

Title: Arithmetic Varieties from String Theory and D-branes

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Abstract:

Arithmetic Varieties from String Theory & D-Branes  
(Work in progress)

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References: RS, JGP 44 (2003)  
RS & S. Underwood, JGP 48 (2003) 169  
M. Lynker & RS, hep-th/0410189  
RS, J. Tyson, to appear

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## The basic picture of string theory

Naive picture:

The 1-dimensional string sweeps out a Riemann surface  $\Sigma_g$  of genus  $g$  — the string world sheet — which is embedded in 10-dimensional spacetime  $M^{1,9}$

$$\begin{aligned}x: \Sigma_g &\longrightarrow M^{1,d} \\(\tau, \sigma) &\mapsto (x^0(\tau, \sigma), \dots, x^d(\tau, \sigma)).\end{aligned}$$

$d$  is a priori an undetermined number. Naive considerations would expect

$$d = 3. \quad (?)$$

( $x$  has (super) symmetry partners carrying half integral representations.)

Spacetime  $M^{1,d}$  comes equipped with different fields:

- Metric  $g$
- Connection  $A$  in a principal bundle with structure group  $E_8 \times E_8$
- 2-Form  $B$

Consistency constraints arise via

- \* supersymmetry (world sheet and spacetime)
- \* conformal invariance (world sheet)



Different solutions of consistency constraints:

- $d = 9$
- Calabi-Yau varieties (geometries)

$$M^{1,9} = M^{1,3} \times X,$$

with  $X$  a compact complex Kähler threefold with vanishing first Chern class.

Set

$$\mathcal{A} = \text{set of CYs } \cup \text{ Special Fanofolds } \cup \dots$$

- Conformal field theories/vertex algebras

$$M^{1,9} = M^{1,3} \times \text{CFT},$$

where CFT is a 2-dimensional field theory on the world sheet Riemann surface which has to satisfy certain constraints.

Set

$$\mathcal{B} = \text{set of CFTs.}$$



**Obvious Questions:**

**Difficult:**

Is it possible to find an underlying string CFT for all CYs, i.e. find a map

$$\tau : \mathcal{A} \longrightarrow \mathcal{B}.$$

(Has been achieved at cohomological level in many examples.)

**Even more difficult:**

Is it possible to find a geometric interpretation for all string theoretic CFTs:

$$\mathcal{C} : \mathcal{B} \longrightarrow \mathcal{A}.$$

(Has not been achieved in a single example in dimensions larger than one).

**String theory induced recurring fantasy:**

- Let  $\mathcal{CY}$  be a hypothetical subcategory in the category of varieties (CYs + Special Fano).
- Let  $\mathcal{CFT}$  be a subcategory in some hypothetical category of CFTs/VOAs
- Then there exists a string theoretic functor  $S$

$$S : \mathcal{CY} \longrightarrow \mathcal{CFT}$$



### **String theory theme: modularity**

Idea for contact between string world sheet and spacetime: **modularity**, one of the important themes in arithmetic geometry.

Hecke (28), Eichler(54), Taniyama(55,57)

Shimura(50s,60s,70s), Weil(67)

Langlands(60s,70s), Deligne(60s,70s), Serre(70s,80s)

Wiles(95), Breuil-Conrad-Diamond-Taylor(01)

Consani, Dieulefait, Hulek, Saito, Scholten, van Geemen, van Straten, Verrill, Yui and others

#### **Observation/Idea:**

Geometry leads to modular forms and vice versa.

**Not good enough –**

**string theory has particular needs:**

Calabi-Yau varieties must lead to particular types of modular forms coming from the string theoretic conformal field theory/affine Lie algebra.

**Task:** Find them.

**Problem:** It is not clear what the best framework is to make this question precise.

**Here:** Consider conformal field theories based on affine Lie algebras.



## String theory and affine Lie algebras

- Want exactly solvable string theory models
  - (simplicity, rich structure, obvious symmetries, computability)
  - (no path integral, no action, no Lagrangian)
- Ingredients: affine Lie algebras/Virasoro algebra
- Physical version: OPEs, correlation functions
- Mathematical versions: VOAs, Operads

Needed in the present context: affine Lie algebra structures

Ingredients:

- Connected compact Lie group  $G$
- Lie algebra  $\mathfrak{g}$
- Loop algebra  
$$L\mathfrak{g} := \mathfrak{g} \otimes \mathbb{C}[t, t^{-1}]$$
- Central extension: exact sequence

$$0 \longrightarrow \mathfrak{a} \longrightarrow \hat{\mathfrak{g}} \longrightarrow \mathfrak{g} \longrightarrow 0$$

with central  $\mathfrak{a}$ , i.e.

$$[a, x] = 0 \quad \forall a \in \mathfrak{a}, x \in \hat{\mathfrak{g}}.$$



- Special case: 1-dimensional extension

$$\mathfrak{a} \cong \mathbb{C} \cdot \mathbf{1}$$

- Affine Lie algebra  $\hat{\mathfrak{g}}$ : central extension of the loop algebra

$$0 \longrightarrow \mathbb{C}K \longrightarrow \hat{\mathfrak{g}} \longrightarrow \text{Lg} \longrightarrow 0$$

with central element  $K$ .

- Virasoro algebra Vir: central extension of the Lie algebra  $\text{Der } \mathcal{K} = \mathbb{C}((t))\partial_t$  of derivatives

$$0 \longrightarrow \mathbb{C}C \longrightarrow \text{Vir} \longrightarrow \text{Der } \mathcal{K} \longrightarrow 0$$

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Example:  $A_1^{(1)}$

Lie algebra:  $A_1$  ( $\text{su}(2), \text{sl}(2)$ )

Cartan matrix:  $A = (2)$

Generalized Cartan matrix:  $\hat{A} = \begin{pmatrix} 2 & -2 \\ -2 & 2 \end{pmatrix}$

Character maps on the upper half-plane

$$\chi_\ell : \mathfrak{H} \times \mathcal{C} \longrightarrow \mathbb{C}$$

defined as

$$\begin{aligned}\chi_\ell^k(\tau, u) &= \text{tr}_{\mathcal{H}_\ell} q^{L_0 - \frac{\ell}{24}} e^{2\pi i u J_0} \\ &= \sum_{\substack{m=-k+1 \\ m=\ell \bmod 2}}^k c_{\ell,m}^k \theta_{m,k}(\tau, u),\end{aligned}$$

with string functions

$$c_{\ell,m}^k(\tau) = \frac{1}{\eta^3(\tau)} \sum_{\substack{-|x| < y \leq |x| \\ (x,y) \text{ or } (\frac{1}{2}-x, \frac{1}{2}+y) \\ \in \mathbb{Z}^2 + \left(\frac{\ell+1}{2(k+2)}, \frac{m}{2k}\right)}} \text{sign}(x) e^{2\pi i \tau((k+2)x^2 - ky^2)}$$

theta functions

$$\theta_{m,n}(\tau, z, u) = e^{-2\pi i u} \sum_{l \in \mathbb{Z} + \frac{m}{2k}} e^{2\pi i \tau n l^2 + 2\pi i z l}$$

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Closed under modular transformations:

$$\tau \mapsto \tau + 1:$$

$$\chi_\ell(\tau + 1, u) = e^{2\pi i (\Delta_\ell^k - \frac{c}{24})} \chi_\ell(\tau, u)$$

$$\tau \mapsto -1/\tau:$$

$$\chi_\ell\left(-\frac{1}{\tau}, \frac{u}{\tau}\right) = e^{\pi i k \tau^2/2} \sum_m S_{\ell m} \chi_m(\tau, u),$$

with modular  $S$ -matrix

$$S_{\ell m} = \sqrt{\frac{2}{k+2}} \sin\left(\frac{(\ell+1)(m+1)\pi}{k+2}\right)$$

and

$$\Delta_\ell^k = \frac{\ell(\ell+2)}{4(k+2)}$$

**Remark.**

The rational numbers given by the level  $k$ , the central charge  $c$ , and the anomalous dimensions  $\Delta_\ell$ , characterize the representations of the model (field content), and determine the structure of the correlation functions (Yukawa couplings).

Hence physically important, and arithmetically boring.

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More interesting parametrization:

(Motivated by 2-dimensional critical systems in 1980s and early 1990s)

Define the quantum dimensions:

$$Q_{\ell m} = \frac{S_{\ell m}}{S_{0m}}$$

These are more interesting numbers number theoretically.

Even better, they know about the anomalous dimensions and the central charge.

Define the Rogers dilogarithm function as

$$L(z) = Li_2(z) + \frac{1}{2} \log(z) \log(1-z),$$

where

$$Li_2(z) = \sum_{n \in \mathbb{N}} \frac{z^n}{n^2}$$

is Euler's dilogarithm.

This function leads to the rather surprising

**Theorem.** (Nahm-Recknagel-Terhoeven, Kuniba-Nakanishi)

$$\frac{1}{L(1)} \sum_{i=1}^k L\left(\frac{1}{Q_{ij}^2}\right) = \frac{3k}{k+2} - 24\Delta_j^{(k)} + 6j.$$

Thus we have two equivalent ways of parametrizing the behavior of the fields at fixed central charge.

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**Questions beyond the spectrum:**

There are a **great** many detailed questions one can ask in the context of the CY/CFT connection and D-branes.

Perhaps the first step would be to understand the characters.

**Problem:**

Understand the geometric origin of CFT characters.

**Observation:**

Characters are modular forms.

**Question:**

Where do modular forms come from?



### 3. Modular forms:

Denote by  $\Gamma_0(N) \subset \text{SL}(2, \mathbb{Z})$  Hecke's congruence group

$$\Gamma_0(N) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \text{SL}_2(\mathbb{Z}) \mid c = 0 \pmod{N} \right\}.$$

**Definition.** A function on the upper half plane  $\mathfrak{H}$  which transforms under

$$g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_0(N) \subset \text{SL}(2, \mathbb{Z})$$

as

$$f(g\tau) = (c\tau + d)^k f(\tau)$$

is called a modular form of weight  $k$  with respect to  $\Gamma$ .

### Hecke construction of modular forms:

Consider L-series

$$L(s) = \sum_n \frac{a_n}{n^s}$$

which satisfy a functional equation

$$L(s) = w L(k-s)$$

for some  $w$ . Then their Mellin transform  $f$  defined via

$$L(f, s) = \frac{(2\pi)^s}{\Gamma(s)} \int_0^{i\infty} (-iz)^s f(z) \frac{dz}{z}$$

defines a modular form of weight  $k$ .

More precisely:



### Converse Theorems:

In practice: one computes L-series via some interesting structure and wants to know whether they are modular.

#### Hecke's Converse Theorem.

Suppose  $N = 1, 2, 3, 4$  and  $L(s) = \sum_n a_n n^{-s}$  converges in some right half plane and continues to an entire function such that  $\Gamma(s)L(s)$  is entire and bounded in vertical strips. If

$$\Lambda(s) := \left( \frac{\sqrt{N}}{2\pi} \right)^s \Gamma(s)L(s)$$

satisfies the functional equation

$$\Lambda(s) = \pm(-1)^{k/2} \Lambda(k-s)$$

then  $f(q) = \sum_n a_n q^n$ ,  $q = e^{2\pi i \tau}$ , is a modular form of weight  $k$  for  $\Gamma_0(N)$ .

#### Generalizations:

- Weil 1967 for any  $N$  via twistings (unpleasant).
- Conrey/Farmer:  $N \leq 15$  (via Euler product).

Q: Where do L-series come from?

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L-functions appear in a great many different contexts:

Themes:

- Number theory (Euler, Riemann, Dirichlet, Dedekind, Hecke)
  - Galois representations (Artin, Tate, Langlands)
  - Algebraic varieties (Hasse-Weil L-functions)
  - Modular forms
- Sometimes all these themes are connected:
- (\*) Abelian varieties with complex multiplication.
  - (\*) Étale cohomology of varieties (schemes).

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### Geometric Zeta-functions

Hasse (1930), Weil (1950s): Consider

$$Z(X/\mathbb{F}_p, t) = \exp \left( \sum_{r \in \mathbb{N}} \#(X/\mathbb{F}_{p^r}) \frac{t^r}{r} \right)$$

with

- For  $p$  prime  $\mathbb{F}_{p^r}$  = extension of

$$\mathbb{F}_p = \{0, 1, 2, \dots, p-1\}$$

of degree  $r$ .

- $t$  is a formal variable, and F.K.Schmidt (1930s) showed rationality:

$$Z(X/\mathbb{F}_p, t) = \frac{\mathcal{P}^{(p)}(t)}{(1-t)(1-pt^2)}$$

with

$$\mathcal{P}^{(p)}(t) = 1 - \beta(p)t + pt^2$$

and

$$\beta(p) = p + 1 - \#(X/\mathbb{F}_p).$$



### Cohomological L-functions:

There still is an arbitrary choice of the prime  $p$  in  $Z(X/\mathbb{F}_p, t)$ . Physically this is unjustified and should be avoided (no 'scale').

Define the Hasse-Weil L-function of a Calabi-Yau  $d$ -fold as the cohomological L-function associated to  $H^d(X)$ .

$$\begin{aligned}L_{HW}(X, s) &= \prod_{\substack{\mathbb{Z}/p \\ p \text{ good prime}}} \frac{1}{\mathcal{P}_d^{(p)}(p^{-s})} \\&= \prod_{\substack{\mathbb{Z}/p \\ p \text{ good prime}}} \frac{1}{\prod_{j=1}^{b_2} \left(1 - \beta_j^{(d)}(p)p^{-s}\right)} \\&=: \sum_n \frac{a_X^d(n)}{n^s},\end{aligned}$$

where the coefficients contain arithmetic/cohomological information about the variety.

### Problem:

Find a string theoretic interpretation of the coefficients  $a_X^d(n)$ .

**Idea:** Look for stringy modular forms.

But: string theory only provides a selection of rather specific modular forms.

Not any old modular form will be useful.



## Geometric Modularity

*Is it possible to understand the modular symmetry of the string geometrically?*

### Simplest framework:

Consider the simplest possible string compactification from 10 to 8 dimensions (unphysical toy models)

$$M^{1,9} \longrightarrow M^{1,7} \times E$$

with elliptic curves  $E$ .

### Examples:

The simplest class of models is provided by the class of elliptic Brieskorn-Pham curves for which physicists have conjectured a relation with exactly solvable conformal field theories.

They are described as polynomials in weighted projective space

$$E^3 = \{(x:y:z) \in \mathbb{P}_2 \mid x^3 + y^3 + z^3 = 0\},$$

$$E^4 = \{(x:y:z) \in \mathbb{P}_{1,1,2} \mid x^4 + y^4 + z^2 = 0\},$$

$$E^6 = \{(x:y:z) \in \mathbb{P}_{1,2,3} \mid x^6 + y^3 + z^2 = 0\}$$

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$$E^5 = \{(x:y:z) \in \mathbb{P}_{1,2,3} \mid x^5 + y^3 + z^2 = 0\}$$



These lead to the Hasse-Weil series

$$L_{HW}(E^3, s) = 1 - \frac{2}{4^s} - \frac{1}{7^s} + \frac{5}{13^s} + \frac{4}{16^s} - \frac{7}{19^s} + \dots,$$

$$L_{HW}(E^4, s) = 1 + \frac{2}{5^s} - \frac{3}{9^s} - \frac{6}{13^s} + \frac{2}{17^s} - \frac{1}{25^s} + \dots,$$

$$L_{HW}(E^6, s) = 1 + \frac{4}{7^s} + \frac{2}{13^s} - \frac{8}{19^s} - \frac{5}{25^s} + \frac{4}{31^s} + \dots,$$

Associated to these are Hasse-Weil forms defined by the Mellin transforms

$$f_{HW}(E^3, q) = q - 2q^4 - q^7 + 5q^{13} + 4q^{16} - 7q^{19} + \dots$$

$$f_{HW}(E^4, q) = q + 2q^5 - 3q^9 - 6q^{13} + 2q^{17} - q^{25} + \dots$$

$$f_{HW}(E^6, q) = q + 4q^7 + 2q^{13} - 8q^{19} - 5q^{25} + 4q^{31} + \dots$$

These forms are normalized cusp Hecke eigenforms of weight 2 and levels 27, 64, 144, respectively.

**Goal:** Derive these modular forms from string theory.

**Specifically:**

We expect these elliptic curves to correspond to affine Lie algebras constructed from two or three copies of the algebra  $A_{1,k}^{(1)}$  at different levels  $k$ . Models:

$$C_n = \left( \bigotimes_{i=1}^{n_f} A_{1,k_i}^{(1)} \right)_{\text{GSO}}$$

These lead to the Hasse-Weil series

$$L_{HW}(E^3, s) = 1 - \frac{2}{4^s} - \frac{1}{7^s} + \frac{5}{13^s} + \frac{4}{16^s} - \frac{7}{19^s} + \dots,$$

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Associated to these are Hasse-Weil forms defined by the Mellin transforms

$$f_{HW}(E^3, q) = q - 2q^4 - q^7 + 5q^{13} + 4q^{16} - 7q^{19} + \dots$$

$$f_{HW}(E^4, q) = q + 2q^5 - 3q^9 - 6q^{13} + 2q^{17} - q^{25} + \dots$$

$$f_{HW}(E^6, q) = q + 4q^7 + 2q^{13} - 8q^{19} - 5q^{25} + 4q^{31} + \dots$$

These forms are normalized cusp Hecke eigenforms of weight 2 and levels 27, 64, 144, respectively.

**Goal:** Derive these modular forms from string theory.

**Specifically:**

We expect these elliptic curves to correspond to affine Lie algebras constructed from two or three copies of the algebra  $A_{1,k}^{(1)}$  at different levels  $k$ . Models:

$$C_n = \left( \bigotimes_{i=1}^{n_f} A_{1,k_i}^{(1)} \right)_{\text{GSO}}$$



Conformal field theoretic models (at  $c = 3$ ):  
(rough indication)

$$\begin{aligned}C_3 &= (A_{1,1}^{(1)})_{\text{GSO}}^{\otimes 3} \\C_4 &= (A_{1,2}^{(1)})_{\text{GSO}}^{\otimes 2} \\C_6 &= (A_{1,1}^{(1)} \otimes A_{1,4}^{(1)})_{\text{GSO}},\end{aligned}$$

where GSO stands for a supersymmetric projection.

These models have a great many  $q$ -expansions in terms of various quantities:

Objects:	Level $k$	Number (rough)
String functions	1	1
	2	3
	4	7
Theta functions	1	1
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#### Remarks.

- It is not obvious from a physical perspective whether this is possible.

String theory provides no guidance what to choose even in this simple case.

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- At  $k = 1$  none of the affine characters  $\chi_\ell$  of  $A_1^{(1)}$  are useful.
- Neither are the  $N = 2$  characters  $\chi_{q,s}^\ell$  or theta functions.
- Nor are the string functions  $c_m^\ell$ , which (essentially) describe the characters of parafermionic theories.
- The weight of  $c_m^\ell$  is  $-\frac{1}{2}$ , hence the weight of

$$\Theta_m^\ell = \eta^3(\tau) c_m^\ell$$

is unity:

- Can make weight 2 with products

$$\Theta_m^\ell(a\tau) \Theta_{m'}^{\ell'}(b\tau)$$

with integers  $a, b$ .

- Insight into  $a, b$  via Weil's conductor/level observation:

in all known examples the level of the modular form agreed with the geometric conductor.

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Compute and compare:

$$\begin{aligned}\Theta_{1,1}^1(q^3)\Theta_{1,1}^1(q^9) &= q - 2q^4 - q^7 + 5q^{13} + 4q^{16} - 7q^{19} + \dots \\ f_{\text{HW}}(E_3, q) &= q - 2q^4 - q^7 + 5q^{13} + 4q^{16} - 7q^{19} + \dots\end{aligned}$$

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The conductor of this curve is 64, hence we would expect the level to be 64, or perhaps a divisor of it.

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**Needed:** new rules.

Need some new gadget that changes sign in just the right way (and some rationale why it didn't appear for  $E_3$ !).

We want to twist the form with a character

$$\chi : \mathbb{Z} \longrightarrow \mathbb{Z}^\times$$

via

$$f_\chi(q) := \sum_n \chi(n) a_n q^n.$$

Class of candidates:

$$\chi_n(p) = \left( \frac{n}{p} \right) = \begin{cases} 1 & n \text{ is a square in } \mathbb{F}_p \\ -1 & n \text{ is not a square in } \mathbb{F}_p \end{cases}$$

What works is  $n = 2$ :

The Hasse-Weil L-function  $L_{HW}(E_4, s)$  of the quartic elliptic curve  $E_4$  factors into the twisted product

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This case is more complicated for two reasons:

- 1) Different levels:  $k = 1, 4$
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$$k = 1 : \Theta_{1,1}^1$$

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It is not obvious which ones to use, but.....

#### Simplifying trick:

Go into homogeneous coordinates. Then the only relevant level is  $k = 1$ .

Hence, consider forms of weight two of the form

$$\Theta_{1,1}^1(q^a)\Theta_{1,1}^1(q^b).$$

The conductor of the Hasse-Weil form is 144, hence the simplest case is

$$\Theta_1^1(q^6)^2 = q - 4q^7 + 2q^{13} + 8q^{19} - 5q^{25} - 4q^{31} + \dots$$

Comparison with the Hasse-Weil modular form

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**Theorem.** *The Mellin transforms of the Hasse-Weil L-functions of the elliptic curves  $E_i$  are modular forms that are determined by Hecke indefinite modular forms of weight one, given as*

$$\begin{aligned} k=1: \quad f_{HW}(E^3, q) &= \Theta_{1,1}^1(q^3)\Theta_{1,1}^1(q^9) \in S_2(\Gamma_0(27)) \\ k=2: \quad f_{HW}(E_4, q) &= \Theta_{1,1}^2(q^4)^2 \otimes \chi_2 \in S_2(\Gamma_0(64)) \\ k=3: \quad f_{HW}(E_6, q) &= \Theta_{1,1}^1(q^6)^2 \otimes \chi_3 \in S_2(\Gamma_0(144)). \end{aligned}$$

**Corollary.** *The Jacobi theta function  $\vartheta_3(\tau)$  at level  $k=2$  is determined by the twisted Hasse-Weil modular form*

$$\vartheta_3(\tau) = \frac{1}{\sqrt{f_{HW}(E_6, q^{1/6})}}.$$

Result: It is possible to provide a string theoretic interpretation of the Hasse-Weil modular forms for the class of Brieskorn-Pham elliptic curves in terms of modular forms derived from Kac-Moody algebras.

Arithmetic geometry

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$$\begin{aligned} k=1: \quad f_{HW}(E^3, q) &= \Theta_{1,1}^1(q^3)\Theta_{1,1}^1(q^9) \in S_2(\Gamma_0(27)) \\ k=2: \quad f_{HW}(E_4, q) &= \Theta_{1,1}^2(q^4)^2 \otimes \chi_2 \in S_2(\Gamma_0(64)) \\ k=3: \quad f_{HW}(E_6, q) &= \Theta_{1,1}^3(q^6)^2 \otimes \chi_3 \in S_2(\Gamma_0(144)). \end{aligned}$$

**Corollary.** *The Jacobi theta function  $\tilde{\phi}_{1,1}(\tau)$  at level  $k=2$  is determined by the twisted Hasse-Weil modular form*

$$\tilde{\phi}_{1,1}(\tau) = \frac{1}{\eta^2(\tau)} \sqrt{f_{HW}(E_4, q^{1/4})}.$$

**Result:** It is possible to provide a string theoretic interpretation of the Hasse-Weil modular forms for the class of Brieskorn-Pham elliptic curves in terms of modular forms derived from Kac-Moody characters.

Arithmetic geometry

at  $k=1$  none of the affine characters  $\chi$  of  $A_1^{(1)}$  are useful.

either are the  $N=2$  characters  $\chi_{E_6}^{\pm}$  or theta functions or are the string functions  $c_m'$  which (essentially) describe the characters of parafermionic theories.

the weight of  $c_m'$  is  $-\frac{1}{2}$ , hence the weight of

$$\Theta_m' = \eta^3(\tau)c_m'$$

unity

can make weight 2 with products

$$\Theta_n'(a\tau)\Theta_m'(b\tau)$$

In integers  $a, b$ ,

insight into  $a, b$  via Weil's conductor/level observation. All known examples the level of the modular form agrees with the geometric conductor.

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**Arithmetic geometry**  $\longleftrightarrow$  **String theory**  
 provides a string-theoretic link between the string theory and the moduli space of elliptic curves.

The weight of  $c_m$  is  $-\frac{1}{2}$ , hence the weight of  $\Theta_m^{\ell}$  is  $\ell$ .  
 $\Theta_m^{\ell} \rightarrow \eta^3$  (the past).

Can make weight 2 with products

$$\Theta_n^{\ell}(a\tau)\Theta_m^{\ell}(b\tau)$$

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**Conductors:** The  $E_i$  have geometric conductors  $N_i$ , 144.

- At  $k = 1$  none of the affine characters  $\chi_\ell$  of  $A_1^{(0)}$  are

• Neither are the  $N = [H : \mathbb{Q}]$  characters  $\chi_{\mathbb{Q}, \ell}$ , or theta functions of the Hecke modular forms.

• Nor are the strong functions, which (essentially) describe factors of polynomial invariants of finite elliptic curves.

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The link

$$\text{Arithmetic geometry} \longrightarrow \text{Affine Lie algebras}$$

provides a stronger link between the string and spacetime than the considerations of the past.

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• Neither are the  $N = 1$  forms  $\chi_{q,\ell}^A$ , or theta functions  $\Theta_{q,\ell}$ .

**Theorem.** For the string functions  $\chi_{q,\ell}$ , which (essentially) determine the elliptic curves  $E_\ell$ , are modular forms that are determined by Hecke and infinite level Hecke actions of polynomial characters of  $SL(2, \mathbb{Z})$ .

**Corollary.** The weight of  $c_{m,n}^\ell$  is  $-1 + m\ell$ . The weight of  $\eta^3(\tau)c_m^\ell$  is unity.

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with integers  $a_{\lambda,b}$ .

• Insight into  $a_{\lambda,b}$  via Weil representation in all known examples ( $\Gamma_0(27)$ ,  $\Gamma_0(54)$ ,  $\Gamma_0(162)$ )

**Conductors:** The  $E_\ell$  have geometric conductors  $N$  –

$$N = \frac{1}{2\pi} \sqrt{\frac{1}{2} \int_{-\infty}^{\infty} L(E_\ell, s)^2 ds}$$

**Result:** It is possible to provide a string theoretic interpretation of the Hasse-Weil modular forms for the class of Brieskorn-Pham elliptic curves in terms of modular forms derived from Kac-Moody characters.

The link

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**Remaining questions:**

- 1) What is the physical meaning meaning of the characters  $\chi_i$  (if any)?
- 2) Why no field at  $k = 1$ ?

**Theorem.** *The twist characters  $\chi_n$  are the quadratic characters associated to the field of quantum dimensions of the underlying affine Lie algebra  $A_1^{(1)}$ .*

**Proof.** The field  $K = \mathbb{Q}(\{Q_{\ell,m}\})$  of quantum dimensions

$$Q_{\ell,m} := \frac{S_{\ell m}}{S_{0m}}$$

arises from the

Here the  $S_{\ell m}$  are elements of the matrix

$$S_{\ell m} = \sqrt{\frac{2}{k+2}} \sin\left(\frac{(\ell+1)(m+1)\pi}{k+2}\right)$$

introduced earlier which describes the modular transformation

$$\tau \mapsto -1/\tau$$

on the  $A_{1,k}^{(1)}$  characters

$$\chi_\ell\left(-\frac{1}{\tau}, \frac{u}{\tau}\right) = e^{\pi i k u^2/2} \sum_m S_{\ell m} \chi_m(\tau, u).$$

Computing these fields explicitly gives the result.  $\square$

**Remark.**

For  $k = 1$  the field is  $K = \mathbb{Q}$ , which is why there is no twist character.



Interesting number fields abound:

- 1) The field of quantum dimensions  $K = \mathbb{Q}(\{Q_{\ell,m}\})$ .
- 2) A priori different are the fields associated to the twist characters  $\chi_n$ .
- 3) The elliptic curves of Brieskorn-Pham type admit complex multiplication, hence there is a CM field.
- 4) The Hecke indefinite modular forms  $\Theta_{\ell,m}^k$  are associated to number fields, hence there are fields  $K_\Theta$ .

Explicitly:

Field	$k = 1$	$k = 2$	$k = 4$
Quantum dimensions	$\mathbb{Q}$	$\mathbb{Q}(\sqrt{2})$	$\mathbb{Q}(\sqrt{3})$
Twist field	$\mathbb{Q}$	$\mathbb{Q}(\sqrt{2})$	$\mathbb{Q}(\sqrt{3})$
CM field	$\mathbb{Q}(\sqrt{-3})$	$\mathbb{Q}(\sqrt{-1})$	$\mathbb{Q}(\sqrt{-3})$
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**String applications:**

- 1) Brieskorn-Pham elliptic curves can be viewed as the simplest compactifications from a 10 dimensional spacetime to 8 dimensions. That's a little boring.
- 2) They also appear as important components of higher dimensional Calabi-Yau varieties, e.g. in elliptic K3 surfaces, such as

$$S_6 = \{(z_0 : z_1 : z_2 : z_3) \in \mathbb{P}_{(1,1,2,2)} \mid z_0^6 + z_1^6 + z_2^3 + z_3^3 = 0\}$$

as well as higher dimensional elliptic fibrations, e.g.

$$X_{12} = \{(z_0 : z_1 : \dots : z_4) \in \mathbb{P}_{(1,1,2,4,4)} \mid z_0^{12} + z_1^{12} + z_2^6 + z_3^3 + z_4^4 = 0\}$$

- 3) They generate factors in the cohomology of higher dimensional varieties via the factorization of Jacobians of higher genus curves a la Faddeev, Gross, Rohrlich and others.

The simplest examples are the Fermat curves of degree 4 and 6.

This establishes the following direction (in this special class)

Spacetime  $\implies$  String Theory

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### From strings to spacetime: elliptic case

#### Two steps:

- 1) Identify 'good' string theoretic modular forms.
- 2) Use ideas of Eichler-Shimura-Deligne-Serre.... to construct varieties.

#### Question:

What are 'good' string theoretic forms?

What are the criteria that single out the three string theoretic forms of theorem 3?

#### Input:

Level $k$	Hecke modular form $\Theta_m^{\epsilon}$
1	$\Theta_{1,1}^1(\tau) = q^{1/12}(1 - 2q - q^2 + 2q^3 + q^4 + 2q^5 + \dots)$
2	$\Theta_{0,0}^2(\tau) = q^{1/16}(1 - 2q + q^3 + q^5 - 2q^6 + 2q^7 + \dots)$ $\Theta_{1,1}^2(\tau) = q^{1/8}(1 - q - 2q^2 + q^3 + 2q^5 + 2q^6 + \dots)$ $\Theta_{2,0}^2(\tau) = q^{9/16}(1 - 2q^2 + 2q^4 + 2q^5 + \dots)$
4	$\Theta_{0,0}^4(\tau) = q^{1/24}(1 - 2q + 2q^3 - q^5 + q^7 + q^{11} + \dots)$ $\Theta_{0,2}^4(\tau) = q^{13/24}(1 - q - q^2 - q^4 + q^5 + q^6 + q^8 + \dots)$ $\Theta_{0,4}^4(\tau) = q^{25/24}(1 - q - q + q^5 + q^7 + q^8 + \dots)$ $\Theta_{1,1}^4(\tau) = q^{5/48}(1 - q - q^2 + q^3 - q^4 + q^5 + q^8 + \dots)$ $\Theta_{1,3}^4(\tau) = q^{23/48}(1 - 2q - q^3 + q^5 + q^6 + 2q^7 + \dots)$ $\Theta_{2,0}^4(\tau) = q^{3/8}(1 - q^2 - q^3 + q^5 + q^8 + q^9 + \dots)$ $\Theta_{2,2}^4(\tau) = q^{1/8}(1 - q - q^3 + 2q^7 + 2q^9 + \dots)$

Hecke indefinite modular forms of weight 1 at conformal levels  
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Hecke indefinite modular forms of weight  $J$  at conformal levels

**Constraint 1.**

Assume building blocks that are cusp forms.

Choice: Hecke indefinite modular forms  $\Theta_{\ell,m}^k$

**Constraint 2.**

Consider only forms of weight 2 of the form

$$\Theta_{\ell_1, m_1}^{k_1}(q^a) \Theta_{\ell_2, m_2}^{k_2}(q^b).$$

**Constraint 3.**

Consider only 'integral' forms.

**Constraint 4.**

Consider only normalized cusp forms.

**Constraint 5.**

Consider only Hecke eigenforms of the above type.

**Constraint 6.**

The level of the modular form must be divisible by the prime factors of  $k+2$ , where  $k$  is the conformal level.

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**Result:** For all three exactly solvable string models there is a **unique** form satisfies all constraints.

Possibly relevant for higher dimensions:

**Constraint 8.**

Consider forms with complex multiplication.



Construction of spacetime from conformal field theory:

Eichler-Shimura Theory

- Suppose we have a cusp form  $f \in S_2(\Gamma_0(N))$  with an expansion

$$f = \sum_n a_n q^n.$$

- Associated to this cusp form is a field

$$K = \mathbb{Q}(\{a_n\}).$$

- The ES construction leads to an abelian variety  $A_f$  of dimension

$$\dim_{\mathbb{C}} A_f = [K : \mathbb{Q}].$$

- $f$  is the Mellin transform of the Hasse-Weil L-function  $L_{HW}(A_f, s)$ .

- Arithmetically, and therefore string theoretically, the variety  $A_f$  carries the important information.

- More generally the idea is to think of the motives  $H^1(A_f)$  as building blocks of string theoretically induced motives.



### Affine Lie algebras and generalized Weierstrass curves

- The interpretation of the Hasse-Weil L-function as the Mellin transform of  $\Theta$ -products is not restricted to the three elliptic Brieskorn-Pham curves.
- The starting point can be more general: consider  $\Theta$ -quotients of

$$\Theta^1 = \Theta_{1,1}^1, \quad \Theta^2 = \Theta_{1,1}^2,$$

which are of weight 2 and Hecke eigenforms.

- Result:

$\Theta$ - Quotient	Conductor	Elliptic curve
$\Theta^1(\tau)\Theta^1(11\tau)$	11	$y^2 + y = x^3 - x^2$
$\Theta^2(\tau)\Theta^2(7\tau)$	14	$y^2 + xy + y = x^3 - x$
$\Theta^1(2\tau)\Theta^1(10\tau)$	20	$y^2 = x^3 + x^2 - x$
$\Theta^2(2\tau)\Theta^2(6\tau)$	24	$y^2 = x^3 + x^2 - x$
$\Theta^1(3\tau)\Theta^1(9\tau)$	27	$y^2 + y = x^3$
$\Theta^2(4\tau)^2$	32	$y^2 = x^3 - x$
$\Theta^1(6\tau)^2$	36	$y^2 = x^3 + 1$
$\frac{\Theta^1(8\tau)^4}{\Theta^1(4\tau)\Theta^1(16\tau)}$	64	$y^2 = x^3 + x$
$\frac{\Theta^1(4\tau)^3\Theta^1(20\tau)^3}{\Theta^1(2\tau)\Theta^1(8\tau)\Theta^1(10\tau)\Theta^1(40\tau)}$	80	$y^2 = x^3 - x^2 - x$
$\frac{\Theta^1(12\tau)^6}{\Theta^1(6\tau)^2\Theta^1(24\tau)^2}$	144	$y^2 = x^3 - 1$

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- The curves of conductor  $N = 27, 64, 144$  in the list are the ones in the theorem:

$$\frac{\Theta^1(8\tau)}{\Theta^1(4\tau)\Theta^1(16\tau)} = \Theta^1(4\tau)^2 \otimes \chi_2$$

$$\frac{\Theta^1(12\tau)}{\Theta^1(6\tau)^2\Theta^1(24\tau)} = \Theta^1(6\tau)^2 \otimes \chi_3.$$

- The curves with  $N = 27, 32, 36, 64, 144$  have complex multiplication and

$$f(E_{64}, q) = f(E_{32}, q) \otimes \chi_2$$

$$f(E_{144}, q) = f(E_{36}, q) \otimes \chi_3$$

- Hence, modulo twists, the only elliptic curves in this list with a known string interpretation are the ones with complex multiplication, confirming our conjecture that exactly solvable Calabi-Yau varieties admit complex multiplication (in an appropriate sense).



## Arithmetic aspects of D-branes

### 1) Exactly solvable D-branes and number fields

General boundary state is expanded in terms of Ishibashi states  $|j\rangle\rangle_{\alpha}$  as

$$|\alpha\rangle_{\alpha} = \sum_{j \in E} \frac{\psi_{\alpha}^j}{\sqrt{S_{0j}}} |j\rangle\rangle_{\alpha},$$

where

$$E = \{j \in \mathcal{T} \mid j = w(\bar{j}), N_{w(j)\bar{j}} \neq 0\}.$$

The coefficients  $\psi_{\alpha}^j$  are determined by the Cardy condition

$$n_{\alpha\beta}^i = \sum_{j \in E} \frac{S_{ij}}{S_{0j}} \psi_{\alpha}^j (\psi_{\beta}^j)^*,$$

where

$$\mathcal{H}_{\alpha\beta} = \bigoplus n_{\alpha\beta}^i \mathcal{H}_i$$

with bulk decomposition

$$\mathcal{H} = \bigoplus_i \mathcal{H}_i.$$

Here the  $S_{ij}$  are elements of the matrix

$$S_{ij} = \sqrt{\frac{2}{k+2}} \sin \left( \frac{(i+1)(j+1)\pi}{k+2} \right)$$

which describes the modular transformation  $\tau \mapsto -1/\tau$  on the  $A_{1,k}^{(1)}$  characters

$$\chi_i \left( -\frac{1}{\tau}, \frac{u}{\tau} \right) = e^{\pi i k u^2/2} \sum_j S_{ij} \chi_j(\tau, u).$$

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