

Title: Motives and Strings. (no audio)

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Abstract: Please Note: There is no audio for most presentations from November 18th due to technical issues. We appologize to presenters and interested parties for this inconvenience.

Perimeter Institute, 19 November 2004  
Workshop on Mirror Symmetry

## Motives and Strings

Jan Stienstra

'Does the *Theory of Everything* include Numbers?'

'The *M* in *M-theory* seems to stand for  
*magic* or *mystery* or *mother of all strings*,  
but could it also mean *Motive*?'



## What is ..... a Motive ?

Barry Mazur in Notices of the AMS november 2004:

Algebraic topology: one cohomology theory, characterized by the Eilenberg-Steenrod axioms, representable by Eilenberg-MacLane spaces and Postnikov towers.

Algebraic geometry: a profusion of different cohomology theories, no axioms, no representing objects:

- Hodge cohomology
- algebraic De Rham cohomology
- crystalline cohomology (for every prime number  $p$ )
- étale  $\ell$ -adic cohomology (for every prime number  $\ell$ )

and comparison isomorphisms between these.

Cohomology theories assign to varieties vector spaces (some with additional structure).

Inverse problem: *given a collection of vector spaces (+ additional structure, comparison isomorphisms), is this produced by the various cohomology theories from one common source?*

*The common source is called a MOTIVE.*

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*The common source is called a MOTIVE.*

More precisely:  
a variety  $X$  is defined over some field  $K$ ,  $\text{char}(K) = 0$ ,  
often with positive transcendence degree over  $\mathbb{Q}$ .

Every embedding  $\sigma$  of  $K$  into  $\mathbb{C}$  leads to a complex variety to which algebraic topology and differential geometry with their cohomology theories apply.

Every embedding  $\sigma$  of  $K$  into a  $p$ -adic field  $\widehat{\mathbb{Q}}_p$  leads to crystalline cohomology spaces.

As the embeddings  $\sigma$  vary, with fixed target  $\mathbb{C}$  or  $\widehat{\mathbb{Q}}_p$ , the variation in the cohomology spaces is described (partly by the action of Galois groups and partly) by *Picard-Fuchs differential equations* and *Gauss-Manin connections*.

*As yet no theory to mix the targets  $\mathbb{C}$  and/or the various  $\widehat{\mathbb{Q}}_p$ !*



Why look for MOTIVE-STRING relation ?

Computations in Type IIb string theory proceed by manipulating solutions of certain differential equations. During the computations there are many denominators. In the end these drop out and true integers remain.

Many differential equations in Type IIb string can be recognized as Picard-Fuchs equations in De Rham cohomology of families of varieties.

The integrality statements can be recognized as consequences of theorems about the crystalline cohomology of families of ordinary varieties.

Challenge for Motive people:  
Crystalline cohomology deals with only one prime  $p$  at a time and puts out statements about  $p$ -adic integrality.

*What mechanism synchronizes the primes  
and leads to true integers?*

Challenge for String people:  
Crystalline cohomology implies extra symmetries in the differential equations.

*Where are these extra symmetries in Nature?*





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## Line & Additive Group



arclength  $\int_a^b dx = b - a$

invariant differential form  $dx$

invariant derivation  $\frac{d}{dx}$

$$dx = dz \Rightarrow z = x + y, \quad y \text{ constant}$$

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Circle



equations for the unit circle in the plane:

$$\begin{array}{ll} r = 1 & \text{polar coordinates} \\ x^2 + y^2 = 1 & \text{Cartesian coordinates} \end{array}$$

arclength  $\int_a^b \frac{dx}{\sqrt{1-x^2}} = \arcsin(b) - \arcsin(a)$

$$\begin{aligned} \arcsin(x) &= \sum_{n \geq 0} (-1)^n \binom{-\frac{1}{2}}{n} \frac{x^{2n+1}}{2n+1} \\ &= 2 \sum_{n \geq 0} \frac{(2n)!}{n!^2} \frac{(x/2)^{2n+1}}{2n+1} \end{aligned}$$

better

$$\frac{1}{2} \arcsin(2x) = \sum_{n \geq 0} \frac{(2n)!}{n!^2} \frac{x^{2n+1}}{2n+1}$$

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$$= \sum_{n=0}^{\infty} \frac{(-2n)!}{n!} \frac{(x^2)^{n+1}}{2n+1}$$

Arxiv

invariant differential form  $\frac{dx}{\sqrt{1-x^2}}$

invariant derivation  $\sqrt{1-x^2} \frac{d}{dx}$

Addition law for trigonometric functions:

$$\frac{dx}{\sqrt{1-x^2}} = \frac{dz}{\sqrt{1-z^2}} \Rightarrow$$

$$z = \sin(\arcsin(x) + \arcsin(y)) \quad y \text{ constant}$$

$$= x\sqrt{1-y^2} + y\sqrt{1-x^2}$$

$$= x + y + \sum_{n \geq 1} (-1)^n \binom{\frac{1}{2}}{n} (xy^{2n} + yx^{2n})$$

Better, with  $z = 2w$ ,  $x = 2u$ ,  $y = 2v$ ,

$$w = u + v - 4 \sum_{n \geq 1} \frac{(2n-3)!}{n!(n-2)!} (uv^{2n} + vu^{2n})$$

The coefficients are integers:

$$-4 \frac{(2n-3)!}{n!(n-2)!} = 2 \binom{2n-2}{n} - 4 \binom{2n-3}{n-1}$$

So, this is a formal group law over  $\mathbb{Z}$ !

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## Multiplicative Group

arclength  $\int_a^b \frac{dx}{x} = \log(b) - \log(a)$

invariant differential form  $\frac{dx}{x}$

invariant derivation  $x \frac{d}{dx}$

$$\frac{dx}{x} = \frac{dz}{z} \Rightarrow z = xy, \quad y \text{ constant}$$

Coordinate change  $z = 1 + w$ ,  $x = 1 + u$ ,  $y = 1 + v$   
gives

$$w = u + v + uv,$$

*the standard multiplicative formal group law over  $\mathbb{Z}$ .*

Applying the chain rule

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Coordinate change  $z = 1 + w, x = 1 + u, y = 1 + v$   
also

$$w = u + v + uv,$$

the standard multiplicative formal group law over  $\mathbb{Z}$ .

## Circle & Multiplicative Group

Substitution  $x \rightarrow ix, y \rightarrow iy, z \rightarrow iz$   
transforms the addition law for trigonometric sine:

$$z = x\sqrt{1-y^2} + y\sqrt{1-x^2}$$

into the addition law for hyperbolic sine:

$$z = x\sqrt{1+y^2} + y\sqrt{1+x^2}$$

i.e.

the two formal group laws are isomorphic over  $\mathbb{Z}[i]$ .

The series

$$\begin{aligned}\sinh(\log(1+u)) &= \frac{1}{2}((1+u) - (1+u)^{-1}) \\ &= u - \frac{1}{2} \sum_{n \geq 2} (-u)^n\end{aligned}$$

establishes an isomorphism, defined over  $\mathbb{Z}[\frac{1}{2}]$ ,  
between the multiplicative group law and the  
addition law for the hyperbolic sine.

*The addition law for the sine function and  
the multiplicative group law are isomorphic  
over  $\mathbb{Z}[i][\frac{1}{2}]$ .*





## Lemniscate & its group law



equations Bernoulli's Lemniscate in the plane:

$$\begin{array}{ll} r^2 = \cos(2\phi) & \text{polar coordinates} \\ (x^2 + y^2)^2 = x^2 - y^2 & \text{Cartesian coordinates} \end{array}$$

$$\text{arclength} \int_a^b \frac{dx}{\sqrt{1-x^4}}$$

$$\begin{aligned} \int \frac{dx}{\sqrt{1-x^4}} &= \sum_{n \geq 0} (-1)^n \binom{-\frac{1}{2}}{n} \frac{x^{4n+1}}{4n+1} \\ &= \sqrt{2} \sum_{n \geq 0} \frac{(2n)!}{n!^2} \frac{(x/\sqrt{2})^{4n+1}}{4n+1} \end{aligned}$$

Coefficients can be made integral by substitution  
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Euler

$$\frac{dx}{\sqrt{1-x^4}} = \frac{dz}{\sqrt{1-z^4}} \Rightarrow$$
$$z = \frac{x\sqrt{1-y^4} + y\sqrt{1-x^4}}{1+x^2y^2}, \quad y \text{ constant}$$

Euler's result was the first example of an addition law for elliptic integrals.

*This marked the beginning of the theory of elliptic curves !*

The elliptic curve in this case is, in homogeneous coordinates in the weighted projective plane  $\mathbb{P}^{[1,1,2]}$ ,

$$X^4 + Y^4 + Z^2 = 0$$

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The elliptic curve in this case is, in homogeneous coordinates in the weighted projective plane  $\mathbb{P}^{1,1,4}$ ,

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One example of a Del Pezzo surface

$dP_7$  is:

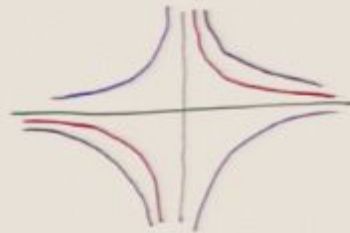
$$W^4 + X^4 + Y^4 + Z^2 = 0 \quad \text{in } \mathbb{P}^{[1,1,1,2]}$$

This is a branched double cover of  $\mathbb{P}^{[1,1,1]}$ ;

On  $\mathbb{P}^{[1,1,1]} = \mathbb{P}^2$ , the usual projective plane,

there is the pencil of conics

$$W^2 - vXY = 0$$



This pencil of conics on  $\mathbb{P}^2$  pulls back to a pencil of elliptic curves on  $dP_7$ :

$$v^2 X^2 Y^2 + X^4 + Y^4 + Z^2 = 0$$

Substitution  $Z \mapsto Z - ivXY$  gives

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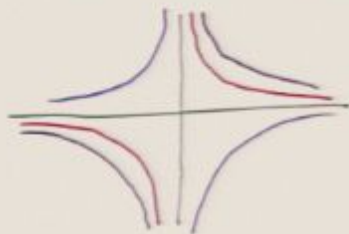
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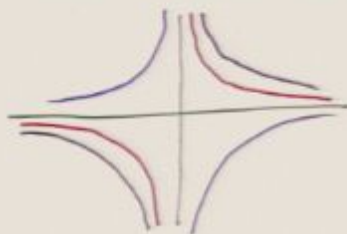
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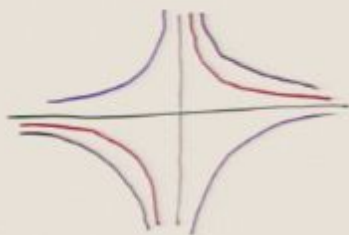
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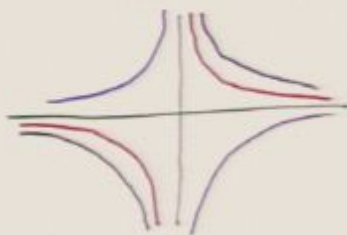
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$$X^4 - Y^4 - Z^2 - tXYZ = 0$$

with  $t$  a complex (deformation) parameter.

Viewed in  $\mathbb{C}^3$  this equation has for  $t^4 \neq 64$  only one singular point, namely  $(0, 0, 0)$ .

*This is a so-called simple elliptic singularity known as the  $\tilde{E}_7$  singularity.*



Viewed in  $\mathbb{P}^{[1,1,2]}$  this equation describes for  $t^4 \neq 64$  a smooth elliptic curve.

*We have a family of elliptic curves with singular fibres at  $t = \pm 2\sqrt{2}, \pm 2\sqrt{-2}, \infty$ .*

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The elliptic curve

$$X^4 + Y^4 + Z^2 - tXYZ = 0$$

a rank 2 *period lattice*.

One period (of a suitably normalized 1-form along a suitable closed path) can be computed via the residue theorem:

with  $u = t^{-1}$  and  $|u|$  sufficiently small:

$$\begin{aligned} f_0 &= \frac{1}{2\pi i} \oint_{\gamma_0} \frac{-t dx}{2z - tx} \Big|_{x^4+1+z^2=txz, |z|<1} \\ &= \left(\frac{1}{2\pi i}\right)^2 \oint \oint \frac{-t dx dz}{x^4 + 1 + z^2 - txz} \\ &= \left(\frac{1}{2\pi i}\right)^2 \oint \oint \frac{1}{1 - ux^{-1}z^{-1}(x^4 + 1 + z^2)} \frac{dx dz}{x z} \\ &= \sum_{n \geq 0} u^n \left(\frac{1}{2\pi i}\right)^2 \oint \oint \left(\frac{x^4 + 1 + z^2}{xz}\right)^n \frac{dx dz}{x z} \\ &= \sum_{m \geq 0} \frac{(4m)!}{m!^2(2m)!} u^{4m} \end{aligned}$$



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A second period

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can be determined from the formula (with  $\epsilon^2 = 0$ )

$$f_0 + f_1\epsilon \equiv \sum_{m \geq 0} \frac{(1+4\epsilon)_{im}}{(1+\epsilon)_m^2 (1+2\epsilon)_{2m}} u^{im+4\epsilon}$$

using the rising Pochhammer symbol: for  $k \geq 0$

$$(a)_k := a(a+1) \cdots (a+k-1)$$

(so  $(1)_k = k!$ )

Note

$$f_1 = 4f_0 \log u + g_1$$

where  $g_1$  is a power series in  $u^{\frac{1}{4}}$  with constant term 0.

Thus if we define  $\tau$  and  $q$  by

$$\tau := \frac{f_1}{f_0}, \quad q := \exp(\tau) = u^{\frac{1}{4}} \exp\left(\frac{g_1}{f_0}\right)$$

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Page 1

1. Let  $V = \mathbb{R}^n$  and  $W = \mathbb{R}^m$ .  
Define a linear map  $T: V \rightarrow W$  by

$T(x_1, x_2, \dots, x_n) = (x_1 + x_2, x_2 + x_3, \dots, x_{n-1} + x_n)$

Find the matrix of  $T$  relative to the standard bases of  $V$  and  $W$ .

Answer:  $\begin{bmatrix} 1 & 1 & 0 & \dots & 0 \\ 0 & 1 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \end{bmatrix}$

$$2\pi i \int_{\gamma_1} \frac{2z - tx}{z^4 + 1 + z^2 - txz}, |z| < 1$$

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$$u^4 = \frac{\eta(q^2)^{14}}{\eta(q)^{24} + 64 \eta(q^2)^{24}}$$

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$$\begin{aligned} \frac{1}{L^4} = u^{-4} &= 64 + \frac{\eta(q)^{24}}{\eta(q^2)^{24}} \\ &= q^{-1} + 40 + 276q + \dots \end{aligned}$$

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More formal group laws.

Put

$$\ell(u) = \int f_0 du = \sum_{m \geq 0} \frac{(4m)!}{m!^2(2m)!} \frac{u^{4m+1}}{4m+1}$$

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$$L(x) = \frac{1}{\sqrt{2}} \int \frac{dx}{\sqrt{1-4x^4}} = \sum_{n \geq 0} \frac{(2n)!}{n!^2} \frac{x^{4n+1}}{4n+1}$$

Then

$$\begin{aligned} \ell^{-1}(\ell(u) + \ell(v)) \\ L^{-1}(L(x) + L(y)) \end{aligned}$$

are two formal group laws over  $\mathbb{Z}$  and they are isomorphic over  $\mathbb{Z}[i]$ .

Moreover  $L^{-1}(L(x) + L(y))$  is the integer version of the addition law for the lemniscate.

i.e. the base of the elliptic pencil

$$X^4 + Y^4 + Z^2 - tXYZ = 0$$

carries in the neighborhood of  $t = \infty$ ,  $u = 0$  a formal group law over  $\mathbb{Z}$ , in the coordinate  $u$ , which is over  $\mathbb{Z}[i]$  isomorphic to the formal group law of the fiber at  $t = 0$ ,  $u = \infty$ .

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## Seiberg-Witten

‘Electric-Magnetic Duality, Monopole Condensation, and Confinement in  $N = 2$  Supersymmetric Yang-Mills Theory’

illustrate their general theory with a first example starting the functions

$$\begin{aligned}\mathbf{a} &= \int f_0 dt \\ \mathbf{a}_D &= \int f_1 dt\end{aligned}$$

They show that  $\mathbf{a}$  and  $\mathbf{a}_D$  give the periods of some meromorphic 1-form without residues on the elliptic curve  $X^2 + Y^2 + Z^2 - tXYZ = 0$ .

Modulo  $\epsilon^2$ :

$$\mathbf{a} + \mathbf{a}_D \epsilon \equiv - \sum_{m \geq 0} \frac{(1 - 4\epsilon)_{4m}}{(1 + \epsilon)_m^2 (1 - 2\epsilon)_{2m}} \frac{u^{4m-1-4\epsilon}}{4m - 1 + 4\epsilon}$$

## Sethna, Witten

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illustrate their general theory with a first example stat-  
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$$\mathfrak{a} = \int \psi, \mathfrak{a}$$
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They show that it includes also the part of  $\mathfrak{a}$  con-  
monopoles,  $\mathfrak{a}$  from without or within  
on the elliptic curve  $\mathfrak{a}^2 = \mathfrak{a}^3 + \mathfrak{a} + \mathfrak{a}^2 + \mathfrak{a}$

Module  $\mathfrak{a}^2$

$$\mathfrak{a} \in \mathfrak{a}^2 = \sum_{k \in \mathbb{Z}} \frac{e^{2\pi i k \mathfrak{a}}}{1 - e^{2\pi i k \mathfrak{a}}} = \sum_{k \in \mathbb{Z}} \frac{e^{2\pi i k \mathfrak{a}}}{1 - e^{2\pi i k \mathfrak{a}}}$$

The functions  $\mathbf{a}$  and  $\mathbf{a}_D$  are used to construct the potential

$$K := \frac{1}{2i}(\bar{\mathbf{a}}\mathbf{a}_D - \mathbf{a}\bar{\mathbf{a}}_D)$$

for a so-called *rigid special Kähler metric* on the base space (Moduli space) of the pencil: the metric is

$$\begin{aligned} \frac{\partial^2 K}{\partial \mathbf{a} \partial \bar{\mathbf{a}}} d\mathbf{a} d\bar{\mathbf{a}} &= \frac{1}{2i} \left( \frac{\partial \mathbf{a}_D}{\partial \mathbf{a}} - \frac{\partial \bar{\mathbf{a}}_D}{\partial \bar{\mathbf{a}}} \right) d\mathbf{a} d\bar{\mathbf{a}} \\ &= \frac{1}{2i} \left( \frac{f_1}{f_0} - \frac{\bar{f}_1}{\bar{f}_0} \right) f_0 \bar{f}_0 dt d\bar{t} \\ &= -2i(\tau - \bar{\tau}) f_0 \bar{f}_0 dt d\bar{t} \end{aligned}$$

Note, with  $\bar{\epsilon} := -\epsilon$ , we can write

$$(\mathbf{a} + \mathbf{a}_D \epsilon) \overline{(\mathbf{a} + \mathbf{a}_D \epsilon)} = \mathbf{a}\bar{\mathbf{a}} + 2iK\epsilon$$

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$$V = \frac{1}{2} B u_x - u u_{xx}$$

It is assumed equal period Euler method on the first square element point, then when the nodes are

$$\frac{\partial V}{\partial u} \frac{du}{dx} = \left( \frac{\partial u_x}{\partial u} - \frac{\partial u_{xx}}{\partial u} \right) \frac{du}{dx}$$

$$= \left( \frac{1}{l} - \frac{1}{l} \right) \frac{1}{l} \frac{du}{dx}$$

$$= 0 \cdot \frac{1}{l} \frac{du}{dx}$$

where  $l$  is the length of the element

$$u = u_1, u = u_2, \dots, u = u_n$$

## A Glimpse of Mirror Symmetry

Golyshev looked at the system of differential equations

$$\frac{d}{dt}\phi = (t - A)^{-1}P\phi$$

where

$$A = \begin{pmatrix} 12 & 552 & 7488 \\ 1 & 40 & 552 \\ 0 & 1 & 12 \end{pmatrix}, \quad P = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

He showed that it has a solution

$$\phi = \begin{pmatrix} * & * & 0 \\ * & * & 0 \\ \mathfrak{a}_D & \mathfrak{a} & 1 \end{pmatrix}$$

and he showed that the entries of  $A$  count how many curves of certain kinds there exist on the DelPezzo surface  $dP_7$ , i.e. the blow up of  $\mathbb{P}^2$  at seven points; more precisely  $A$  is the matrix for multiplication with the anti-canonical class in the quantumcohomology of  $dP_7$ .

Note: classical theory relates this DelPezzo to the root system  $E_7$ .

Notice the coordinate translation  $t \mapsto t - \lambda$  has the same effect as replacing  $A$  by  $A - \lambda$ .

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Note: classical theory relates this DelPezzo to the root system  $E_7$ .

Notice the coordinate translation  $t \mapsto t + \lambda$  has the same effect as replacing  $A$  by  $A - \lambda$ .



## A Glimpse of Mirror Symmetry

Golyshev looked at the system of differential equations

$$\frac{d}{dt}\phi = (t - A)^{-1}P\phi$$

where

$$A = \begin{pmatrix} 12 & 552 & 7488 \\ 1 & 40 & 552 \\ 0 & 1 & 12 \end{pmatrix}, \quad P = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

He showed that it has a solution

$$\phi = \begin{pmatrix} \alpha & \alpha & 0 \\ \alpha & \alpha & 0 \\ \alpha_D & \alpha & 1 \end{pmatrix}$$

and he showed that the entries of  $A$  count how many curves of certain kinds there exist on the DelPezzo surface  $dP_7$ , i.e. the blow up of  $\mathbb{P}^2$  at seven points; more precisely  $A$  is the matrix for multiplication with the anti-canonical class in the quantumcohomology of  $dP_7$ .

Note: classical theory relates this DelPezzo to the root system  $E_7$ .

Notice the coordinate translation  $t \mapsto t + A$  has the

## A Glimpse of Mirror Symmetry

Gibson looked at the system of differential equations

$$\frac{d}{dt}\Phi = t^{-1}A^{-1}t^2\Phi$$

where

$$A = \begin{pmatrix} 12 & 752 & 7188 \\ 1 & 40 & 752 \\ 0 & 1 & 12 \end{pmatrix} \quad P = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

He showed that it has a solution

$$\Phi = \begin{pmatrix} x + 0 \\ x + 0 \\ a_1 + a_2 t \end{pmatrix}$$

and he showed that the entries of  $A$  count how many curves of certain kinds there exist on the Fano variety surface  $\mathbb{P}^2$ , i.e. the blow up of  $\mathbb{P}^2$  at seven points, more precisely  $A$  is the matrix for multiplication with the unit canonical class in the quantum cohomology of  $\mathbb{P}^2$ .

Note: classical theory relates this Fano variety to the rest system  $E_7$ .

$(A, B, \lambda)$

1.  $(432, 0, 60)$

$A$

$$\begin{pmatrix} 60 & 20520 & 1339200 \\ 1 & 312 & 20520 \\ 0 & 1 & 60 \end{pmatrix}$$

2.  $(64, 0, 12)$

$$\begin{pmatrix} 12 & 552 & 7488 \\ 1 & 40 & 552 \\ 0 & 1 & 12 \end{pmatrix}$$

3.  $(27, 0, 6)$

$$\begin{pmatrix} 6 & 108 & 756 \\ 1 & 15 & 108 \\ 0 & 1 & 6 \end{pmatrix}$$

4.  $(16, 0, 4)$

$$\begin{pmatrix} 4 & 40 & 192 \\ 1 & 8 & 40 \\ 0 & 1 & 4 \end{pmatrix}$$

5.  $(11, -1, 3)$

$$\begin{pmatrix} 3 & 20 & 75 \\ 1 & 5 & 20 \\ 0 & 1 & 3 \end{pmatrix}$$

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$$1. \quad \frac{u_m}{m!(2m)!(3m)!}$$

$$c_n = u_m \quad \text{if } 6|n \quad m = \frac{n}{6}$$

$$2. \quad \frac{(4m)!}{m!^2(2m)!}$$

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$$5. \quad \sum_{k=0}^m \binom{m}{k}^2 \binom{m+k}{k}$$

$$m = n$$

$$6. \quad \sum_{k=0}^m \binom{m}{k}^3$$

$$m = n$$

$c_n = 0$  otherwise

$c_n =$  constant term of  $p(x,y)^n$

Recurrence relation

$$(m+1)^2 u_{m+1} - (Am^2 + Am + \lambda) u_m + Bm^2 u_{m-1} = 0$$

$$\begin{array}{ll}
 1. & \frac{u_m}{m!(2m)!(3m)!} \quad \text{if } 6|n \quad m = \frac{n}{6} \\
 2. & \frac{(4m)!}{m!^2(2m)!} \quad \text{if } 4|n \quad m = \frac{n}{4} \\
 3. & \frac{(3m)!}{m!^3} \quad \text{if } 3|n \quad m = \frac{n}{3} \\
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 \end{array}$$

$$5. \quad \sum_{k=0}^m \binom{m}{k}^2 \binom{m+k}{k} \quad m=n$$

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Recurrence relation

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$$(m+1)^2 u_{m+1} - (Am^2 + Am + \lambda) u_m + Bm^2 u_{m-1} = 0$$

$$u_0 = 1$$





1.  $X^2 + Y^3 + Z^6$

2.  $X^2 + Y^4 + Z^4$

3.  $X^3 + Y^3 + Z^3$

4.  $(X+Y)(XY+Z^2)$

5.  $(X+Y+Z)(X+Z)(Y+Z)$

6.  $(X+Y)(Y+Z)(Z+X)$

$\underbrace{\hspace{10em}}$   
 $P(X, Y, Z)$

$$p(x, y) = x^{-1} y^{-1} P(x, y, 1)$$

$$p: \mathbb{C}^* \times \mathbb{C}^* \rightarrow \mathbb{C}$$

level sets are elliptic curves  
(minus a few points)<sup>+</sup>

critical values  $\Leftrightarrow$  singular fibres

physics:  $p$  is Landau - Ginzburg potential

interested in critical points of  $P$



More formal group laws.

Recall

$$\alpha = \int f_0 dt = - \int f_0 \frac{du}{u^2} = - \sum_{m \geq 0} \frac{(4m)!}{m!^2(2m)!} \frac{u^{4m-1}}{4m-1}$$

Note that the coefficients

$$\begin{aligned} \frac{(4m)!}{m!^2(2m)!} \frac{1}{4m-1} &= 4 \frac{(4m-3)!}{m!^2(2m-2)!} \\ &= -2 \frac{(4m-2)!}{m!^2(2m-2)!} + 8 \frac{(4m-3)!}{m!(m-1)!(2m-2)!} \end{aligned}$$

are sums of multinomial coefficients and, hence, are integers.

This can also be stated as: *The function*

$$\frac{1}{\alpha} = u + 4u^5 + 76u^9 + 2224u^{13} + \dots$$

*is the logarithm of a formal group law over  $\mathbb{Z}$  in the coordinate  $u$  which is over  $\mathbb{Z}$  isomorphic to the additive group law.*

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Chloroformol group has

$n = 0$

$$n = \int_0^{\infty} \frac{1}{x} dx = \int_0^{\infty} x^{-1} dx = \sum_{k=1}^{\infty} \frac{1}{k} \left( \frac{1}{k} - \frac{1}{k+1} \right)$$

or  $n = 1 - \ln 2$

Char $^n$  =  $\frac{1}{x} = \frac{1}{x} \cdot \frac{1}{1-x} = \frac{1}{x} \sum_{k=0}^{\infty} x^k$

$$= \sum_{k=0}^{\infty} x^{k-1} = \sum_{k=-1}^{\infty} x^k \quad \text{Char } n = 1$$

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If the characteristic of a ring is  $n$ , then  $n \cdot 1 = 0$ .  
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Lerche, Mayr, Warner  
'Non-critical Strings, Del Pezzo Singularities  
and Seiberg-Witten curves'

work in one of their examples with the functions

$$\mathfrak{b} = \frac{1}{4} \int f_0 \frac{du}{u}$$
$$\mathfrak{b}_D = \frac{1}{4} \int f_1 \frac{du}{u}$$

They relate these to periods of *multivalued 1-forms* on the elliptic curve  $X^4 + Y^4 + Z^2 - tXYZ = 0$ .

They also define a function  $\mathcal{F}$  so that

$$\mathfrak{b}_D = \frac{d\mathcal{F}}{d\mathfrak{b}}.$$

The functions  $\mathfrak{b}$  and  $\mathcal{F}$  are then used to construct the potential

$$K = -\log(\operatorname{Im}\mathcal{F} - \operatorname{Im}\mathfrak{b} \operatorname{Re}\mathfrak{b}_D)$$

for a so-called *local special Kähler metric* on the base space (Moduli space) of the pencil.



### Lemma: Moyal-Wigner

Consider an atom with  $\hat{H}(x, p) = \frac{p^2}{2m} + V(x)$  and ordering Weyl  $\rho = \rho_W$

with  $m = m_0$ . Then  $\rho$  couple with the function

$$W = \int \frac{dx}{2\pi} e^{-ix} \hat{H}(x, p)$$

$$W = \int \frac{dx}{2\pi} e^{-ix} \left( \frac{p^2}{2m} + V(x) \right)$$

The  $\rho$  obeys the  $\rho$ - $\rho$  Poisson equation  $\rho \{ \rho, \hat{H} \} = -\rho \hat{H}$  on the algebra  $\rho = \hat{X}^2 + \hat{P}^2 + \dots + \hat{X}\hat{P} + \dots = 0$

The  $\rho$  obeys the  $\rho$ - $\rho$  Poisson equation

$$\rho \{ \rho, W \} = -\rho W$$

The function  $W$  and  $\rho$  are related by  $\rho = \rho_W$  and the potential

$$V(x) = \int \frac{dx}{2\pi} e^{-ix} (W - \frac{p^2}{2m})$$

for  $\rho = \rho_W$  with  $\rho = \rho_W$  and  $\rho = \rho_W$  in the case  $\rho = \rho_W$  and  $\rho = \rho_W$

## From curves to surfaces to threefolds.

The total space of our elliptic pencil can be embedded as the surface

$$\mathcal{S}: \quad (X^4 + Y^4 + Z^2)U - XYZT = 0$$

of bidegree  $(4, 1)$  in  $\mathbb{P}^{[1,1,2]} \times \mathbb{P}^{[1,1]}$ .

The elliptic pencil arises via the projection from  $\mathcal{S}$  onto the projective line  $\mathbb{P}^{[1,1]}$ . Projection onto the weighted projective plane  $\mathbb{P}^{[1,1,2]}$  shows  $\mathcal{S}$  as the blow up of this plane in 8 points.

One DelPezzo surface  $d\mathbb{P}_7$  is given by

$$\mathcal{P}: \quad W^4 + X^4 + Y^4 + Z^2 = 0$$

in weighted projective 3-space  $\mathbb{P}^{[1,1,1,2]}$ .

The elliptic pencil is obtained by intersecting  $\mathcal{P}$  with the pencil of quartics  $W^4U - XYZT = 0$ .

Lerche, Mayr, Warner model their work on the family of non-compact Calabi-Yau threefolds

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$$K = -\log(\operatorname{Im}\mathcal{F} - \operatorname{Im}\mathfrak{b} \operatorname{Re}\mathfrak{b}_D)$$

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$$\begin{aligned}
 & \dots \\
 & \dots \\
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From

$$f_0 + f_1 \epsilon = \sum_{m \geq 0} \frac{(1 + 4\epsilon)_{4m}}{(1 + \epsilon)_{2m}^2 (1 + 2\epsilon)_{2m}} u^{4m+4\epsilon} + \mathcal{O}(\epsilon^2)$$

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Then

$$\mathcal{F} = \frac{1}{2} (\mathfrak{b}\mathfrak{b}_D - \mathfrak{c})$$

and thus

$$\begin{aligned} K &= -\log(\operatorname{Im} \mathcal{F} - \operatorname{Im} \mathfrak{b} \operatorname{Re} \mathfrak{b}_D) \\ &= -\log \left( \frac{i}{4} (\mathfrak{c} - \bar{\mathfrak{c}} + \mathfrak{b}\bar{\mathfrak{b}}_D - \bar{\mathfrak{b}}\mathfrak{b}_D) \right) \\ &= -\log \frac{i}{4} (\Omega | \bar{\Omega}) \end{aligned}$$

where

$$\begin{aligned} \Omega &= e_0 + \mathfrak{b}e_1 + \mathfrak{b}_D e_2 + \mathfrak{c}e_3 \\ \langle e_n | e_m \rangle &= (-1)^m \delta_{n+m,4} \end{aligned}$$



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$$\langle e_n | e_m \rangle = (-1)^m \delta_{n-m, 3}$$

From

$$h(x) = \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \frac{d^k}{dx^k} (a^x - 1) = \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} a^x \ln^k a$$

$$h(x) = \frac{1}{a} \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a$$

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and thus we can define  $h(x) = \frac{1}{a} \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a$

$$f(x) = \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a = a^x \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a$$

Then

$$f(x) = \frac{1}{a} \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a$$

and thus

$$h(x) = \frac{1}{a} \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a = \frac{1}{a} \sum_{k=0}^{\infty} \frac{(x+b)^k}{k!} \ln^k a$$

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## A Glimpse of Mirror Symmetry

The function  $\mathcal{F} - \frac{b^3}{6}$  is a function of  $e^b$ . It can be written as a series

$$\mathcal{F} = \frac{1}{6}b^3 + \sum_{k \geq 1} a_k \text{Li}_3(e^{kb})$$

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$\text{Li}_3$  is the *trilogarithm function*:

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$$\mathcal{F} = \frac{1}{6} \mathfrak{b}^3 + \sum_{k \geq 1} a_k \text{Li}_3(e^{k\mathfrak{b}})$$

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Hence, setting  $Q := e^{\mathfrak{b}}$ , we find

$$q := \exp(\tau) = Q \prod_{k \geq 1} (1 - Q^k)^{-k^2 a_k}$$

So, the numbers  $a_k$  can be recovered directly from the relation between  $q$  and  $Q$ .

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Handwritten text, possibly bleed-through from the reverse side of the page. The text is faint and difficult to decipher but appears to include several lines of writing.

More formal group laws.

Recall

$$\mathfrak{b} = 4 \int f_0 \frac{du}{u} = \log u^4 + \sum_{m \geq 1} \frac{(4m)!}{m!^2(2m)!} \frac{u^{4m}}{m}$$

It can be shown that the series expansion of  $Q$

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is a power series in  $u^4$  with all its coefficients in  $\mathbb{Z}$ .

So if we define  $Q(x) = x \exp \left( \sum_{m \geq 1} \frac{(4m)!}{m!^2(2m)!} \frac{x^m}{m} \right)$

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$$Q^{-1}(Q(x)Q(y)) \in \mathbb{Z}[[x, y]]$$

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Mean normal distribution

mean

$$\mu = \frac{1}{n} \sum_{i=1}^n x_i$$

standard deviation

$$\sigma = \sqrt{\frac{1}{n} \sum_{i=1}^n (x_i - \mu)^2}$$

variance

$$\sigma^2 = \frac{1}{n} \sum_{i=1}^n (x_i - \mu)^2$$

then

$$\sigma^2 = \frac{1}{n} \sum_{i=1}^n x_i^2 - \mu^2$$

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Recall

$$\begin{aligned} \mathfrak{b} &= 4 \int f_0 \frac{du}{u} \\ &= 4 \left( \frac{1}{2\pi i} \right)^2 \int \oint \oint \frac{dx dz dt}{x^3 + 1 + z^2 - txz} \\ &= \frac{1}{\pi^2} \oint \oint \log \left[ \frac{txz - (x^3 + 1 + z^2)}{xz} \right] \frac{dx dz}{x z} \end{aligned}$$

So  $-\frac{1}{\pi} \operatorname{Re} \mathfrak{b}$  is the *logarithmic Mahler measure* of the Laurent polynomial

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$$m(F) := \frac{1}{(2\pi i)^2} \iint_{|x|=|z|=1} \log |F(x, z)| \frac{dx dz}{x z}.$$

C. Smyth, D. Boyd and others found many examples of Laurent polynomials  $F$  for which the (logarithmic) Mahler measure equals up to a rational factor and to many decimal places the value at  $s = 0$  of the derivative of the L-function of the zero locus of  $F$  (suitably compactified)

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with integer values for the parameter  $t$ , is not in their lists: probably one did not yet look for it.

*Thus, special values of L-functions, the main enumerative problem about Motives, appear alongside with instanton counts, the main enumerative problem about Strings*

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1. The first part of the document discusses the importance of maintaining accurate records of all transactions.

2. It also outlines the various methods used to collect and analyze data, including surveys, interviews, and focus groups.

3. The results of the study are presented in a series of tables and graphs, showing a clear trend of increasing participation over time.

4. Finally, the document concludes with a series of recommendations for future research and implementation of the findings.

5. The overall goal of the study was to provide a comprehensive overview of the current state of the field and to identify areas for further investigation.

Beukers - Stienstra (Math. Ann. 271 (1985))  
On the Picard-Fuchs equation and the formal  
Brauer group of certain elliptic K3-surfaces

gives examples in which the L-function of  
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If the same ideas can be applied to our  
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this would show that

$$f_0 \frac{1}{\sqrt{u}} \frac{du}{dt}$$

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## Modular Forms

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$f_0(u(\tau))$  is a modular form of weight 1

$\frac{du}{d\tau}$  is a modular form of weight 2

All integrands

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