

Title: Mixed-state entanglement in the light of pure-state entanglement constrained by superselection rules

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Abstract: Quantum Information Theory

Mixed-state entanglement in the light of pure-state entanglement constrained by superselection rules

Stephen Bartlett



THE UNIVERSITY
OF QUEENSLAND
AUSTRALIA

Collaborators:
Robert Spekkens
Howard Wiseman

A generic quantum system

 ρ

Density matrix

 $\mathcal{E} : \rho \rightarrow \rho'$

Trace-preserving CP maps

- Consider all states, all maps, all measurements
- Dimension (number of degrees of freedom) determines everything

A generic quantum system

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Density matrix trace-preserving CPTP maps
BORING!

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1. Consider an interesting *subset* of operations

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1. Consider an interesting *subset* of operations
2. Explore the resulting information theory

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Interesting restrictions

Split the system up into two parts

- From a quantum info point of view, things get interesting
- bipartite: restrict to LOCC
- *entanglement* → Bell inequalities, quantum cryptography

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Easier than
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A guide

Can pure state entanglement with SSRs
teach us about mixed state entanglement?

Mixed State Entanglement 101

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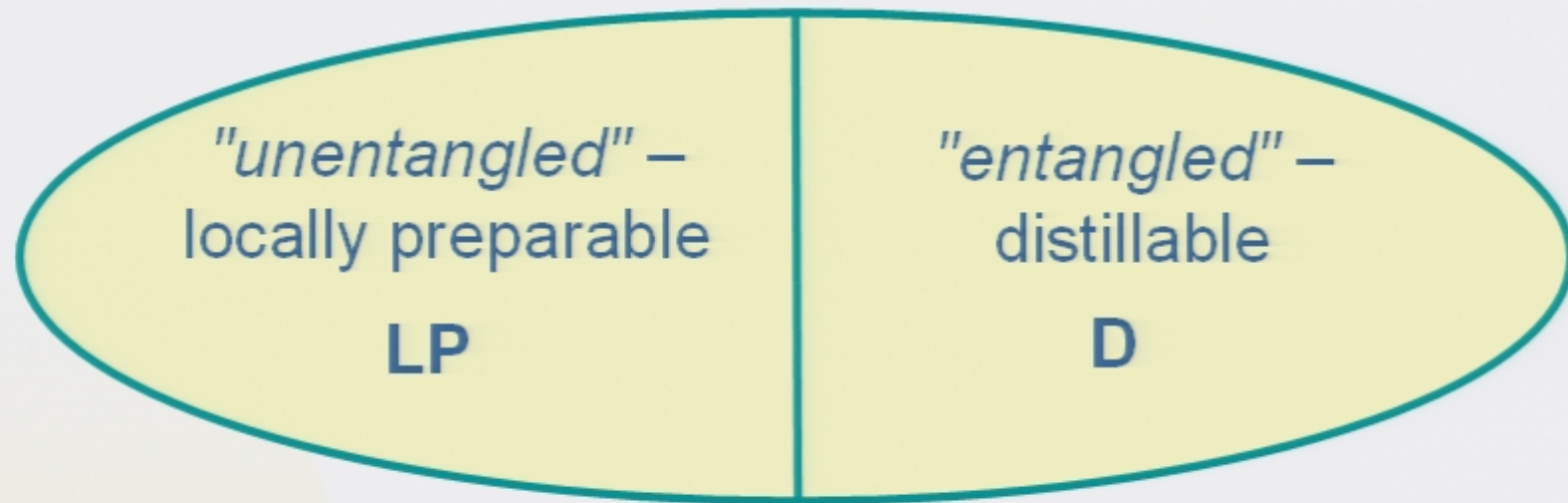
Good references:

Eggeling, Vollbrecht, Werner and Wolf,
Phys. Rev. Lett. **87**, 257902 (2001)

Horodecki (x3), arXiv:quant-ph/0109124

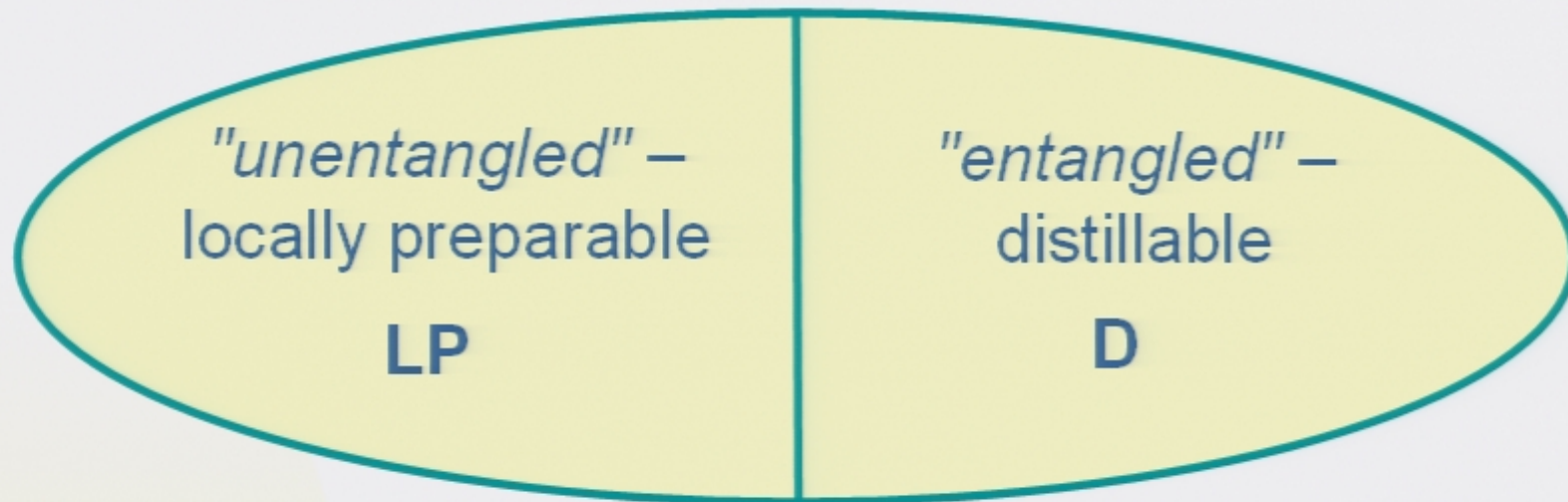
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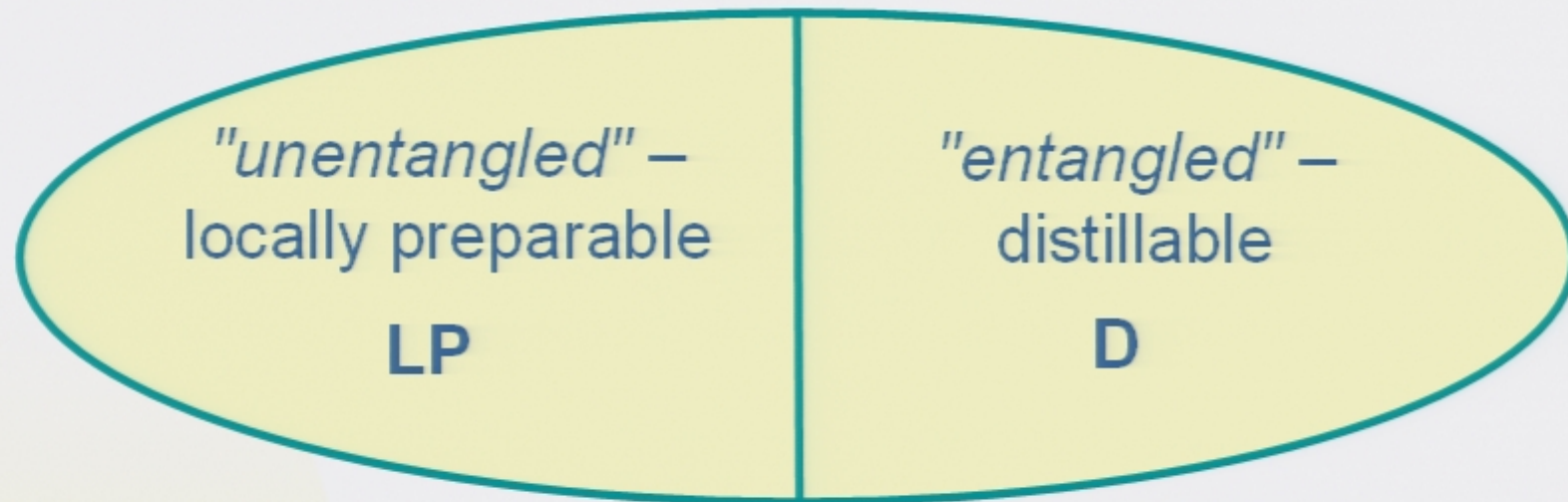


"Unentangled"

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preparable with LOCC
can't make entanglement
- *Mathematically:*
product states

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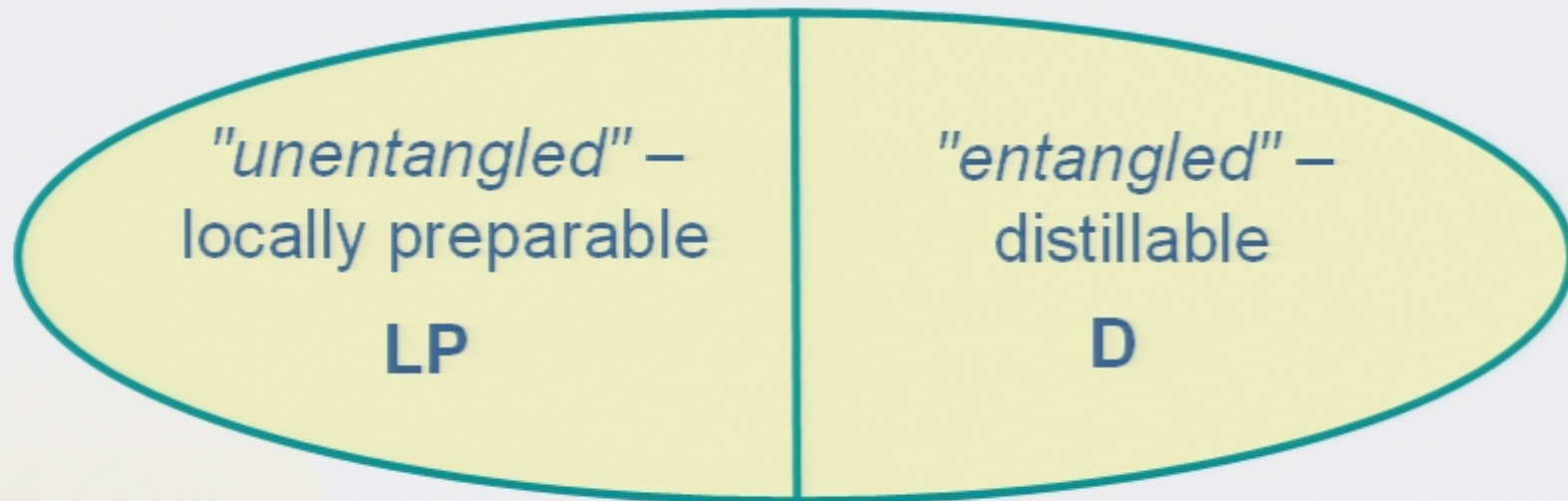
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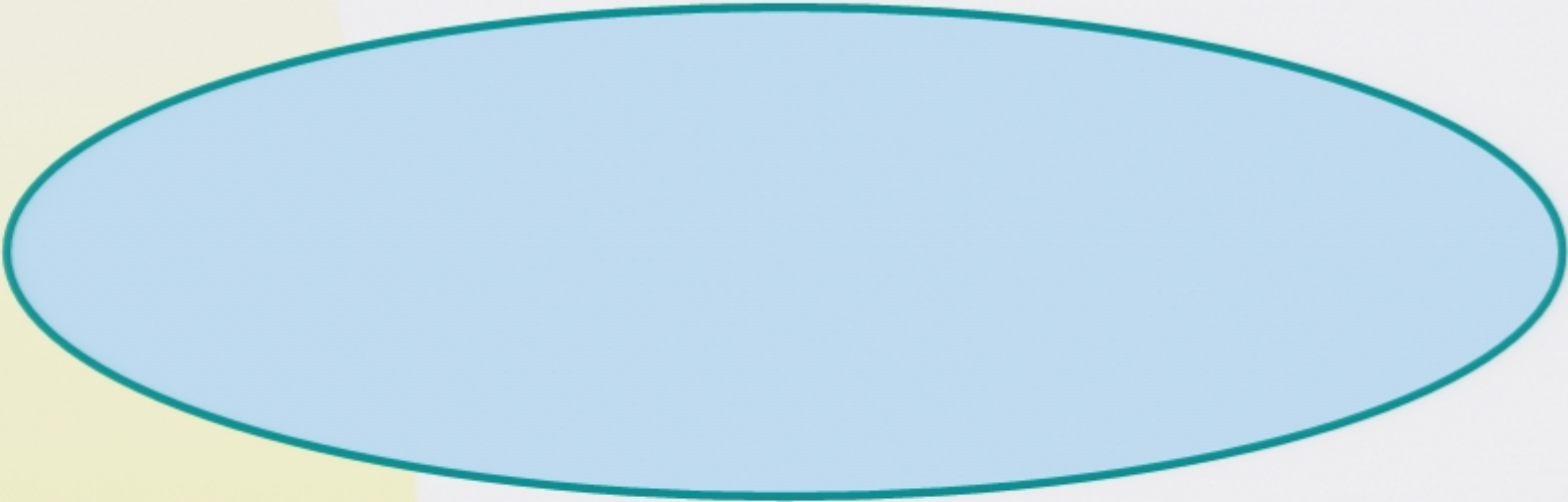
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$$LP \cap D = 0$$

Entanglement as a resource

Mixed state entanglement

All mixed states

A large, light blue oval with a dark blue outline is positioned in the lower half of the slide. It is centered horizontally and occupies a significant portion of the lower half of the frame. The text "All mixed states" is centered above the oval.

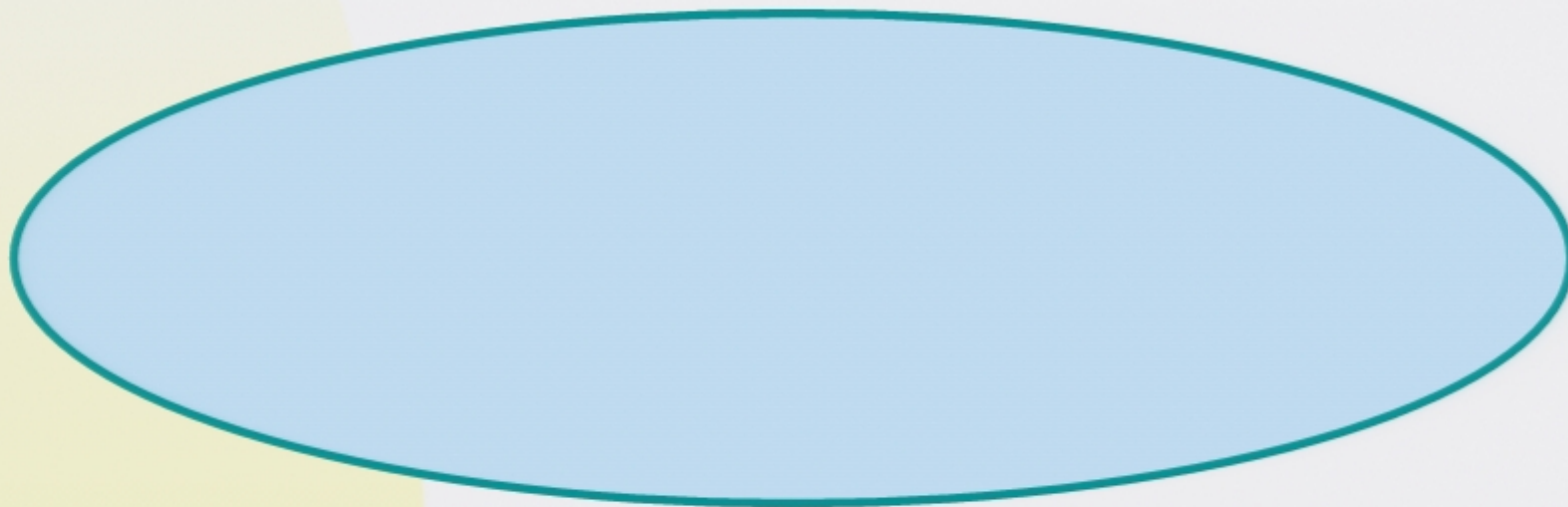
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$$\rho = \sum_i p_i \sigma_i^A \otimes \sigma_i^B$$

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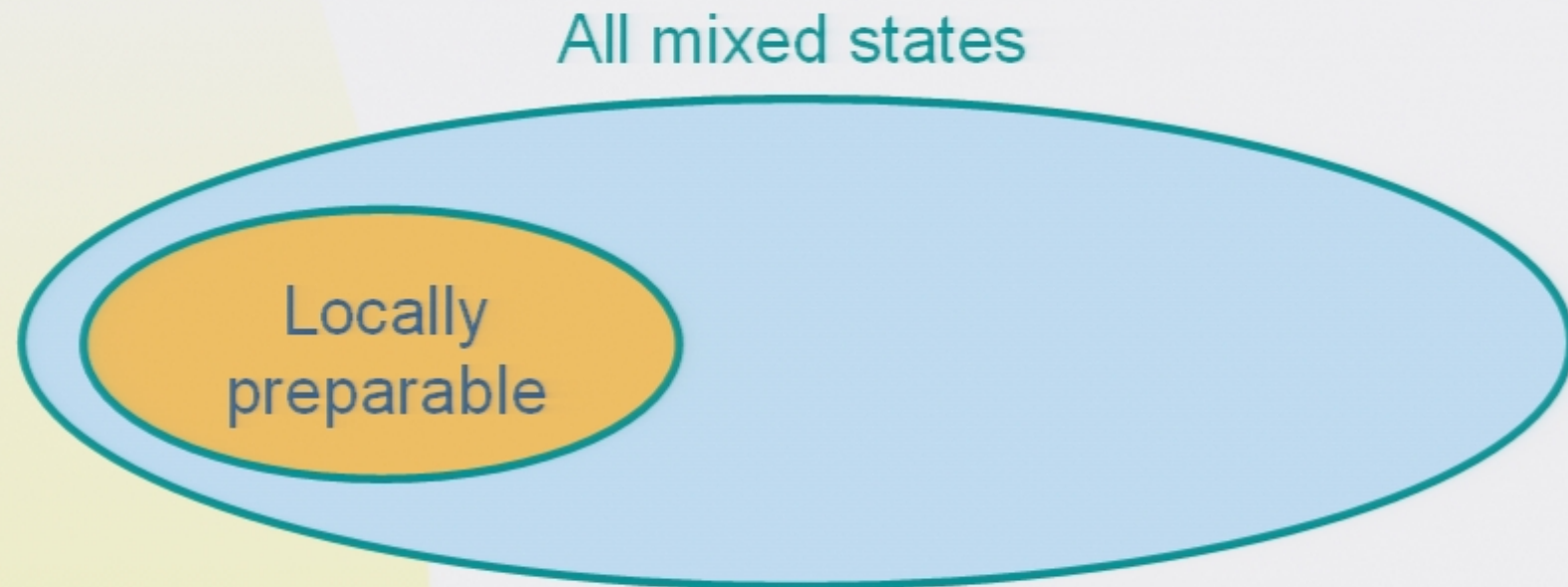


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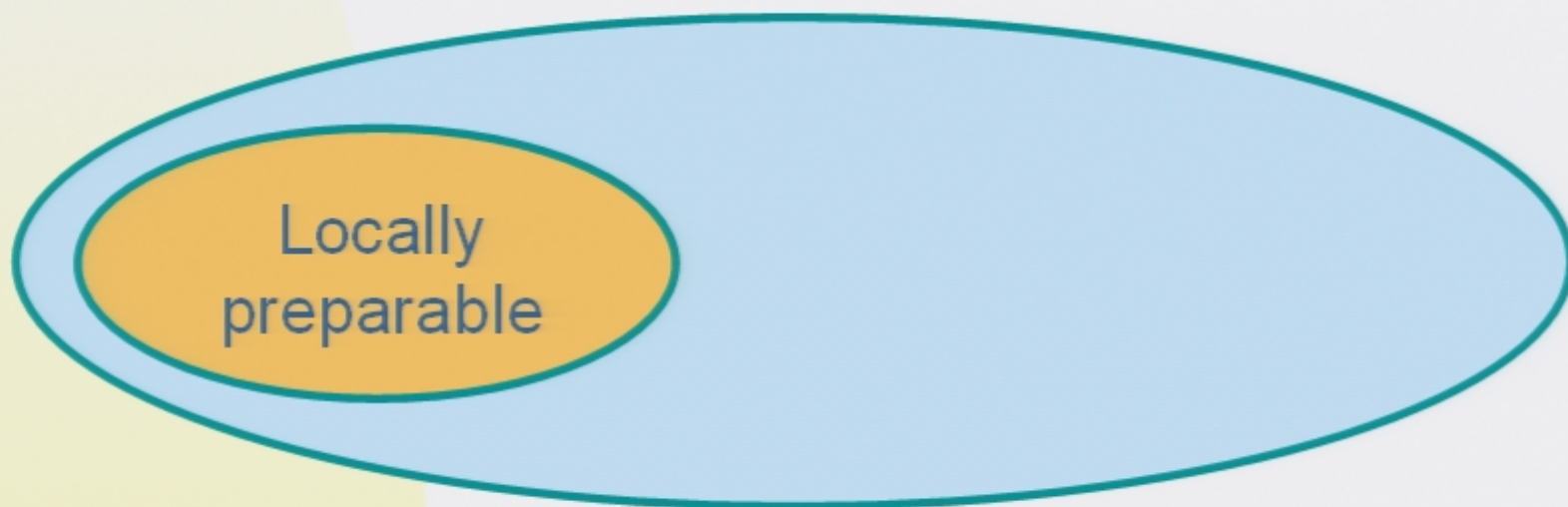
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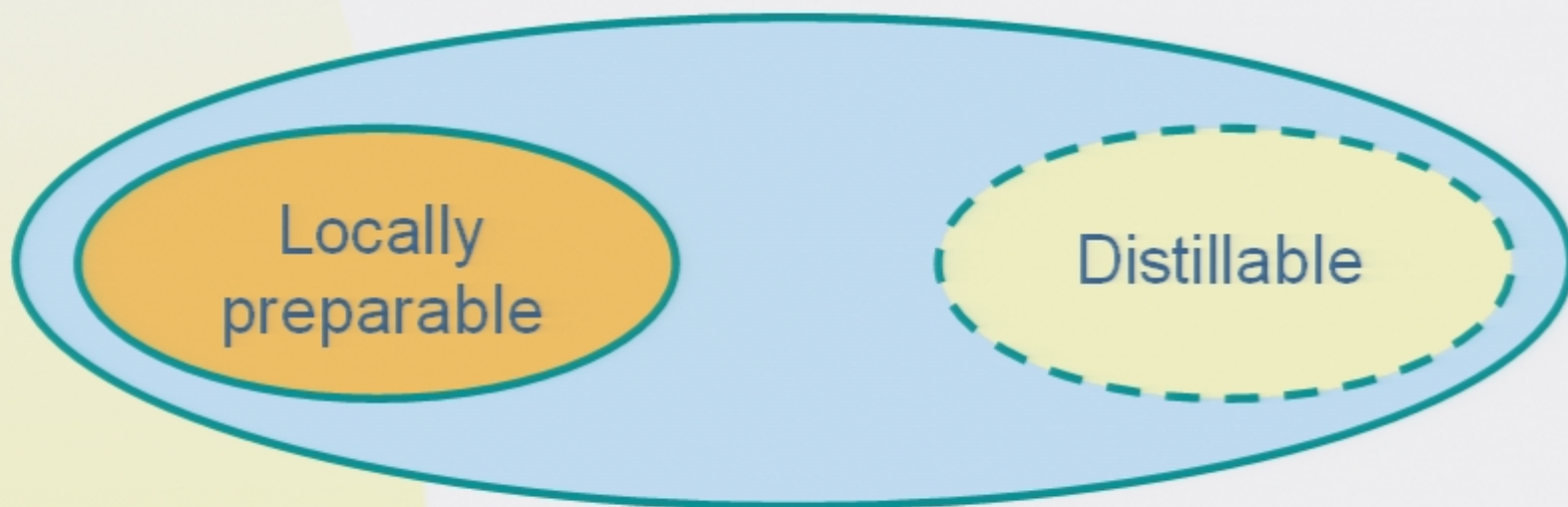
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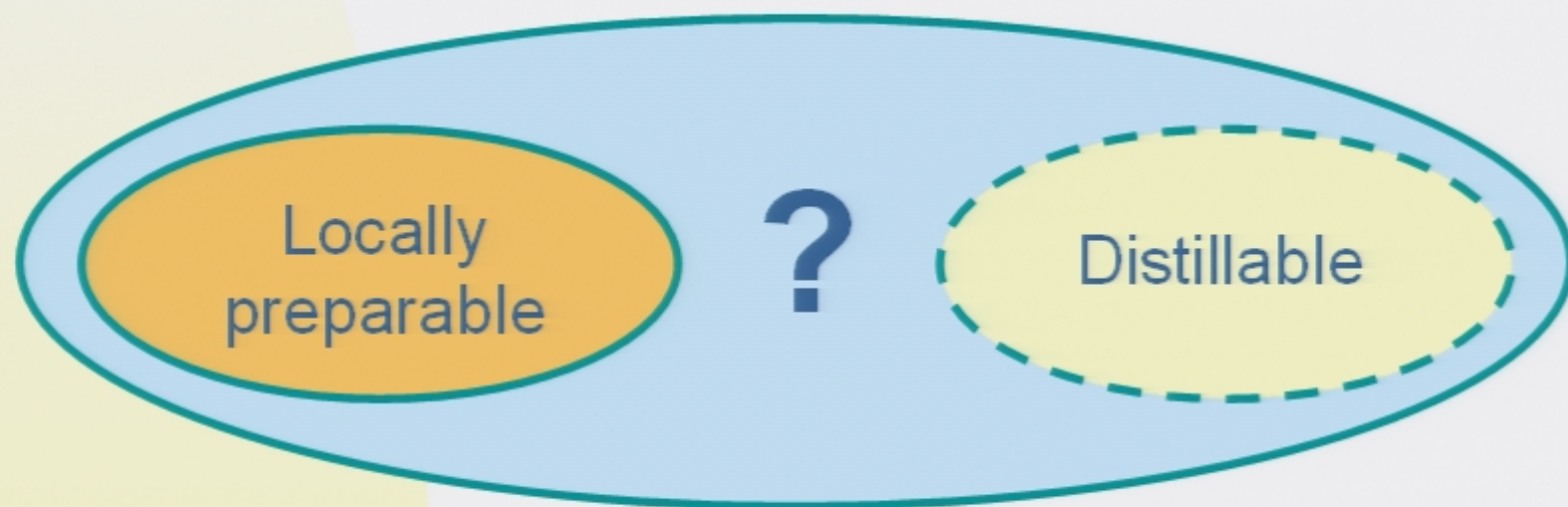
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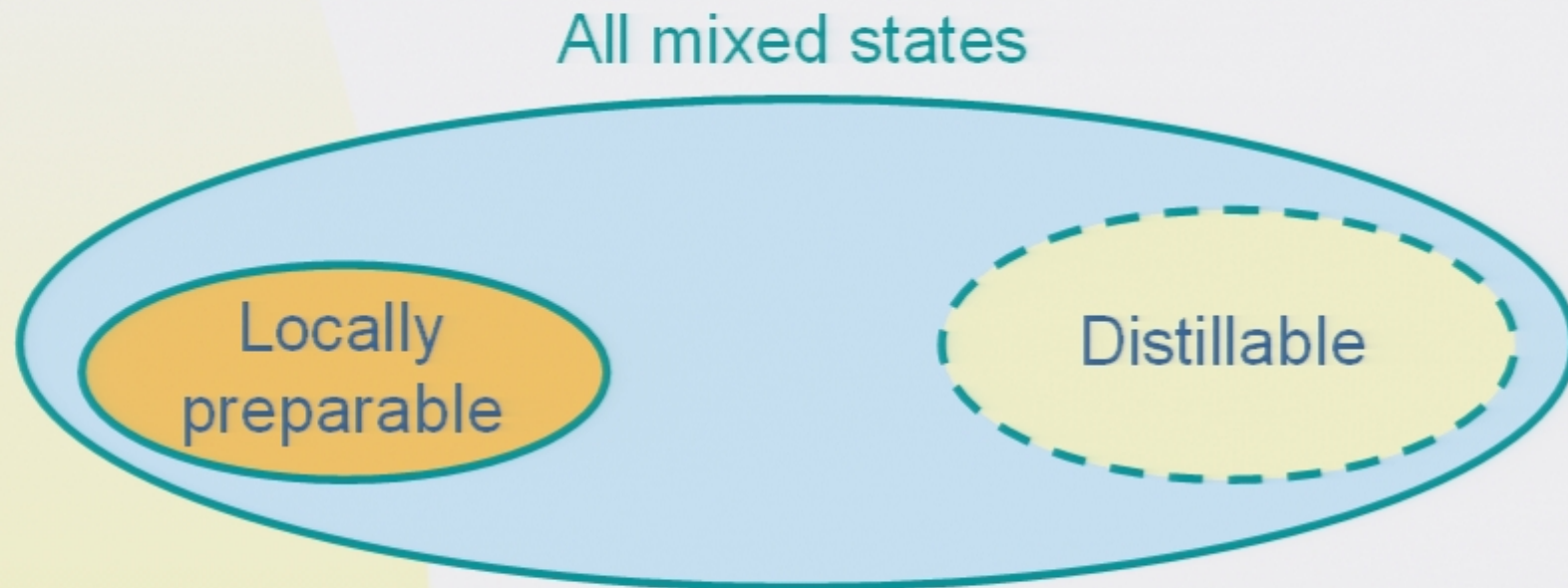


Do there exist states that are not locally preparable and not distillable?

Bound entanglement

1998: Horodecki (x3)

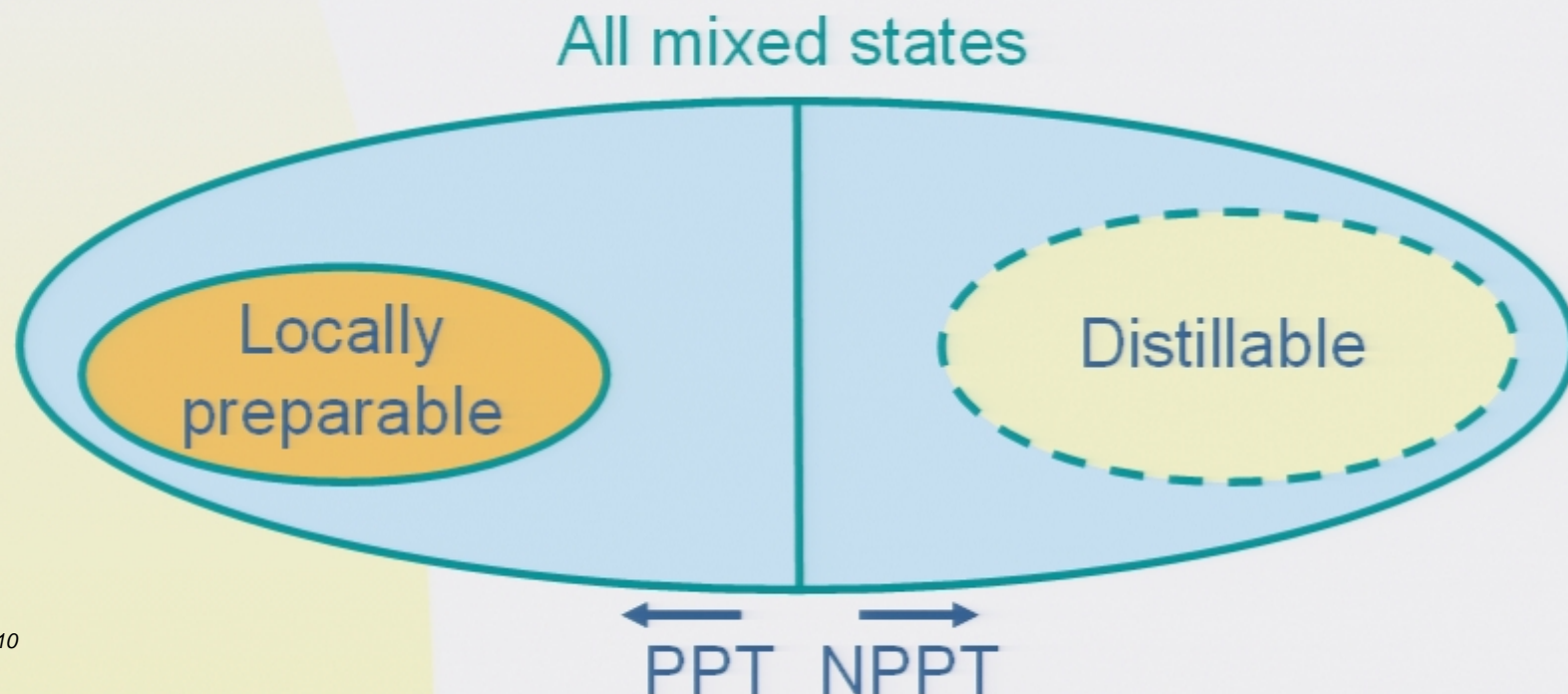
- LOCC operations: preserve positivity of partial transpose (PPT)
- Distillation is the creation of pure Bell states, which are NPPT (*not* PPT)
- \therefore PPT states that are not locally preparable cannot be distilled under LOCC



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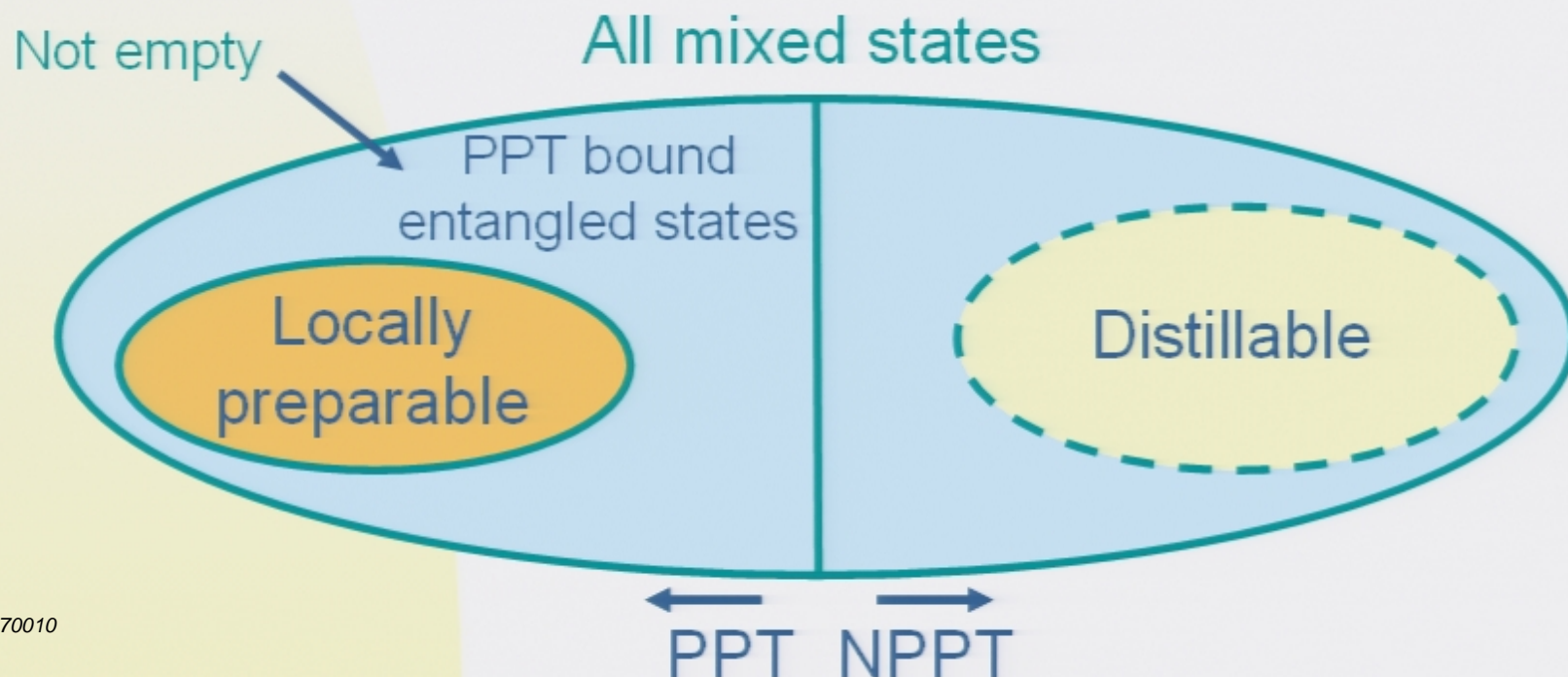
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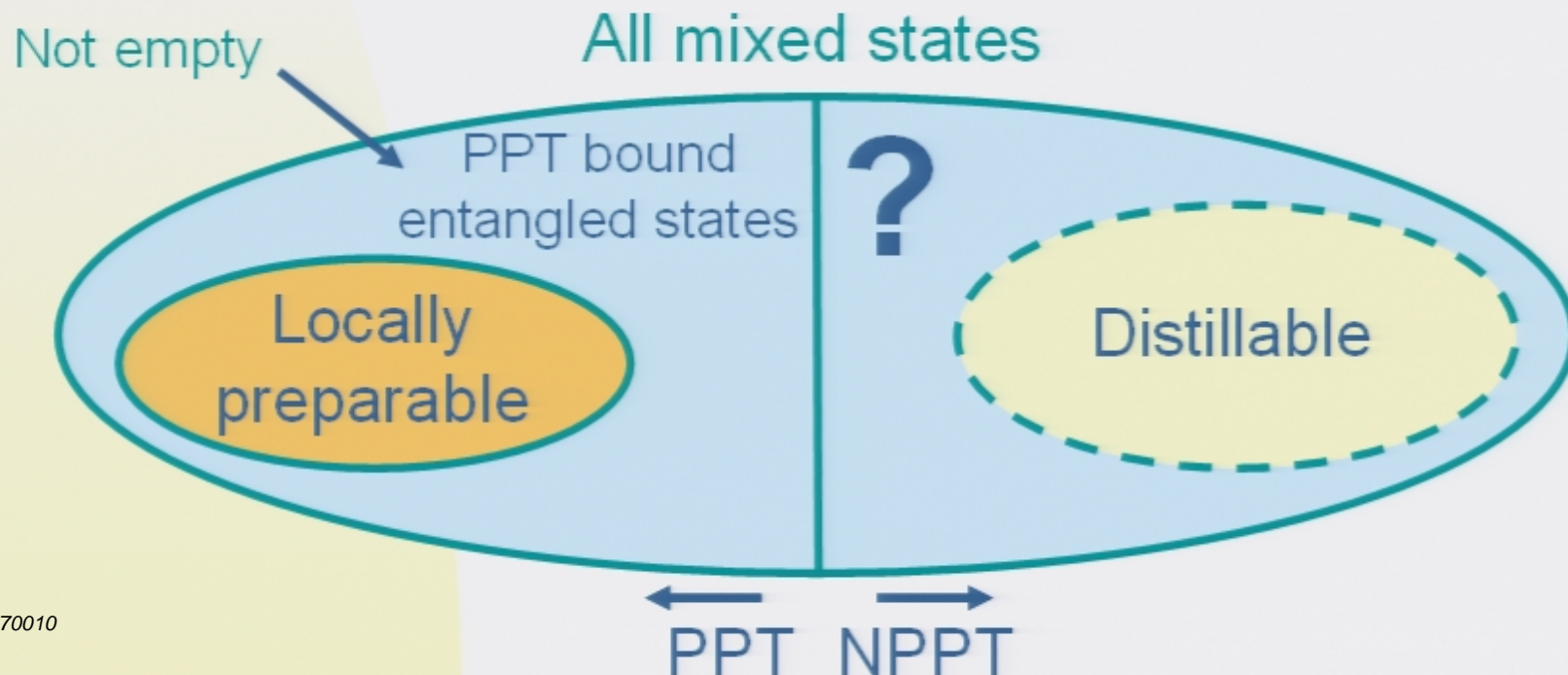
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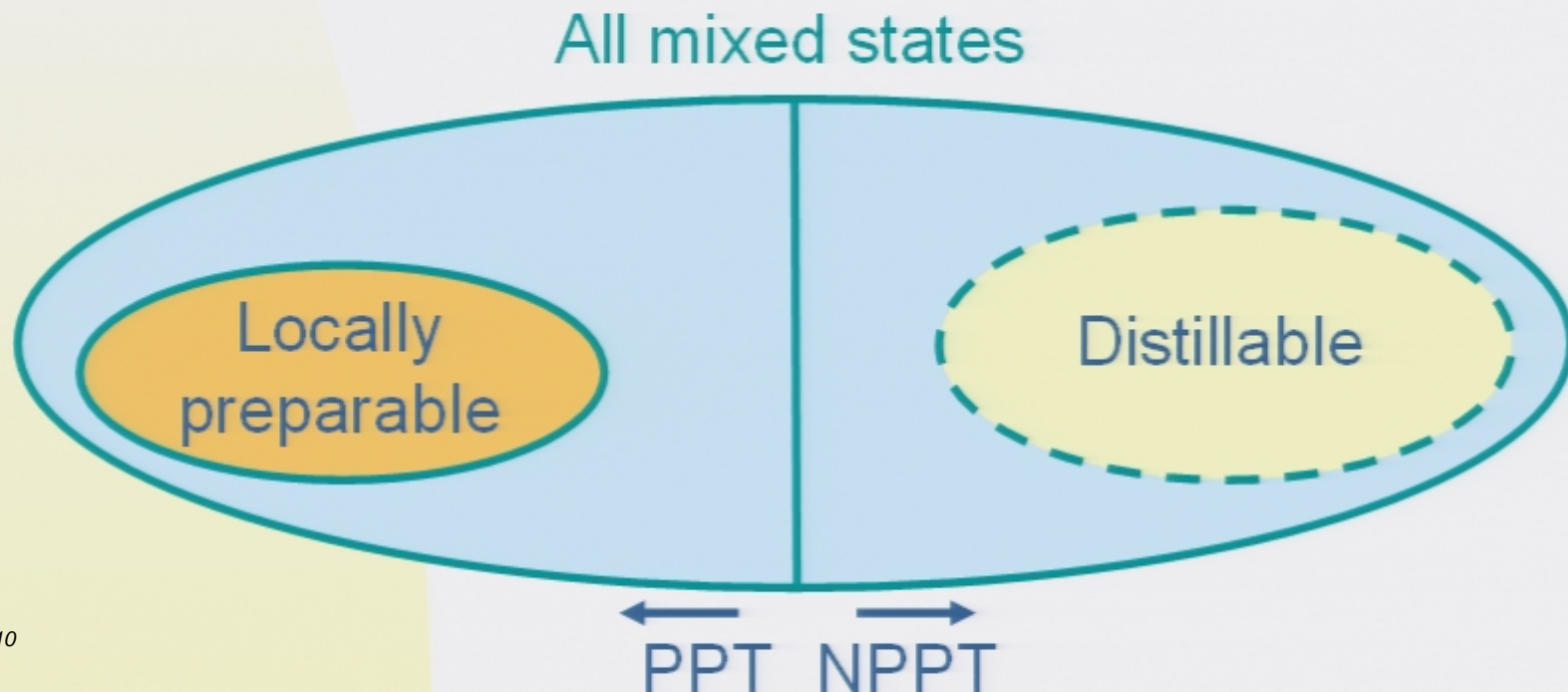


1-distillable

- "Distillable" is too hard to characterise
- Define *1-distillable*: one copy can, with some finite probability, be projected to an entangled 2-qubit state

$$\exists |\psi\rangle \text{ s.t. } \langle \psi | \rho^{T_A} | \psi \rangle < 0$$

- Easy to classify

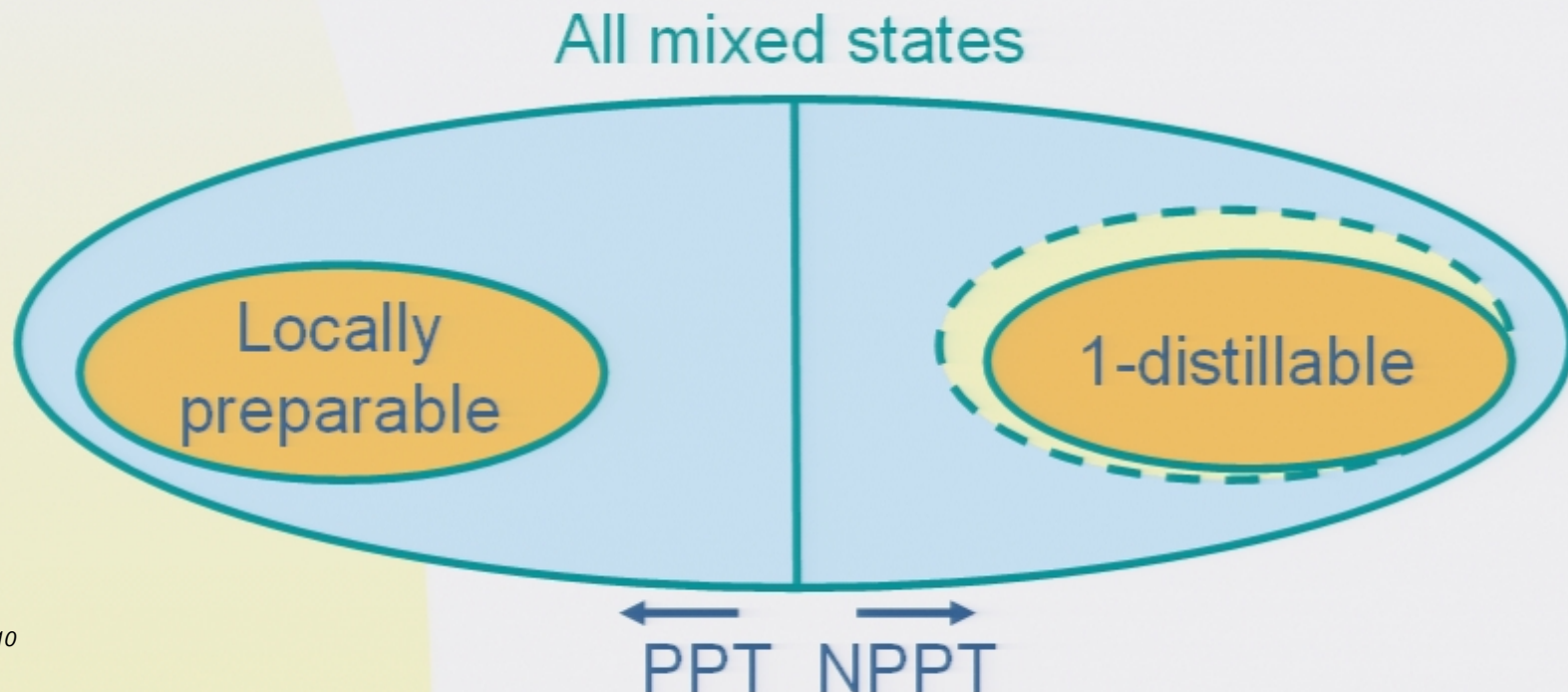


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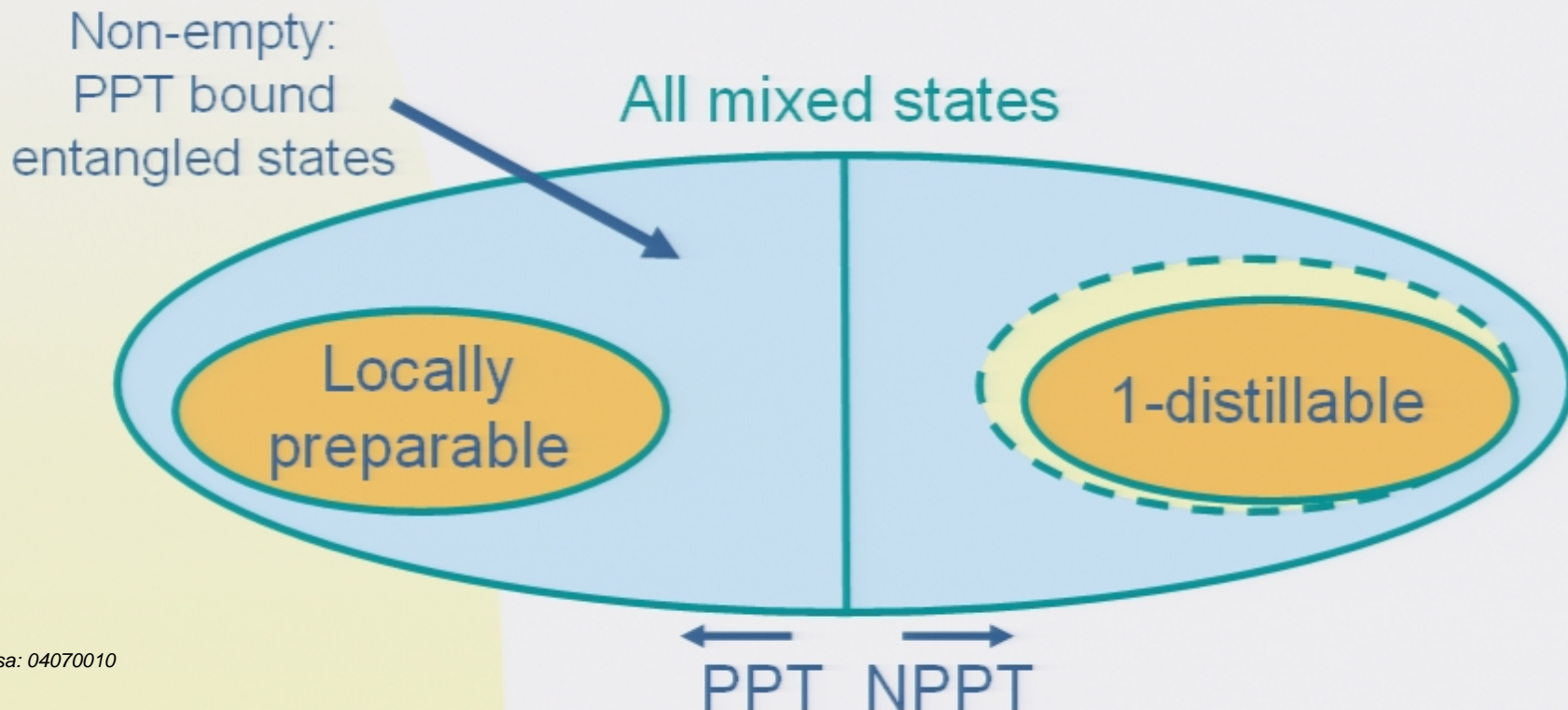


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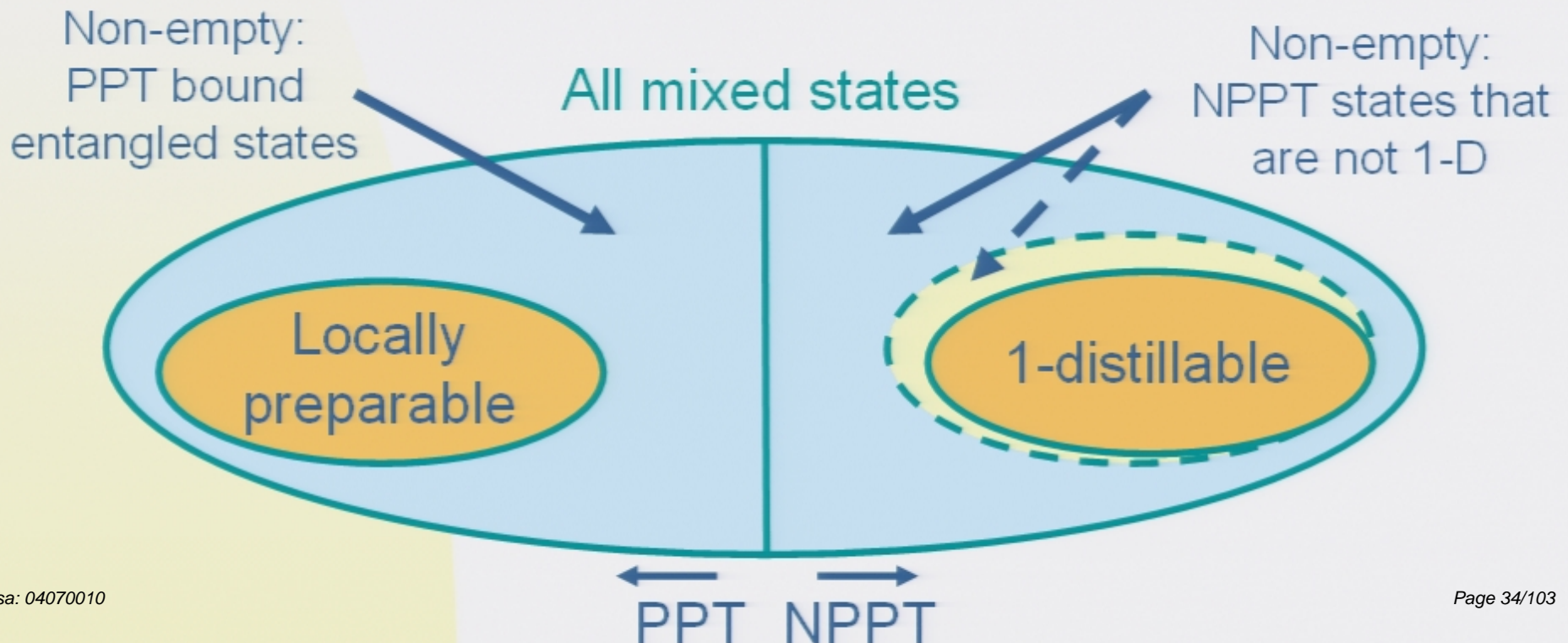


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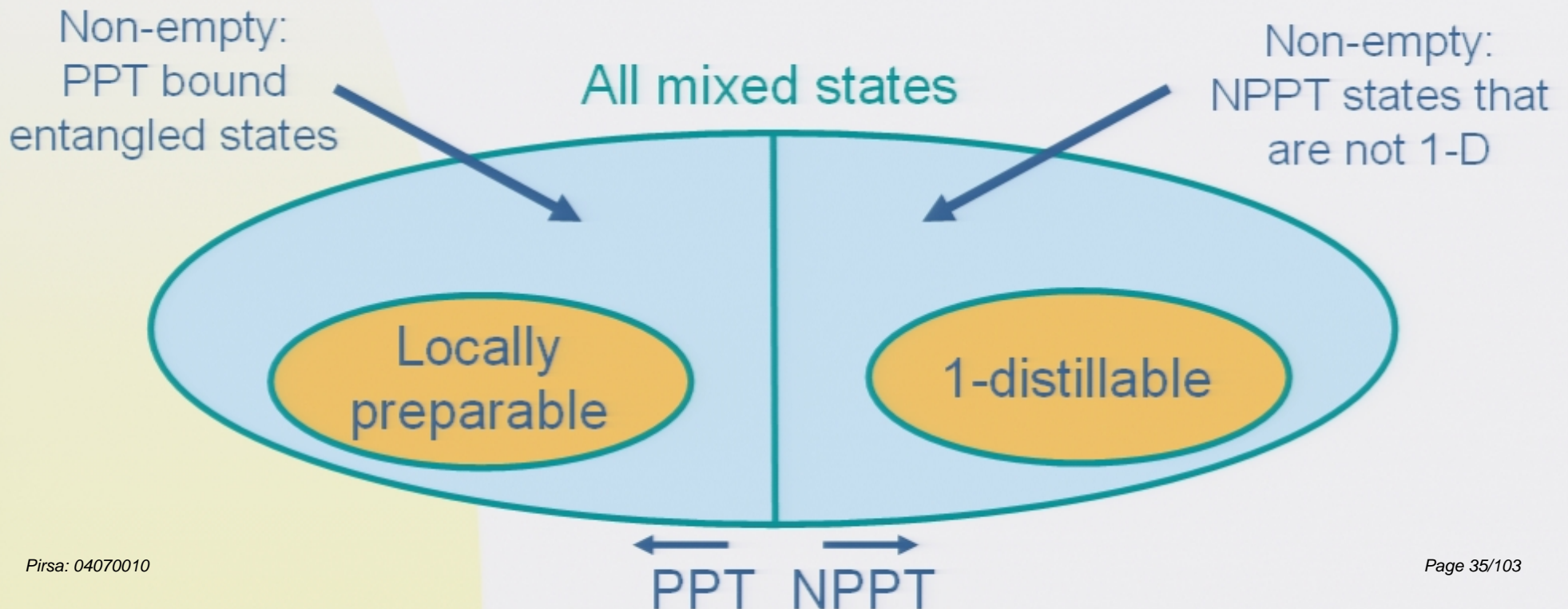
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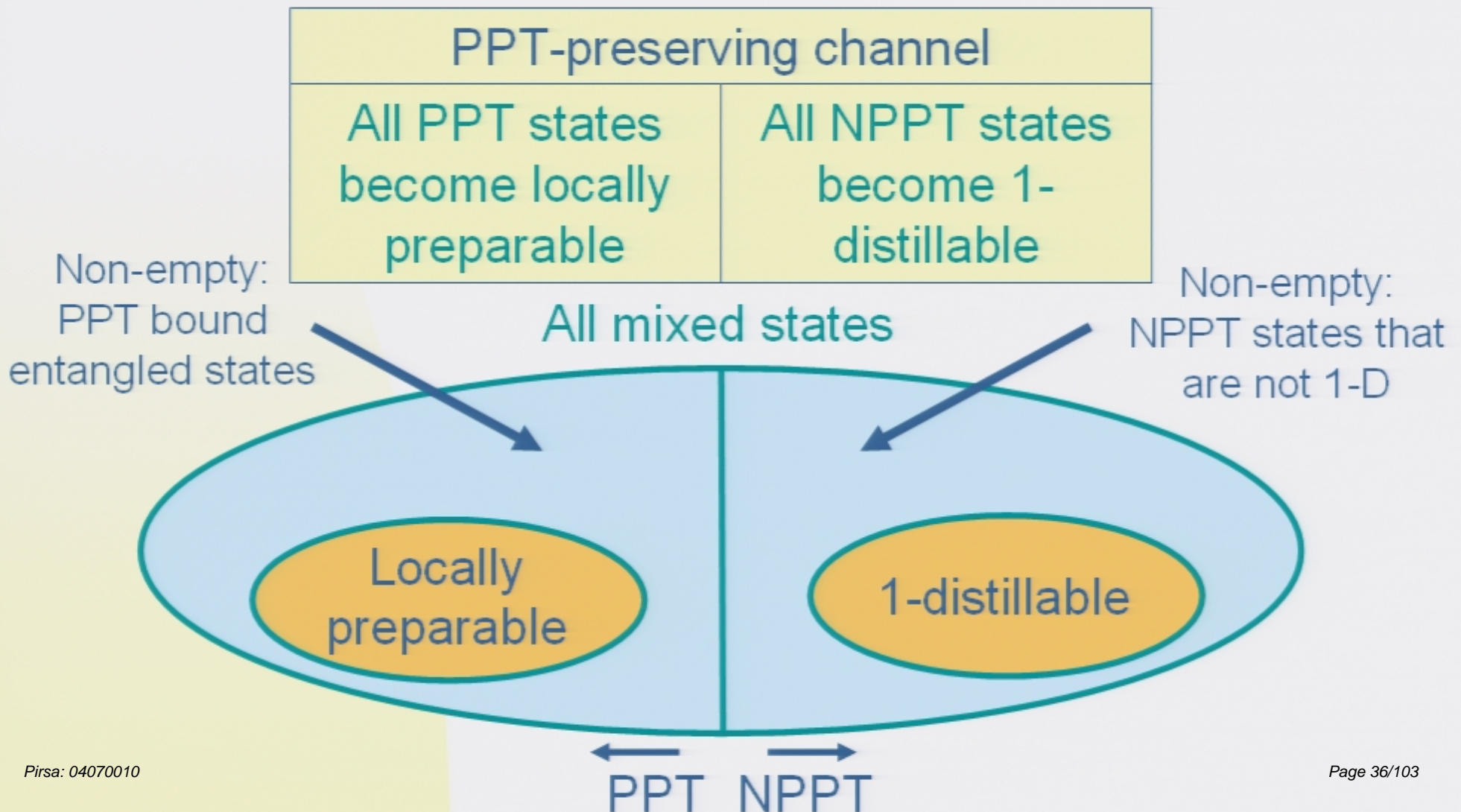
A helpful channel

- Channel C: distributes PPT states (PPT-preserving channel)
- A resource in addition to LOCC – can do more



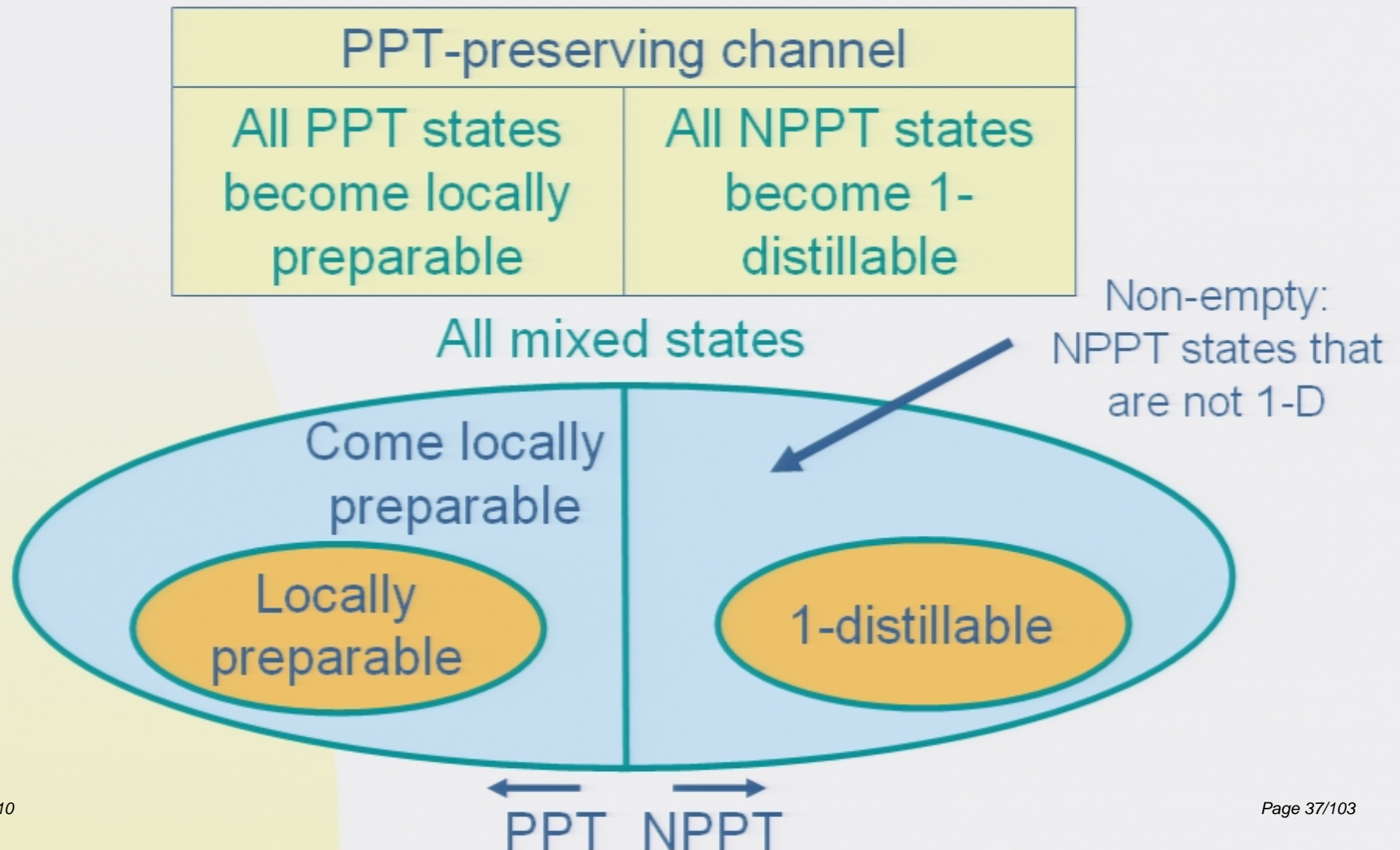
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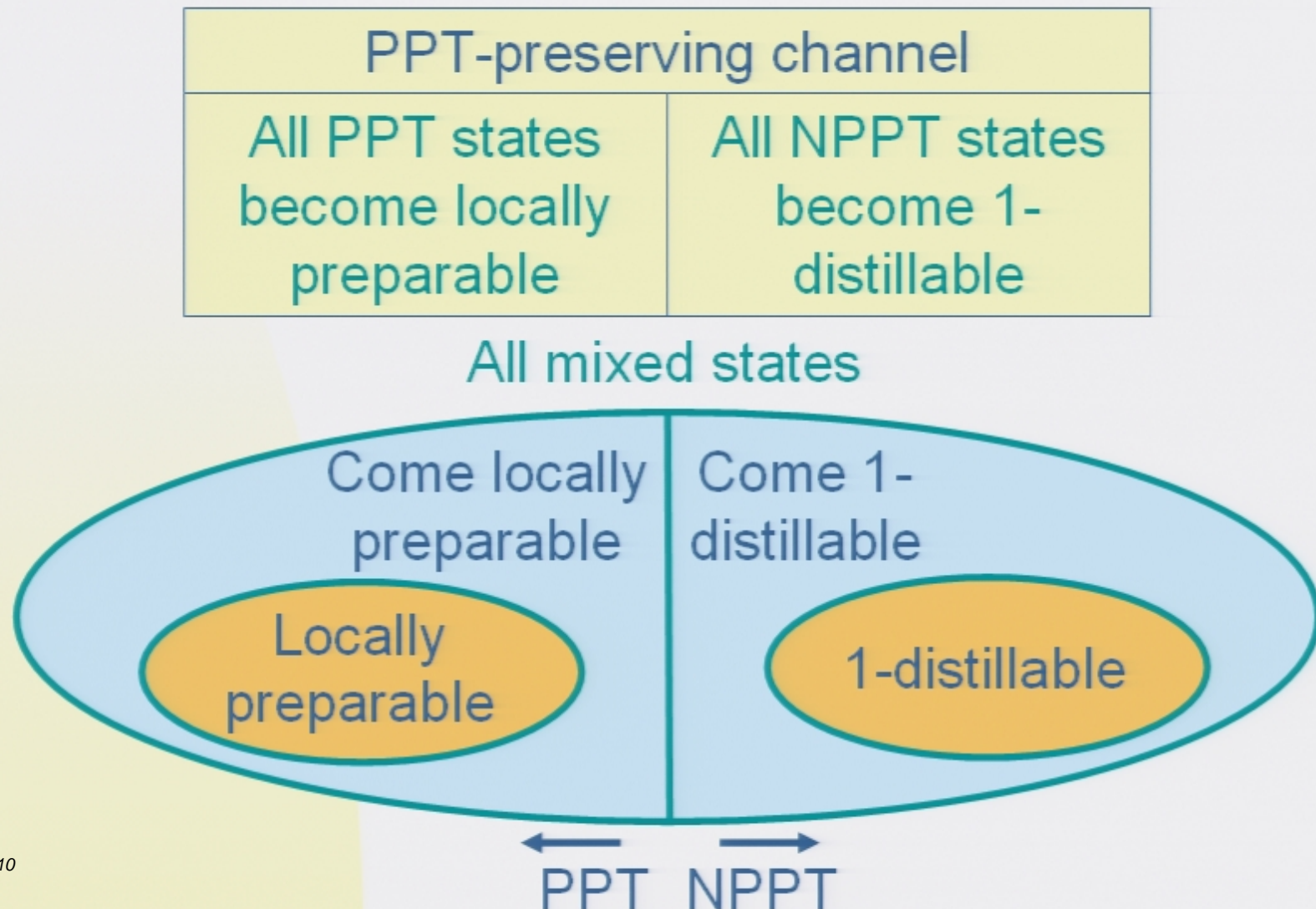
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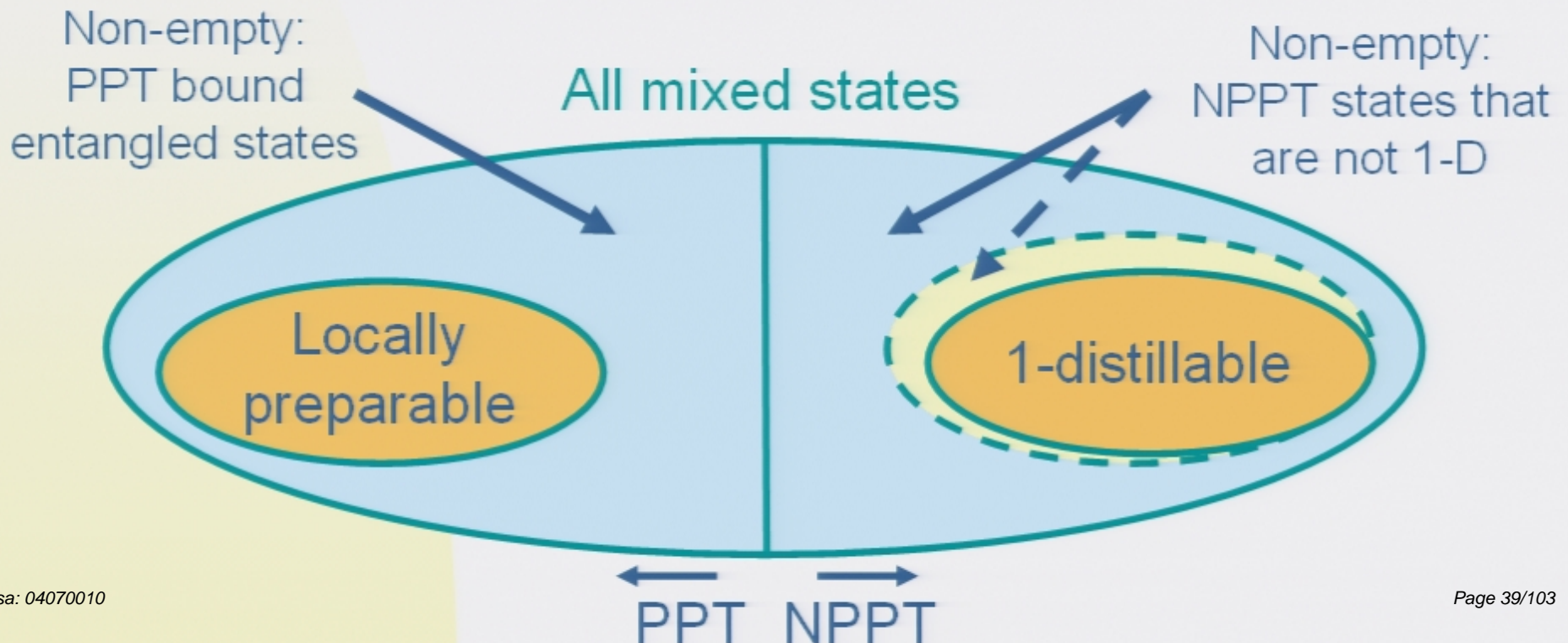


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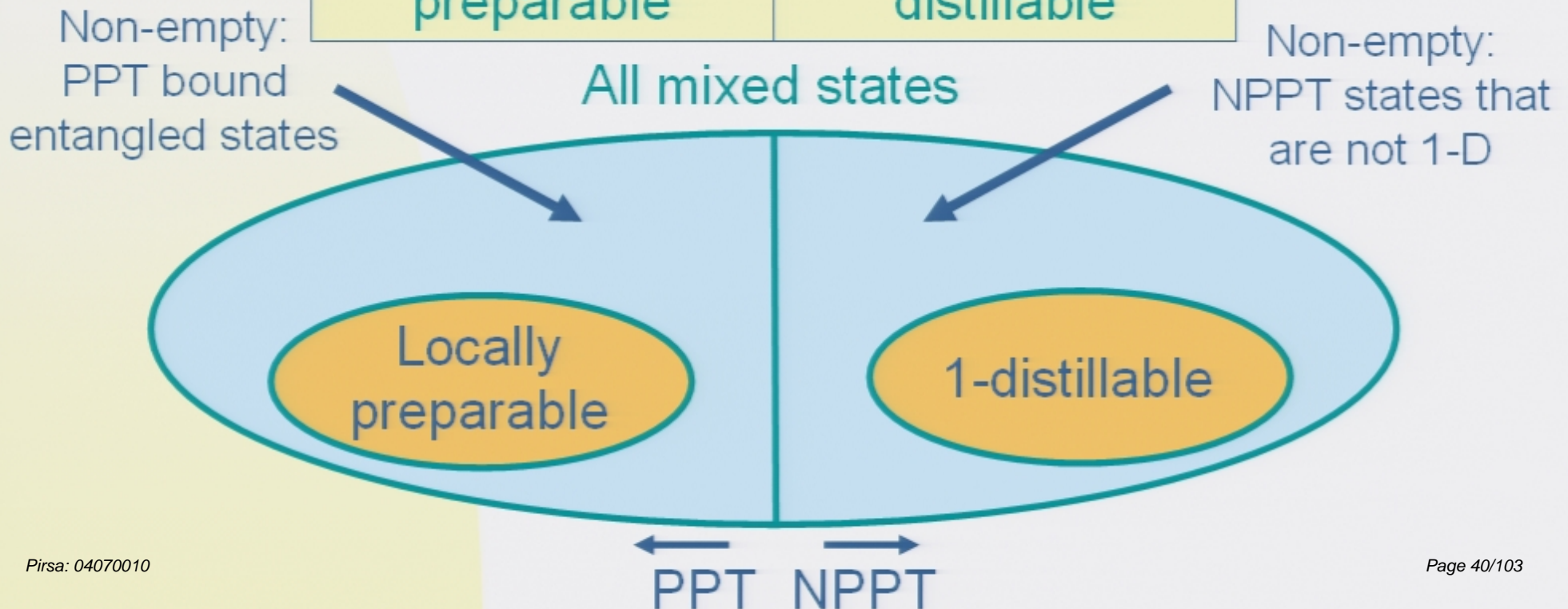
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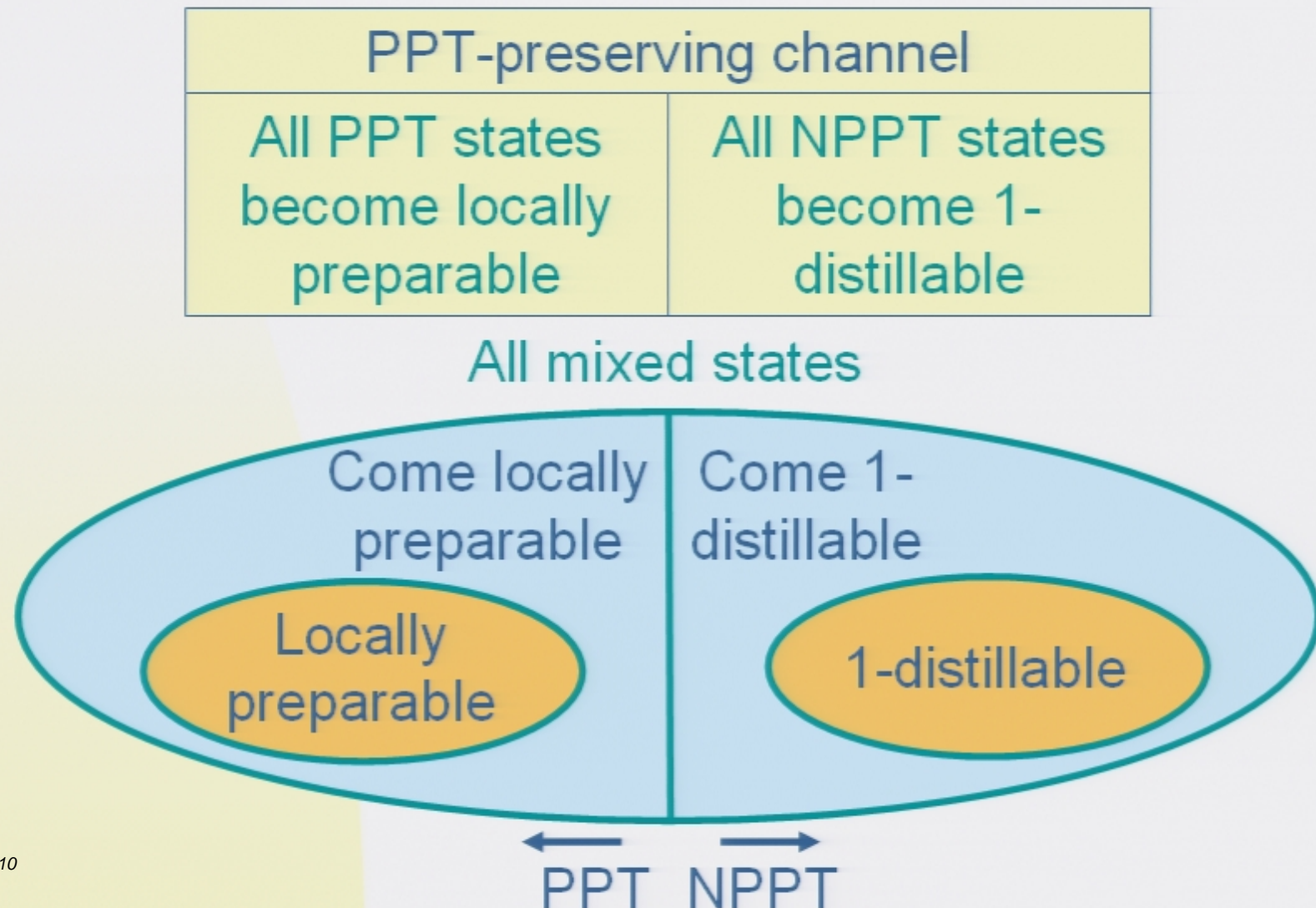
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PPT-preserving channel	
All PPT states become locally preparable	All NPPT states become 1-distillable



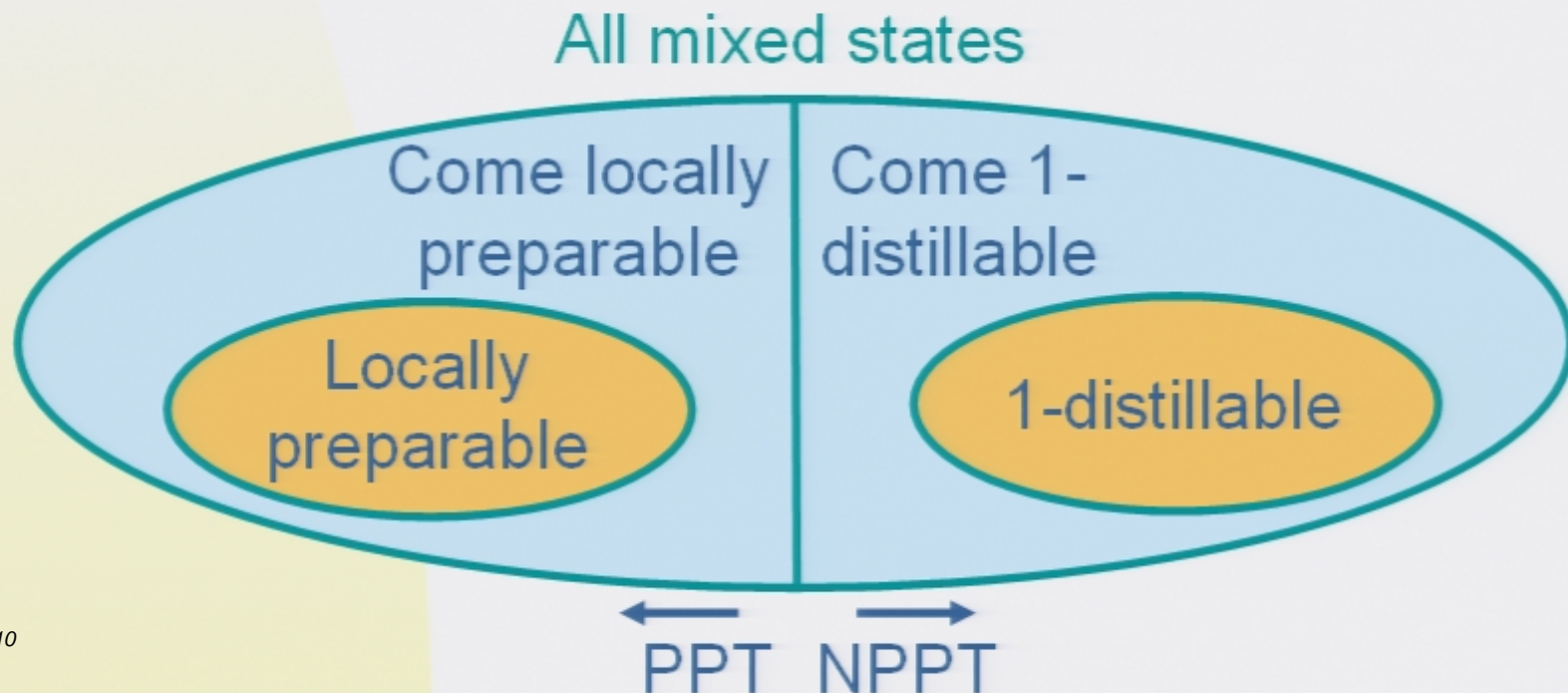
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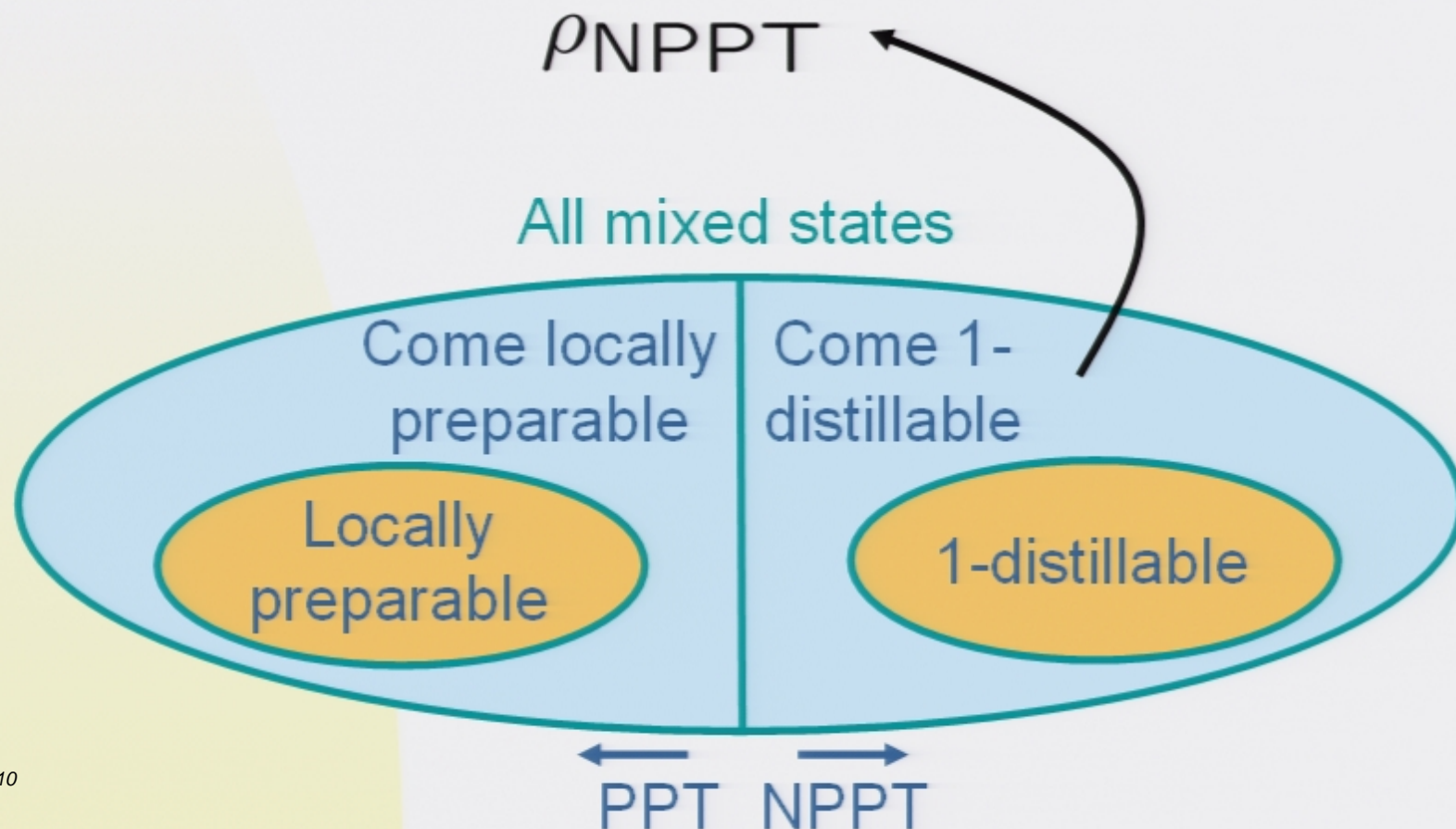
Activation

- Channel C: implemented probabilistically using PPT states (Jamiołkowski isomorphism)
- For every "come 1-distillable" state, there exists a "come locally preparable" state that *activates* the 1-distillability



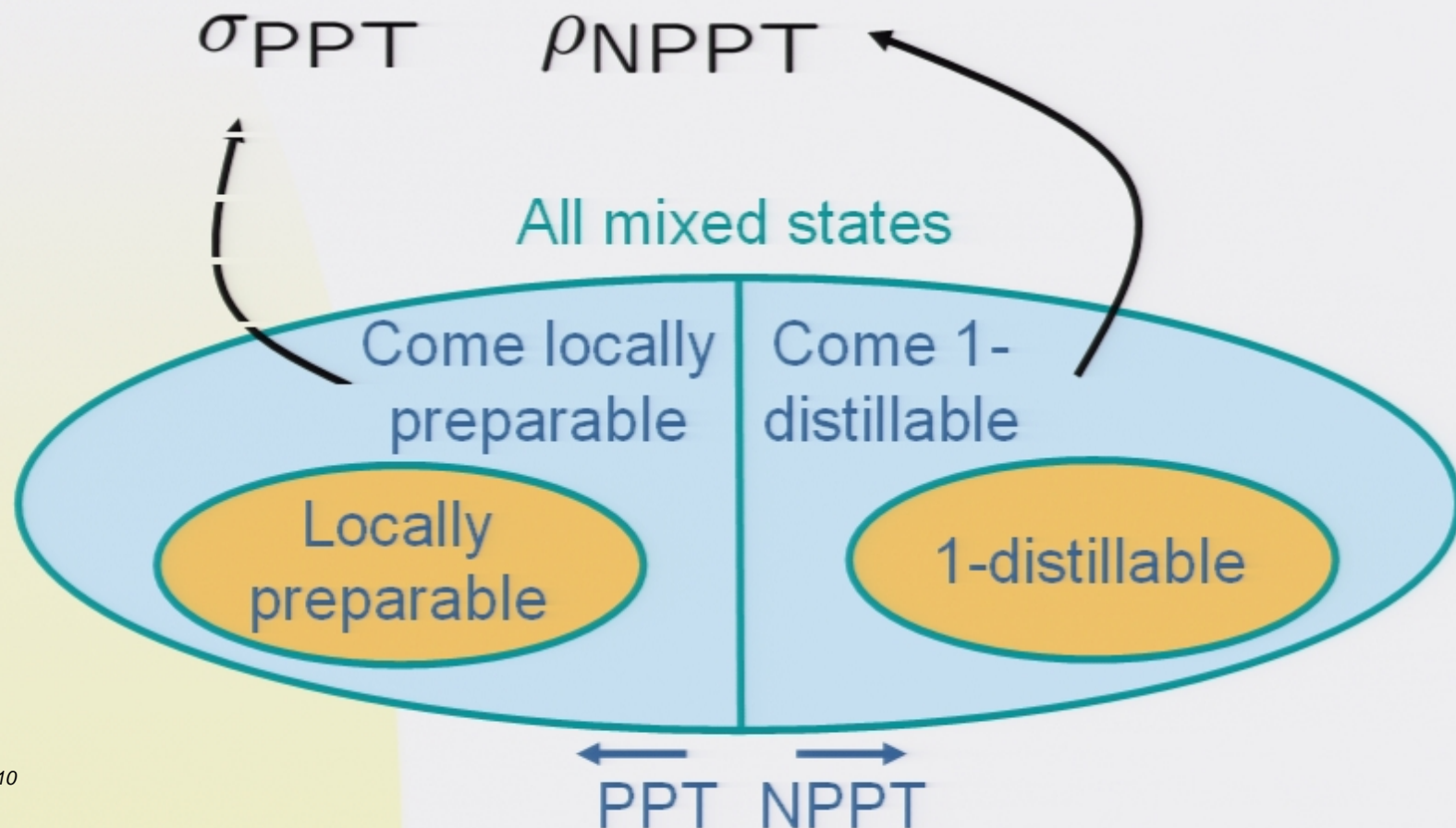
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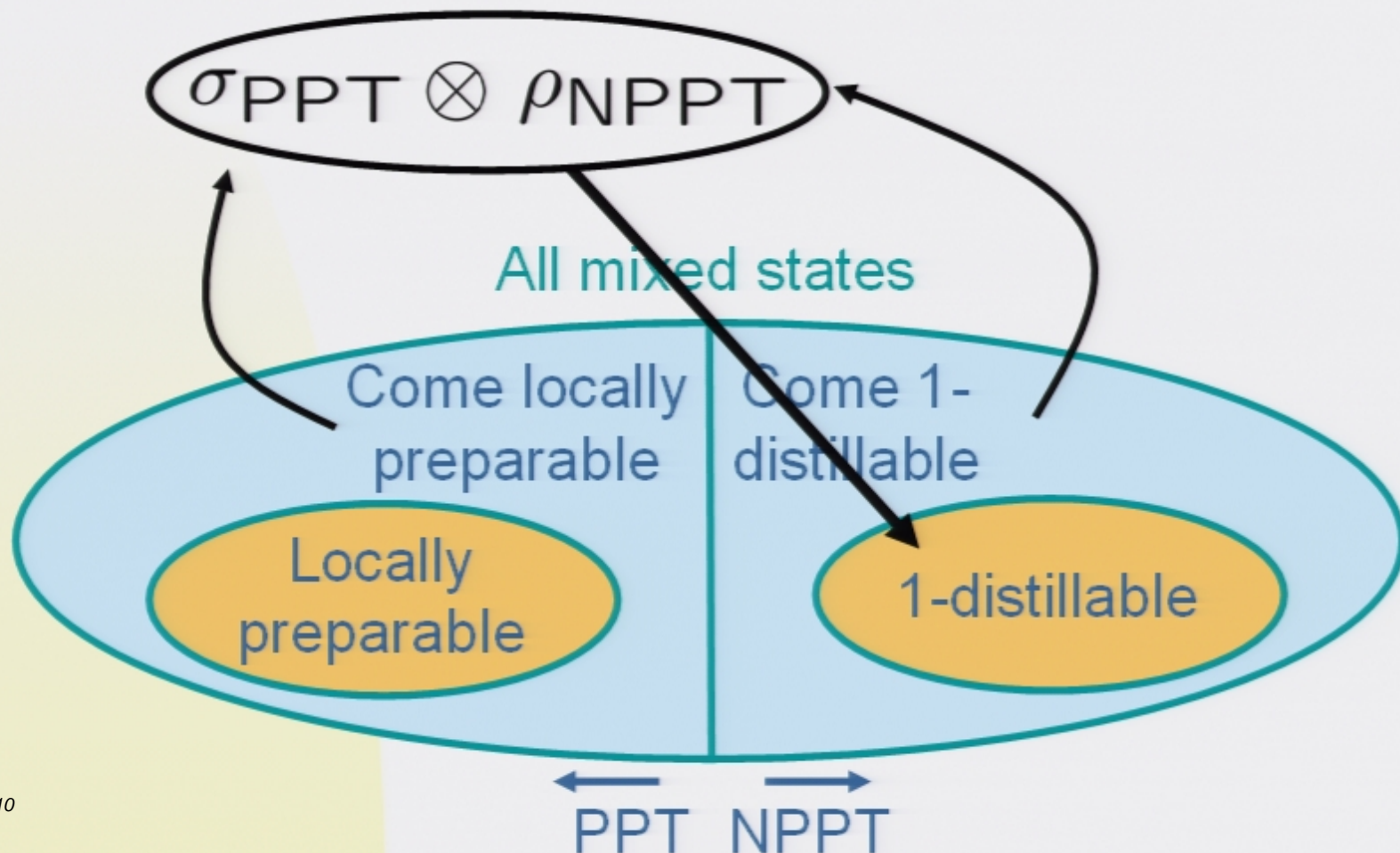
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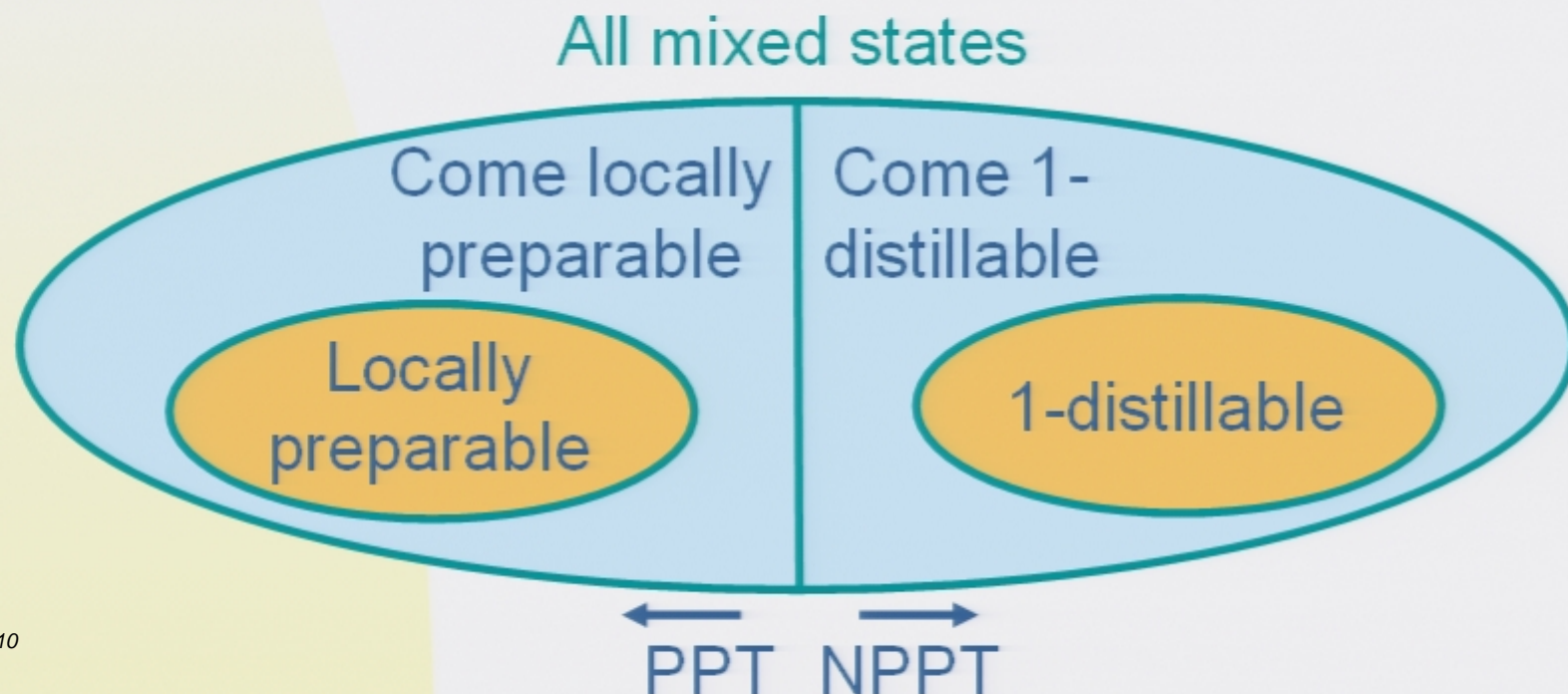
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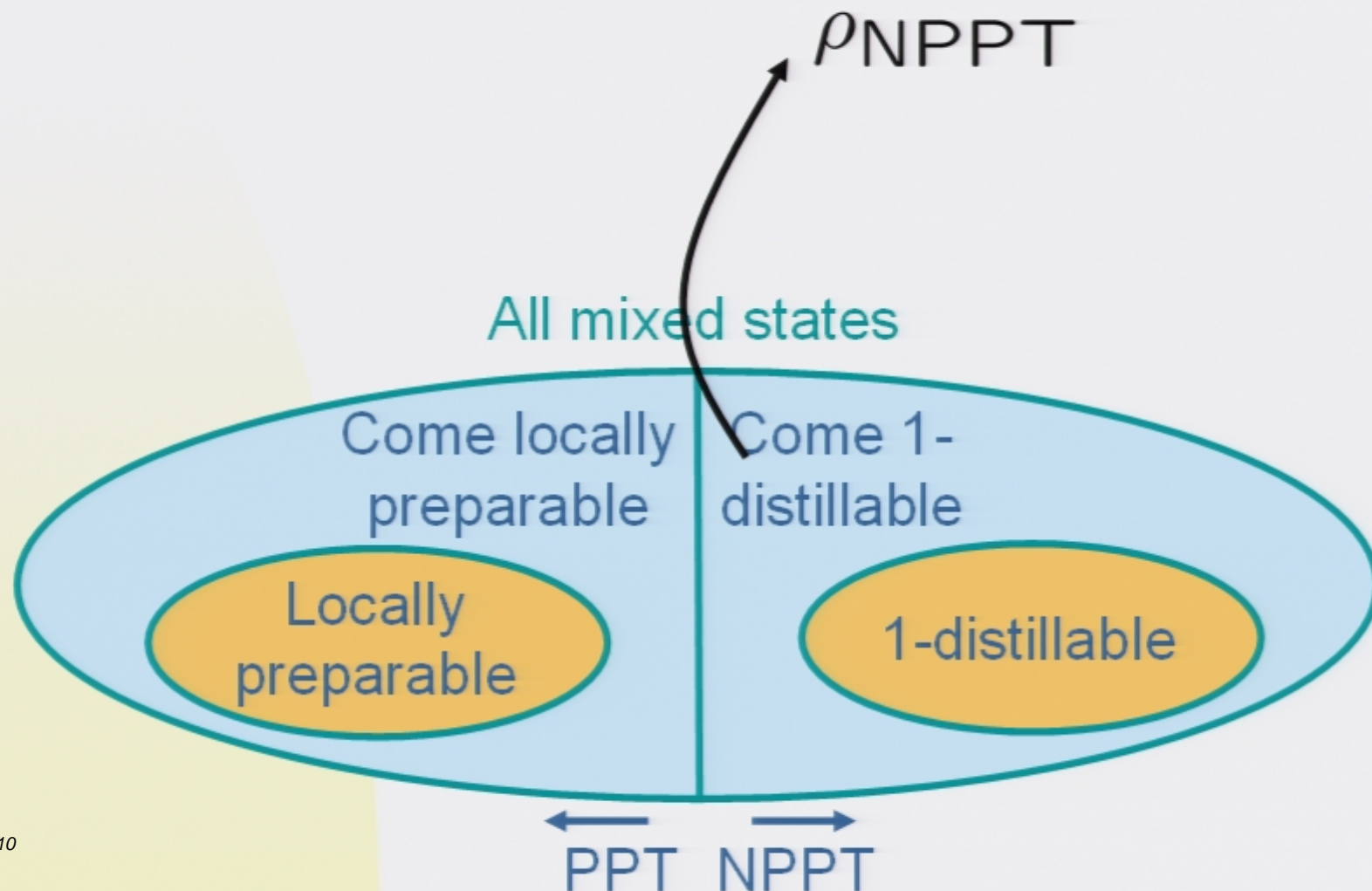
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- Some "come 1-distillable" states, although not 1-distillable, become 1-distillable when given 2 (or N) copies



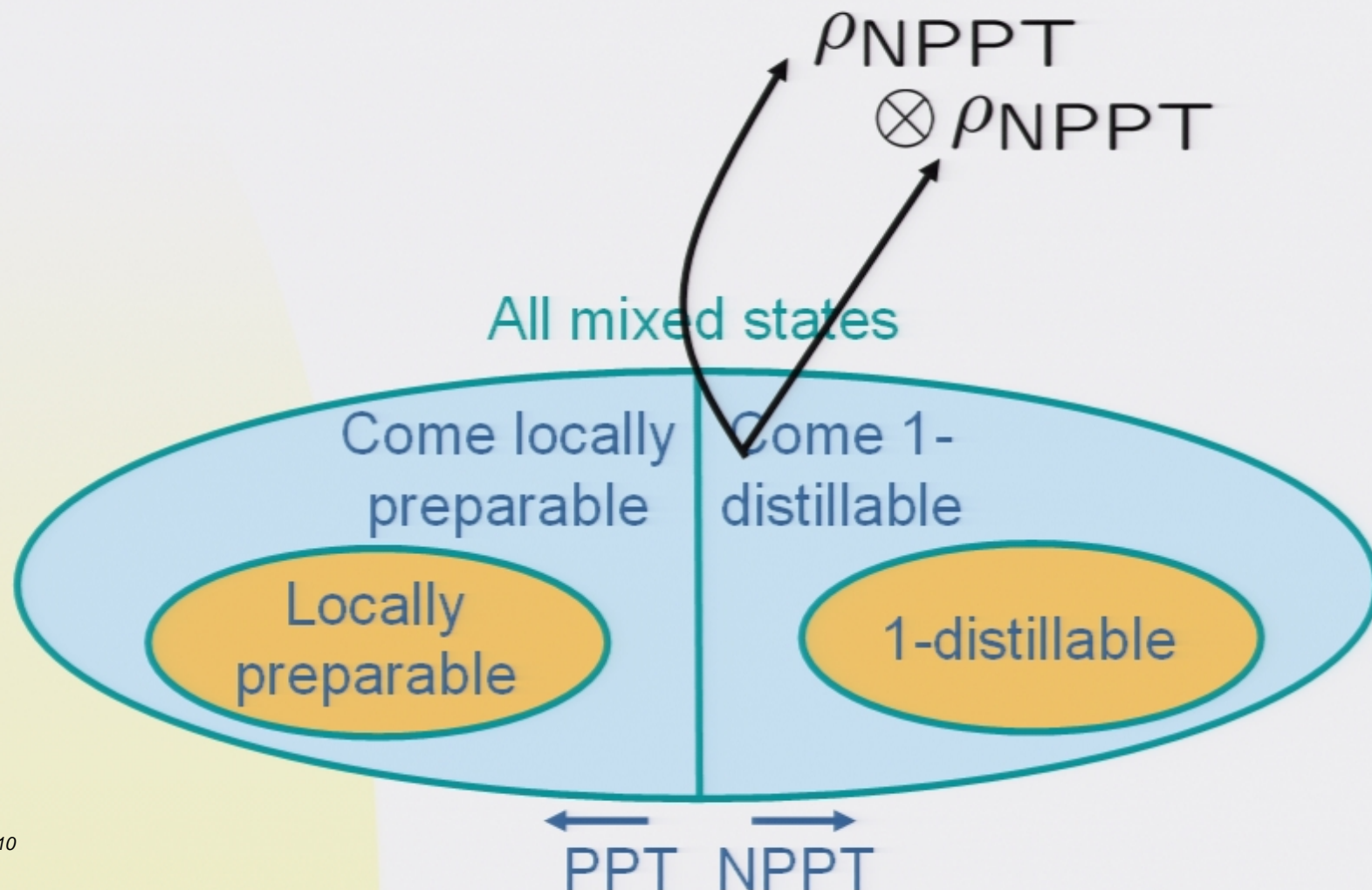
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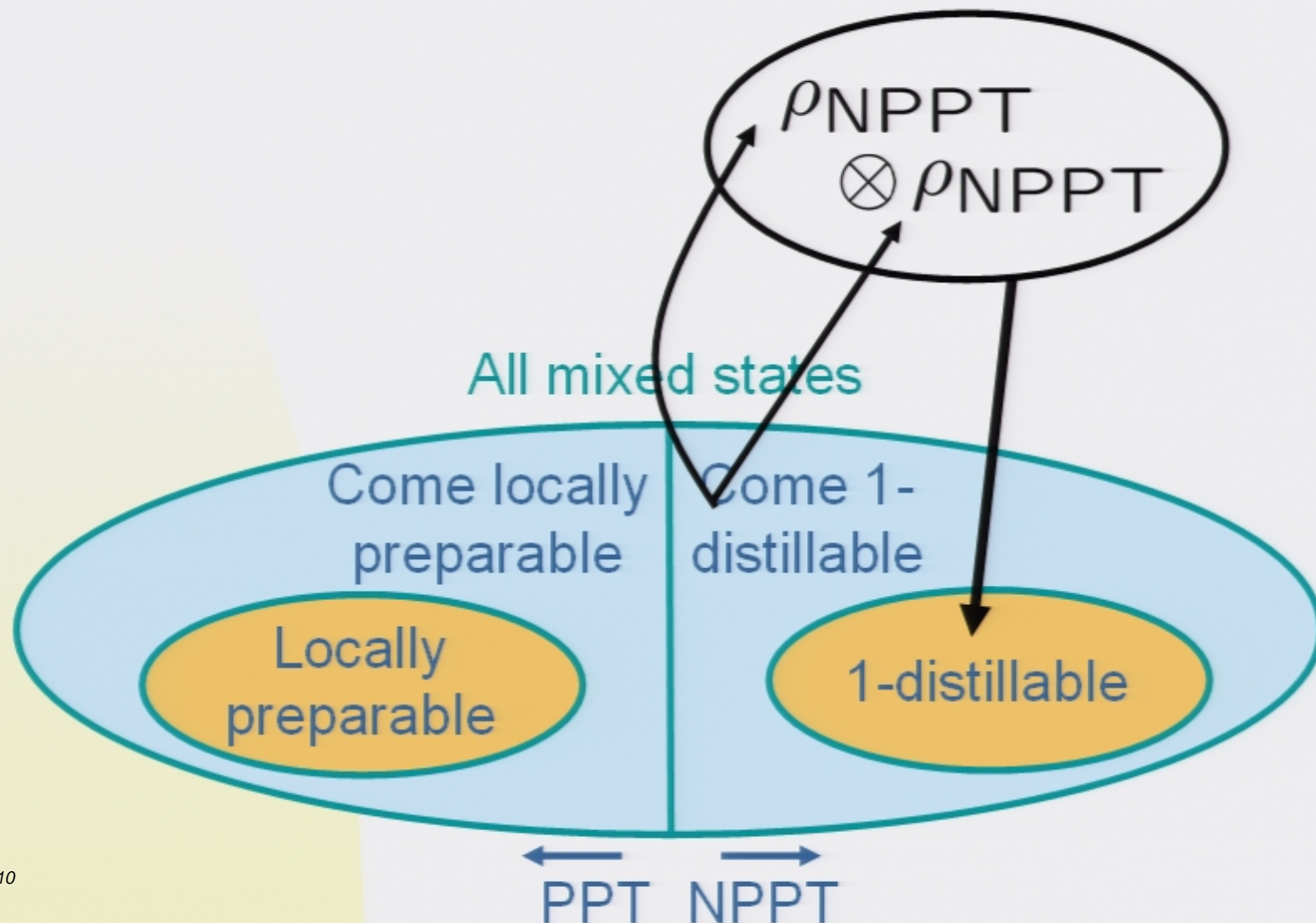
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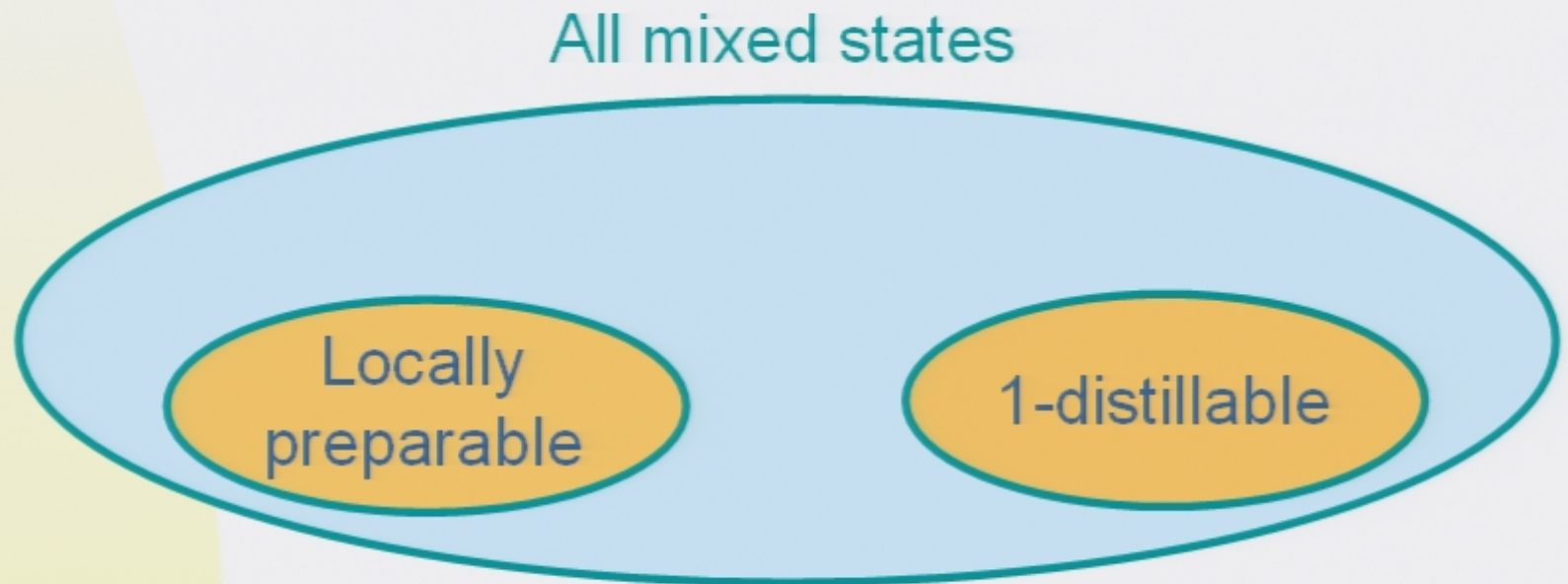
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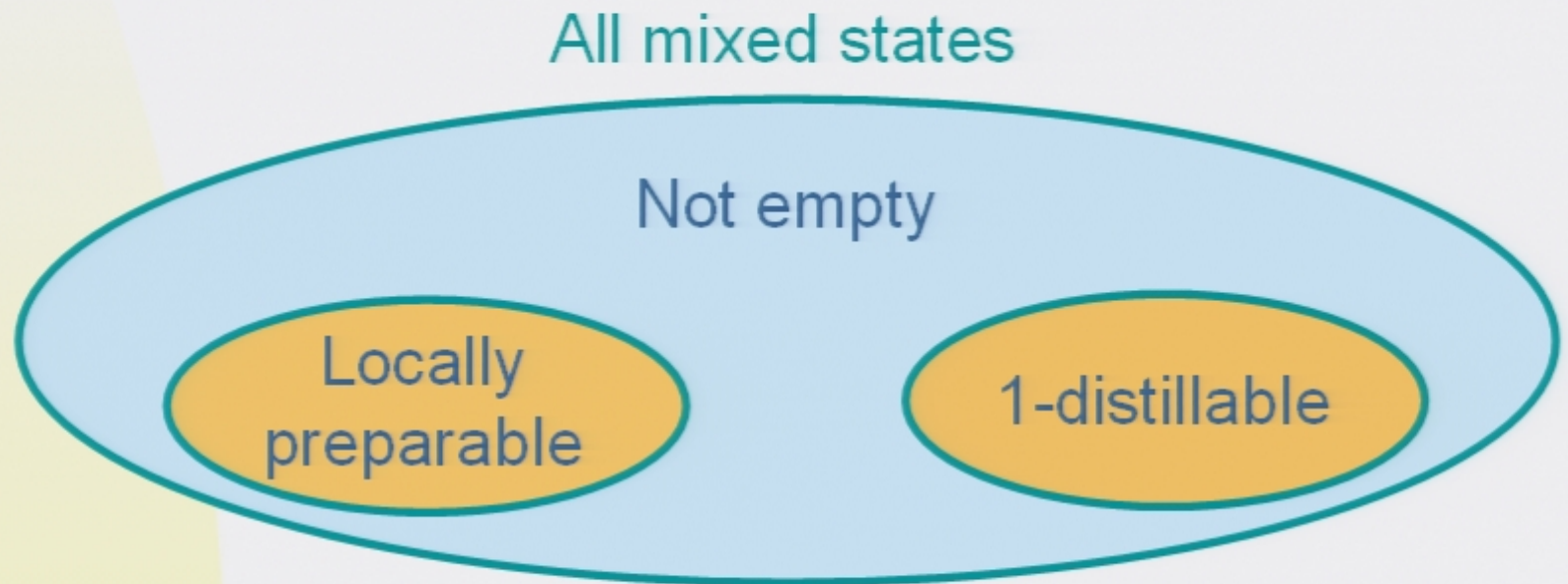
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- Locally-preparable and 1-distillable states



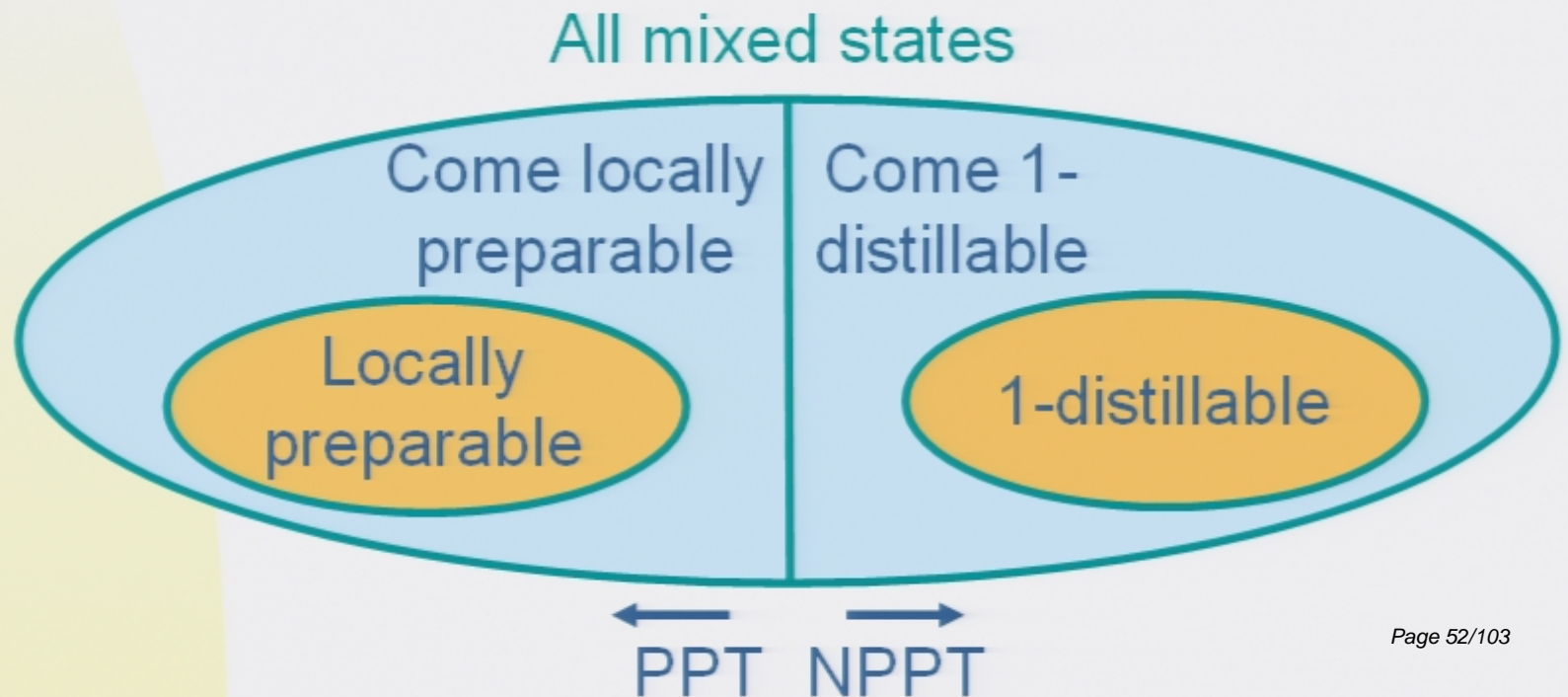
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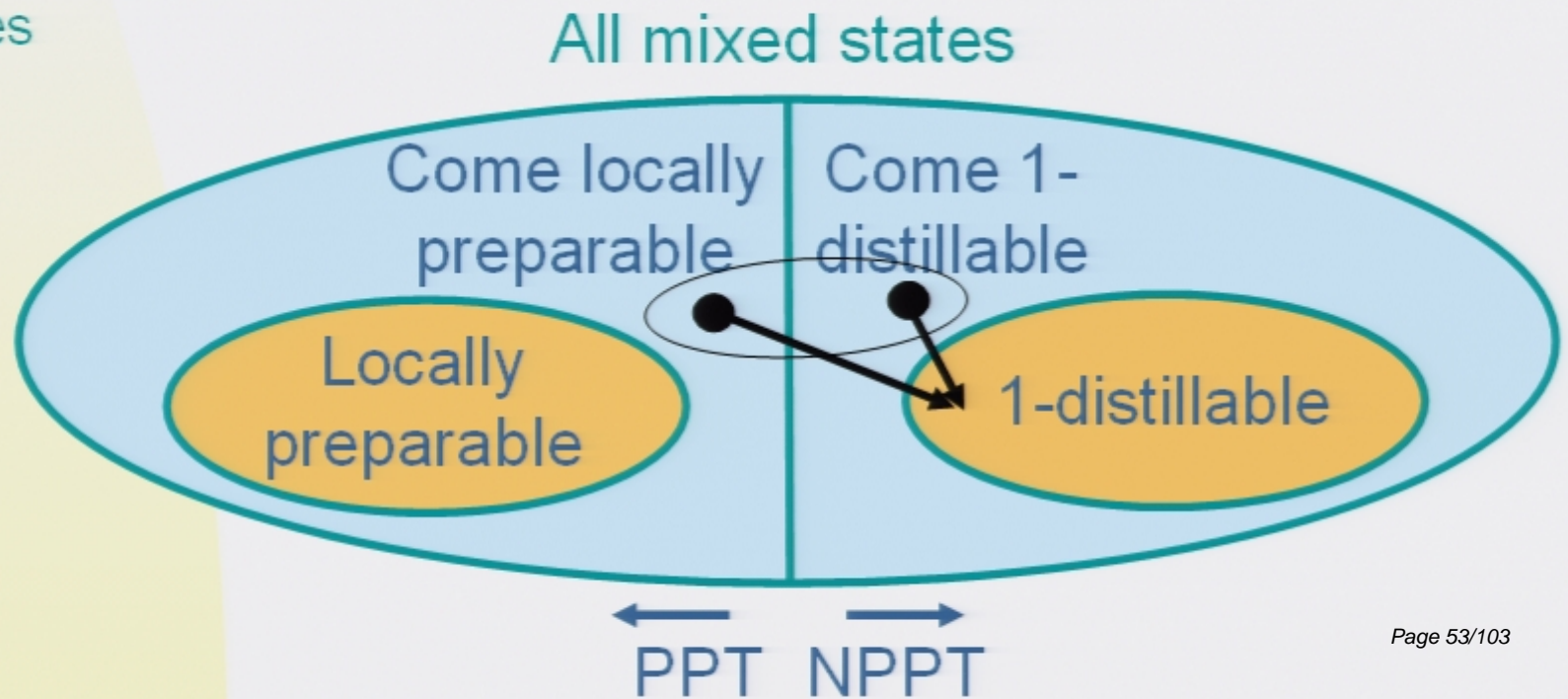
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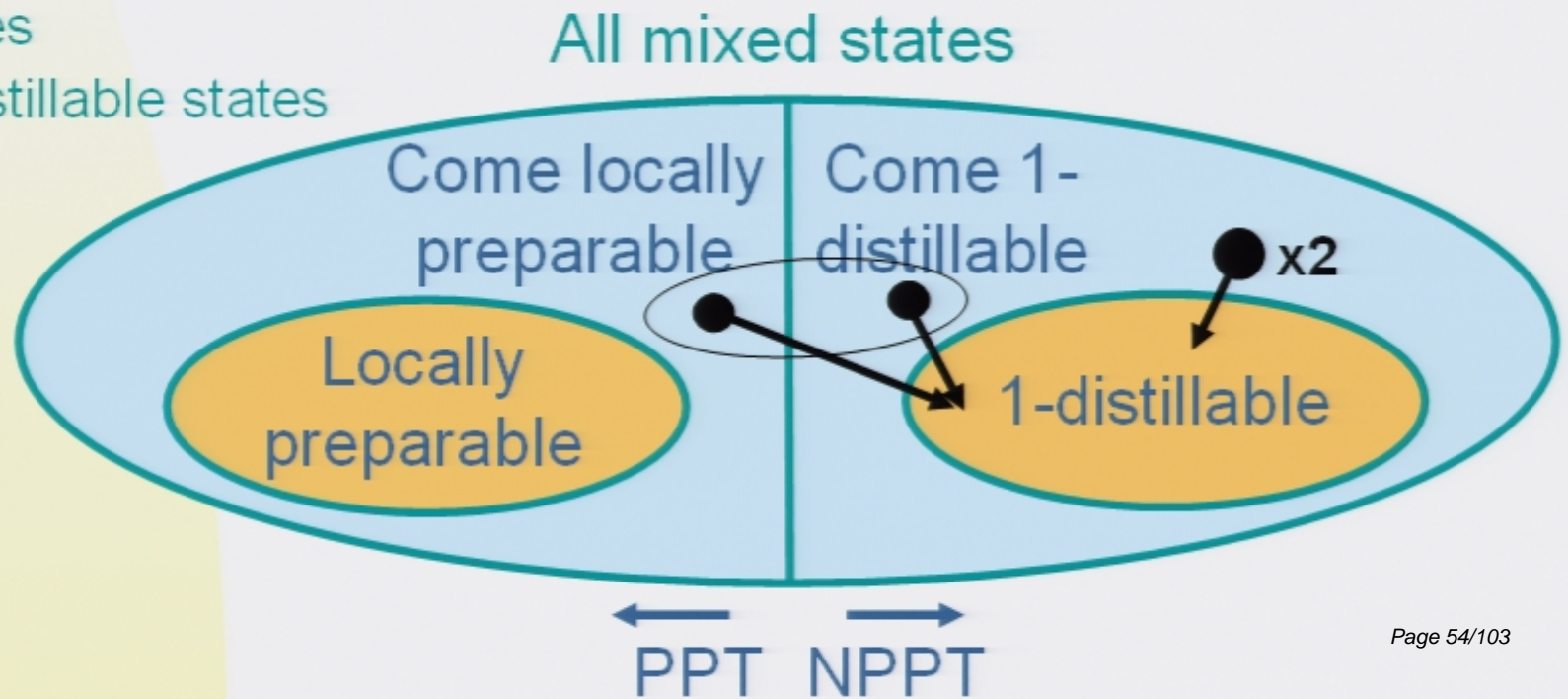
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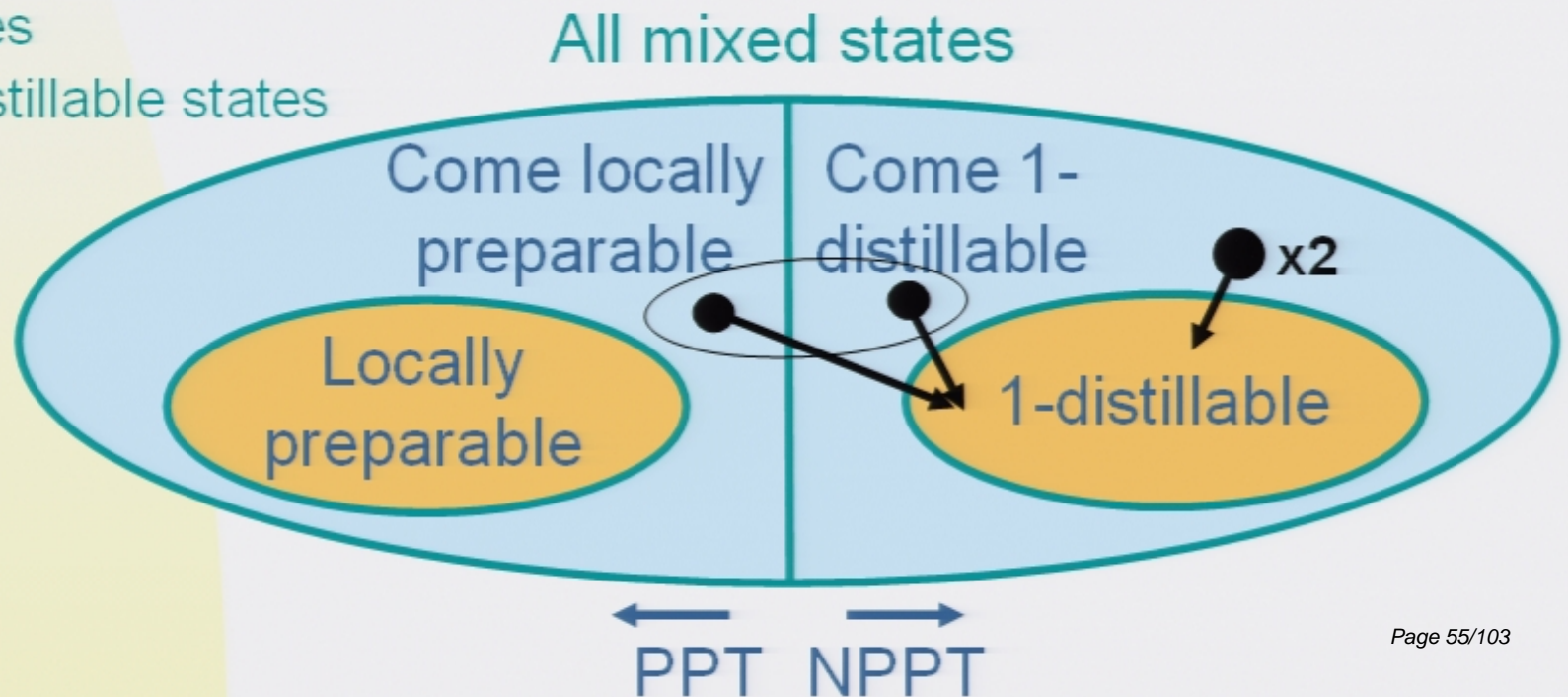


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Big open question:

Are all C1-D states distillable?



Entanglement constrained by superselection rules

Operational SSRs

A superselection rule is an additional *restriction*

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G-SSR: Restriction on allowed operations; CP maps \mathcal{O} are *G-covariant*

$$\begin{aligned}\mathcal{O}[T(g)\rho T^\dagger(g)] \\ &= T(g)\mathcal{O}[\rho]T^\dagger(g) \\ &\quad \forall g \in G, \rho\end{aligned}$$

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■ U(1)-SSR

- ◆ Generated by a Hermitian operator \hat{Q}

$$T(\xi) = \exp(i\xi\hat{Q})$$

- ◆ All operations must commute with \hat{Q}
- ◆ Describes SSRs for charge, particle number
- ◆ Operations block-diagonal wrt eigenspaces of \hat{Q}

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■ SU(2)-SSR

- ◆ Generated by angular momentum operators
$$\{\hat{L}_x, \hat{L}_y, \hat{L}_z\}$$
- ◆ All operations are rotationally invariant
- ◆ No direction axis to prepare, measure spins

Effect of G -covariance

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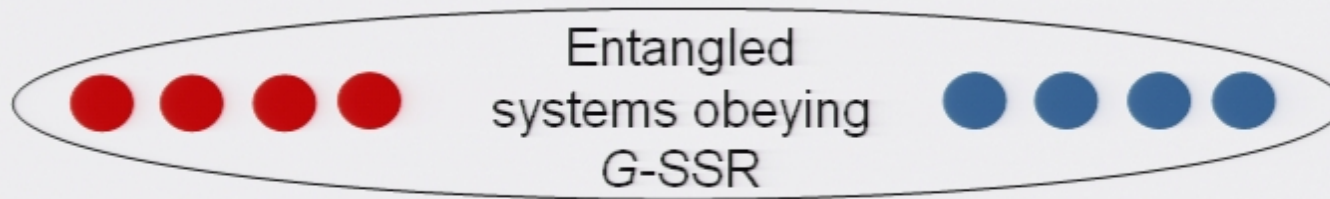
- U(1) example:

$$\frac{1}{\sqrt{2}}(|q_1\rangle + |q_2\rangle) \leftrightarrow \frac{1}{2}(|q_1\rangle\langle q_1| + |q_2\rangle\langle q_2|)$$

Entanglement with SSRs

Alice

Bob



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- All operations (including preparations, transformations and measurements) must be G-covariant in the presence of a G-SSR
- Any state ρ is indistinguishable from $T(g)\rho T^\dagger(g)$
- Operationally equivalent to describe the state as

$$\mathcal{G}[\rho] = \int_G dv(g) T(g) \rho T^\dagger(g)$$

- SSR as *decoherence*:

~~Environment~~

Indistinguishability

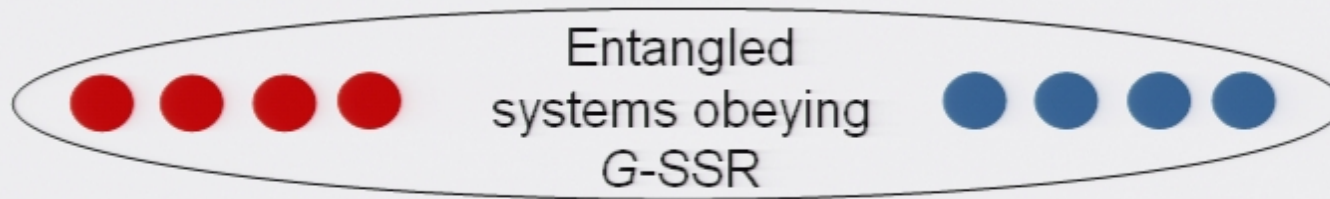
- U(1) example:

$$\frac{1}{\sqrt{2}}(|q_1\rangle + |q_2\rangle) \leftrightarrow \frac{1}{2}(|q_1\rangle\langle q_1| + |q_2\rangle\langle q_2|)$$

Entanglement with SSRs

Alice

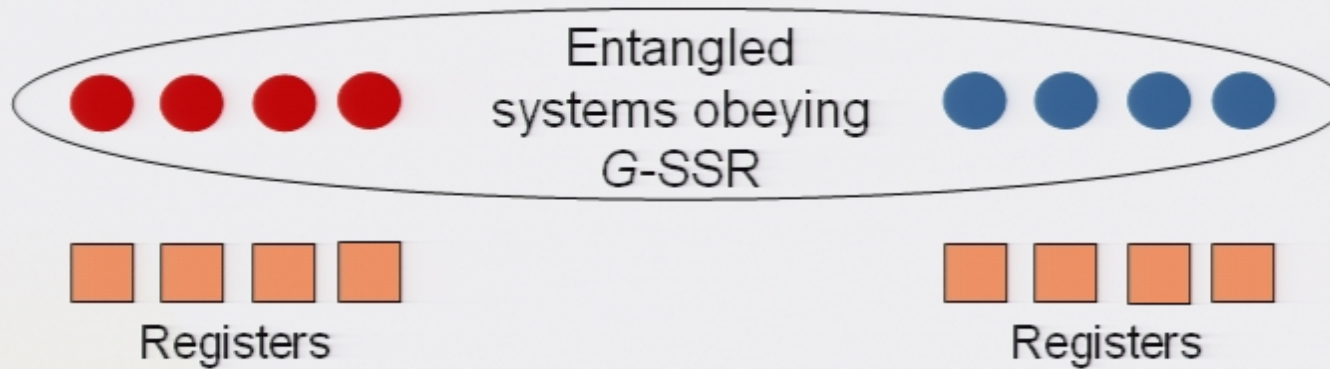
Bob



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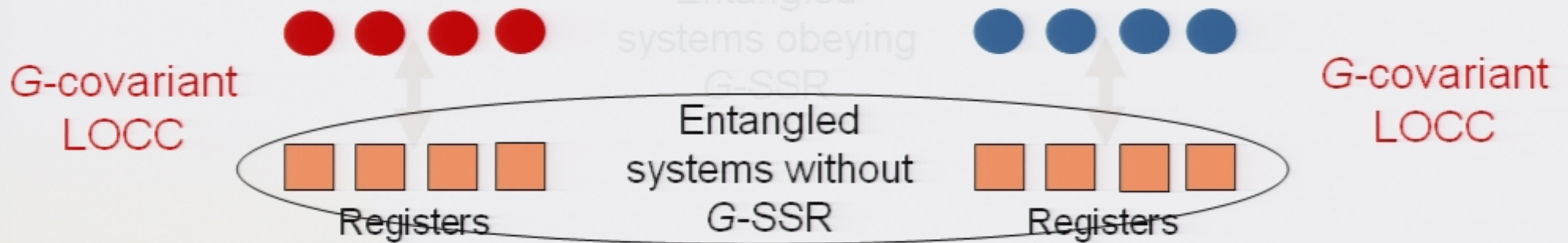
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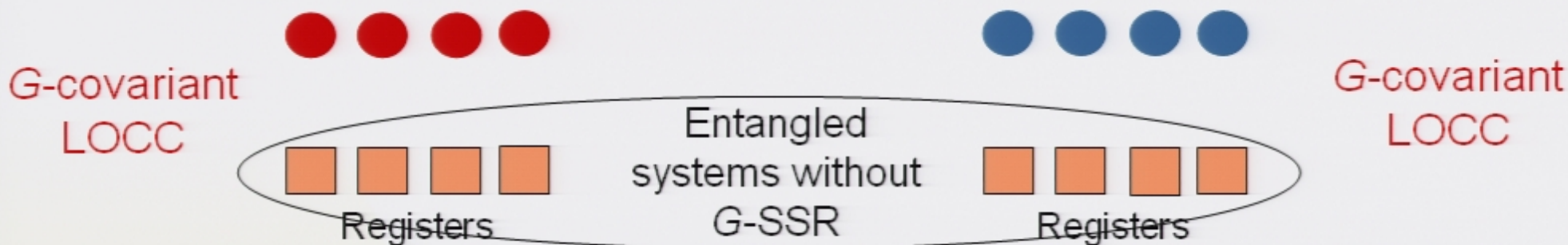
Bob



Entanglement with SSRs

Alice

Bob



- Quantify entanglement transferable to "standard quantum registers" using LOCC obeying the SSR

Amount of entanglement present in the state ρ in the presence of G-SSR

=

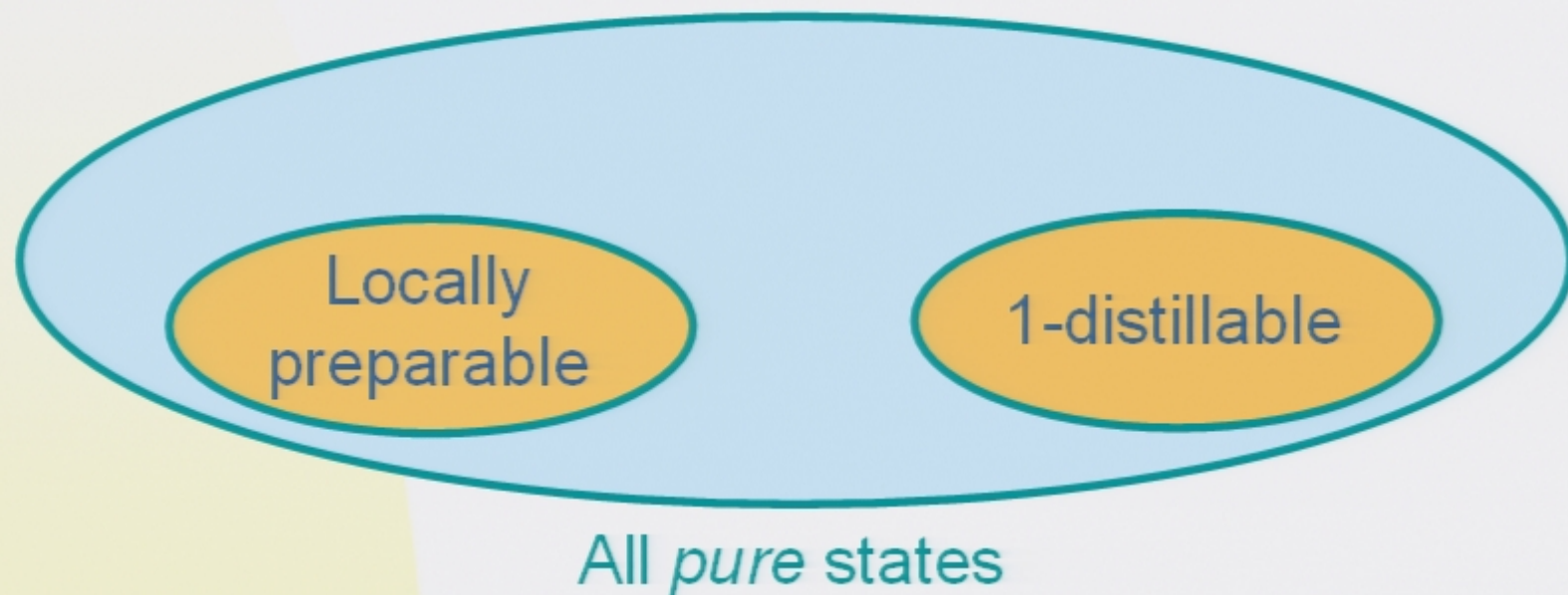
Standard entanglement in $(\mathcal{G}_A \otimes \mathcal{G}_B)[\rho]$

Wiseman and Vaccaro, *Phys. Rev. Lett.* **91**, 097902 (2003)

Bartlett and Wiseman, *Phys. Rev. Lett.* **91**, 097903 (2003)

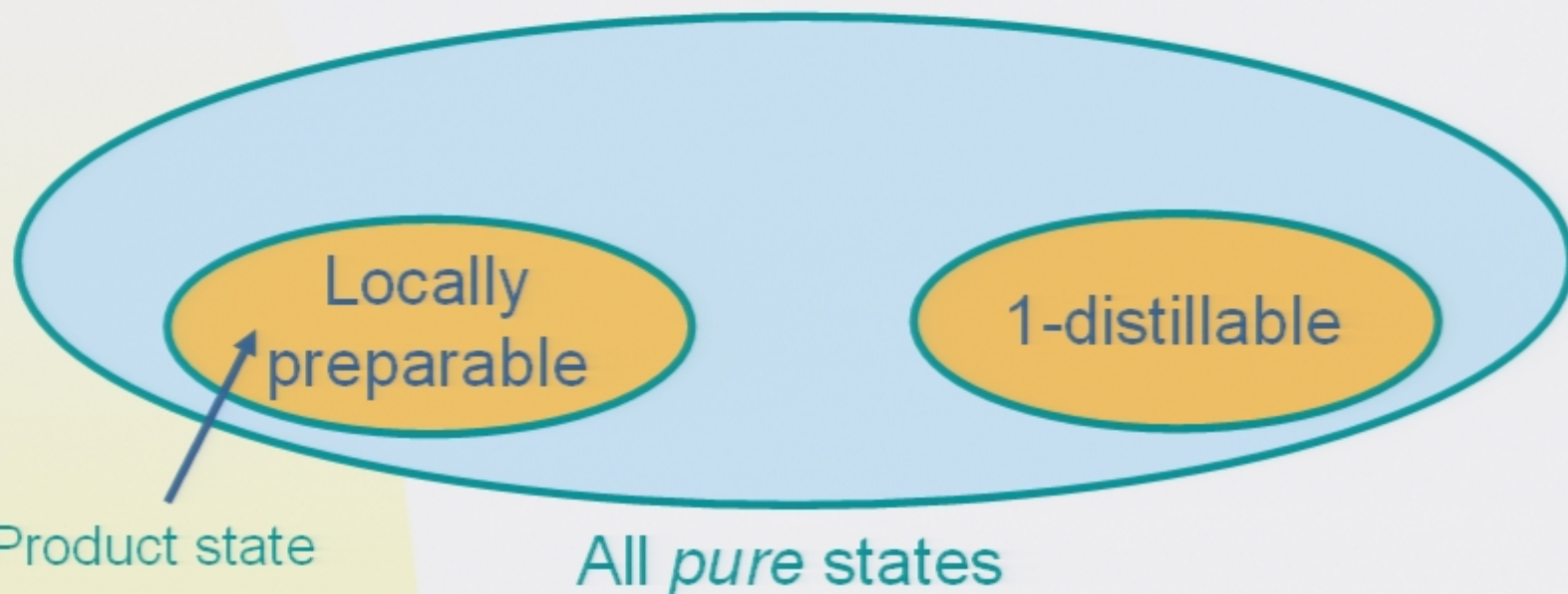
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- Classification mimics the structure of mixed state entanglement



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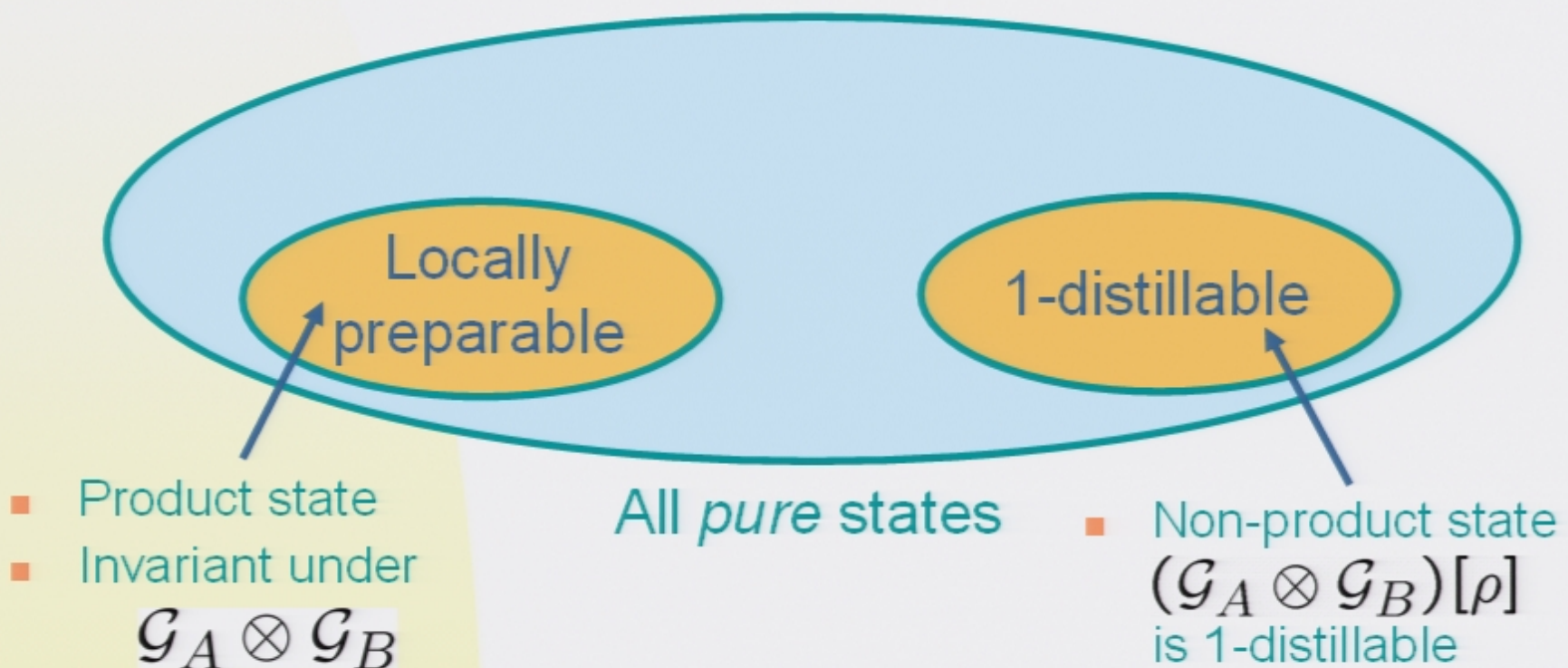


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- Invariant under

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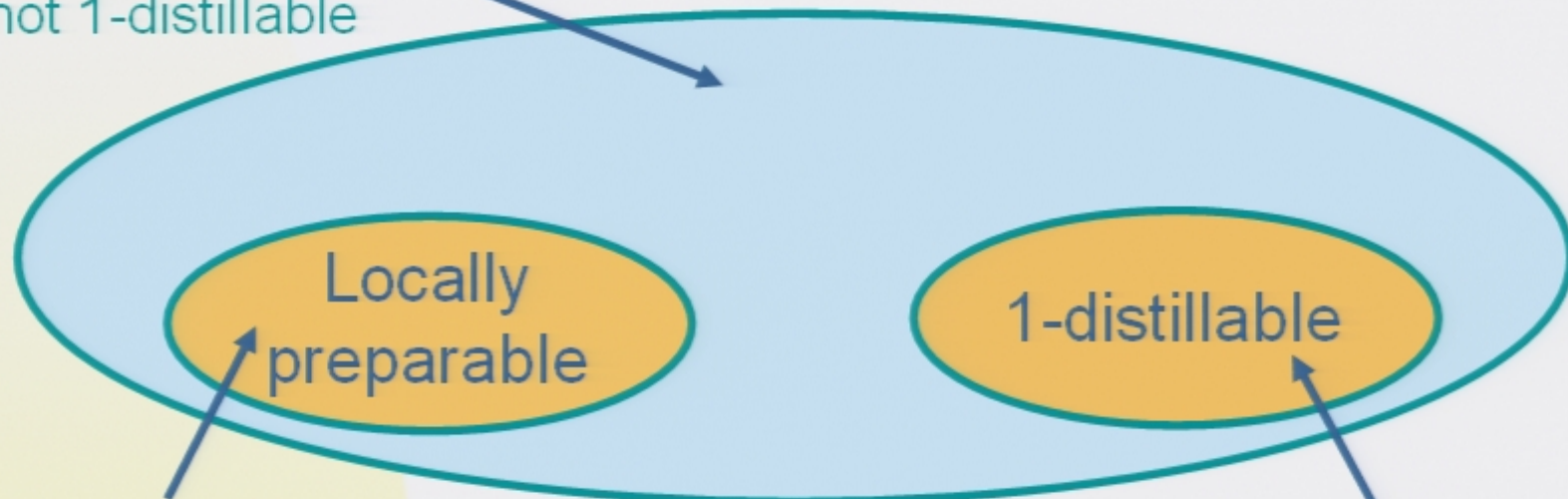
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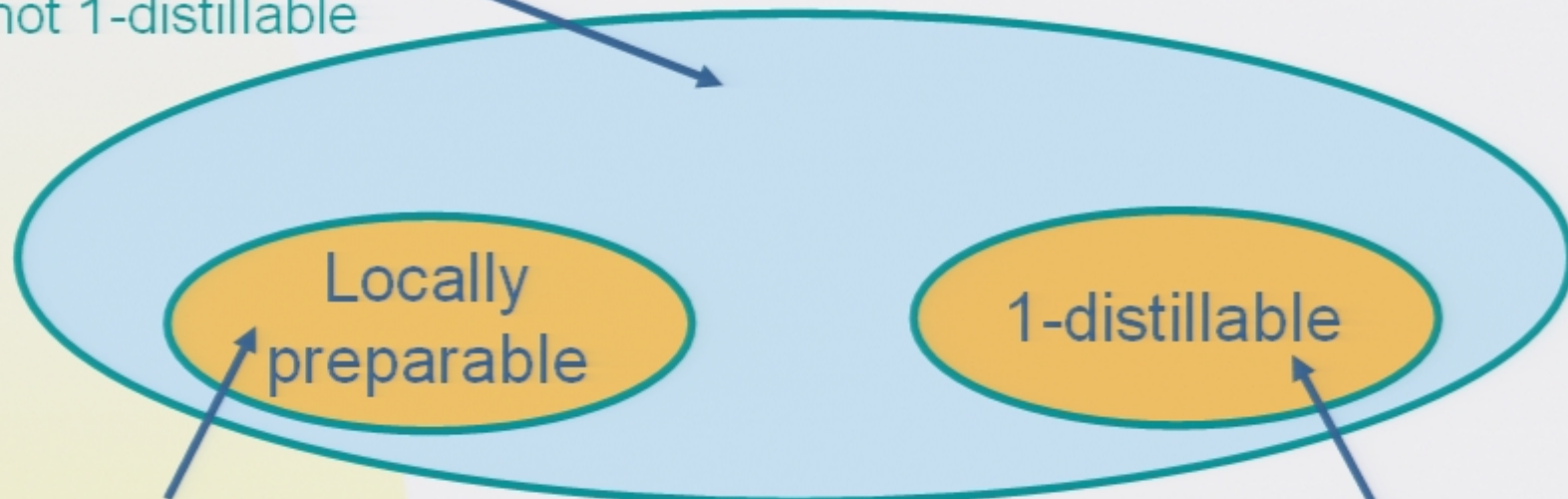
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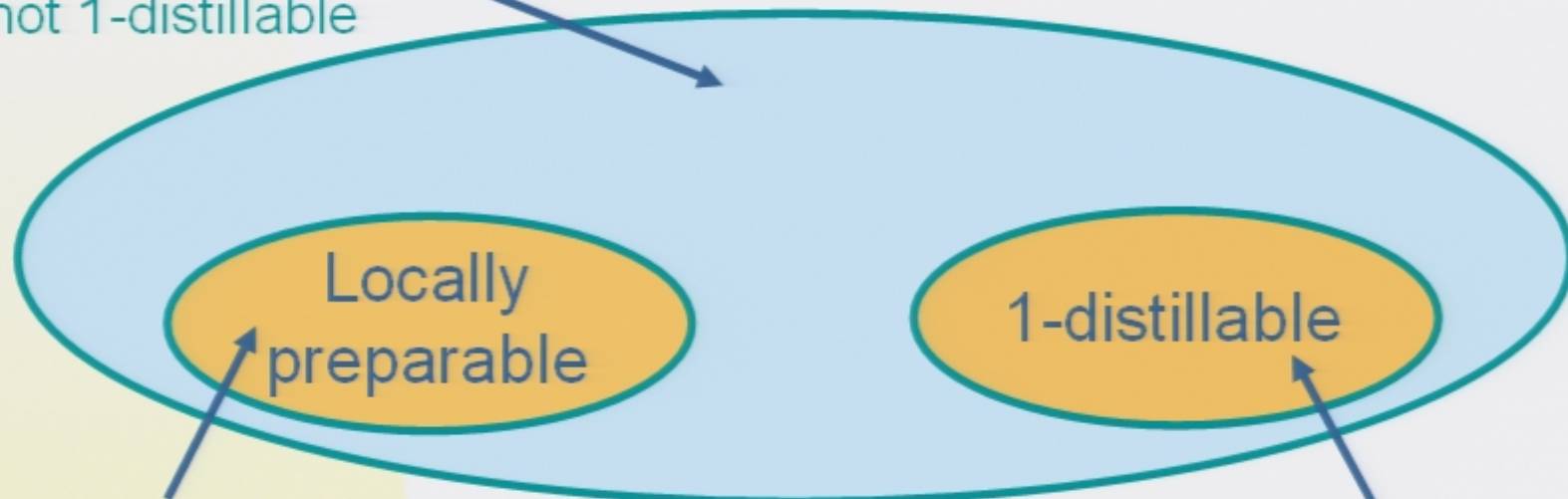
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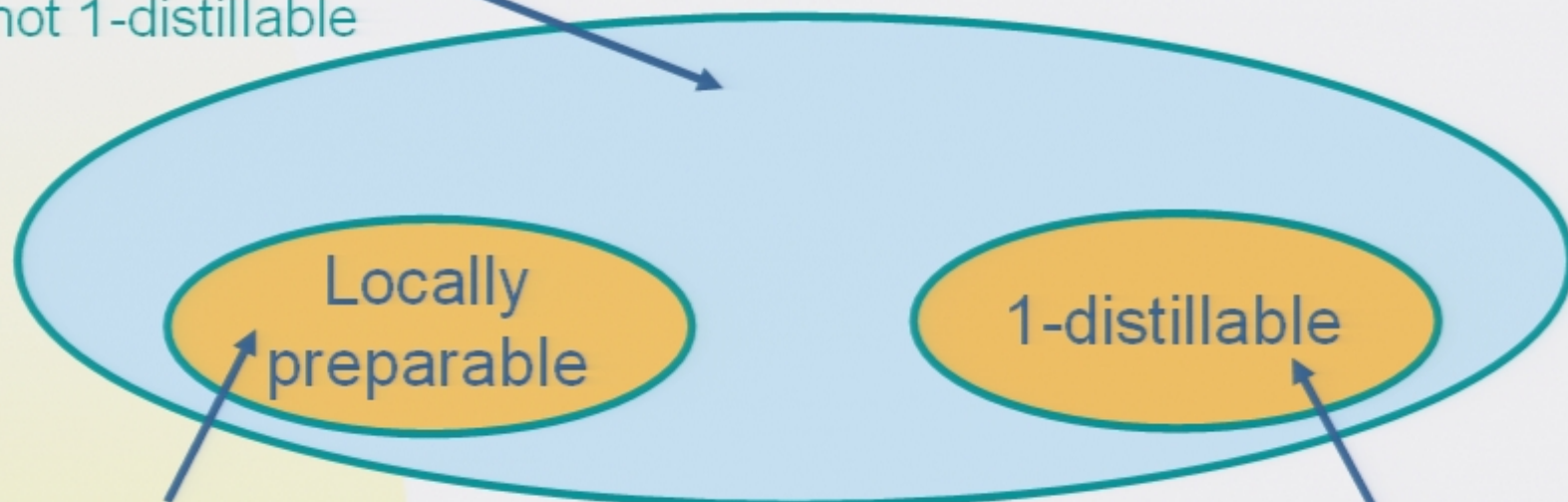
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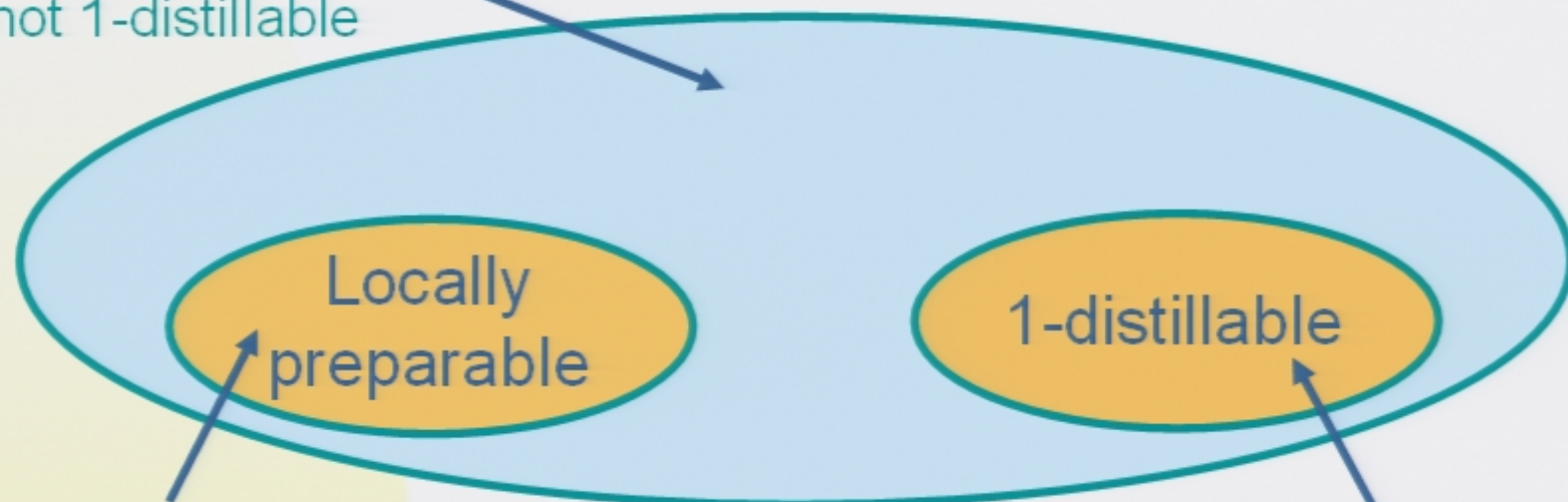
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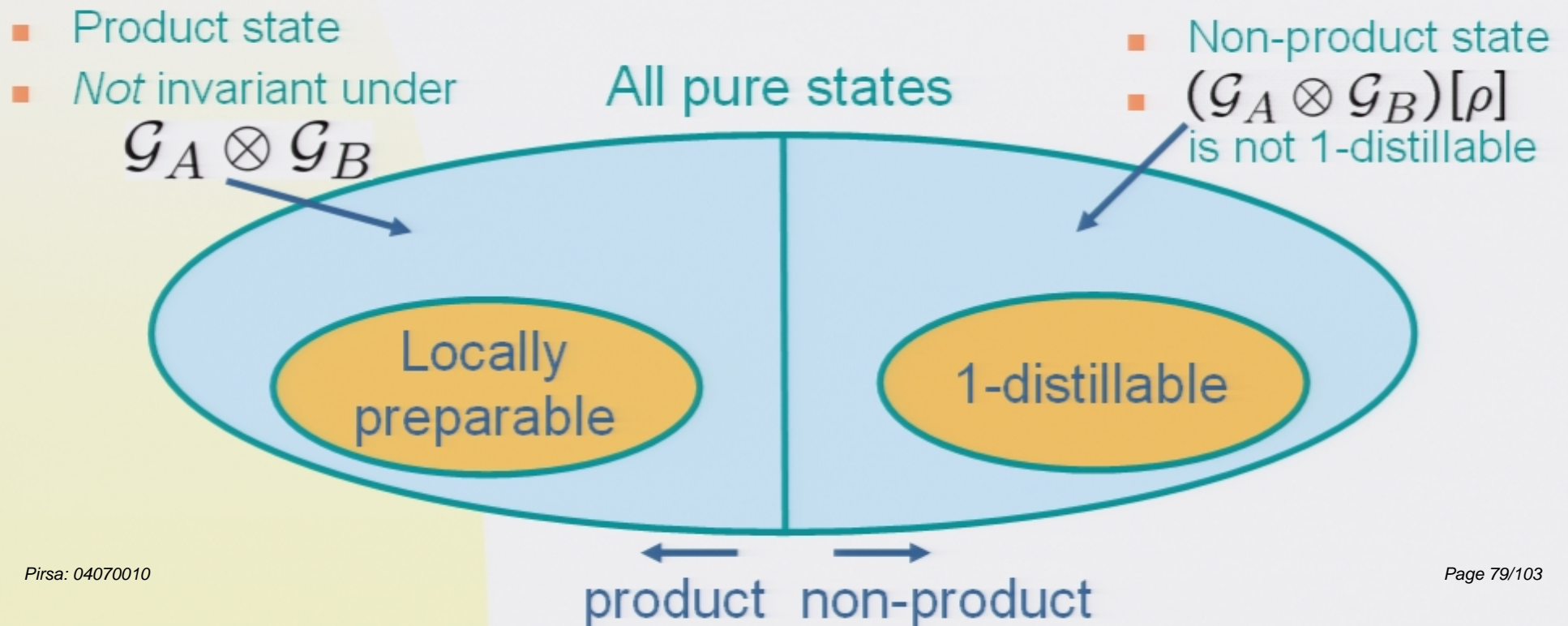
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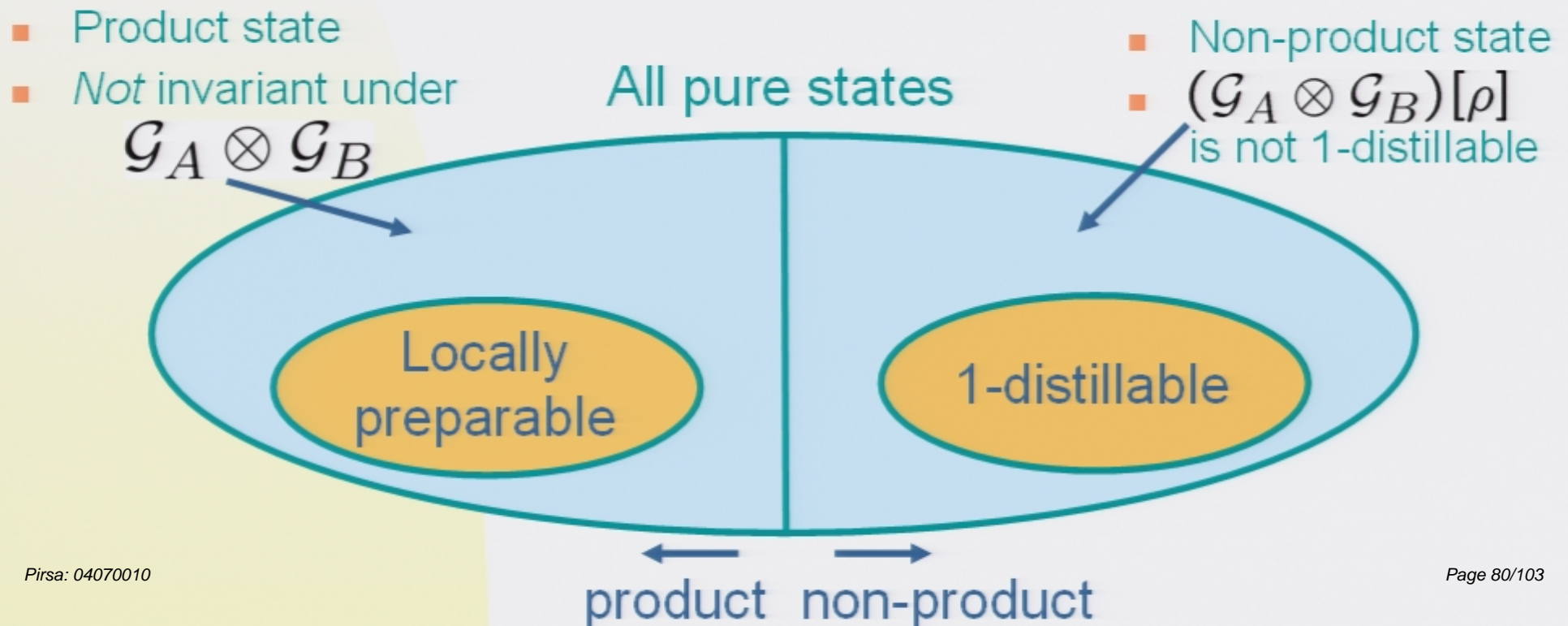
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A helpful channel – shared reference frame



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- Channel C: a shared reference frame, obviating the SSR
- Can distribute any product state (even non-invariant ones)

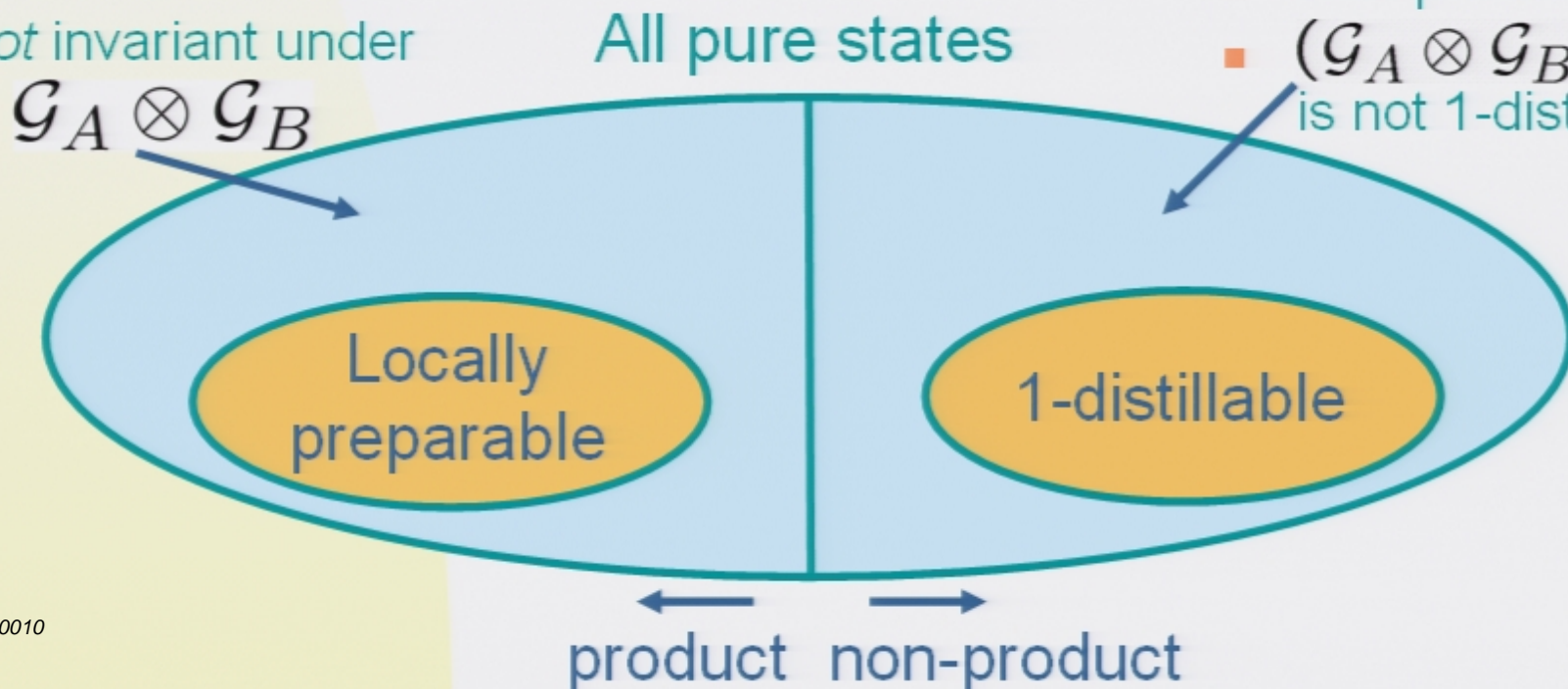


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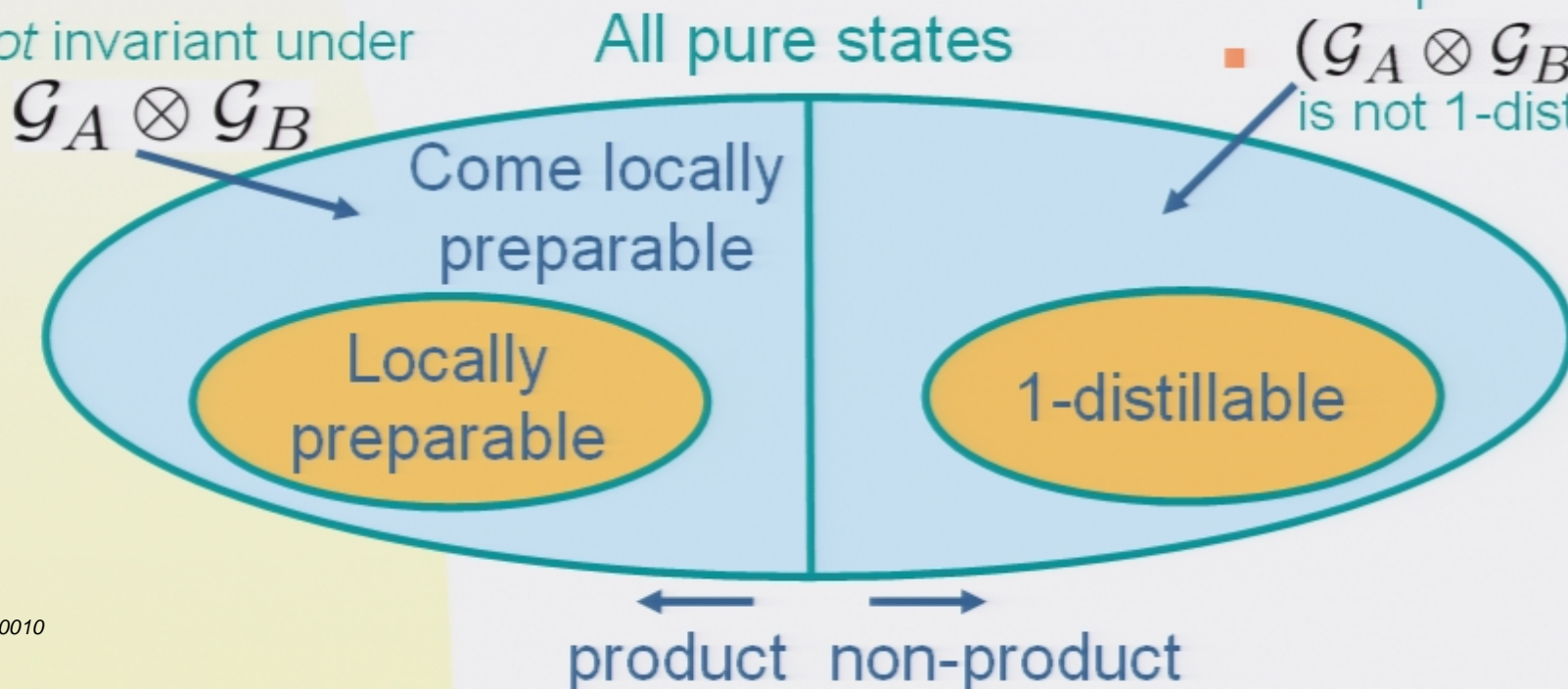


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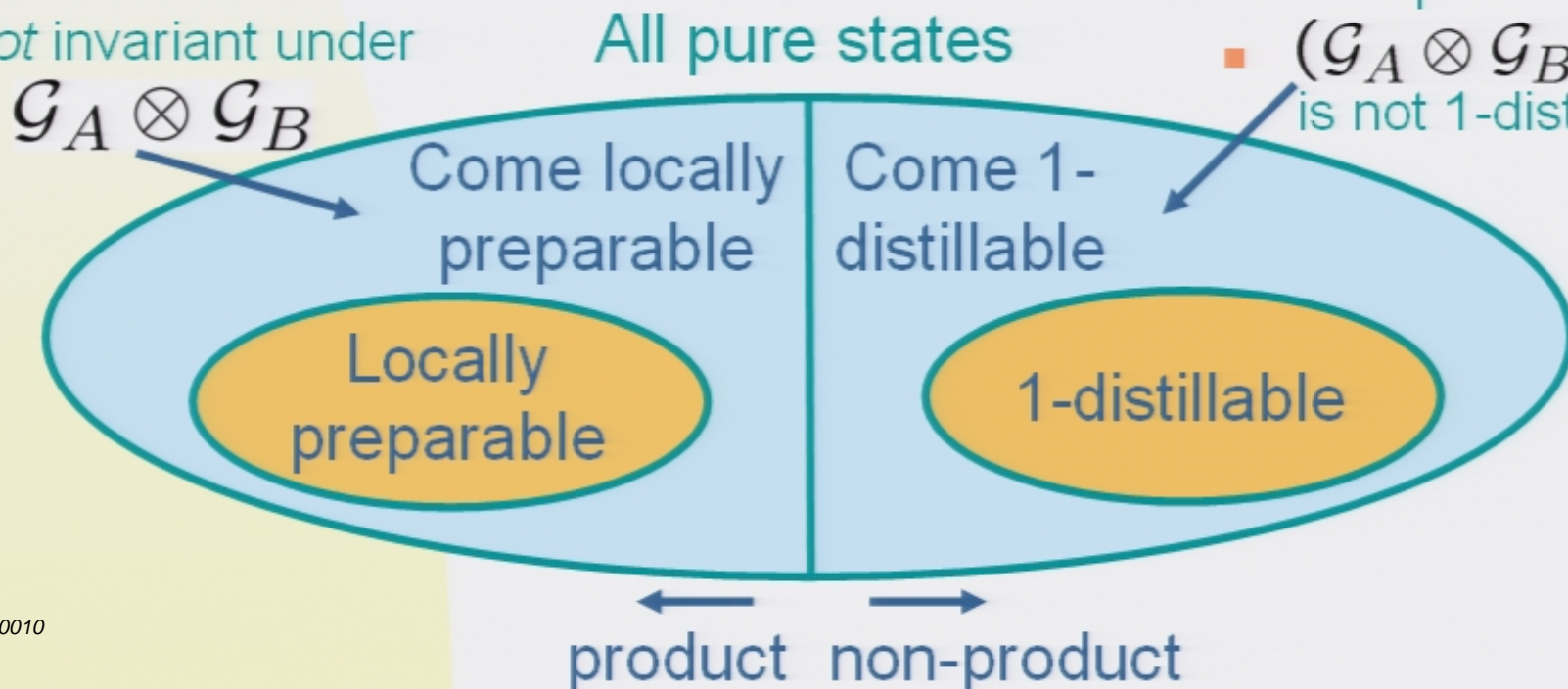


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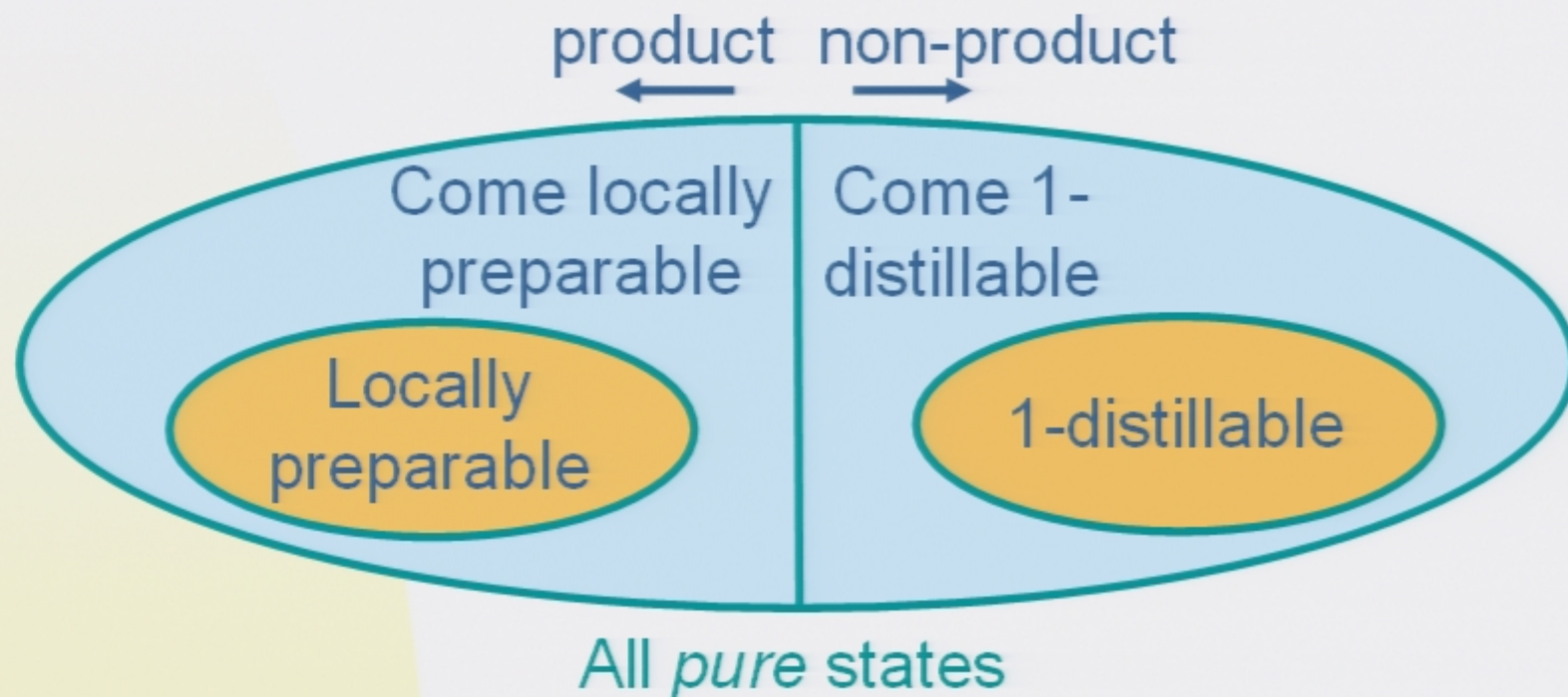
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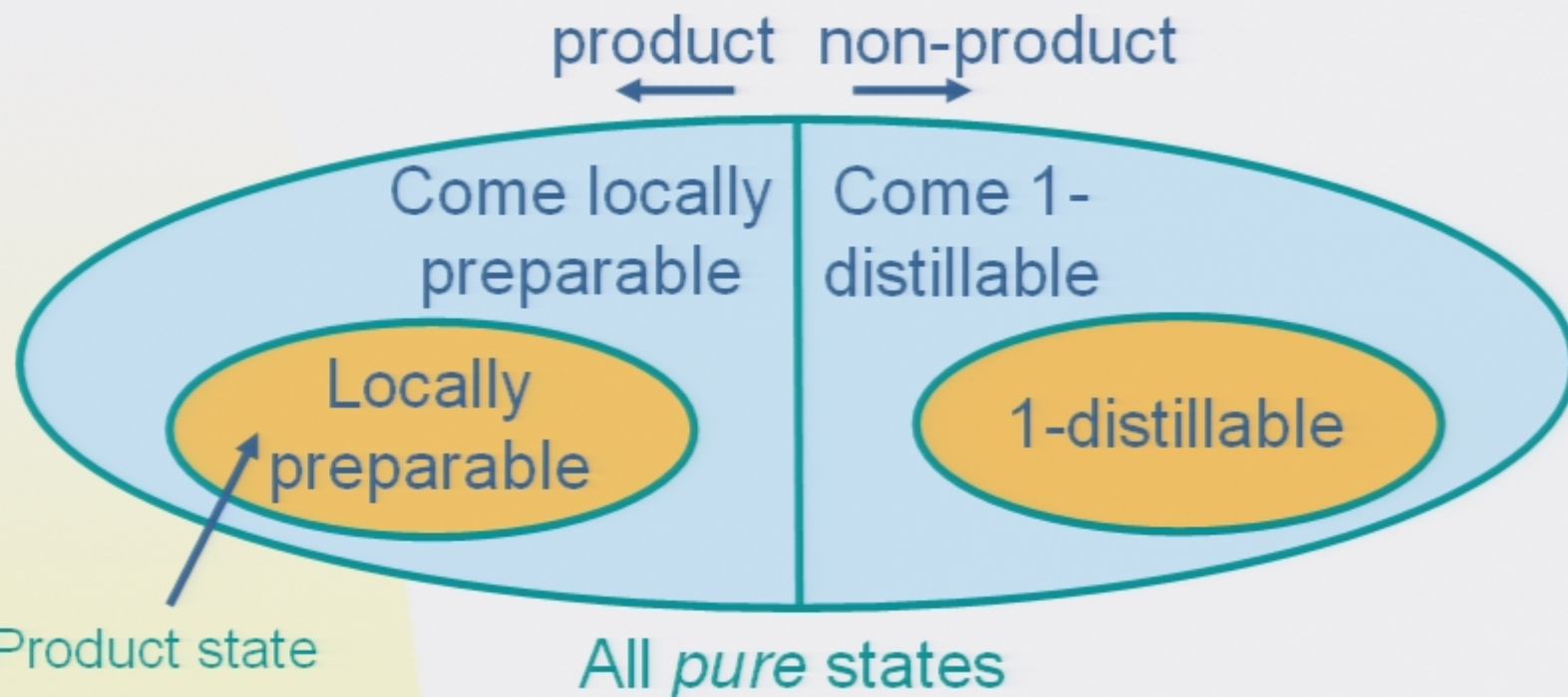
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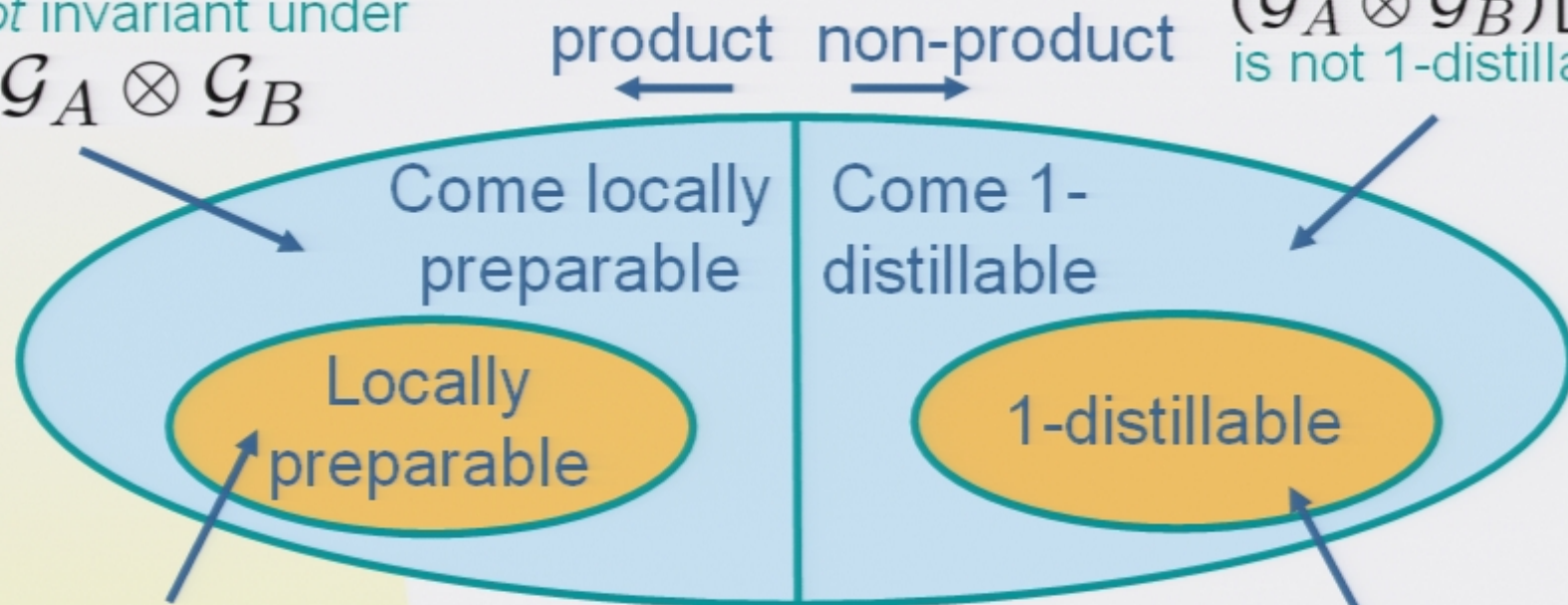
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Abelian SSR example

- U(1) SSR for particle number, Fock states $|n\rangle_{A,B}$

$$(|0\rangle + |1\rangle)_A \otimes (|0\rangle + |1\rangle)_B$$

$$|1\rangle_A|0\rangle_B + |0\rangle_A|1\rangle_B$$

product non-product
 \longleftrightarrow \longrightarrow

Come locally
preparable

Come 1-
distillable

Locally
preparable

1-distillable

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Activation

- Channel C : implemented probabilistically using non-invariant product states (Jamiołkowski isomorphism)
- For every "come 1-distillable" state, there exists a "come locally preparable" state that serves as a finite shared reference frame and *activates* the 1-distillability

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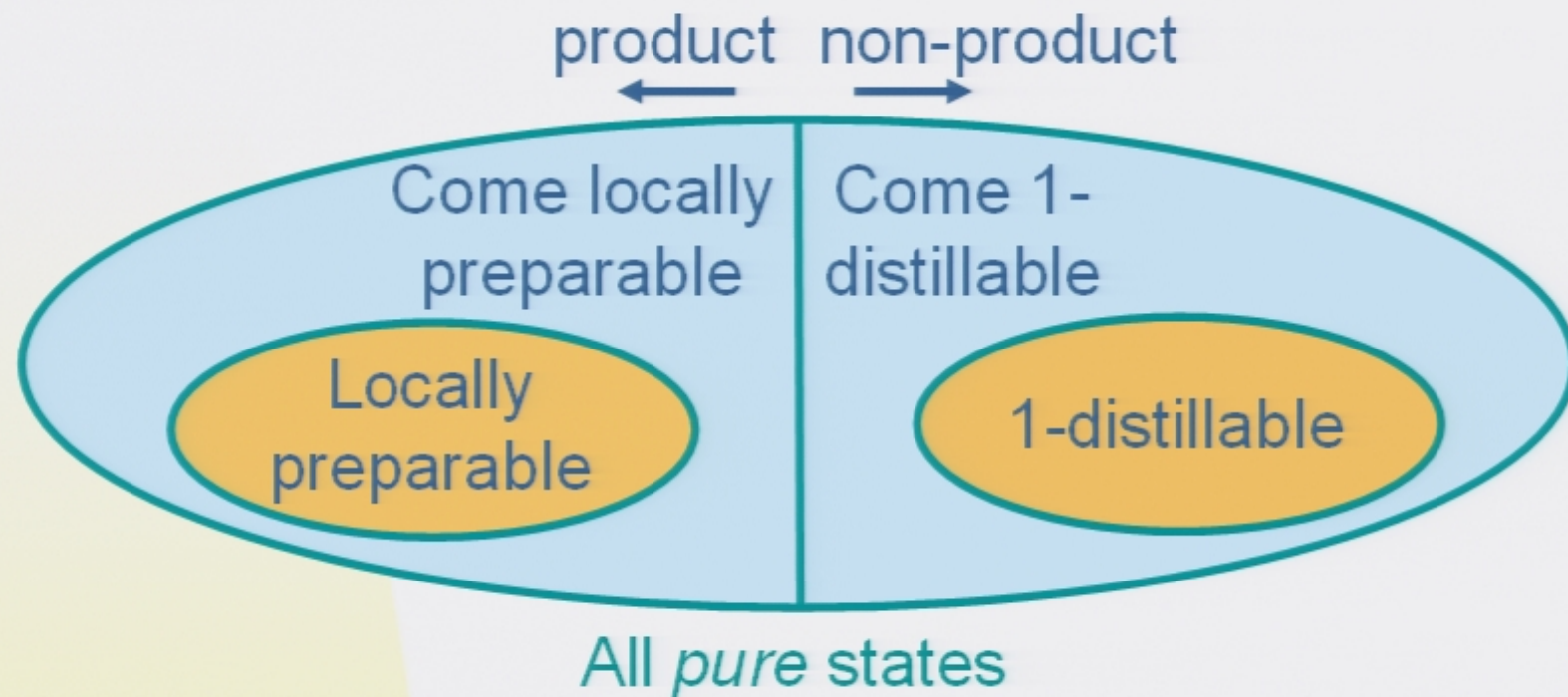
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- Alice and Bob each measure total local particle number
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Resulting state: $|0, 1\rangle_A |0, 1\rangle_B + |1, 0\rangle_A |1, 0\rangle_B$

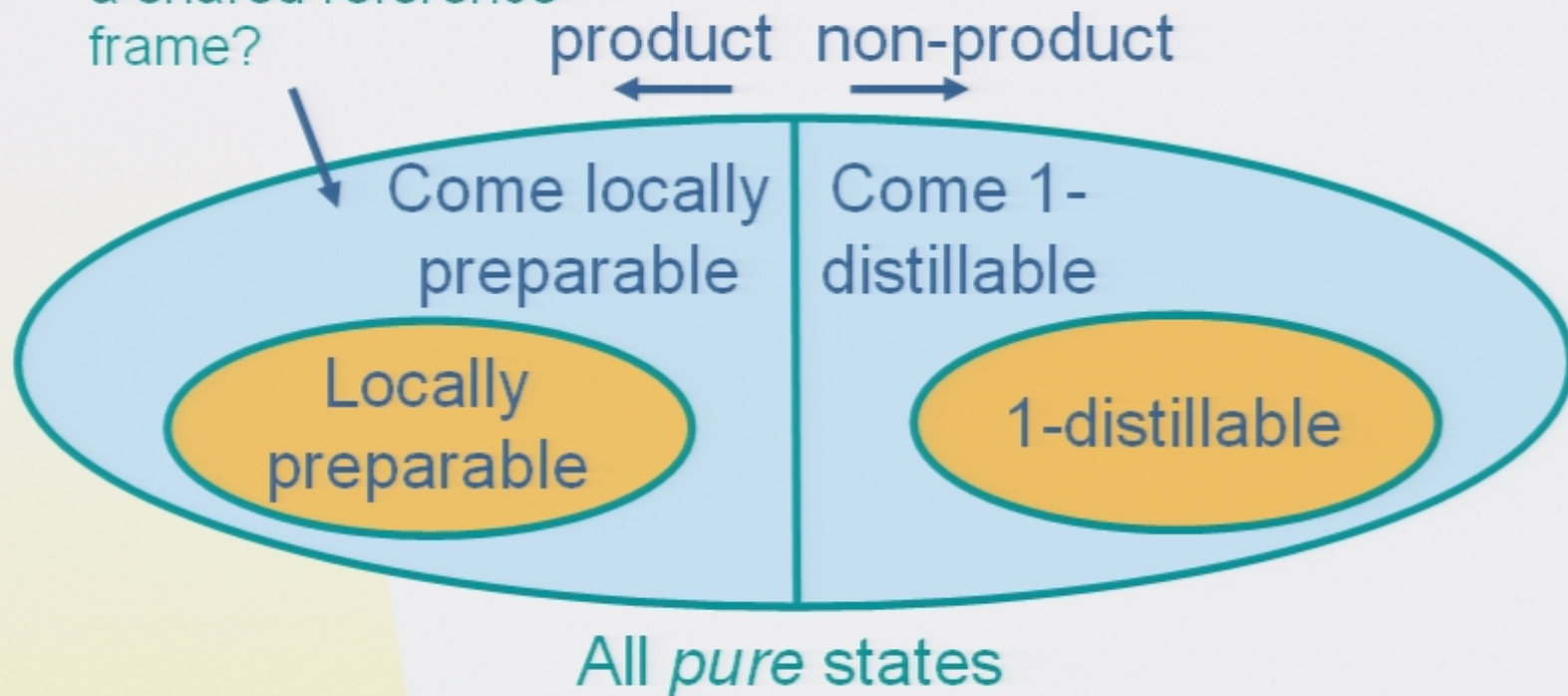
1-distillable!

What states are good reference frames?



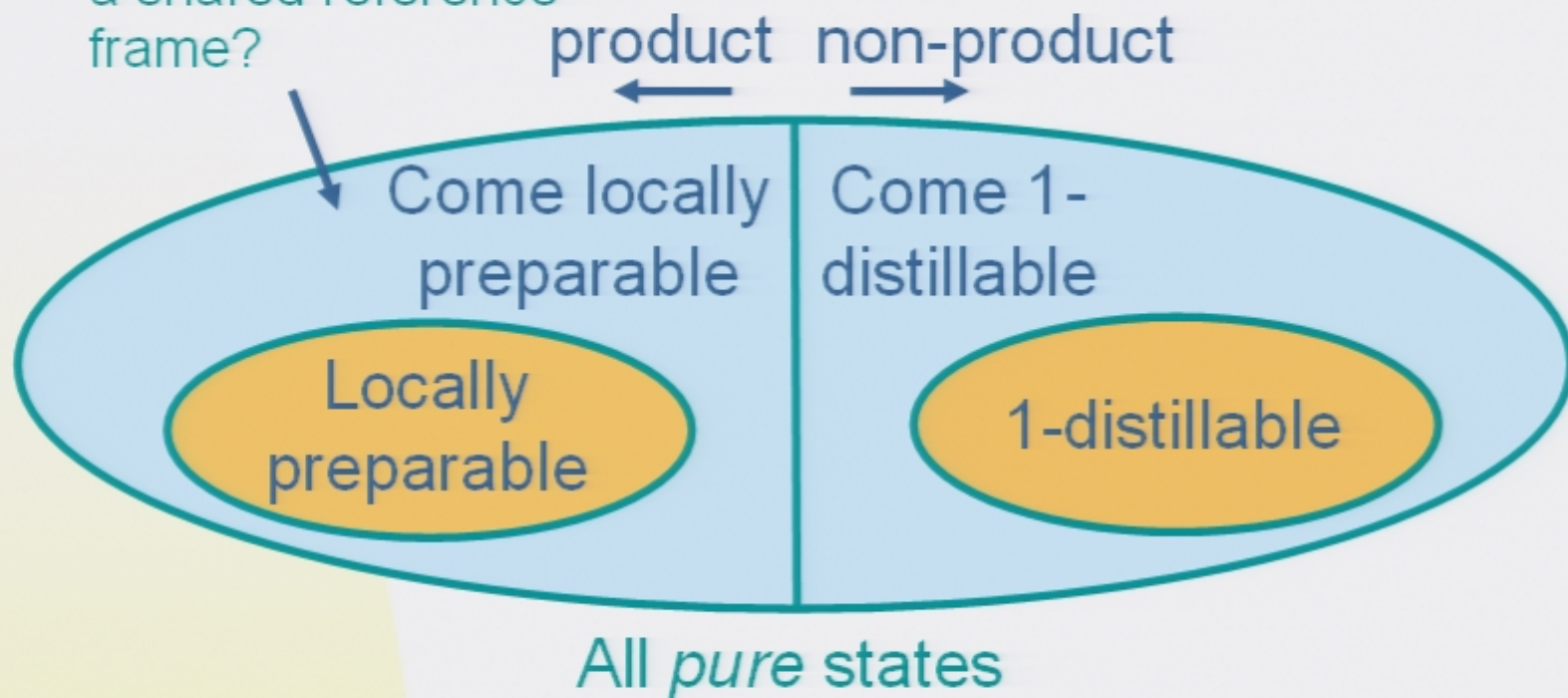
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- Is it only the "come locally preparable" states that can serve as a shared reference frame?



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No! Some C1-D states can activate entanglement in other states, and themselves...

Activation / 2-distillability

- C1-D states can possess some "reference frameness"
- Entanglement in one copy may be "1-bound", but two copies may be distilled: *2-distillable*

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 $= |1, 1\rangle_A|0, 0\rangle_B + |1, 0\rangle_A|0, 1\rangle_B$
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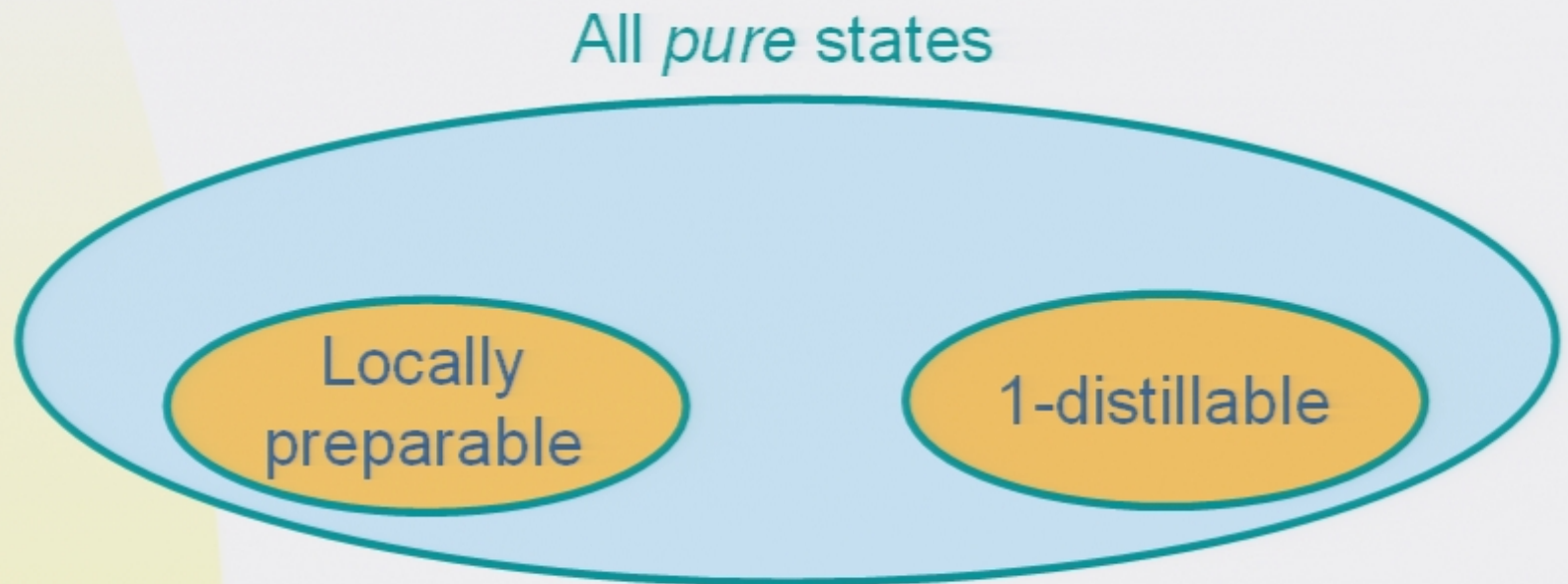
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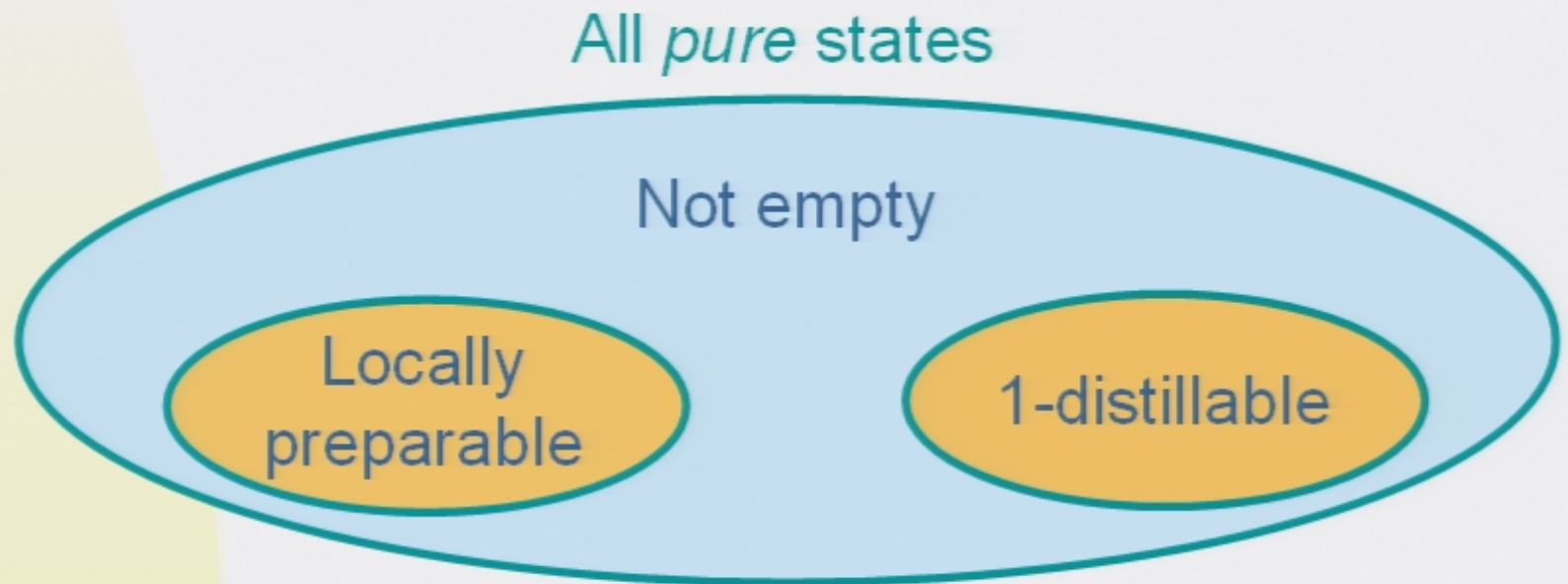
Summary

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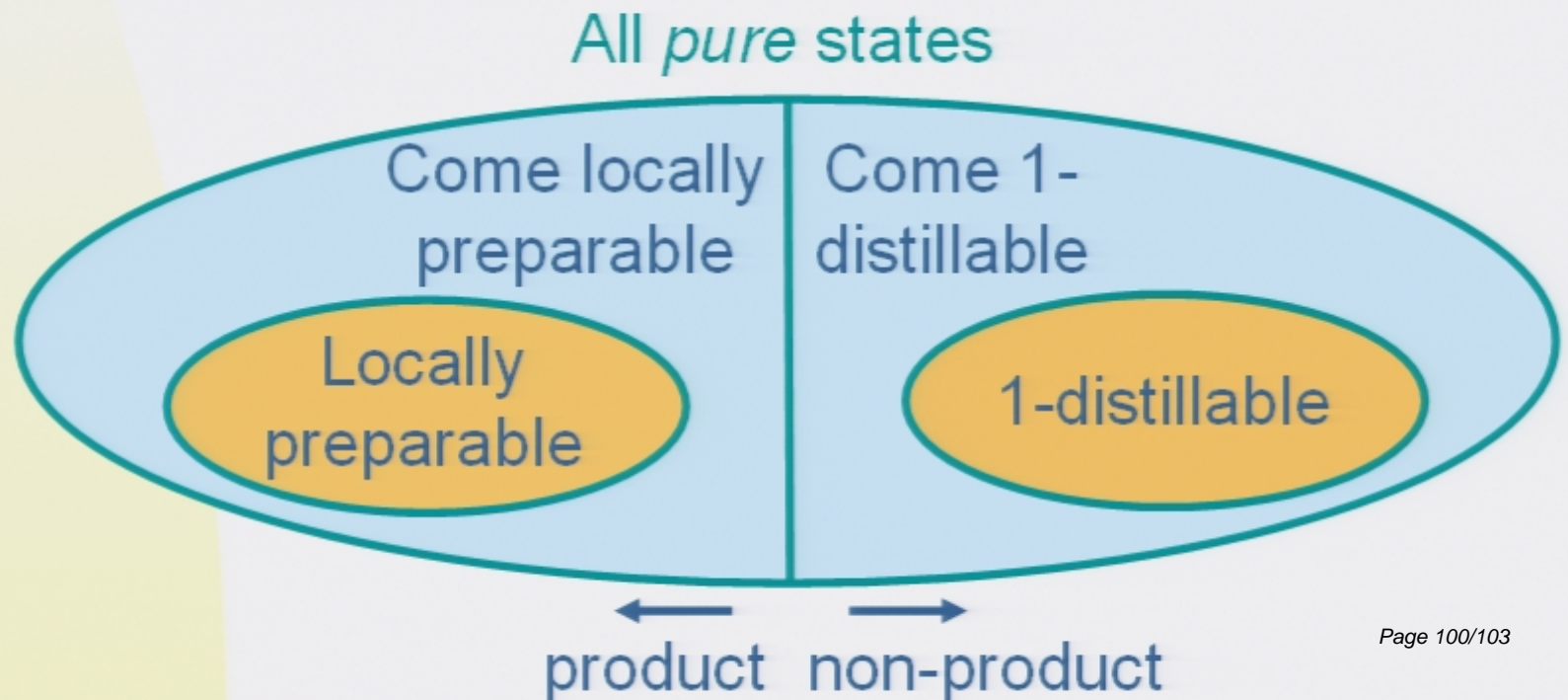
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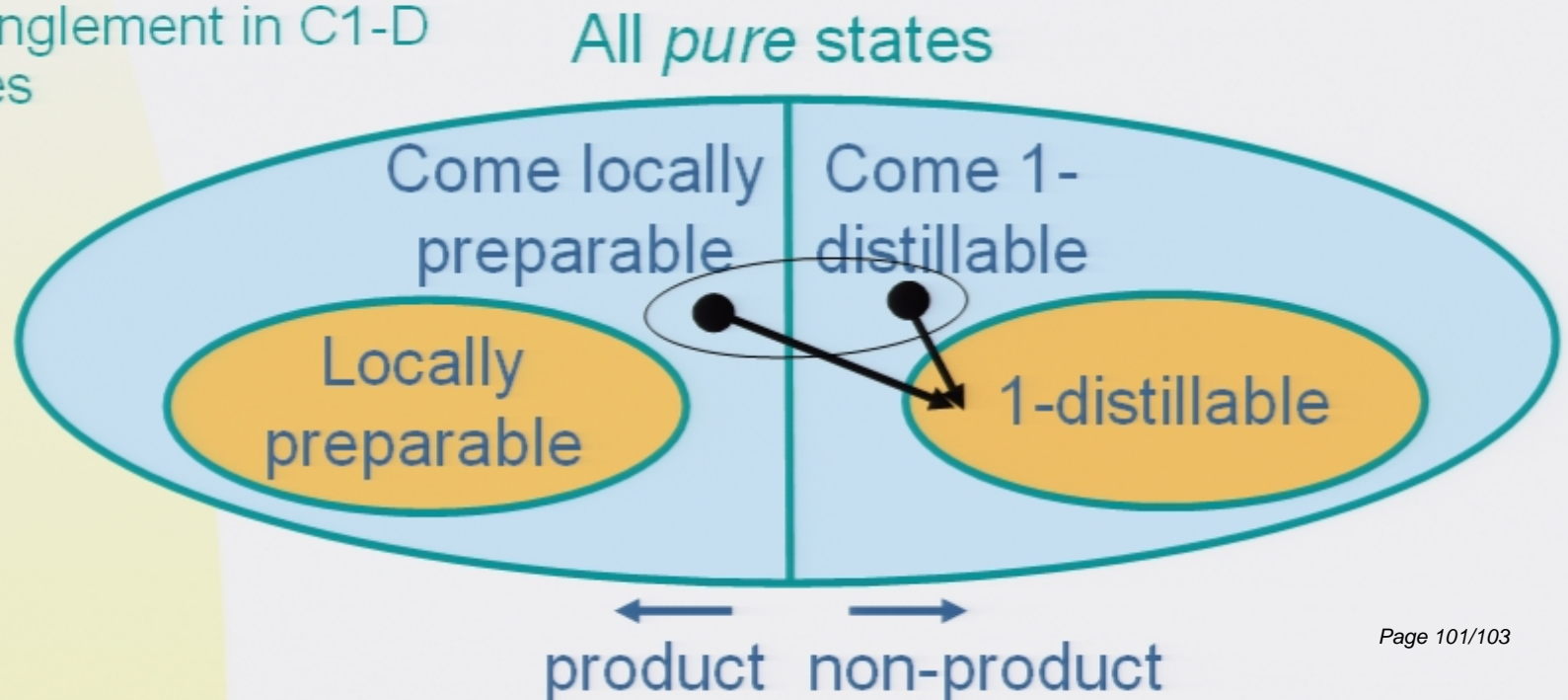
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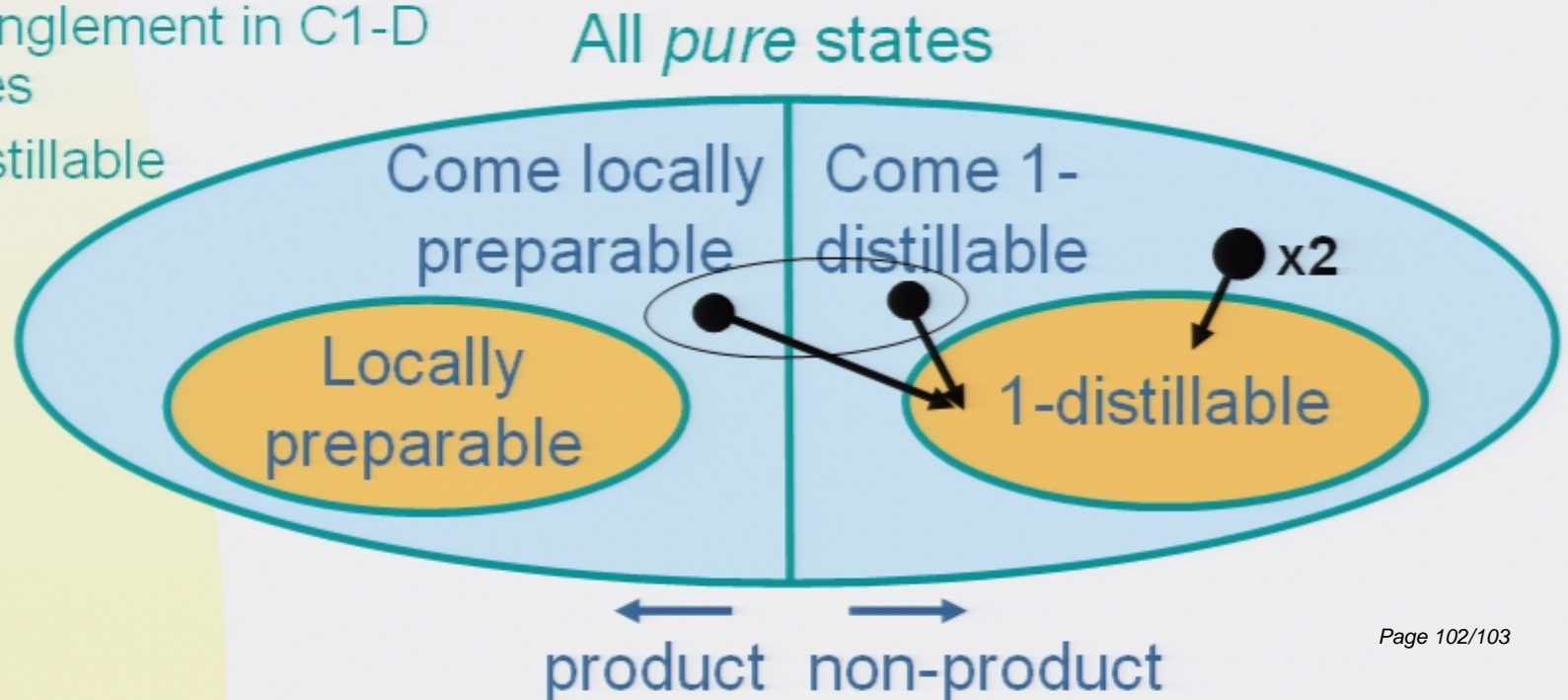
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Useful to think of bound entanglement in terms of restrictions and reference systems

