

Title: Supersymmetric large extra dimensions and dark energy.

Date: Oct 07, 2004 02:00 PM

URL: <http://pirsa.org/04100010>

Abstract:

OUTLINE

- 6D SUGRA AND THE COSMOLOGICAL CONSTANT

- * Classical Scale Invariance and Self Tuning (Weinberg's No Go)

- Loops, SUSY and LEDs

- BUT WHAT ABOUT:

- * Hidden Fine Tunings; UV-sensitive Loops;

- Brane Phase Transitions; New Forces;

- Dynamical Warping; Radius Stabilization;

- Time-Varying G ; ED Energy Loss.....

- ARE WE MISLED?

- * NOT the MSSM

NATURALNESS + DARK ENERGY

- COSMOLOGICAL CONSTANT: ($\lambda_{\text{obs}} \sim (10^{-3} \text{eV})^4$)

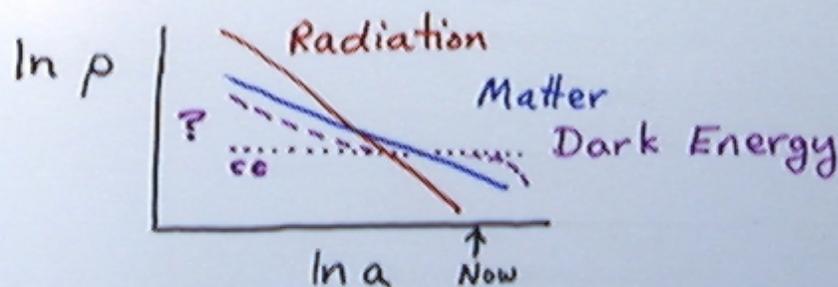
Why doesn't the electron generate too large a c.c.?

$$\delta\lambda \sim m^4 \quad (\text{non supersymmetric})$$

$$\delta\lambda \sim m_{\text{SS}}^2 M^2 \quad (\text{supersymmetric})$$

or m_{SS}^4

- DARK ENERGY: ($m_\phi \lesssim H_0 \sim 10^{-33} \text{eV}$)



How can such a light scalar be stable against loop corrections?

Even if so, why isn't there an observable force?

WISH LIST

(1) IN THE FUNDAMENTAL MICROSCOPIC THEORY, WHY IS λ_0 SMALL?

(2) GIVEN THAT λ_0 IS SMALL, WHY DOES IT STAY SMALL AS ONE INTEGRATES OUT THE SCALES TO WHERE λ IS MEASURED?

ALL KNOWN HIERARCHIES SATISFY BOTH
(eg atom/nucleus ; superconductor/electron..)

★ Solutions to (1) may need to await discovery of microscopic fundamental theory. (string theory?...)

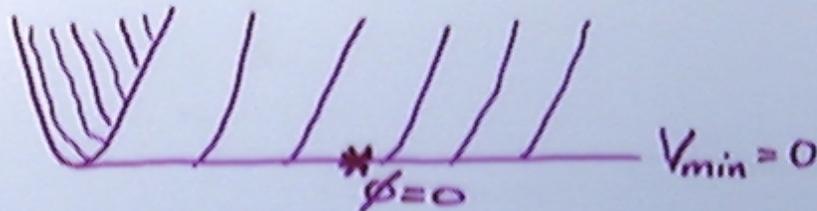
★ Solutions to (2) must change
LOW ENERGY physics we think

WEINBERG'S No-Go

- DYNAMICAL RELAXATION + SPONTANEOUSLY BROKEN SCALE INVARIANCE.

$$g^{\mu\nu} \frac{\delta S}{\delta g^{\mu\nu}} \propto \frac{\delta S}{\delta \phi} \iff \hat{g}_{\mu\nu} = \phi g_{\mu\nu}$$

- ex $V_{\text{eff}} = \lambda_{ijkl} \phi^i \phi^j \phi^k \phi^l$ with flat dirⁿ.



- ▶ Scale invariant point ($\phi=0$): $V_{\text{min}}=0$
- ▶ Scale invariance of $V \Rightarrow V_{\text{trough}}(\phi) = 0$.

SLED

- SUPPOSE PHYSICS IS EXTRA DIMENSIONAL ABOVE THE CC SCALE:

- Experimentally possible if $r < 100 \mu\text{m}$
(so $1/r > 0.01 \text{ eV}$) provided:

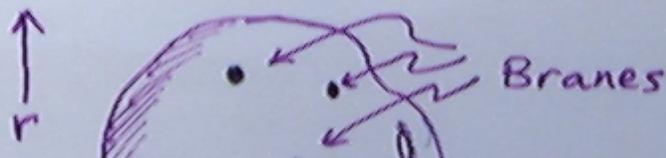
* 6D World at these scales

* $M_g^2 \sim \frac{M_p^2}{r}$ ADD

- IN STRING VERSION WOULD EXPECT SUPERSYMMETRY, BUT:

- Must be broken at scale M_g
on the branes

- Implies susy breaking scale in
bulk is of order $m_{sb} \sim \frac{M_g^2}{M_p} \sim \frac{1}{r}$



SLM

- SUPPOSE PHYSICS IS EXTRA DIMENSIONAL ABOVE THE CC SCALE:

- Experimentally possible if $r < 100 \mu\text{m}$
(so $1/r > 0.01 \text{ eV}$) provided:

* 6D World at these scales

* $M_g^2 \sim \frac{M_p^2}{r} \sim (10^4 \text{ GeV})^2$ ADD

- IN STRING VERSION WOULD EXPECT SUPERSYMMETRY, BUT:

- Must be broken at scale M_g
on the branes

- Implies susy breaking scale in
bulk is of order $m_{sb} \sim \frac{M_g^2}{M_p} \sim \frac{1}{r}$



CC IN 6 DIMENSIONS

- SUSY FORBIDS A C.C. IN 6D

$$\mathcal{L} = -\frac{1}{2}R - \frac{1}{2}(\partial\phi)^2 - \frac{1}{12}e^{-2\phi}G_3^2 - \frac{1}{4}e^{-\phi}F_2^2 \oplus^* g^2 e^\phi + \dots$$

* + → ROMAN'S SUGRA (NON CHIRAL)

- → NISHINO-SERGIN SUGRA (CHIRAL)

- INTEGRATING OUT BRANE PARTICLES GIVES
 $\sim O(M_p^4)$ CONTRIBUTION TO BRANE TENSION.

→ a local energy distribution: $T_i \delta^3(x-x_i)$

- CLASSICAL BULK RESPONSE: $R_{sing} = -2\sum_i T_i \delta^3(x-x_i)$
 + smooth parts

- EFFECTIVE 4D VACUUM ENERGY AFTER
CLASSICALLY INTEGRATING OUT BULK:

$$\lambda_{eff} = \sum_i T_i + \frac{1}{2} \int d^3y [R + (\partial\phi)^2 + \dots] = 0$$

↑
 $\sum T_i$ cancels R_{sing}

Chen, Luhn, Ponton

e.g.: SS SOLUTION

- Fluxes (Freund-Rubin ansatz): ^{-Salam, Sezgin} PLB 147 (84)

$$F_{mn} = f \epsilon_{mn} \quad (n=1 \text{ monopole})$$

$$ds^2 = g_{\mu\nu}(x) dx^\mu dx^\nu + g_{mn}(y) dy^m dy^n$$

(maximal symmetry)

- STABILIZATION WITH A SURPRISE → stabilization, minimum action principle (compactification)

$$\gamma \left[-\frac{1}{2} R - \frac{1}{4g^2} e^{-\phi} F^2 - 2g^2 e^{\phi} \right]$$



$$V_{\text{eff}} \sim \frac{1}{r^2} \left[-\frac{2}{r^2} + \frac{e^{-\phi}}{g^2 r^4} + g^2 e^{\phi} \right]$$



$$= \frac{g^2 e^{\phi}}{r^2} \left[1 - \frac{1}{g^2 r^2 e^{\phi}} \right]^2$$

* $t = g^2 r^2 e^{\phi}$ fixed

* $s = \frac{r^2}{g^2 e^{\phi}}$ flat

* flat 4D space!



- PRESERVES $N=1$ $D=4$ SUSY

$$f_{ab} = SS_{ab} + \dots \quad K = -\ln(S+S^*) - \ln(T+T^*+V)$$

$$\sum_{\text{FI}} = \text{const}$$

Aghababaei, CB, Parameswaran

GENERAL SOLUTIONS

- MOST GENERAL SOLUTION TO CHIRAL GAUGED SUGRA Gibbons, Gaiens + Pope th/0307238

- * 4D Maximal Symmetry
- * Axisymmetry in Extra Dimensions



- * ALL of these solutions have **FLAT 4D** *
- * Tensions can be positive

- WHY ARE THEY ALL FLAT?

ABCFQTZ

th/0308064

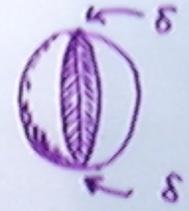
- * Classical Scale Invariance!

$$S = \int d^6x e^{-2\phi} \left[-\frac{1}{2}R + \dots \right] + \sum_i \int d^4x e^{-2\phi} \left[T_i + \dots \right]$$

SELF TUNING?

● RUGBY BALL TENSIONS MUST BE EQUAL:

WHAT HAPPENS IF ONE
BRANE TENSION CHANGES?



* Time-dependent transients must preserve topological constraints + conservation laws.

* **CANNOT** be addressed using truncated 4D theory: $\phi(x)$, $ds^2 = a(x)ds_4^2 + b(x)ds_2^2$

MASS OF $t = r^2 e^\phi$ SAME AS KK MASS

Garriga + Porrati

● IS EVOLUTION TOWARDS ANOTHER FLAT GGP SOLUTION? Yes, for small enough δT

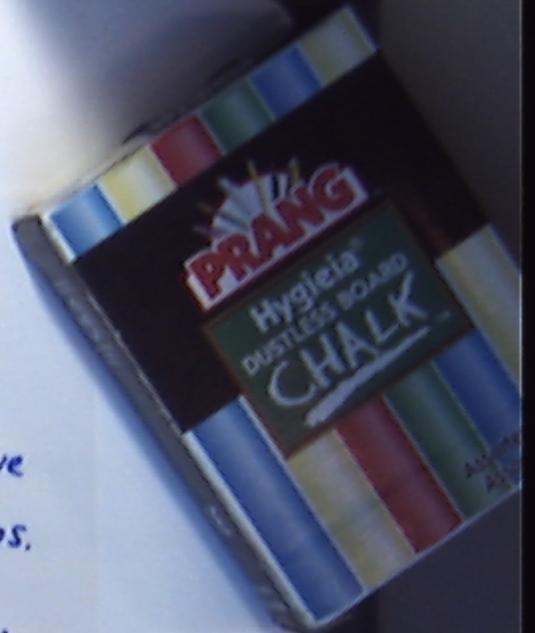
* All conical singularity solutions satisfy:

$$\frac{T_+ - T_-}{2\pi} + \frac{1}{g} \left(1 + \frac{T_+}{2\pi} \right) \left(1 + \frac{T_-}{2\pi} \right) = 0$$

BQTZ
#10408109

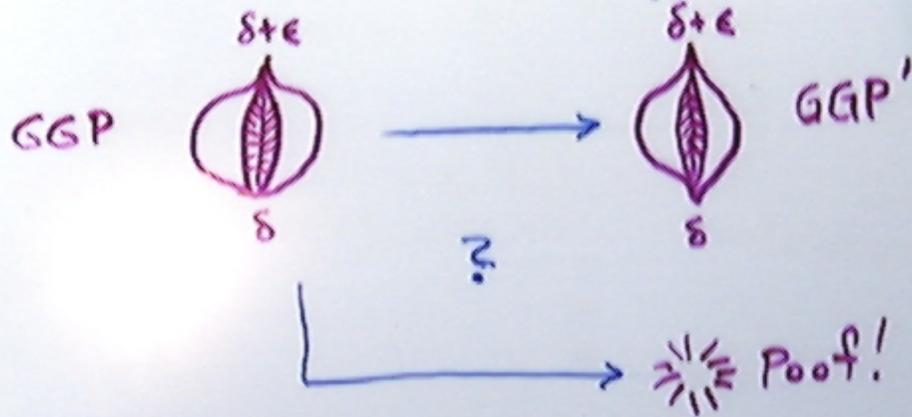
ABCFQTZ

#10308064

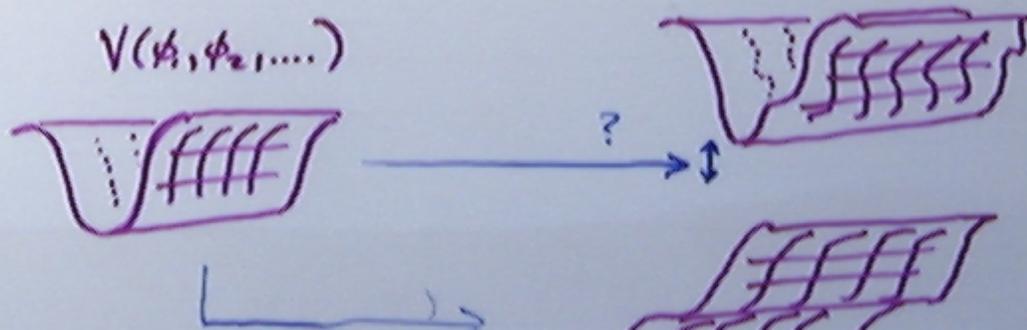


GGP SOLUTIONS + STABILITY

- Since GGP assume maximal 4D symmetry (Mink, AdS, dS), what does it say about self tuning?



- Regarded as 4D theory with many KK modes:



BULK LOOPS

- INTEGRATING OUT BULK MASSLESS 6D FIELDS: (1-Loop)

* On spheres and rugby balls: $V(r) = \frac{\alpha}{r^4} \left[\ln\left(\frac{r}{r_0}\right) + \beta \right]$
 \uparrow UV cutoff

* On torii and orbifolds: $V(r, U) = \frac{W(U, U^*)}{r^4}$
 \uparrow shape moduli
 Poppite Ponton
 Poppita Pelosa
 BAHQ

- ULTRAVIOLET DIVERGENCES: (1-Loop)

* In bulk (dim reg):

$$S_{\infty} = \int d^6x \left[c_1 R^3 + c_2 e^{\phi} R^2 + c_3 e^{2\phi} R + c_4 e^{3\phi} + \dots \right]$$

* On branes (dim reg):

$$S_{\infty} = \int d^4x \left[b_1 R^2 + b_2 e^{\phi} R + b_3 e^{2\phi} + \dots \right]$$

Recall: $R \sim \frac{1}{r^2}$ $\int d^3x \sim r^3$ $e^{\phi} \sim \frac{1}{r^2}$



BULK LOOPS (2)

- INTEGRATE OUT MASSIVE 6D MODES:
(eg: 10D-6D KK modes;) $M \sim \text{TeV}$

* Sensitivity to M arises for local effective terms in 6D and on branes

$$S_{\text{bulk}} = \int d^6x \left[M^6 + M^4(R + e^{\phi}) + M^2(R^2 + e^{\phi}R + \dots) \right]$$

$$S_{\text{brane}} = \int d^4x \left[M^4 + M^2(R + e^{\phi} + \dots) \right]$$

BULK LOOPS (2)

- INTEGRATE OUT MASSIVE 6D MODES
(eg: 10D-6D KK modes, ...) $M = \text{TeV}$

* Sensitivity to M arises for local effective terms in 6D and on branes

$$S_{\text{brane}} = \int d^6x \left[\cancel{M^6} + M^6 \underbrace{(R + \overset{\circlearrowleft}{e^2})}_{!} + M^6 \underbrace{(R^2 + e^2 R + \dots)}_{!} \right]$$

$$S_{\text{brane}} = \int d^6x \left[M^6 + \underbrace{M^6 (R + e^2 + \dots)}_{!} \right]$$

⊗: 6D SUSY

↗: Renormalizations of classical action

!: Dangerous terms: $\delta V(r) = \frac{M^6}{r^2} \underbrace{[e^2 + \dots]}_{\text{higher loops}}$

BULK LOOPS (2)

- INTEGRATE OUT MASSIVE 6D MODES
(eg: 10D-6D KK modes;) $M \sim \text{TeV}$

* Sensitivity to M arises for local effective terms in 6D and on branes

$$S_{\text{bulk}} = \int d^6x \left[\cancel{M^6} + \overset{\nearrow}{M^4} (R + e^\phi \dots) + \underbrace{M^2 (R^2 + e^\phi R + \dots)}_{!} \right]$$

$$S_{\text{brane}} = \int d^4x \left[\overset{\nearrow}{M^4} + \underbrace{M^2 (R + e^\phi + \dots)}_{!} \right]$$

⊗: 6D SUSY

↗: Renormalizations of classical action

!: Dangerous terms: $\delta V(r) \sim \frac{M^2}{r^2} \left[\underbrace{1}_{!} e^\phi + \dots \right]$
higher loops

BULK LOOPS (2)

● INTEGRATE OUT MASSIVE 6D MODES:
 (eg: 10D-6D KK modes;) $M \sim \text{TeV}$

* Sensitivity to M arises for local effective terms in 6D and on branes

$$S = \int d^6x \left[\cancel{M^6} + M^4 \overbrace{(R + e^\phi + \dots)}^{\nearrow} + M^2 \overbrace{(R^2 + e^\phi R + \dots)}^{\nearrow} \right]$$

$$S_{\text{brane}} = \int d^4x \left[M^4 + \overbrace{M^2 (R + e^\phi + \dots)}^{\nearrow} \right] !$$

⊗: 6D SUSY

↗: Renormalizations of classical action

!: Dangerous terms: $\delta V(r) \sim \frac{M^2}{r^2} \left[1 + \overbrace{e^\phi + \dots}^{\text{higher loops}} \right]$



TOPOLOGY VS TUNING

- Topological constraints must be satisfied in order for solutions to classical eqns to exist.
eg: $\sum_i Q_i = 0$ for Coulomb's Law in compact space

- Once existence is ensured, general solution shows it is necessarily flat in 4D.

- Easy to confuse this with fine tuning of c.c. IF internal dimensions are flat:

$$\triangleright \frac{1}{4\pi} \int d^2y \sqrt{g} R = \sum_i \frac{\delta_i}{2\pi} = \chi = \begin{cases} 2 & (\text{sphere}) \\ 0 & (\text{torus}) \\ i & \end{cases}$$

$$\lambda_{\text{eff}} = \sum_i T_i - \frac{1}{2} \int d^2y \sqrt{g} R = \sum_i T_i - 2\pi \chi$$

$$= \sum_i T_i \quad (\text{torus}) \quad \text{many} \dots$$

$$= \sum_i T_i - 4\pi \quad (\text{sphere}) \quad \text{Arkani-Hamed, Hall, Smith + Weiner}$$

CONCLUSIONS

- SLED HAS PROMISE AS AN UNDERSTANDING OF DARK ENERGY IN A NATURAL WAY.

- ★ Hidden Tunings of Couplings?
- ★ Too Large UV Loops?
- ★ Response to tension Changes?
- ★ Why Don't 2D Warp Strongly?
Why Are 2D So Large?
- ★ Sensible Post BBN Cosmology?
- ★ New Light Scatter Bounds?
- ★ SN + Phenomenology Bounds?

- NOVEL REALIZATION OF WEAK SCALE SUSY BREAKING:

- ★ POSSIBLY NO SUSY PARTNERS EXCEPT FOR THE GRAVITON!
- ★ LED AT COLLIDERS AND MORE!

- ENORMOUSLY PREDICTIVE, WITH MANY



CONCLUSIONS

- SLED HAS PROMISE AS AN UNDERSTANDING OF DARK ENERGY IN A NATURAL WAY.

- * Hidden Tunings of Couplings? 😊
- * Too Large UV Loops? 😊
- * Response to tension Changes? ?
- * Why Don't 2D Warp Strongly? ?
- * Why Are 2D So Large? 😊
- * Sensible Post BBN Cosmology? 😊
- * New light Scalar Bounds?
- * SN + Phenomenology Bounds?

- NOVEL REALIZATION OF WEAK SCALE SUSY BREAKING:

* POSSIBLY NO SUSY PARTNERS EXCEPT FOR THE GRAVITON!

* LED AT COLLIDERS AND MORE!