gauge field part = Aharonov-Bohm phase
- scalar field part = W-boson mass
- unique N=4 supersymmetric structure

- heavy meson

$$X^0 = \tau, \quad \mathbf{X}_{\parallel}(\sigma) = \sigma \hat{\mathbf{r}}_3, \quad U = U(\sigma)$$

- U -eqn of motion

$$-2Z^{-1}U'' + (\partial_U Z^{-1})(2U'^2 + Z^{-1}) = 0$$

- first integral of motion:

$$Z^2 U'^2 + Z = \frac{g^2 N}{U_*^4} \quad \leftarrow \text{int. const.}$$

- solution $U(\sigma)$ via implicit function form:
easy! mechanical particle analog

$$q(t) := U_*/U(\sigma) \quad \rightarrow \text{bounded orbit}$$

- heavy quark

$$X^0 = \tau, \quad X_{\parallel}(\sigma) = 0, \quad \frac{1}{2\pi\alpha'} U(\sigma) = \sigma$$

- string energy = quark mass

$$E_{\text{string}} = \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{U'^2} = \Delta\sigma$$

- interpretation: radial distance = energy scale

$$r(\text{gravity}) \leftrightarrow \frac{1}{\lambda}(\text{gauge})$$

geometric UV-IR duality, holography

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