

Title: Semi-Discrete Solution to the Dynamics of Loop Quantum Gravity

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Abstract:

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# Consistent discretization as a dynamics for loop quantum geometry

*Jorge Pullin*  
*Horace Hearne Laboratory*  
*for Theoretical Physics*  
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With Rodolfo Gambini



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The mechanism can eliminate the black hole information paradox

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*“While being a fascinating possibility, such a procedure would be a rather drastic step in the sense that it would render most results of LQG obtained so far obsolete”* T. Thiemann *“The Phoenix Project”*  
gr-qc/0305080

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- Therefore it can be consistently imposed as a constraint.
- The resulting theory, with this extra constraint, has the exact same kinematics as LQG but an explicit dynamics (no Hamiltonian constraint).

We start with the usual action,

$$S = \int dt d^3x \left( \tilde{P}_i^a F_{0a}^i - N^a C_a - NC \right)$$

$$C^a = \tilde{P}_i^a F_{ab}^i$$

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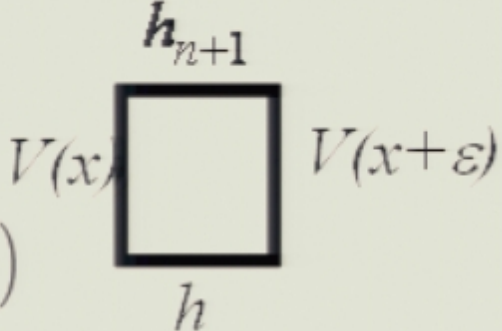
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Where  $V$  is a “vertical” parallel transport (in the time direction). To get this expression one first discretizes in time and space, replaces the  $F_{ab}$  with a holonomy and then takes the limit in which the spatial links go to zero. The last term just enforces that the  $V$ 's be in  $SU(2)$ .

$SU(2)$  gauge invariance is exactly preserved.

If now one performs an infinitesimal spatial transformation  $x'^a = x^a + v^a$  one immediately sees the action is invariant. The associated conserved quantity through Noether's theorem is,

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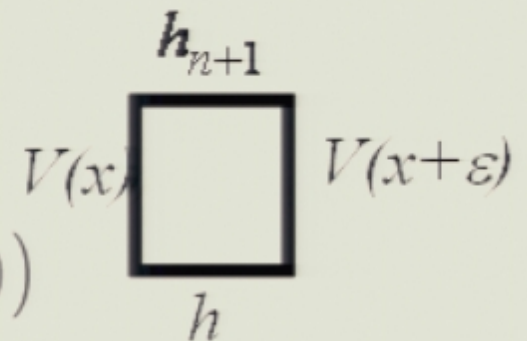
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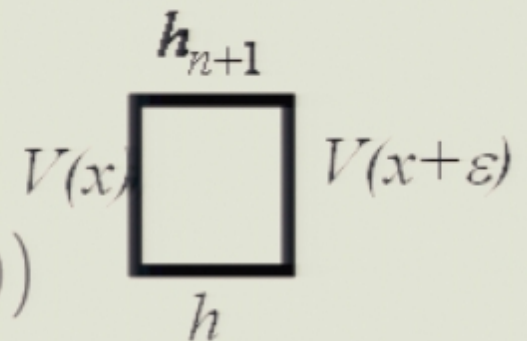
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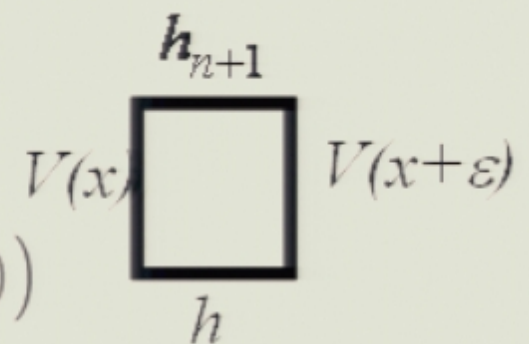
Quantization is straightforward! (no constraints) but computationally complex for the case of full GR. The quantum kinematics is the usual of LQG

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## An explicit example: BF theory

$$L(n, n + 1) = \text{Tr} \left\{ B_0(x) F_{12}(x) + B_1(x) \left( A_2(x) - V(x) A_{n+1,2}(x) V^{-1}(x) + \partial_2(V(x)) V^{-1}(x) \right) \right. \\ \left. + B_2(x) \left( V(x) A_{n+1,1}(x) V^{-1}(x) + V(x) \partial_1 V^{-1}(x) - A_1(x) \right) + \mu \left( V(x) V^\dagger(x) - 1 \right) \right\}$$

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The momentum of  $V$  also vanishes and yields Gauss' law.

Quantization is achieved by implementing the evolution equations as a unitary transformation. This can be immediately done:

$$U(A', A) = \langle A', n + 1 | A, n \rangle$$

$$= \delta(A'_1 - V^{-1}A_1V + \partial_1(V^{-1})V) \delta(A'_2 - V^{-1}A_2V + \partial_2(V^{-1})V) \exp\left(\text{Tr} \int B_0 F\right),$$

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One has to impose the constraint  $F_{12}=0$ , which yields the usual space of states of BF theory (Ooguri, Noui-Perez).

The unitary evolution depends on two free functions  $B_0$  and  $V$ , but it is clear that on the space of physical states the evolution is independent on them.

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- This opens a concrete avenue for numerical quantum gravity.