

Title: The Problem of Dynamics in Quantum Gravity

Date: Oct 29, 2004 10:00 AM

URL: <http://pirsa.org/04100021>

Abstract:

Physical Ramifications of Quantum Geometry

Dual goal:

- Introduction / Context for several presentations in the workshop
- Discuss a few recent developments

Emphasis: Can unravel a lot of interesting physics even though many issues open concerning dynamics of the full theory.

Organization

- Quantum geometry Fleischhack, LOST
- Quantum cosmology Bojowald, Singh, Vandersloot, Willis, Novi, Perez, AA, ...
- Quantum black holes Domagala, Lewandowski, Meisner, Engle, Van Den Broeck, Bojowald, AA, ...

Quantum Geometry

• spinors; tetrads

$$S_P(E, \omega) = \frac{i}{2\pi} \int \text{tr} \omega e^{\hat{A}} e^{\hat{A}^\dagger} d^4x$$

$\rightarrow (E_a^i, K_a^i) \in \Gamma_{\text{grav}} = \text{Geometrodynamics}$

• other forces of nature: Connection Dynamics

$$S_G(E, \omega) = S_P - \frac{i}{2\pi} \int \text{tr} \omega e^{\hat{A}} e^{\hat{A}^\dagger} d^4x$$

$\rightarrow (A_a^i, P_a^i) \in \Gamma_{\text{grav}} \quad (\text{Yang-Mills connection})$

• Relation: $E_a^i = \text{tr} \omega P_a^i$; Riemann Geom

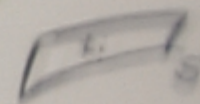
$$\text{rk}_a^i = A_a^i - \Gamma_a^i$$

• Techniques from gauge theories:

Holonomies: $h_\alpha(A)$

'Electric fluxes': $\Phi_{S,L} = \int_S t_i^* P^i$

elementary
variables



Quantum Geometry

- Spinors; tetrads

$$S_P(e, \omega) = \frac{1}{4k} \int \epsilon_{IJKL} e^I \wedge e^J \wedge \Omega^{KL}$$

$$\leadsto (E^a_i, k_a^i) \in \Gamma_{\text{Grav}} : \text{Geometrodynamics}$$

- Other forces of Nature: Connection dynamics

$$S_H(e, \omega) = S_P - \frac{1}{2k\hbar} \int e^I \wedge e^J \wedge \Omega_{IJ}$$

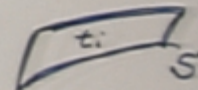
$$\leadsto (A_a^i, P^a_i) \in \Gamma_{\text{Grav}} \quad (\text{SU}(2)\text{-Connection Theory})$$

- Relation: $E^a_i = 8\pi\hbar G P^a_i$; Riem Geo
 $rk_a^i = A_a^i - \Gamma_a^i$

- Techniques from gauge theories:

elementary variables {

$$\begin{aligned} &\text{Holonomies : } h_e(A) \\ &\text{'Electric fluxes' : } P_{S,t} = \int_S t_i \star P^i \end{aligned}$$

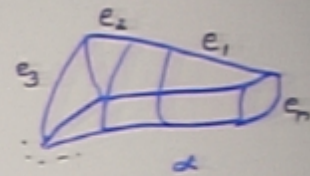


Quantum kinematics:

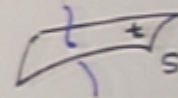
- configuration variables

$$F_\alpha(A) = f(h_{e_1}(A), \dots, h_{e_n}(A))$$

Holonomy (unital C^*) algebra : $\overline{\text{cyl}}$.



- Momenta: $P_{S,t}$ define derivations on $\overline{\text{cyl}}$
($\sim$ Action of VFS)



Fundamental assumption of LQG:
Quantization based on this type of Algebra

Very surprising recent result:

- Quantum algebra $\hat{\mathcal{A}}$ generated by $\hat{F}_\alpha, \hat{P}_{S,t}$ admits a unique diff.-invariant state.
(Lewandowski, Okolow, Sahlmann, Thiemann, Fleischhack)
- Chronology
- Contrast with reps of CCR in Minkowskian QFTs. Diff invariance qualitatively different
↓
Absence of background geometry

Fleischhack

LOST

(Lewandowski - Okolow - Sahlmann - Thiemann)

Piecewise C^ω

Piecewise C^ω

Sub-analytic, Finite-Widely
triangulated

same

Stratified
Analytic

same

Holonomies

Holonomies

Exponentiated electric fluxes

Electric Fluxes

* Algebra, C^* ;

"Mixed" Algebra, $\ast$ -Alg

$e^{i\hat{\mu}P}$ in QM]

[Ex: $e^{i\hat{\lambda}X}$, $\hat{p}$ in QM]

↳ (Weak-continuity)

Diff-inv. state s.t. on

Unique Diff invariant state.

acting rep, Diff acts

Rep obtained by GNS-
construction:

Common domain for all
Flux operators

Fleischhack

holonomies:

Edges

Piece wise C^∞

Momentum Fluxes:
Surfaces

Sub-analytic, Finite-Widely
triangulated

Symmetries:
Diffeos

Stratified
Analytic

Algebra

Holonomies

Exponentiated electric fluxes

"Weyl" Algebra, C^* ;

[Ex: $e^{\hat{\lambda}x}$, $e^{\hat{\mu}p}$ in QM]
↳ (Weak-continuity)

Reps: Unique Diff-inv. state s.t. on
the resulting rep, Diff acts
'naturally'.

LOS

(Lewandowski - ok)

Pieces

Elect

Mixed

[Ex: $e^{\hat{\lambda}x}$]

Unique Diff
Rep obtained
construction
Common domain
Flux operators

Leischhack

LOST

(Lewandowski - Okolow - Sahlmann -
Thiemann)

Piecewise C^ω

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Sub-analytic, Finite-Widely
triangulated

: same :

Stratified
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Holonomies

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Differentiated electric fluxes

Electric Fluxes

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Inv-state st on

Unique Diff invariant state.

g rep, Diff acts

Rep obtained by GNS-
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Common domain for all
Flux operators

2. Application: Quantum Cosmology (Bojowald)

- Symmetry reduction (Ex: Hom Isotropic)

$$A_a^i = c \, \hat{e}_a^i$$

$$P^a_i = p \sqrt{q^0} \, \hat{e}^a_i$$

$$\text{scale factor: } a^2 = |P| \quad ; \quad \text{curv} \sim 1/a^2$$

- Follow procedure from full theory

$$\hat{h}_\mu = \widehat{\exp i \mu c} \quad \text{well-defined but not } \hat{c}$$

$$\mathcal{H} = L^2(\bar{R}_{\text{Bohr}}, d\mu_0)$$

(AA, Bojowald, Lewandowski)

New quantum mechanics
inequivalent to std QGD
Dramatic consequences

- Curvature unbounded classically

$$\text{QGD: } \hat{a} \text{ SA; } \hat{a}^{-1} \text{ SA unbounded}$$

$$\text{Here: } \hat{p} = \frac{1}{\ell_P^2} \sum_\mu |\mu\rangle \mu \langle \mu| \quad ; \quad \langle \mu | \mu' \rangle = \delta_{\mu' \mu}$$

$1/\mu$ not a measurable fn on $\text{SP}(\hat{p})$

Naive $\hat{p}^{-1}$ not SA (not even dd)

- $\hat{p}^{-1}$ defined using strategy from full theory (Thiemann)

Properties:

- Commutates with $\hat{p}$
- Eigenvalues $\sim 1/\mu$ when $a \geq 10^2 \ell_P$
- Important deviations in the deep Planck regime: **operator is bounded above!**

$$\text{Bound} = \left(\frac{12}{8\pi\gamma} \right) \frac{1}{G\hbar} \sim 10^{55} \times \text{curv at the horizon of } 1M_\odot \text{ BH}$$

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- Recall H-atom

$$E_0 = -\frac{me^4}{2 \hbar^2}$$

- Q-Einstein Eq : Hamiltonian Constraint

$$|\Psi\rangle = \sum_{\mu} \underbrace{\psi_{\mu}(\phi)}_{\text{Matter}} \underbrace{|\mu\rangle}_{\text{geometry}}$$

$$C_+ \psi_{\mu+\mu_0}(\phi) + C \psi_{\mu}(\phi) + C_- \psi_{\mu-\mu_0}(\phi) = \hat{H}_{\phi} \psi_{\mu}(\phi)$$

can be interpreted as an evolution equation in discrete time steps.

$C_{\pm}(\mu, \mu_0)$ } such that evolution does not stop @ $\mu=0$
singularity resolved, but classical space-time
"dissolves" in the deep Planck regime. only
quantum state meaningful.

- 2.2 • Semi-classical corrections to Friedman Eq
(AA, Bojowald, Willis)

- Non-trivial issue: Semi-classical behavior at arbitrarily large late times
- Conceptually important test of Thiemann-type regularization (in the limited context of quantum cosmology). For example: There is propagation
- Einstein's eqs do get corrections in LQG

Quantum Evolution

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2.3 : New issues raised by Green & Unruh "Difficulties with closed isotropic LQC"

Issue:

- No 'problem' with the open case.
- Spatial topology S^3 ($k=1$), Massless scalar field

Solns to the Hamiltonian constraint which are also eigenfunctions of scalar field energy (numerically) diverge after a finite no of steps.

Interesting issues raised but no obvious 'difficulty' so far (because of our incompleteness)

- Physical states in cyl^* . Solns may be in cyl^* i.e. genuinely distributional.
- Not using self-adjoint constraint, cannot ask if there are 'physical' infinities.
- Semi-classical states should be constrained but unclear if eigenstates of energy are such states.

BUT burden on us to show there is no problem!
Need to complete the quantization program.

STATUS:

- 2.4 • Willis (+AA+MB): Natural SA constraint does exist. Descends from full theory. Again singularity resolved for std matter.

open issues:

compact case, explicit solutions?
Group averaging a physical Hilbert space?

- 2.5 • Vandersloot (+Novi+Perez): One example.
open model, no scalar field but Λ .
No true degrees of freedom $\rightarrow$ Phys. H.S. 1dim?

Yes! Group Averaging $\checkmark$
Physical state $\in \mathcal{H}_{\text{kin}}$!
Nothing diverges

(Caution:

Factor ordering arrived at using Euclidean Plebanski action. unclear if it can be easily extended to genuine matter in $-\rightarrow\rightarrow\rightarrow$ domain)

3. Quantum BHs

3.1 : Correction of the ABK counting .

- Horizon quantum geometry ✓
- Idea of CS surface states $\rightsquigarrow$ Entropy $\checkmark$
 $= \ln N$
- ABK lower bound on N fine but upper bound incorrect. Thus ABK undercounted $N \Rightarrow$ B-I parameter γ changed from 0.1274... to 0.2375...
(Domagala & Lewandowski; Meissner)

Wheeler's
It from Bit
Only partial answer
($J > \frac{1}{2}$ contributions)

BH evaporation
Transitions much
richer, not just a
equidistant line
spectrum

Quasi-
normal
Scenario
even less
plausible
physically.

• Alternative new possibility (AA-Meissner)

ensemble of structures

Meissner's algebraic equation for γ
interpreted as vanishing of the
chemical potential

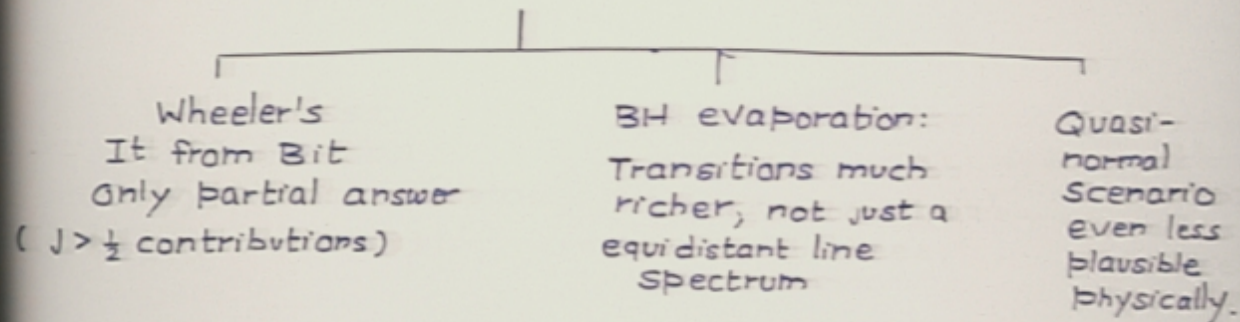
structures $\longleftrightarrow$ photons
water quanta

conversion

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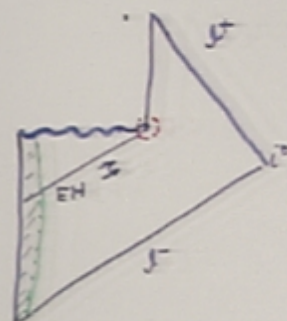
- Attractive new possibility (AA + Meissner)

BH $\longleftrightarrow$ Ensemble of punctures

Meissner's algebraic equation for τ interpreted as vanishing of the chemical potential

punctures $\longleftrightarrow$ photons
Matter quanta

Conversion



standard scenario

- QFT in CST + Heuristics of back reaction

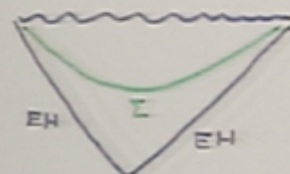
Limitations:

Cannot trust this diagram

- near the entire singularity
- role of event horizon: Global concept

Quantum geometry effects near singularity

Hint: Schwarzschild interior (KS minisuperspace)



$\Sigma: S^2 \times \mathbb{R}$, homogeneous

2 DoF : radius of S^2 : μ
anisotropy : ϕ

Results (AA + Bojowald, Modesto)

- co-triad operator bounded above
- 'Evolution' : μ as 'time'

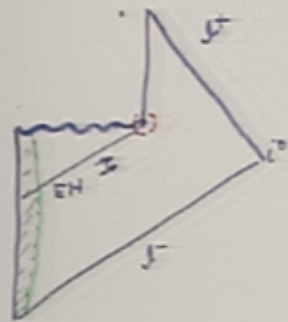
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Hamiltonian constraint $\rightarrow$ difference eq. on $\psi_{\mu}(\phi)$

"Time evolution in discrete steps"

coefficients such that **singularity resolved**.
But classical space-time "dissolves"

Issue of Information LOSS (AA+Bojowald)



standard scenario

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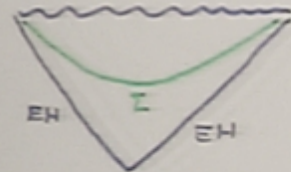
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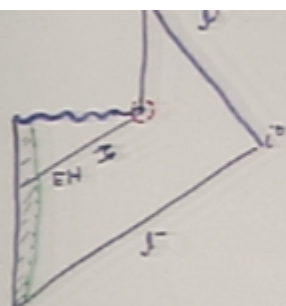
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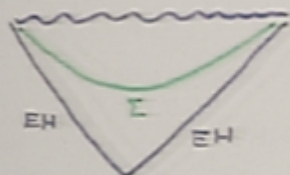
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- 'Evolution': μ as 'time'

$$|\Psi\rangle = \sum_{\mu} \underbrace{\gamma_{\mu}(\phi)}_{\downarrow} |\mu\rangle$$

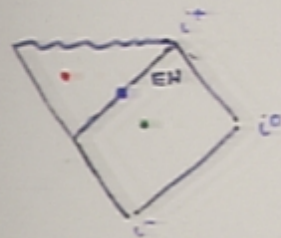
Hamiltonian constraint $\rightarrow$ difference eqⁿ on $\gamma_{\mu}(\phi)$

"Time evolution in discrete steps"

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Dynamical horizons

AA + krishnan

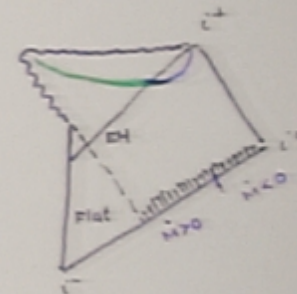
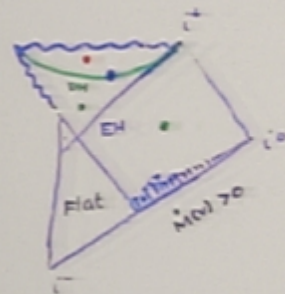


Eternal black hole

- untrapped
- marginally trapped
- Trapped

Dynamical situations : A DH is $S^2 \times \mathbb{R}$, foliated by MT_S ($\Theta_{(l)} = 0$, $\Theta_{(n)} < 0$).

Ex: Vaidya Solutions (spherical; null fluid)



DH: space-like

Time-like

Area of MT_S Increases

Decreases

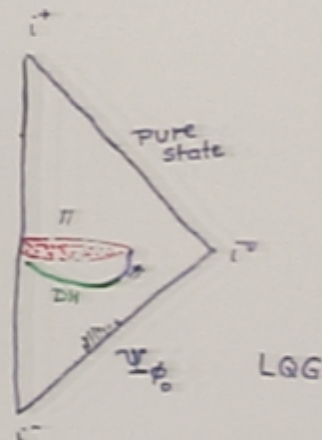
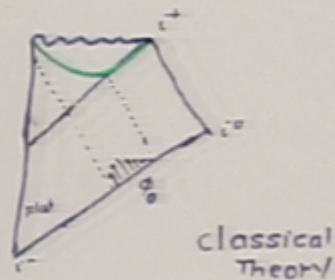
Energy flux +ve

-ve

Integral generalization of first law holds on general DH_S

BH evaporation: A possible, consistent paradigm

- Gather these precise results to develop a qualitative heuristic picture within LQG.
- Spherical collapse of a scalar field: Large BH

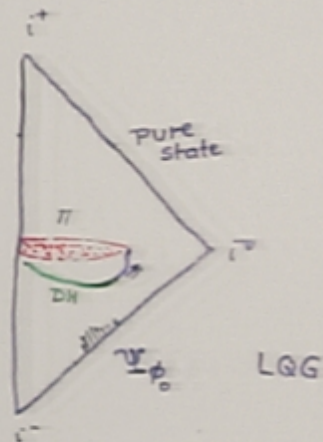
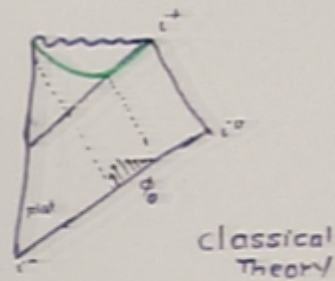


Deep Planck regime
No classical space-time

Assumptions:

- 1) Full quantum state becomes semi-classical again

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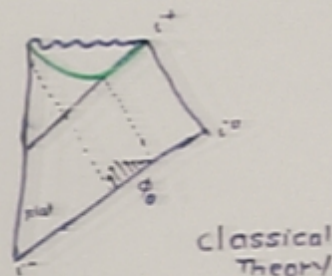
Deep Planck regime
No classical space-time

Assumptions:

- 1) Full quantum state becomes semi-classical again
- ii) $QDH \sim QIH$ when area $\gg \ell_P^2$
(slow evaporation regime)

- Gather these precise results to develop a qualitative heuristic picture within LQG.

- Spherical collapse of a scalar field: Large BH



- Space-time does not end in a singularity $\rightarrow$ No sink of information.

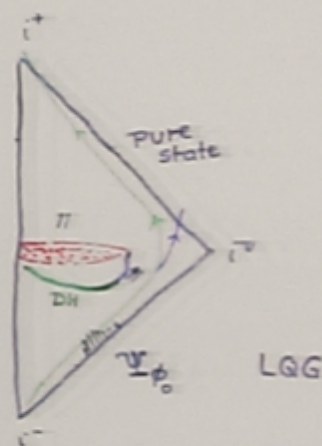
- pure states on $\mathcal{H}^- \rightarrow$ pure states on $\mathcal{H}^+$

- 'Event horizon': no longer a useful notion

- DHs form in semi-classical region and evaporate.

$$\frac{dR}{dt} = -8\pi G R^2 \cdot T_{ab} \ell^a \ell^b$$

Area Quanta $\rightarrow$ ϕ quanta



- Recovery of correlations (patience!)

Deep Planck regime
No classical space-time

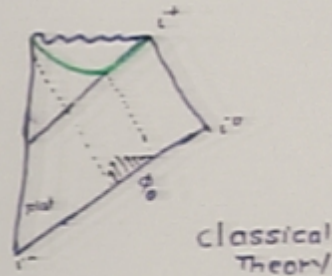
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- i) Full quantum state becomes semi-classical again

- ii) QDH $\sim$ QIH when area $\gg \ell_p^2$
(slow evaporation regime)

- Gather these precise results to develop a qualitative heuristic picture within LQG.

- spherical collapse of a scalar field: Large BH



- Space-time does not end in a singularity $\rightarrow$ No sink of information.

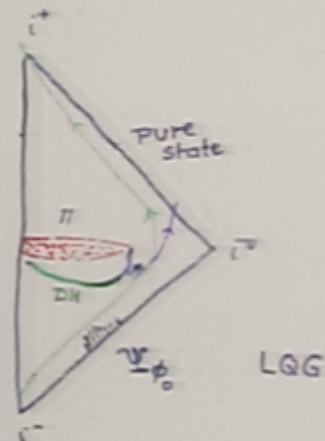
- pure states on $\mathcal{H}^- \rightarrow$ pure states on $\mathcal{H}^+$

- 'Event horizon': no longer a useful notion

- DHs form in semi-classical region and evaporate.

$$\frac{dR}{dt} = -8\pi G R^2 \cdot T_{ab} \hat{r}^a \hat{r}^b$$

Area quanta $\rightarrow$ ϕ quanta



- Recovery of correlations (patience!)

Deep Planck regime
No classical space-time

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Summary:

Even though we do not have a satisfactory treatment of full quantum dynamics, we can address a number of important conceptual issues.

Progress since the last conference in this series 9 months ago is substantial. Solutions to problems which were open & brand new developments

These in turn provide new challenges and great opportunities in the form of concrete new problems as in quantum cosmology as well as brand new directions as in quantum black holes.

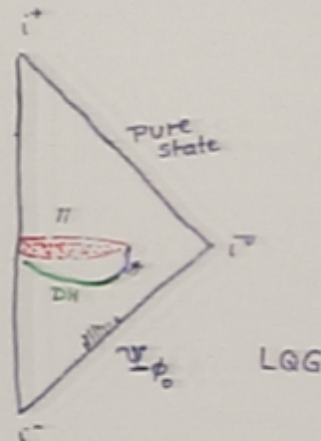
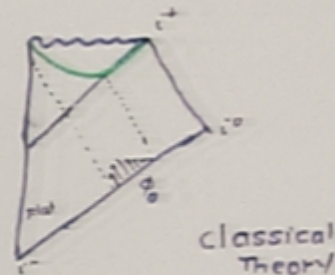
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