

Title: Entanglement, Critical Phenomena & RG Flows

Date: Jan 19, 2005 04:00 PM

URL: <http://pirsa.org/05010010>

Abstract:

Entanglement, Critical Phenomena & RG Flows

Perimeter Institute

19/01/05

Enrique Rico Ortega

@

Universitat Barcelona



Entanglement, Critical Phenomena & RG Flows

Perimeter Institute

19/01/05

Enrique Rico Ortega

@

Universitat Barcelona



UNIVERSITAT DE BARCELONA



Entanglement, Critical Phenomena & RG Flows

- Entanglement in Quantum Critical Phenomena.

G. Vidal, J.I. Latorre, E. Rico, A. Kitaev. Phys.Rev.Lett. 90 (2003) 227902

- Ground State Entanglement in Quantum Spin Chains.

J.I. Latorre, E. Rico, G. Vidal. Quant. Inf. & Comp. Vol.4 no.1 (2004) pp.048-092

- Fine-Grained Entanglement Loss along Renormalization Group Flows.

J.I. Latorre, C.A. Lütken, E. Rico, G. Vidal. quant-ph/0404120

- Renormalization Group Transformation on quantum states.

F. Verstraete, J.I. Cirac, J.I. Latorre, E. Rico, M.M. Wolf, quant-ph/0410227



Pirsa: 05010010

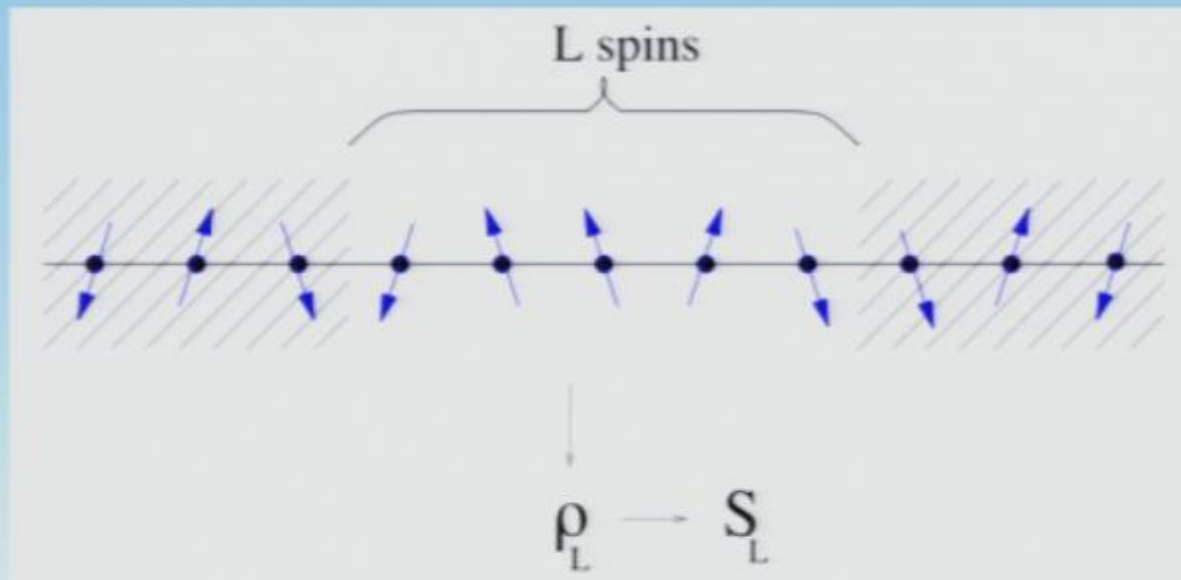


Page 4/55

Outline

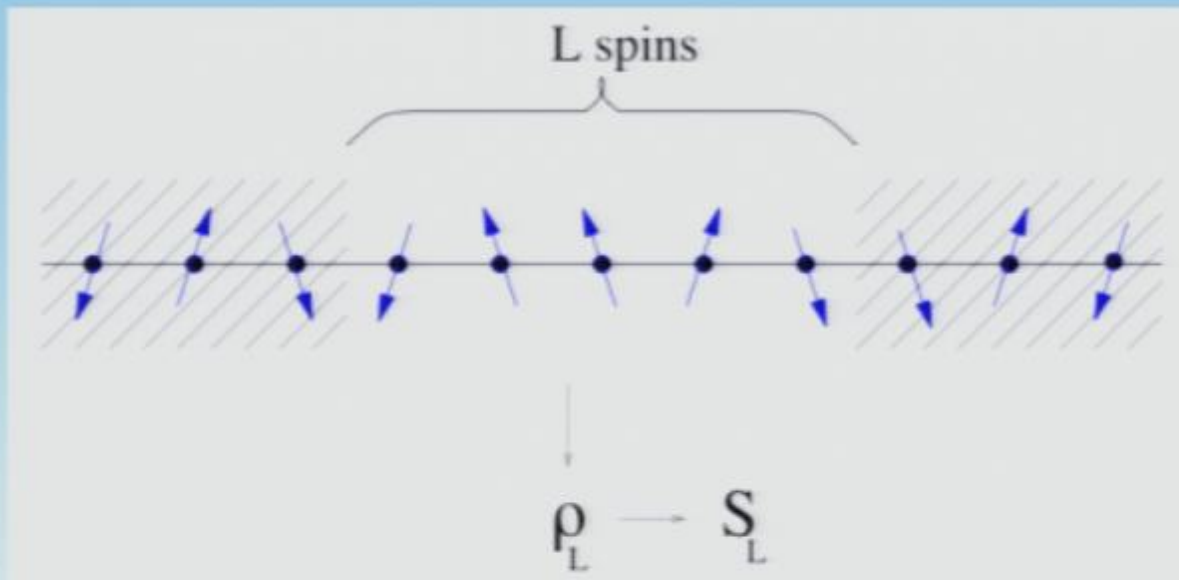
- Introduction (Motivation & Entanglement measures)
- Entanglement in Quantum Spin Models
 - Heisenberg and XY model
 - Analysis of the Ground State Entanglement
 - Connection with higher Dimension
 - Connection with Conformal Field Theory
- Entanglement loss along RG Flows
- Matrix Product States
 - Definition
 - Fixed points in Scale Transformation

Introduction & Motivation



$$\langle \mathbf{O}_i \cdot \mathbf{O}_j \rangle - \langle \mathbf{O}_i \rangle \langle \mathbf{O}_j \rangle \rightarrow \begin{cases} e^{-|i-j|/\xi} \Rightarrow \text{Finite Correlation} \\ \frac{1}{|i-j|^{\nu}} \Rightarrow \text{Phase Transition} \end{cases}$$

Introduction & Motivation



$$\langle \mathbf{O}_i \cdot \mathbf{O}_j \rangle - \langle \mathbf{O}_i \rangle \langle \mathbf{O}_j \rangle \rightarrow \begin{cases} e^{-|i-j|/\xi} \Rightarrow \text{Finite Correlation} \\ \frac{1}{|i-j|^v} \Rightarrow \text{Phase Transition} \end{cases}$$

$$\rho_L = \text{Tr}_{N-L}(|\Psi_G\rangle\langle\Psi_G|) \Rightarrow S_L = -\text{Tr}(\rho_L \log_2 \rho_L)$$

$$\rho_L(\lambda, \gamma) \succ \rho_{L'}(\lambda', \gamma')$$

Introduction & Entanglement measures

Schmidt Rank.- $|\psi\rangle = \sum_{i=1}^D \sqrt{\lambda_i} |i_A\rangle |i_B\rangle$ $\rho_A = \text{tr}_B(|\psi\rangle\langle\psi|) = \sum_{i=1}^D \lambda_i |i_A\rangle\langle i_A|$

Introduction & Entanglement measures

Schmidt Rank.- $|\psi\rangle = \sum_{i=1}^D \sqrt{\lambda_i} |i_A\rangle |i_B\rangle$ $\rho_A = \text{tr}_B(|\psi\rangle\langle\psi|) = \sum_{i=1}^D \lambda_i |i_A\rangle\langle i_A|$

Entropy of Entanglement.- $E(|\psi\rangle) = S(\rho_A) = -\sum_{i=1}^D \lambda_i \log_2 \lambda_i$

$$|\psi\rangle^{\otimes N} \xrightarrow{LOCC} |EPR\rangle^{\otimes M} \Leftrightarrow E(|\psi\rangle) = \frac{M}{N} \quad \text{Bennett et al (1996)}$$

Introduction & Entanglement measures

Schmidt Rank. - $|\psi\rangle = \sum_{i=1}^D \sqrt{\lambda_i} |i_A\rangle |i_B\rangle$ $\rho_A = \text{tr}_B(|\psi\rangle\langle\psi|) = \sum_{i=1}^D \lambda_i |i_A\rangle\langle i_A|$

Entropy of Entanglement. - $E(|\psi\rangle) = S(\rho_A) = -\sum_{i=1}^D \lambda_i \log_2 \lambda_i$

$$|\psi\rangle^{\otimes N} \xrightarrow{LOCC} |EPR\rangle^{\otimes M} \Leftrightarrow E(|\psi\rangle) = \frac{M}{N} \quad \text{Bennett et al (1996)}$$

Majorization. - $|\psi\rangle \xrightarrow{LOCC} |\tilde{\psi}\rangle \Leftrightarrow |\psi\rangle \prec |\tilde{\psi}\rangle$

$$\lambda_1^\downarrow \leq \tilde{\lambda}_1^\downarrow$$

Nielsen, Vidal (2000)

$$\lambda_1^\downarrow + \lambda_2^\downarrow \leq \tilde{\lambda}_1^\downarrow + \tilde{\lambda}_2^\downarrow$$

$\vdots$

$$\lambda_1^\downarrow + \lambda_2^\downarrow + \dots + \lambda_n^\downarrow = \tilde{\lambda}_1^\downarrow + \tilde{\lambda}_2^\downarrow + \dots + \tilde{\lambda}_n^\downarrow$$

Quantum Spin Models

Heisenberg model

$$H_{\text{XXZ}} = \sum_n \left[\frac{1}{2} (\sigma_n^x \sigma_{n+1}^x + \sigma_n^y \sigma_{n+1}^y + \Delta \sigma_n^z \sigma_{n+1}^z) - \lambda \sigma_n^z \right] = H_0 + \Delta H_1 + \lambda H_2$$

Quantum Spin Models

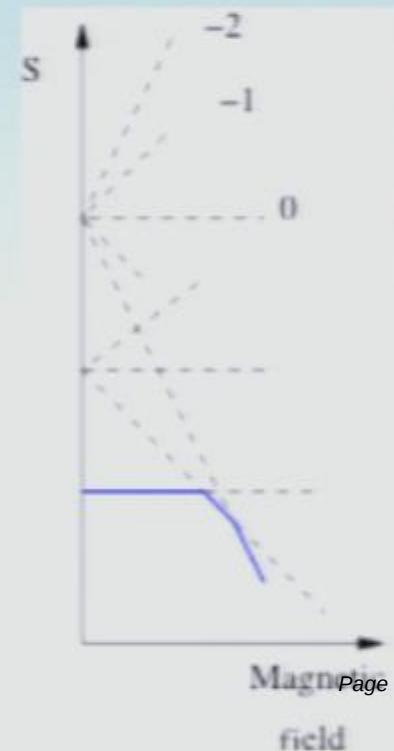
Heisenberg model

$$H_{XXZ} = \sum_n \left[\frac{1}{2} (\sigma_n^x \sigma_{n+1}^x + \sigma_n^y \sigma_{n+1}^y) + \Delta \sigma_n^z \sigma_{n+1}^z \right] - \lambda \sigma_n^z = H_0 + \Delta H_1 + \lambda H_2$$

Solution of the model.- **Bethe Ansatz**

- Rotational symmetry
about the z-axis in the spins space
- Translational invariance
by any number of lattice spacing

$$|F\rangle = |\uparrow\uparrow\uparrow\cdots\uparrow\uparrow\rangle \Rightarrow |\psi\rangle = \sum_n e^{ikn} \sigma_n^- |F\rangle$$



Quantum Spin Models

XY model

$$H_{XY} = \sum_{l=1}^N \left(\frac{1+\gamma}{2} \sigma_l^x \sigma_{l+1}^x + \frac{1-\gamma}{2} \sigma_l^y \sigma_{l+1}^y - \lambda \sigma_l^z \right)$$

Solution of the model.- **Spinless fermions** Lieb, Schultz, Mattis (1961)

Quantum Spin Models

XY model

$$H_{XY} = \sum_{l=1}^N \left(\frac{1+\gamma}{2} \sigma_l^x \sigma_{l+1}^x + \frac{1-\gamma}{2} \sigma_l^y \sigma_{l+1}^y - \lambda \sigma_l^z \right)$$

Solution of the model.- **Spinless fermions** Lieb, Schultz, Mattis (1961)

Canonical Transformations:

- **Jordan-Wigner**

$$a_l = \left(\prod_{m<l} \sigma_m^z \right) \frac{\sigma_l^x - i \sigma_l^y}{2} \quad \{a_l, a_m^+\} = \delta_{l,m}$$

- **Fourier**

$$d_k = \frac{1}{\sqrt{N}} \sum_{l=-\frac{N}{2}+1}^{\frac{N}{2}-1} a_l \text{Exp} \left[i \frac{2\pi}{N} kl \right] \quad \{d_k, d_p^+\} = \delta_{k,p}$$

- **Bogoliubov**

$$b_k^+ = u_k d_k^+ + i v_k d_{-k} \quad \{b_k, b_p^+\} = \delta_{k,p}$$

Quantum Spin Models

XY model

$$H_{XY} = \sum_{l=1}^N \left(\frac{1+\gamma}{2} \sigma_l^x \sigma_{l+1}^x + \frac{1-\gamma}{2} \sigma_l^y \sigma_{l+1}^y - \lambda \sigma_l^z \right)$$

$$\xrightarrow{N \rightarrow \infty} \int_{-\pi}^{\pi} \frac{d\phi}{2\pi} \sqrt{(\lambda - \cos \phi)^2 + \gamma^2 \sin^2 \phi} b_{\phi}^{\dagger} b_{\phi}$$

$$\xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{d\phi}{2\pi} \sum_{s=R,L} \sqrt{m^2 + c^2 \phi_s^2} b_{s,\phi}^{\dagger} b_{s,\phi}$$

Quantum Spin Models

XY model

Ground State:

$$b_\phi |\Psi_G\rangle = 0, \forall \phi \Rightarrow n_\phi = \langle b_\phi^\dagger b_\phi \rangle = 0 \quad \rho = \prod_{\otimes} \rho_\phi \Rightarrow \rho_\phi = \begin{pmatrix} n_\phi & 0 \\ 0 & 1 - n_\phi \end{pmatrix}$$
$$\Gamma_{mn} = \frac{1}{2\pi} \int_{-\pi}^{\pi} d\phi \cdot \text{Exp}[-i\phi(m-n)] U_\phi \rho_\phi U_\phi^\dagger$$

Quantum Spin Models

XY model

Ground State:

$$b_\phi |\Psi_G\rangle = 0, \forall \phi \Rightarrow n_\phi = \langle b_\phi^\dagger b_\phi \rangle = 0 \quad \rho = \prod_{\otimes} \rho_\phi \Rightarrow \rho_\phi = \begin{pmatrix} n_\phi & 0 \\ 0 & 1 - n_\phi \end{pmatrix}$$

Wick's Theorem:

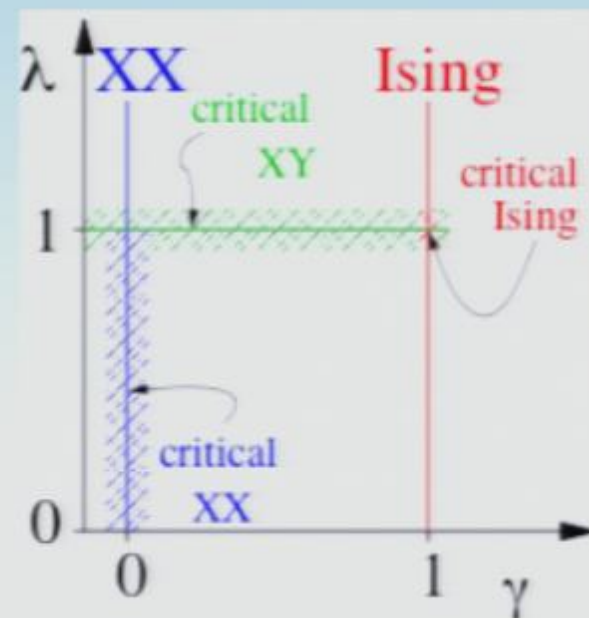
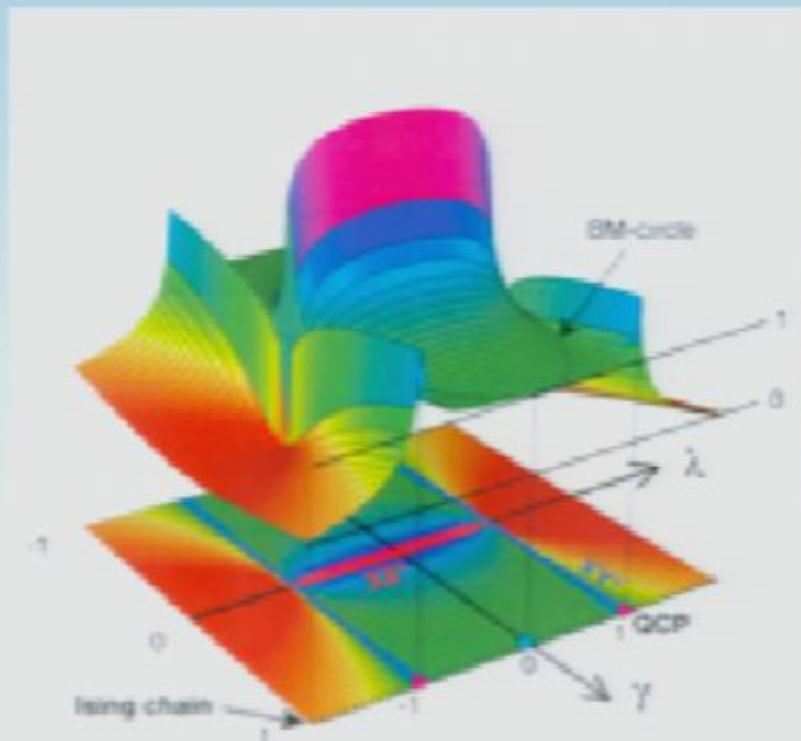
$$\Gamma_{mn} = \frac{1}{2\pi} \int_{-\pi}^{\pi} d\phi \cdot \text{Exp}[-i\phi(m-n)] U_\phi \rho_\phi U_\phi^\dagger$$

$$\langle a_{m1} a_{m2} a_{m3}^\dagger a_{m4}^\dagger \rangle = \langle a_{m1} a_{m2} \rangle \langle a_{m3}^\dagger a_{m4}^\dagger \rangle - \langle a_{m1} a_{m3}^\dagger \rangle \langle a_{m2} a_{m4}^\dagger \rangle + \langle a_{m1} a_{m4}^\dagger \rangle \langle a_{m2} a_{m3}^\dagger \rangle$$

$$\rho_L = \lim_{N \rightarrow \infty} \text{Tr}_{N-L} |\Psi_G\rangle \langle \Psi_G| = \prod_{\otimes} \begin{pmatrix} v_\phi & 0 \\ 0 & 1 - v_\phi \end{pmatrix}$$

Peschel(2002)

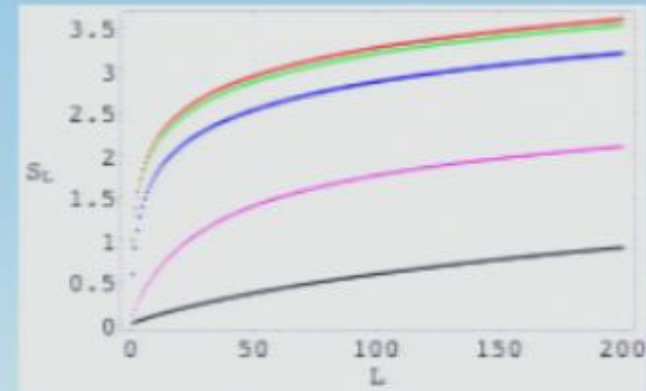
Analysis of the Ground State Entanglement



Analysis of the Ground State Entanglement

XX Model:

$$S_L^{XX} = \frac{1}{3} \log_2 L + \frac{1}{6} \log_2 (1 - \lambda^2)$$



Analysis of the Ground State Entanglement

Heisenberg Model:



$$S_L \cong \frac{1}{3} \log_2 L$$

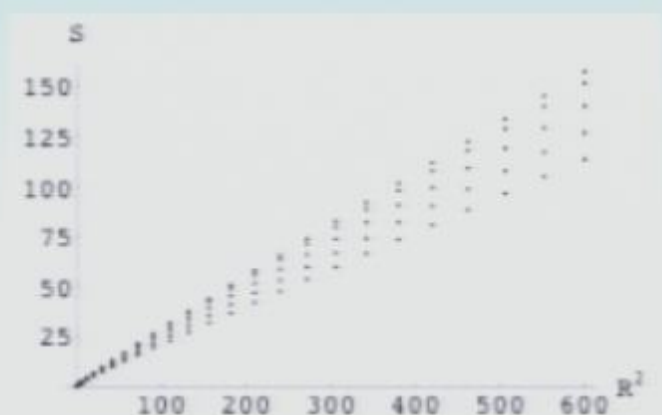
Entanglement in higher Dimensions

Srednicki (1993):
$$H = \frac{1}{2} \int d^3x \left(\Pi^2(x) + |\nabla \phi(x)|^2 + m^2 |\phi(x)|^2 \right)$$

Entanglement in higher Dimensions

Srednicki (1993):
$$H = \frac{1}{2} \int d^3x \left(\Pi^2(x) + |\nabla \phi(x)|^2 + m^2 |\phi(x)|^2 \right)$$

$$\phi = \sum \mu \cdot \phi_{in} \cdot \phi_{out}$$



$$S_{d=3} \propto R^2 \Rightarrow S_d \propto R^{d-1}$$

Connection with Conformal Field Theory

Connection with Conformal Field Theory

Wilczek, Callan, Holzhey, Larsen (1994):

Calabrese, Cardy (2004)

$$S_L = \frac{c + \bar{c}}{6} \log_2 L$$

**Heisenberg
& XX model:**

$$c = 1$$

Ising model:

$$c = \frac{1}{2}$$

Entanglement Loss along RG-Flows

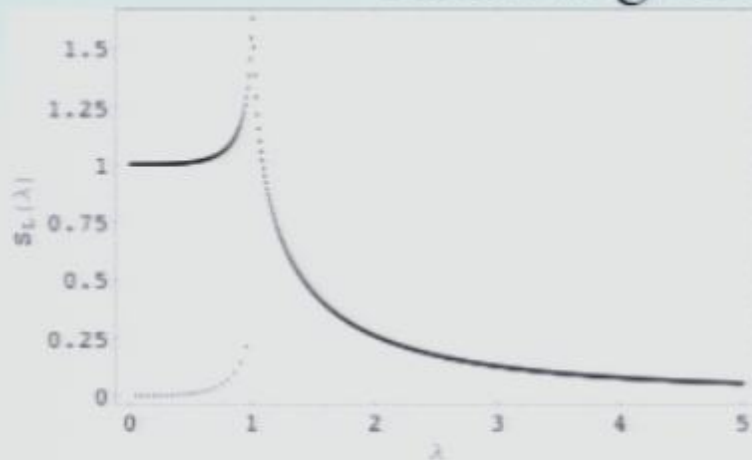
$$I \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^\dagger b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^\dagger b_k$$

Entanglement Loss along RG-Flows

$$l \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^+ b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^+ b_k$$

Renormalization Group Transformation

- Integration of the high-energy modes
- Rescaling of the parameters



$$\begin{aligned} a &\rightarrow a e^{-l} \\ m &\rightarrow m e^l \Rightarrow \begin{aligned} \frac{\partial a}{\partial l} &= -a \\ \frac{\partial \lambda}{\partial l} &= |1 - \lambda| \\ \frac{\partial c}{\partial l} &= 0 \end{aligned} \end{aligned}$$

Entanglement Loss along RG-Flows

•**Global loss.**- Direct consequence of universal scaling of entanglement in 1+1D systems and conformal symmetry

c-theorem: There is a monotonic nonincreasing function C along the Renormalization Group Flow.

Zamolodchikov (1986)

$$S = \frac{c + \bar{c}}{6} \log_2 L$$

$$c^{UV} > c^{IR} \Rightarrow S^{UV} > S^{IR}$$

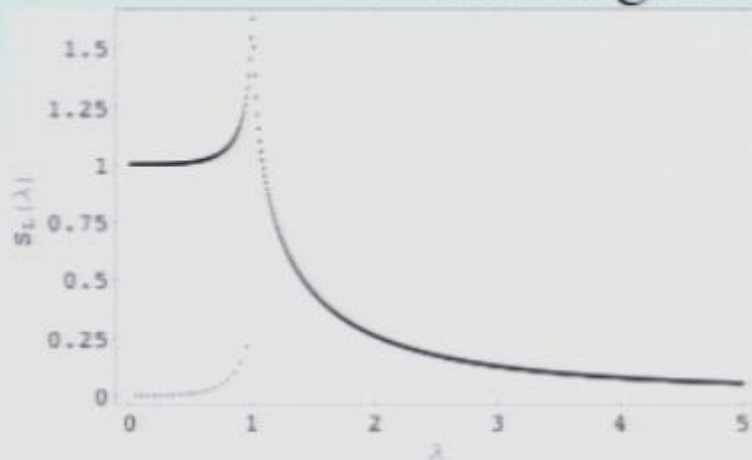
•**Monotonic loss of entanglement.**- Entanglement is nonincreasing step by step under renormalization group flows.

Entanglement Loss along RG-Flows

$$l \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^+ b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^+ b_k$$

Renormalization Group Transformation

- Integration of the high-energy modes
- Rescaling of the parameters



$$\begin{aligned} a &\rightarrow a e^{-l} \\ m &\rightarrow m e^l \Rightarrow \begin{aligned} \frac{\partial a}{\partial l} &= -a \\ \frac{\partial \lambda}{\partial l} &= |1 - \lambda| \\ \frac{\partial c}{\partial l} &= 0 \end{aligned} \end{aligned}$$

$$H = \sigma_i^z \sigma_{i+1}^z + \lambda \sigma_i^x + \epsilon \mathbb{O}$$

$$|\psi\rangle \approx |\uparrow\uparrow\uparrow\rangle \quad \text{top} \quad \text{bottom}$$

$$\sigma \rightarrow \sigma \quad |\psi\rangle \approx \text{down}$$

$$H = \sigma_i^z \sigma_{i+1}^z + \lambda \sigma_i^x + \epsilon \approx 0$$

$$|\psi\rangle \approx |\uparrow\uparrow\rangle \quad \uparrow \rightarrow \text{up} \quad \downarrow \rightarrow \text{down}$$

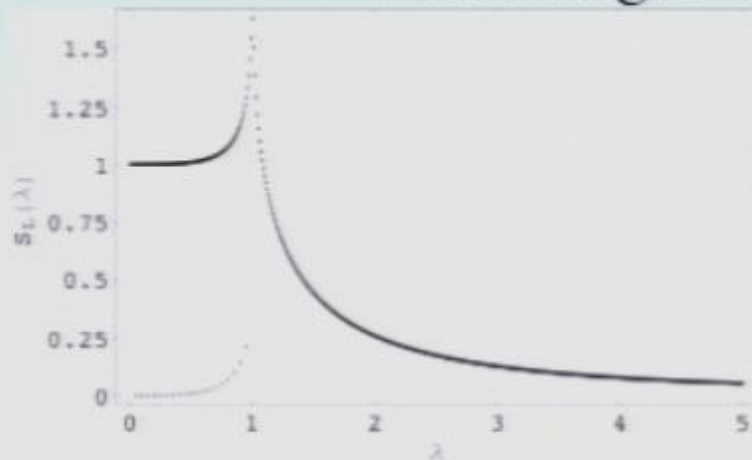
$$\psi \rightarrow 0 \quad |\psi\rangle \approx \text{down}$$

Entanglement Loss along RG-Flows

$$l \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^+ b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^+ b_k$$

Renormalization Group Transformation

- Integration of the high-energy modes
- Rescaling of the parameters



$$\begin{aligned} a &\rightarrow a e^{-l} \\ m &\rightarrow m e^l \Rightarrow \begin{aligned} \frac{\partial a}{\partial l} &= -a \\ \frac{\partial \lambda}{\partial l} &= |1 - \lambda| \\ \frac{\partial c}{\partial l} &= 0 \end{aligned} \end{aligned}$$

$$H = \sigma_i^z \sigma_{i+1}^z + \lambda \sigma_i^x + \epsilon \in O$$

$$|\psi\rangle \sim (|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle)$$

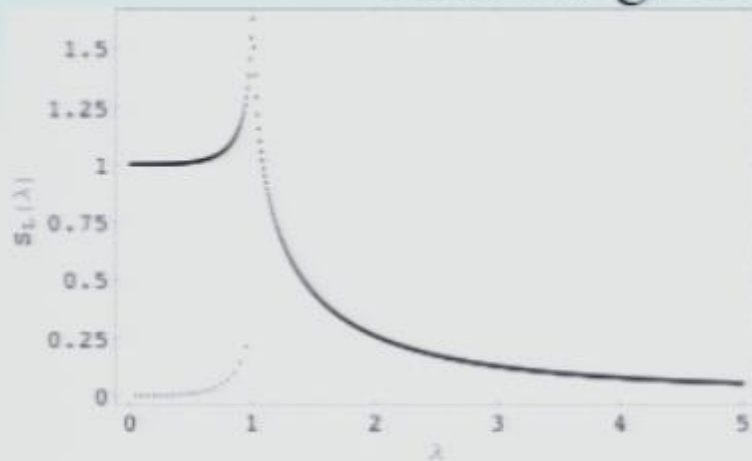
$$\sigma \rightarrow 2$$

Entanglement Loss along RG-Flows

$$l \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^+ b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^+ b_k$$

Renormalization Group Transformation

- Integration of the high-energy modes
- Rescaling of the parameters



$$\begin{aligned} a &\rightarrow a e^{-l} \\ m &\rightarrow m e^l \Rightarrow \begin{aligned} \frac{\partial a}{\partial l} &= -a \\ \frac{\partial \lambda}{\partial l} &= |1 - \lambda| \\ \frac{\partial c}{\partial l} &= 0 \end{aligned} \end{aligned}$$

Entanglement Loss along RG-Flows

- Global loss.**- Direct consequence of universal scaling of entanglement in 1+1D systems and conformal symmetry

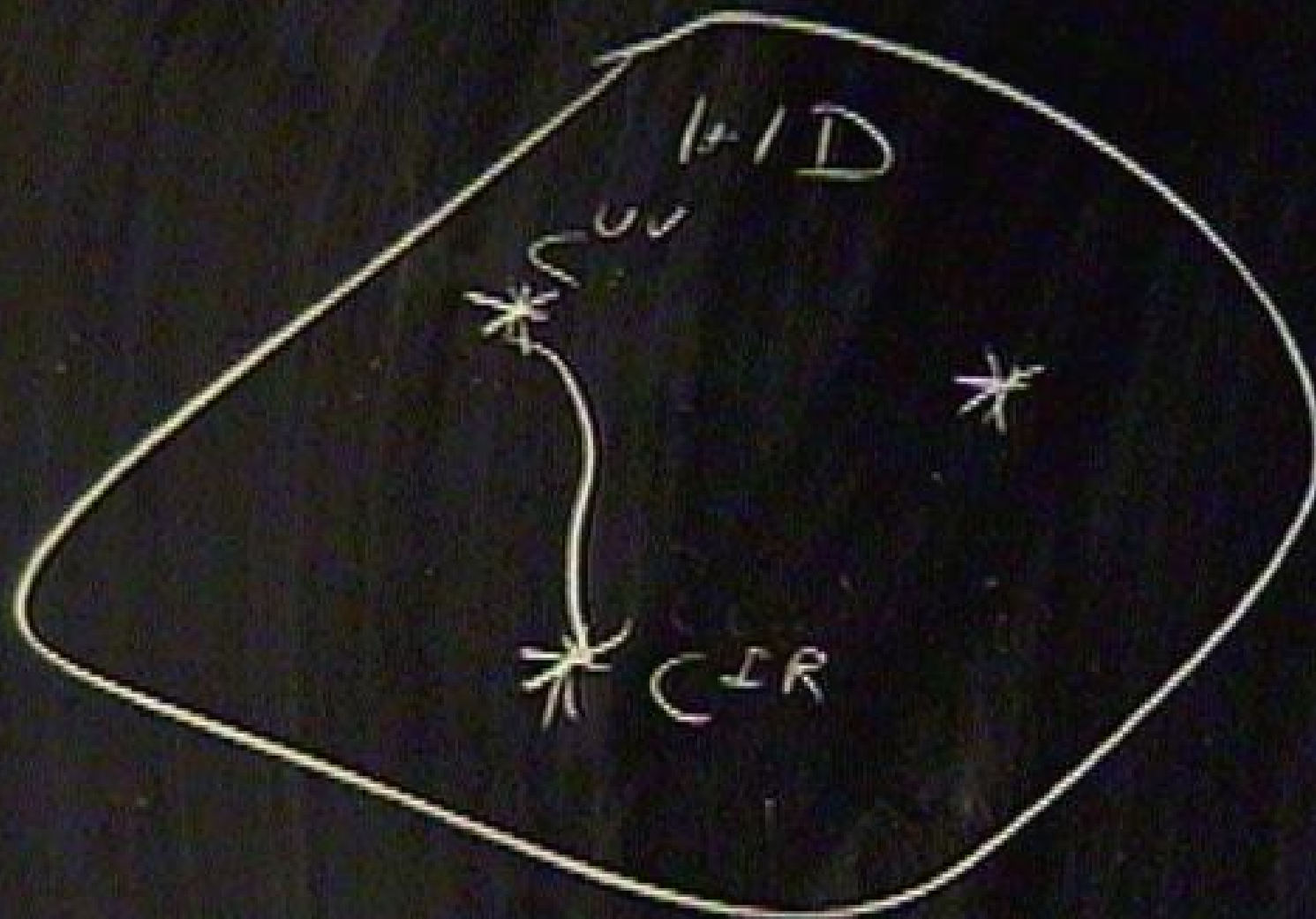
c-theorem: There is a monotonic nonincreasing function C along the Renormalization Group Flow.

Zamolodchikov (1986)

$$S = \frac{c + \bar{c}}{6} \log_2 L$$

$$c^{UV} > c^{IR} \Rightarrow S^{UV} > S^{IR}$$

- Monotonic loss of entanglement.**- Entanglement is nonincreasing step by step under renormalization group flows.



Entanglement Loss along RG-Flows

- Global loss.**- Direct consequence of universal scaling of entanglement in 1+1D systems and conformal symmetry

c-theorem: There is a monotonic nonincreasing function C along the Renormalization Group Flow.

Zamolodchikov (1986)

$$S = \frac{c + \bar{c}}{6} \log_2 L$$

$$c^{UV} > c^{IR} \Rightarrow S^{UV} > S^{IR}$$

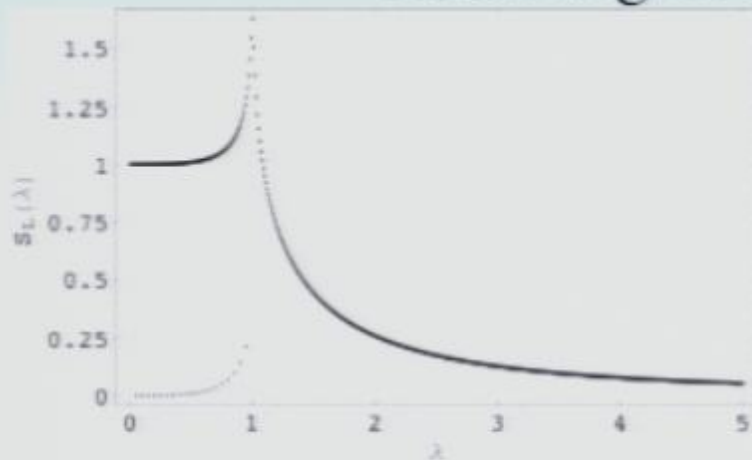
- Monotonic loss of entanglement.**- Entanglement is nonincreasing step by step under renormalization group flows.

Entanglement Loss along RG-Flows

$$l \sin g \xrightarrow{\text{low-energy}} \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left| \frac{1-\lambda}{a} \right|^2 + c^2 k^2} b_k^+ b_k = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \sqrt{\left(\frac{m}{a} \right)^2 + c^2 k^2} b_k^+ b_k$$

Renormalization Group Transformation

- Integration of the high-energy modes
- Rescaling of the parameters



$$\begin{aligned} a &\rightarrow a e^{-l} \\ m &\rightarrow m e^l \Rightarrow \begin{aligned} \frac{\partial a}{\partial l} &= -a \\ \frac{\partial \lambda}{\partial l} &= |1 - \lambda| \\ \frac{\partial c}{\partial l} &= 0 \end{aligned} \end{aligned}$$

Entanglement Loss along RG-Flows

- Global loss.**- Direct consequence of universal scaling of entanglement in 1+1D systems and conformal symmetry

c-theorem: There is a monotonic nonincreasing function C along the Renormalization Group Flow.

Zamolodchikov (1986)

$$S = \frac{c + \bar{c}}{6} \log_2 L$$

$$c^{UV} > c^{IR} \Rightarrow S^{UV} > S^{IR}$$

- Monotonic loss of entanglement.**- Entanglement is nonincreasing step by step under renormalization group flows.



$$E(14) > E(RG(14))$$

$$S(p) > S(p')$$

* CIR

$$\begin{aligned} E(14 >) &> E(RG(14 >)) \\ S(p) &> S(p') \\ 14 > &< 14' > \end{aligned}$$

Entanglement Loss along RG-Flows

Majorization.-Partial order relation $\vec{p} \prec \vec{q}$

$$\vec{p} = D\vec{q} = \left(\sum_j r_j \Pi_j \right) \vec{q}$$

$$\sum_j r_j = 1$$

$$p_1 \leq q_1$$

$$p_1 + p_2 \leq q_1 + q_2$$

.....

$$\sum_i p_i \leq \sum_i q_i$$

•**Fine-grained Loss of Entanglement.**-Majorization at any infinitesimal step in the RG-Flow.

$$(\lambda, \gamma) \xrightarrow{\text{RG-flow}} (\lambda', \gamma')$$

$$\rho(\lambda, \gamma) \prec \rho(\lambda', \gamma') \Rightarrow S(\rho(\lambda, \gamma)) > S(\rho(\lambda', \gamma'))$$

Entanglement Loss along RG-Flows

- **Fine-grained Loss of Entanglement.-**

**Heisenberg
& XY model:**

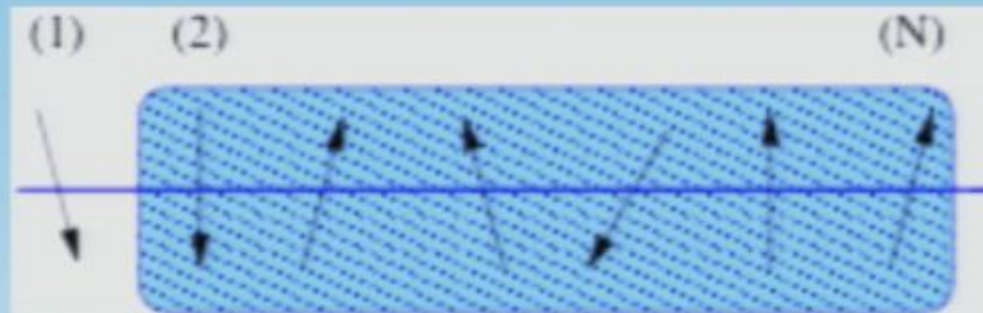
$$\rho(L, \gamma, \lambda) = \prod_{\otimes k} \rho_k(L, \gamma, \lambda)$$

$$\rho_k[L, \gamma, \lambda] = \frac{\exp(-n_k \varepsilon_k[L, \gamma, \lambda])}{Z_k[L, \gamma, \lambda]} \quad \text{Peschel (2004)}$$

$$\rho_1 \prec \rho'_1 \Rightarrow \rho_1 = D\rho'_1 \quad \rho_2 \prec \rho'_2 \Rightarrow \rho_2 = \tilde{D}\rho'_2$$

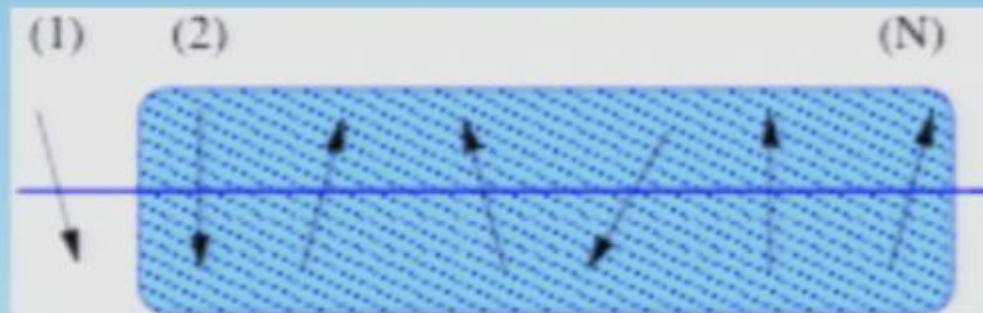
$$(\rho_1 \otimes \rho_2) = (D \otimes \tilde{D})(\rho'_1 \otimes \rho'_2) \Rightarrow \rho_1 \otimes \rho_2 \prec \rho'_1 \otimes \rho'_2$$

Matrix Product States



$$|\psi\rangle = \sum_{a=1}^D \lambda_a^{(1)} |\phi_a^{(1)}\rangle |\phi_a^{(2,\dots,N)}\rangle = \sum_{a=1}^D \sum_{s_1=1}^d \lambda_a^{(1)} |s_1\rangle \langle s_1 | \phi_a^{(1)} \rangle |\phi_a^{(2,\dots,N)}\rangle$$

Matrix Product States



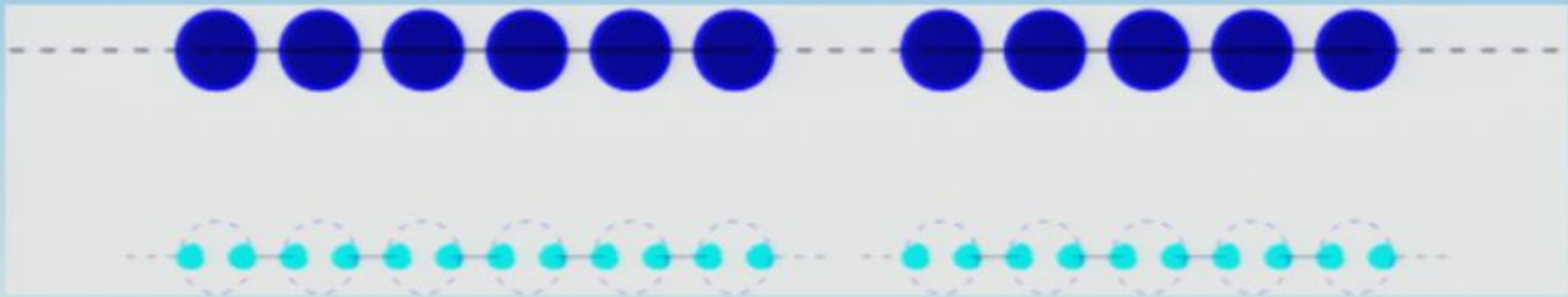
$$|\psi\rangle = \sum_{a=1}^D \lambda_a^{(1)} |\phi_a^{(1)}\rangle |\phi_a^{(2,\dots,N)}\rangle = \sum_{a=1}^D \sum_{s_1=1}^d \lambda_a^{(1)} |s_1\rangle \langle s_1 | \phi_a^{(1)} \rangle |\phi_a^{(2,\dots,N)}\rangle$$

$$= \sum_{a=1}^D \sum_{s_1=1}^d |s_1\rangle A_a[s_1] |\phi_a^{(2,\dots,N)}\rangle = \sum_{a,b=1}^D \sum_{s_1, s_2=1}^d |s_1\rangle A_a[s_1] |s_2\rangle A_{ab}[s_2] |\phi_b^{(3,\dots,N)}\rangle$$

$$= \dots = \sum_{\{s_j\}=1}^d \text{tr}(A[s_1] \cdots A[s_N]) |s_1\rangle \dots |s_N\rangle$$

Matrix Product States

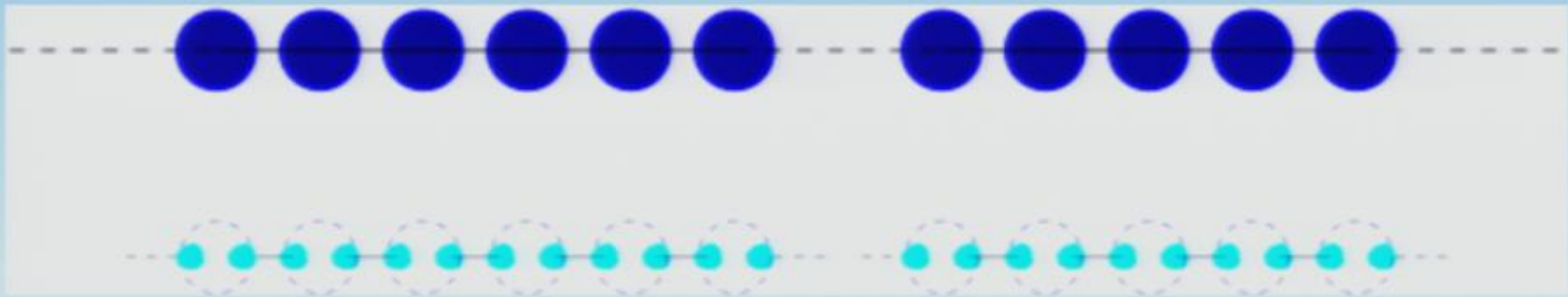
$$A_{D \times D}[s] \left| s \right\rangle_{s=1}^d$$



$$C^D \otimes C^D \xrightarrow{A} C^d$$

Matrix Product States

$$A_{D \times D}[s] \left| s \right\rangle_{s=1}^d$$



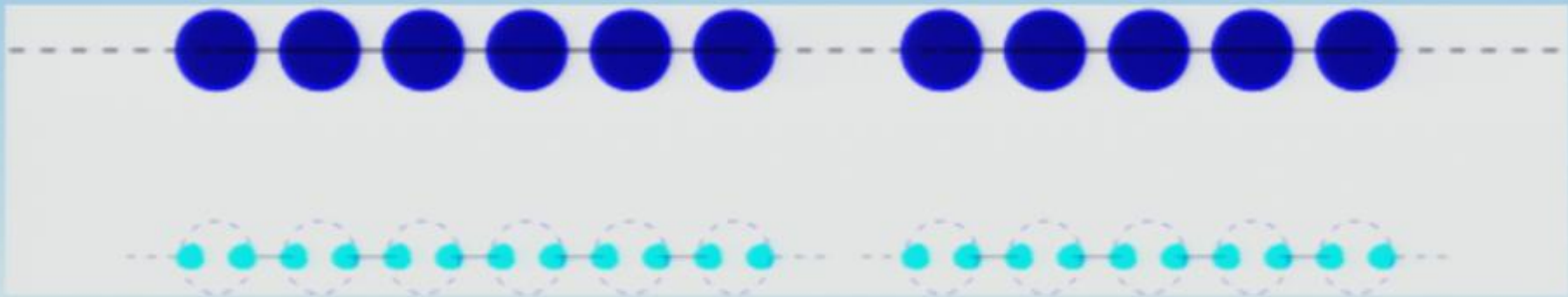
$$C^D \otimes C^D \xrightarrow{A} C^d$$

Properties of MPS by definition

variant

Matrix Product States

$$A_{D \times D}[s] \left| s \right\rangle_{s=1}^d$$



$$C^D \otimes C^D \xrightarrow{A} C^d$$

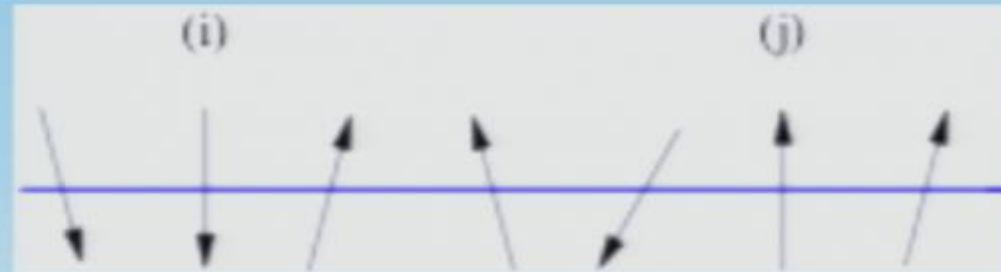
Properties of MPS by definition

1. Translational invariant

picture)



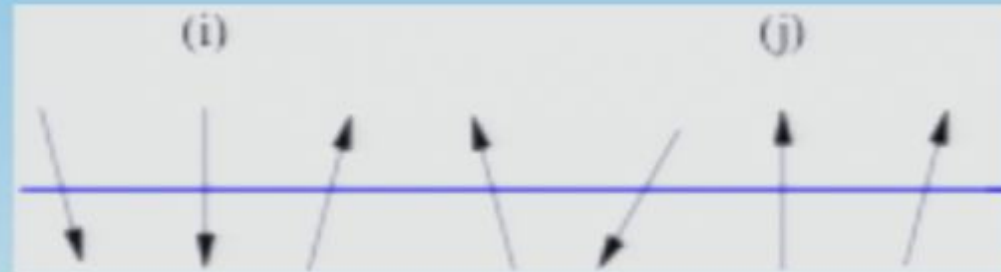
Fixed Points in Matrix Product States



$$\langle O_i O_j \rangle = \text{tr} \left(E^{N-j+i-2} \tilde{O}_i E^{j-i} \tilde{O}_i \right)$$

$$E = \sum_{s=1}^d A^*[s] \otimes A[s] \quad \tilde{O} = \sum_{s,s'=1}^d A^*[s] \otimes A[s'] \langle s|O|s' \rangle$$

Fixed Points in Matrix Product States



$$\langle O_i O_j \rangle = \text{tr} \left(E^{N-j+i-2} \tilde{O}_i E^{j-i} \tilde{O}_i \right)$$

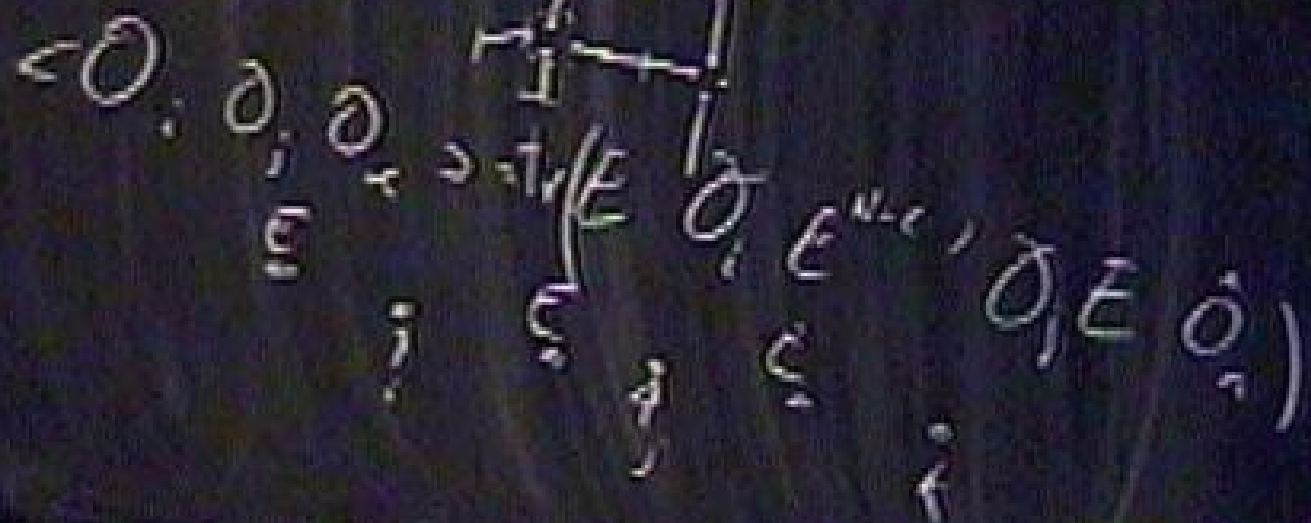
$$E = \sum_{s=1}^d A^*[s] \otimes A[s] \quad \tilde{O} = \sum_{s,s'=1}^d A^*[s] \otimes A[s'] \langle s|O|s' \rangle$$

**Scale
Transformation**

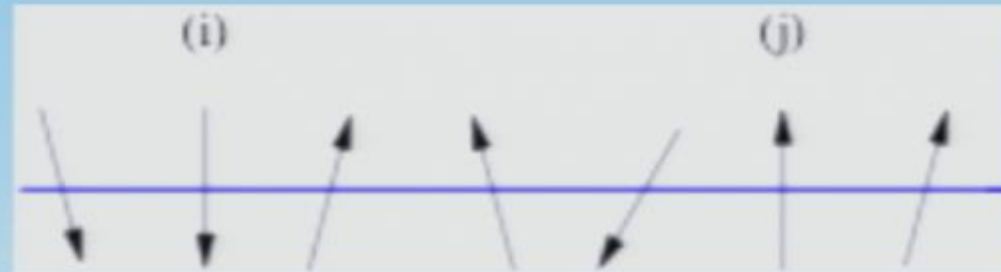
$$E \rightarrow E' = E \cdot E = \sum_{s'=1}^{d'} A'^*[s'] \otimes A'[s']$$

Fixed Points

$$E = E' = E \cdot E$$



Fixed Points in Matrix Product States



$$\langle O_i O_j \rangle = \text{tr} \left(E^{N-j+i-2} \tilde{O}_i E^{j-i} \tilde{O}_i \right)$$

$$E = \sum_{s=1}^d A^*[s] \otimes A[s] \quad \tilde{O} = \sum_{s,s'=1}^d A^*[s] \otimes A[s'] \langle s|O|s' \rangle$$

**Scale
Transformation**

$$E \rightarrow E' = E \cdot E = \sum_{s'=1}^{d'} A'^*[s'] \otimes A'[s']$$

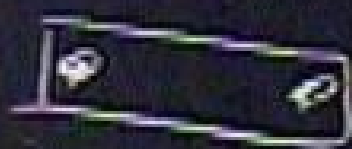
Fixed Points

$$E = E' = E \cdot E$$

$$H = \sigma_z^1 \sigma_z^2 + \sigma_z^1 + \sigma_z^2 + \epsilon 0$$

$$|\psi\rangle = |\uparrow\uparrow\rangle_{\text{top}} + |\downarrow\downarrow\rangle_{\text{bottom}}$$

$$C \rightarrow 0 \quad |\psi\rangle = |\downarrow\downarrow\rangle_{\text{top}}$$



E'

E'

E'

Conclusions

- Universal Scaling Law of Entanglement in 1+1 D Systems in the critical intervals
- Entanglement behaviour is controlled by Conformal Symmetry
- Away from the critical points Entanglement is saturated by a “mass” scale
- DMRG (Density Matrix Renormalization Group) failure in higher Dimensions
- Entanglement is no increasing under RG Flows
- Relation between Majorization and C-functions
- Description of Real Space Renormalization Group Transformations within Matrix Product States

Entanglement, Critical Phenomena & RG Flows

- Entanglement in Quantum Critical Phenomena.

G. Vidal, J.I. Latorre, E. Rico, A. Kitaev. Phys.Rev.Lett. 90 (2003) 227902

- Ground State Entanglement in Quantum Spin Chains.

J.I. Latorre, E. Rico, G. Vidal. Quant. Inf. & Comp. Vol.4 no.1 (2004) pp.048-092

- Fine-Grained Entanglement Loss along Renormalization Group Flows.

J.I. Latorre, C.A. Lütken, E. Rico, G. Vidal. quant-ph/0404120

- Renormalization Group Transformation on quantum states.

F. Verstraete, J.I. Cirac, J.I. Latorre, E. Rico, M.M. Wolf, quant-ph/0410227



Pirsa: 05010010



UNIVERSITAT DE BARCELONA



UNIVERSITY
OF OSLO



Page 55/55

MPO