

Title: Interpretation of Quantum Theory: Lecture 17

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URL: <http://pirsa.org/05030100>

Abstract:

## Consistent/Decoherent Histories

- Conceptual difficulties in QM come from introducing probabilities in the wrong way
- Histories approach: consistent introduction of probabilities eliminates difficulties and resolves (tames) paradoxes
- History of histories:
  - Griffiths 1984
  - Omnès 1987
  - Gell-Mann and Hartle 1990
  - Many subsequent papers, books
  - Griffiths, CONSISTENT QUANTUM THEORY (Cambridge 2002)

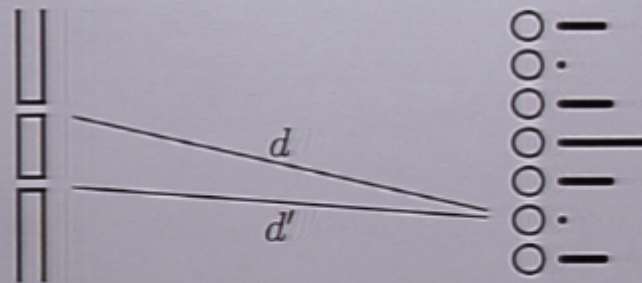


first 12 chapters at  
<http://quantum.phys.cmu.edu>

## Histories and Paradoxes

- Paradoxes that are resolved/tamed using histories:
  - Einstein, Podolsky, Rosen
  - Double slit
  - Bell, Kochen, Specker
  - Greenberger, Horne, Zeilinger
  - Hardy
  - Aharonov and Vaidman multiple box
  - Wheeler delayed choice
  - Elitzur and Vaidman noninteracting measurement
  - ...
- Paradoxes that are not resolved using histories:
  - Any suggestions?

## Double Slit I

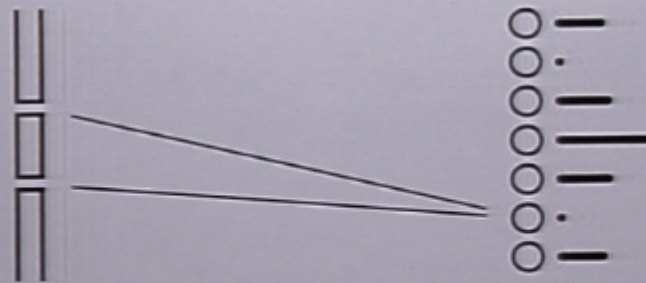


- Slit system, detectors in interference region
  - Horizontal bars: counting rates
  - Interference depends on *difference*  $d - d'$ , so
    - particles pass through slits *coherently*
  - Particles arrive *randomly* at detectors
- Consistent histories
  - Randomness an intrinsic part of nature
  - Anti-Einstein. There are no hidden variables

## Histories and Measurements

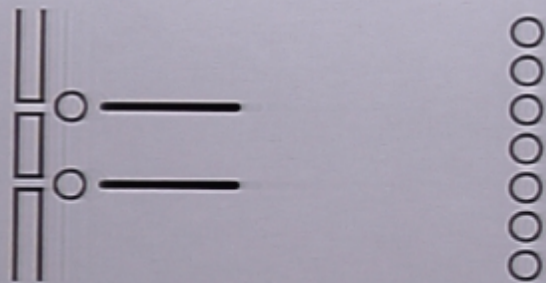
- Textbook QM:
  - Randomness arises through *measurements*
- Histories:
  - Randomness *intrinsic* in QM
  - Measurements are examples of physical processes
  - Same quantum principles govern *all* processes
  - There is *no* classical world, apparatus
  - *Sometimes* classical mechanics is a good approx
  - Quantum principles determine those circumstances

## Double Slit II



- Experimentalist:
  - Detector triggers *because* particle arrives
  - Just *before* detection particle was near detector, on its way to detector
- Historian:
  - Good experimentalists know what they're doing
  - Triggered detector indicates arrival of particle
  - QM justifies this talk; indicates its limitations
  - Theorists should not bully competent people!

### Double Slit III



- Detectors directly behind slits
  - Particles arrive at random
  - Total counting rate same as before
  - *One* detector, not both, detects each particle

- Explanations

- Experimentalist:

Particle came through slit preceding detector—

Collimators work this way

- Textbook:

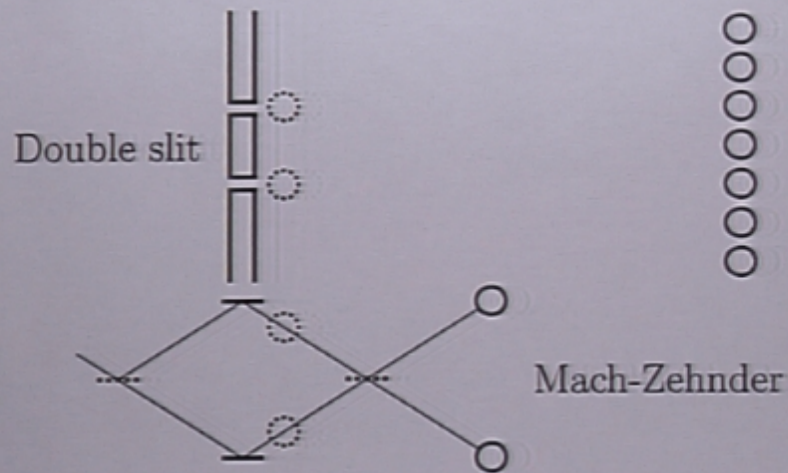
Cannot discuss what happened before measurement

“Great Smoky Dragon”

- Historian:

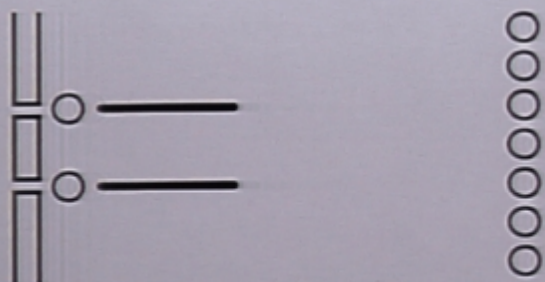
QM supports experimentalist account

## Double Slit + Mach-Zehnder



- Correspondences:
  - Which slit?  $\leftrightarrow$  Which arm?
  - Detectors behind slits  $\leftrightarrow$  inside interferometer
  - In interference zone  $\leftrightarrow$  following 2d beam splitter
- For precise description, use Mach-Zehnder
  - Basic idea applies to double slit

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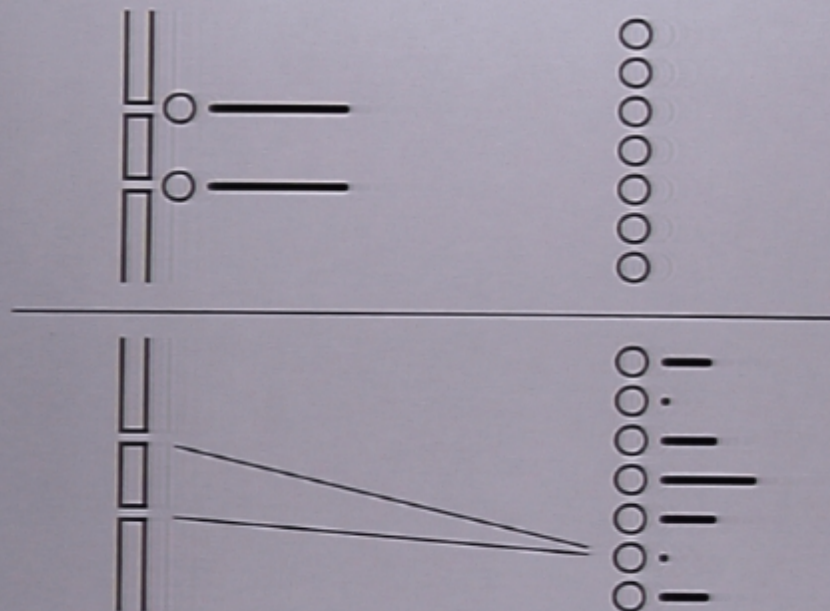
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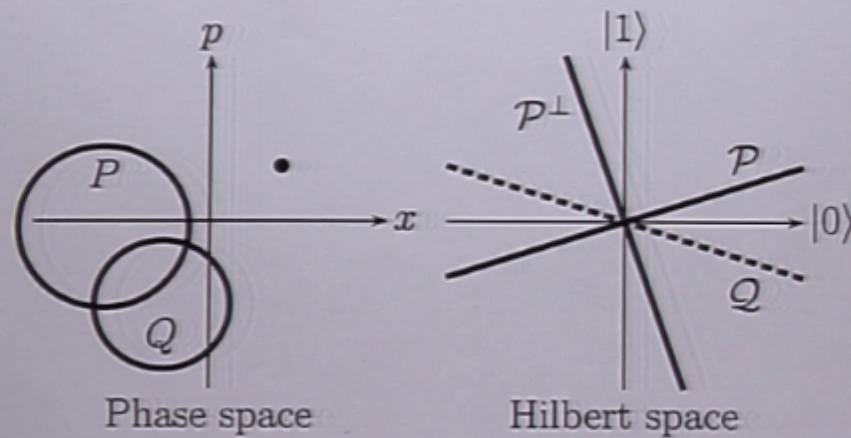
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## Double Slit IV



- Detectors behind slits removed at the last moment
  - Detectors remain:
    - particle came through definite slit
  - Detectors removed:
    - particle passed through slits coherently
- Particle could enter slit system before decision to remove detectors was made! (Wheeler delayed choice)
  - Does the future influence the past?

## Phase Space and Hilbert Space



|                | Classical       | Quantum                            |
|----------------|-----------------|------------------------------------|
| Physical state | Point           | Ray                                |
| Property $P$   | Subset $P$      | Subspace $\mathcal{P}$             |
| NOT $P$        | Compl. $\sim P$ | Orthog. compl. $\mathcal{P}^\perp$ |
| $P$ AND $Q$    | $P \cap Q$      | ?                                  |

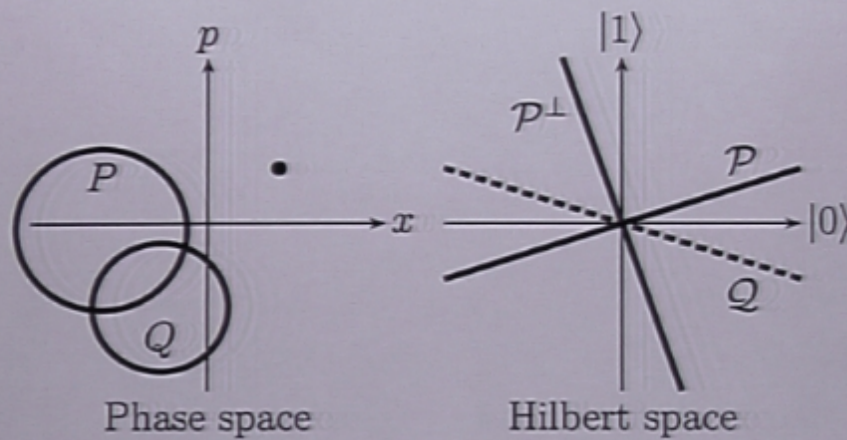
## Spin Half Particle

- $S_z = +1/2$  is a physical state
  - Ray in Hilbert space. Point on Bloch sphere
- $S_z = -1/2$  is negation of  $S_z = +1/2$ 
  - Orthogonal ray. Antipode on Bloch sphere
- For any spin-half particle,
  - Either  $S_z = +1/2$  or  $S_z = -1/2$ , not both
  - Stern-Gerlach measurement shows which is the case

---
- Nothing special about  $z$ . The  $x$  axis is just as good.
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- $S_z = +1/2$  AND  $S_x = +1/2$  is *meaningless*:
  - Hilbert space QM assigns it no meaning
    - No corresponding ray in the Hilbert space
  - No experiment which can measure it
    - Because it is not there!

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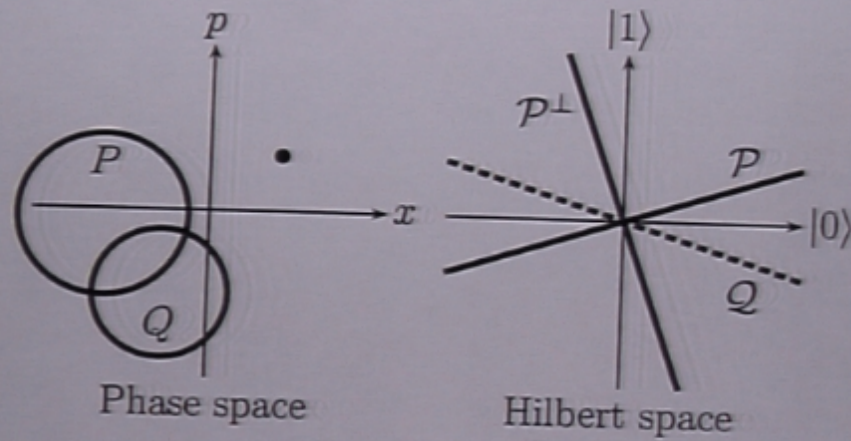
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## Logic of Quantum Properties

- False statement is one whose negation is true
  - “Pennsylvania is a Canadian province”
- Meaningless statement: not formed according to rules governing proper use of the language
  - Example: “ $P \wedge \vee Q$ ”
  - Negation of meaningless statement is meaningless
- Classical physical system:
  - Meaningful to combine two properties with AND
  - “The position is...” AND “The momentum is...”
- Quantum physical system:
  - Use AND only with *compatible* properties
  - Compatible: projectors commute:  $PQ = QP$
  - $S_z = +1/2$ ,  $S_x = +1/2$  are *incompatible*

## Quantum Logic

- George Birkhoff and John von Neumann, Ann. Math. 37 (1936) 823, "The Logic of Quantum Mechanics"
  - $S_z = +1/2$  AND  $S_x = +1/2$  is *meaningful, false*
  - Must modify rules of logic:
$$A \vee (B \wedge C) \neq (A \vee B) \wedge (A \vee C)$$
$$A \wedge (B \vee C) \neq (A \wedge B) \vee (A \wedge C)$$
- Consistent histories recognizes logical problem
  - But solves it in a different way
  - $S_z = +1/2$  AND  $S_x = +1/2$  is *meaningless*
- Rules of logic remain unchanged, but one must
  - Recognize and exclude meaningless statements
- Single framework rule:
  - Meaningful quantum descriptions use a *single collection* of mutually compatible properties
  - Incompatible descriptions cannot be combined!
- Spin half
  - Can discuss  $S_z$ , which is  $+1/2$  or  $-1/2$
  - Can discuss  $S_x$ , which is  $+1/2$  or  $-1/2$
  - Cannot *combine* these discussions
    - Doing so makes no sense in Hilbert space QM

## Probabilities I

- Standard (textbook) probability theory:  $(\mathcal{S}, \mathcal{E}, \text{Pr})$
  - Sample space  $\mathcal{S}$  of *mutually-exclusive possibilities*
    - One and only one occurs in a given experiment
    - Examples:
      - $\{H, T\}$  for coin toss
      - $\{1, 2, 3, 4, 5, 6\}$  for roll of die
  - Event algebra  $\mathcal{E}$ .
    - Assume  $\mathcal{S}$  discrete;  $\mathcal{E} =$  all subsets of  $\mathcal{S}$
  - Probability distribution  $\text{Pr}$ 
    - To each  $s_i$  in  $\mathcal{S}$  assign  $p_i = \text{Pr}(s_i) \geq 0$ ;  $\sum_i p_i = 1$ .
- 
- Quantum mechanics: three options for probabilities
    - (i) Use standard theory; (ii) Invent new one;
    - (iii) Become confused (very popular option)
  - Consistent histories uses *standard* probability theory
    - There are two tasks:
      - Define quantum sample space  $\mathcal{S}$
      - Introduce probabilities  $\text{Pr}$

## Probabilities II

- Example of spin half
- $S_z = +1/2, -1/2$  are mutually exclusive possibilities
  - If one is true, the other is false
  - One, only one occurs in Stern-Gerlach experiment
  - They constitute the  $S_z$  sample space
- Likewise,  $S_x = +1/2, -1/2$  constitute  $S_x$  sample space
- $S_z$  and  $S_x$  sample spaces are *incompatible*
  - Events cannot be combined
  - Probabilistic inference cannot be combined
- WARNING!
  - *Incompatible* is a *quantum* concept
  - *Mutually Exclusive* is classical or quantum
  - Do not confuse the two!

## Probabilities III

- General structure of quantum sample spaces
- Decomposition of the identity in projectors  $\{P^j\}$ 
  - (Superscript is label, not power)
  - $P^j = (P^j)^\dagger$ ,  $P^j P^k = \delta_{jk} P^j$ ,  $I = \sum_j P^j$
  - Each  $P^j \leftrightarrow$  physical property (Hilbert subspace)
  - $j \neq k \Rightarrow P^j P^k = 0$ : mutually exclusive properties
  - $\sum_j P^j = 1$ : at least one property is true.
- Event algebra  $\mathcal{E}$  consists of all projectors of type
$$P = \sum_j \pi_j P^j, \quad \pi_j = 0 \text{ or } 1$$
- Example: Orthonormal basis  $\{|\phi^j\rangle\}$ ;  $P^j = |\phi^j\rangle\langle\phi^j|$ .
- $S_z$  sample space for spin half:  $I = [z^+] + [z^-]$ 
  - Use  $[\psi]$  as abbreviation for dyad  $|\psi\rangle\langle\psi|$ .
- Before discussing quantum probabilities, make sure sample space exists! Many quantum paradoxes and other confusion can be traced to nonexistent sample spaces!

## Born Rule I

- Time development of *closed* or *isolated* physical system
  - Open system: make it part of larger closed system
  - Use Schrödinger Eqn to compute probabilities
  - Born rule is first (but not last!) step
- Unitary time development operator  $T(t, t')$ 
  - Comes from solving Schrödinger's equation
  - Time-independent  $H$ :  $T(t, t') = e^{-i(t-t')H/\hbar}$
- Assume  $|\psi_0\rangle$  at  $t_0$ 
  - Sample space  $\mathcal{S}$ : basis  $\{|\phi_1^k\rangle\}$  at  $t_1$
- Born probabilities:

$$\Pr(\phi_1^k) = \Pr(\phi_1^k | \psi_0) = |\langle \phi_1^k | T(t_1, t_0) | \psi_0 \rangle|^2$$

- $\Pr(\phi_1^k)$  = prob of  $[\phi_1^k]$ , *not measurement* of  $\phi_1^k$ .
  - Good measurements reveal pre-existing properties.
  - Use *quantum* description *including* apparatus to discuss measurements

## Born Rule II

- Born probabilities *depend on basis*  $\{|\phi_1^k\rangle\}$
- Example. Spin half,  $|\psi_0\rangle = |z^+\rangle$ ,  $H = 0$ ,  $T(t, t') = I$
- $S_z$  basis  $\{|z^+\rangle, |z^-\rangle\}$  at  $t_1$ :
  - $\text{Pr}(z_1^+) = 1$ ,  $\text{Pr}(z_1^-) = 0$
  - Subscript 1 indicates time  $t_1$ .
- $S_x$  basis  $\{|x^+\rangle, |x^-\rangle\}$  at  $t_1$ :
  - $\text{Pr}(x_1^+) = 1/2 = \text{Pr}(x_1^-)$
- Probabilities refer to properties of particle!
  - Bases incompatible; *cannot* assign probability to  $S_z = +1/2$  AND  $S_x = -1/2$  at time  $t_1$
- “Gyroscope with axis in  $z$  direction” is misleading
  - $S_z = +1/2$  at  $t_0$ ,  $S_x = -1/2$  at  $t_1$  does *not* mean change in direction of axis!
  - Better picture: gyroscope axis in random direction
  - Given  $z$  component at  $t_0$ , what is probability of  $x$  component at  $t_1$ ?

## Pre-Probability $|\psi_t\rangle$

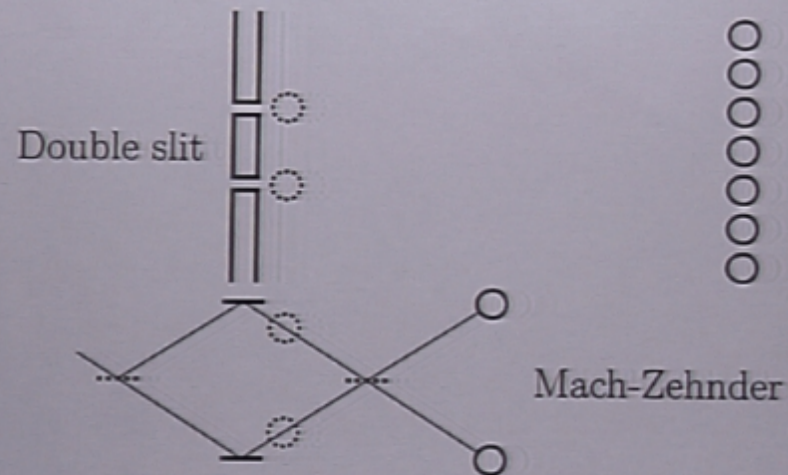
- Born probability

$$\text{Pr}(\phi_1^k) = \text{Pr}(\phi_1^k | \psi_0) = |\langle \phi_1^k | T(t_1, t_0) | \psi_0 \rangle|^2$$

can be calculated in different ways.

- 1. Integrate Schrödinger Eqn from  $t_0$  to  $t_1$ 
  - $|\psi_1\rangle = T(t_1, t_0) |\psi_0\rangle$
  - $\text{Pr}(\phi_1^k | \psi_0) = |\langle \phi_1^k | \psi_1 \rangle|^2$
- 2. Integrate Schrödinger Eqn from  $t_1$  to  $t_0$ 
  - $|\phi_0^k\rangle = T(t_0, t_1) |\phi_1^k\rangle$
  - $\text{Pr}(\phi_1^k | \psi_0) = |\langle \phi_0^k | \psi_0 \rangle|^2$
- Approaches 1 and 2 equally good
  - Compare E&M: same result using different gauge
- Physical reality:  $|\psi_0\rangle$  and the  $\{|\phi_1^k\rangle\}$ ;  
however,  $|\psi_1\rangle$  and  $\{|\phi_0^k\rangle\}$  are *pre-probabilities*:  
tools for computing probabilities, *not* physical reality!
- “Wave function of universe”  $|\psi_t\rangle = T(t, t_0) |\psi_0\rangle$ 
  - Everett:  $|\psi_t\rangle$  represents physical reality
  - Histories:  $|\psi_t\rangle$  is pre-probability: useful for finding Born probabilities; inadequate for others

## Double Slit + Mach-Zehnder



- Correspondences:
  - Which slit?  $\leftrightarrow$  Which arm?
  - Detectors behind slits  $\leftrightarrow$  inside interferometer
  - In interference zone  $\leftrightarrow$  following 2d beam splitter
- For precise description, use Mach-Zehnder
  - Basic idea applies to double slit