

Title: Efficient Alignment of Reference Frames and without Shared Entanglement

Date: Jul 13, 2004 10:45 AM

URL: <http://pirsa.org/04070005>

Abstract: Quantum Information Theory

Efficient alignment of reference frames with and without entanglement



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In collaboration with: M. Baig & R. Muñoz-Tapia

Ariadne, Theseus, the Minotaur

[Theseus] got from Ariadne, who had fallen in love with him, the famous thread, and that having been instructed by her how to make his way through the intricacies of the Labyrinth, he slew the Minotaur and sailed off with Ariadne and the youths. " —Plutarch, *Theseus* 19.1



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Outline

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 - A related problem: estimation of one-qubit gates

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- Saving resources

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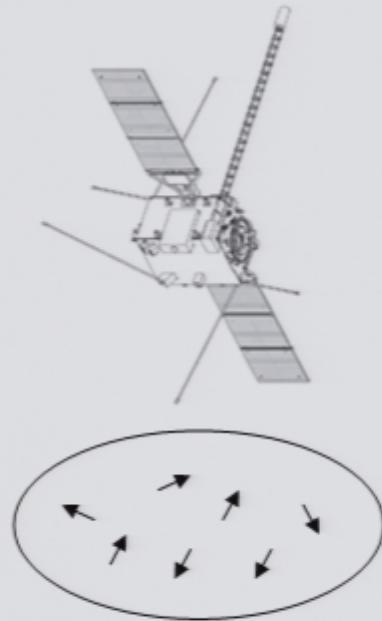
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- Saving resources
- Summary and conclusions

The problem

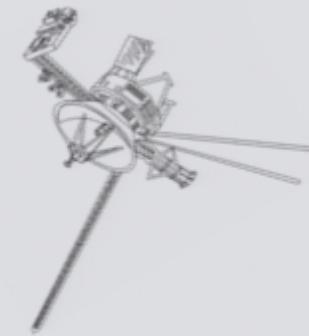
Point of view 1

A communicates a spatial frame to B by sending N spin $\frac{1}{2}$ particles. B is allowed to perform **generalized** measurements.
•No classical communication allowed.

A



B

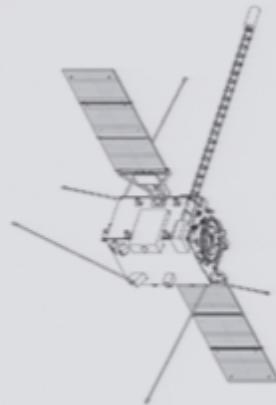


The problem

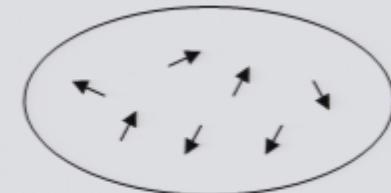
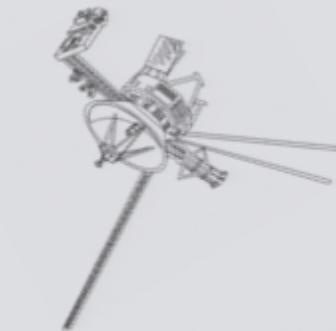
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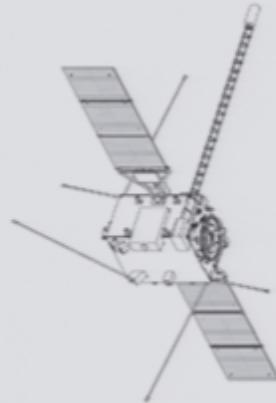


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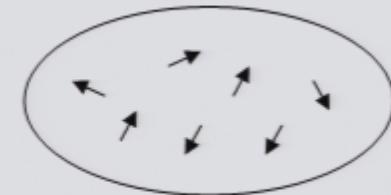
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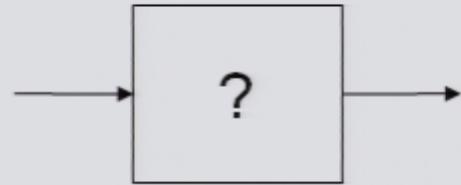


B



The problem

Point of view 2 (*Reverse engineering*): estimation of unknown 1-qubit gates

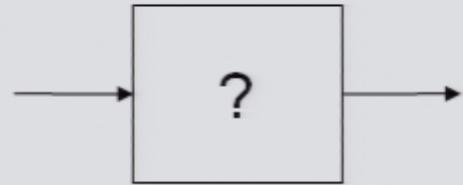


We are allowed to **make N calls** to the unknown gate (but input state made out of as many qubits as we need). We may use **generalized** measurements.

The problem

Point of view 2

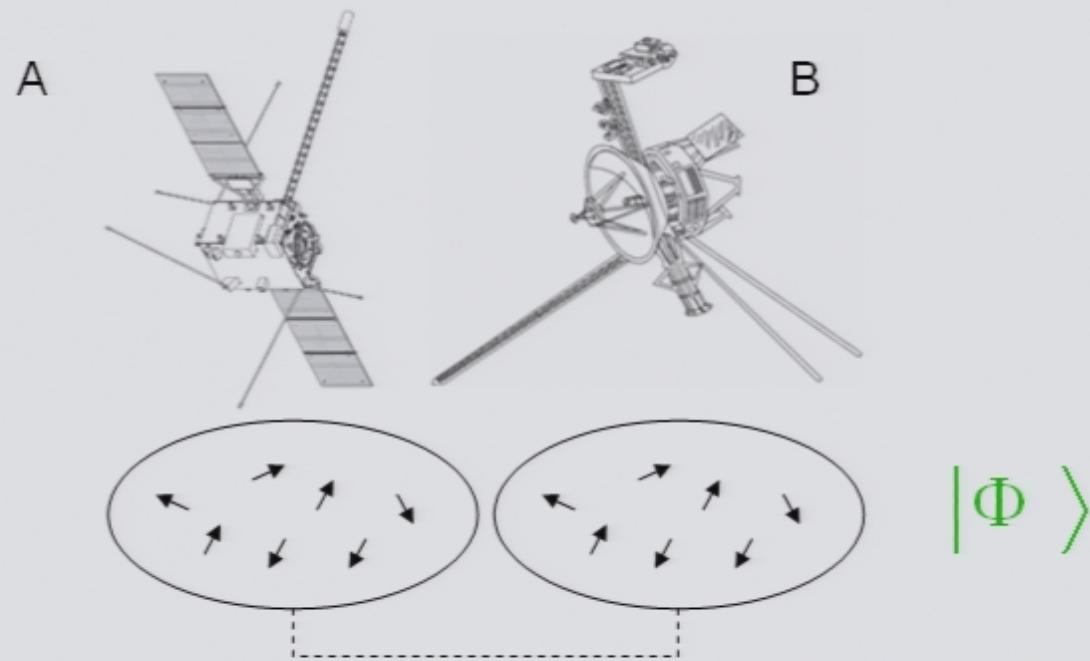
(Reverse engineering): estimation of unknown 1-qubit gates



We are allowed to make N calls to the unknown gate (but input state made out of as many qubits as we need). We may use generalized measurements.

The optimal scheme (dense covariant coding)

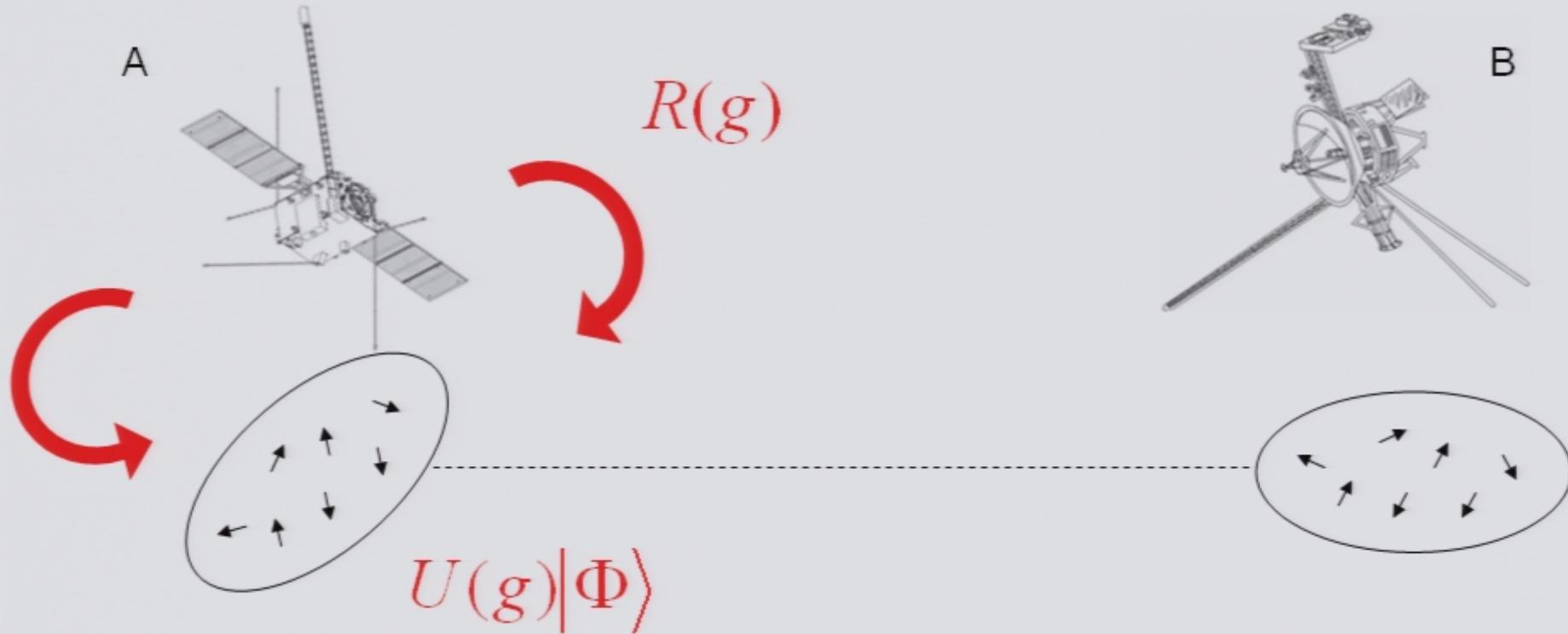
Point of view 1: Preparation



Additional resources: *entanglement + N extra spins*

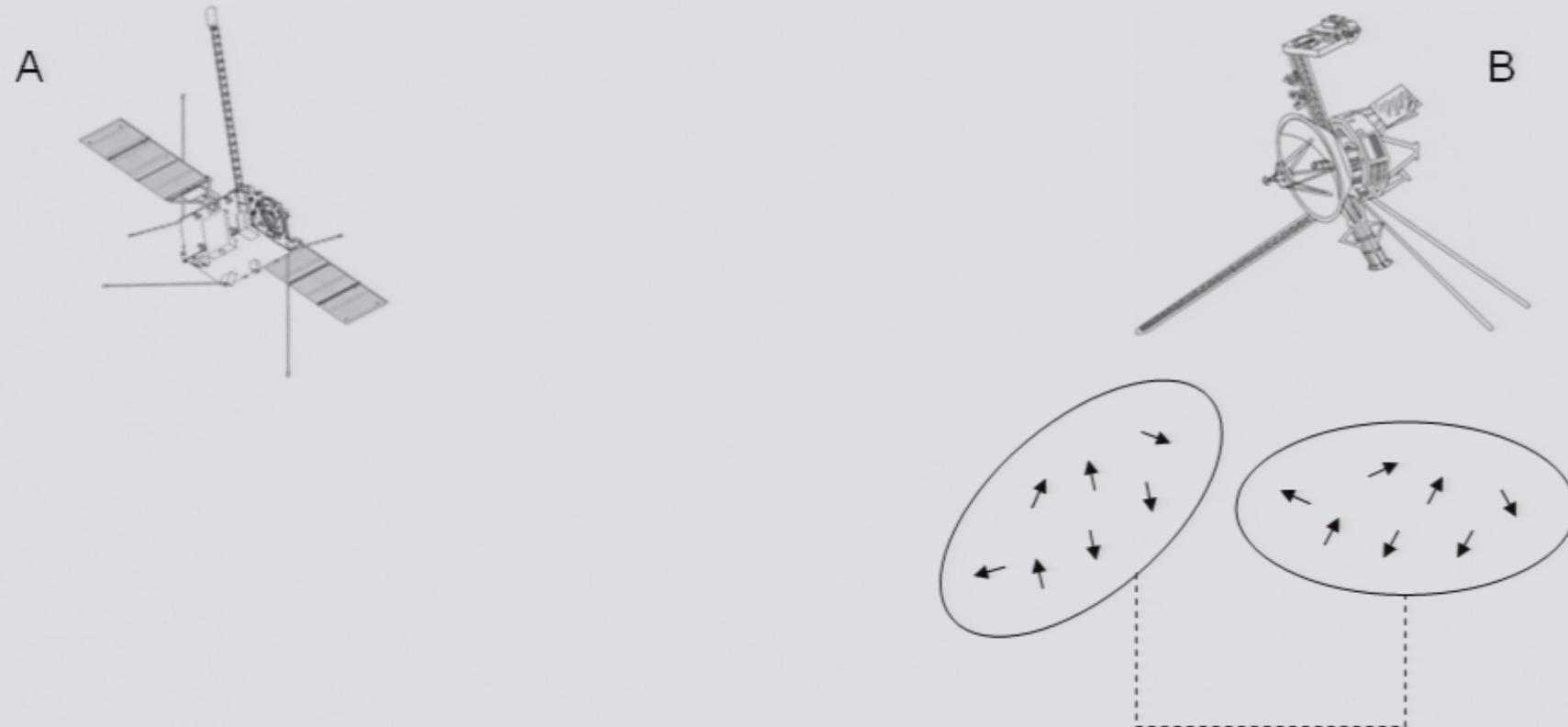
The optimal scheme (dense covariant coding)

Point of view 1: The journey



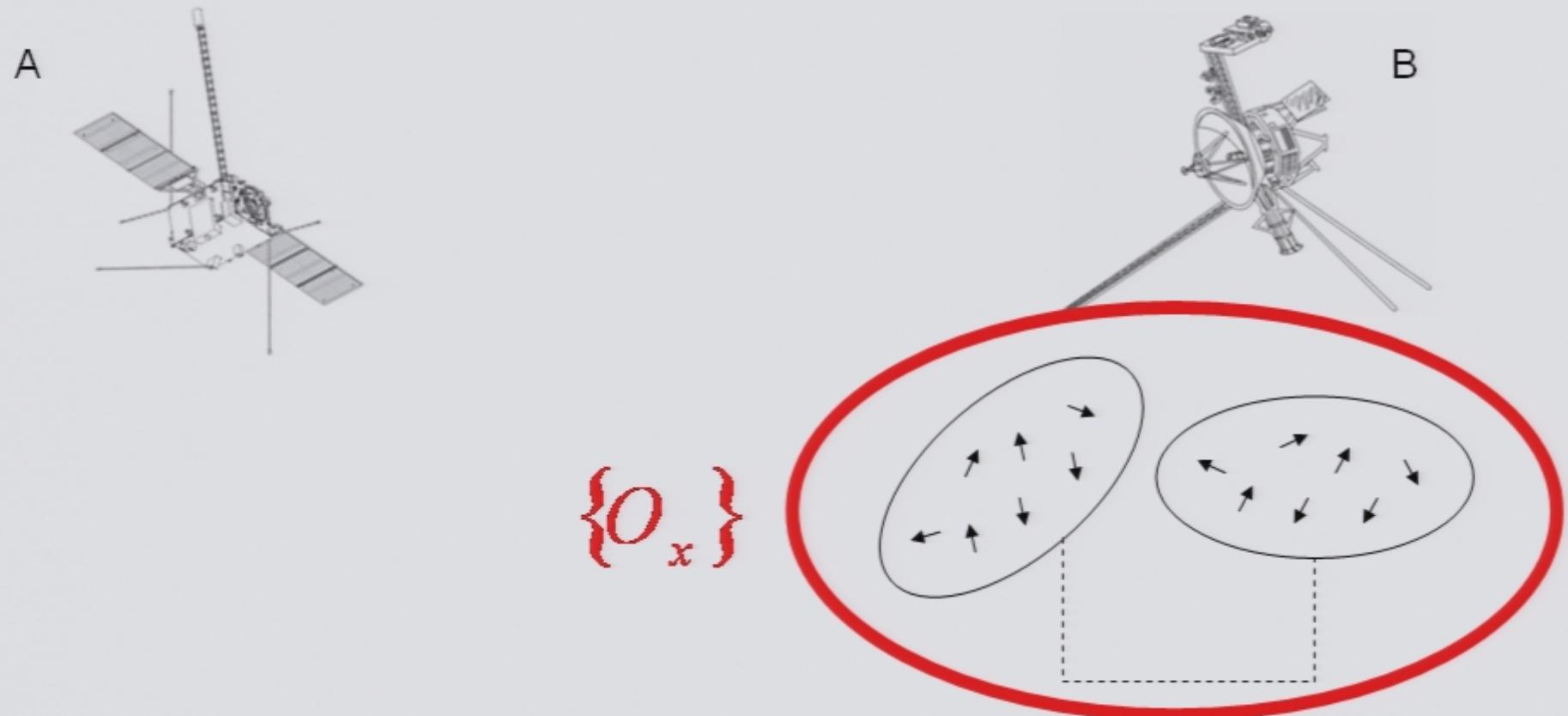
The optimal scheme *(dense covariant coding)*

Point of view 1: communication, measurement and alignment



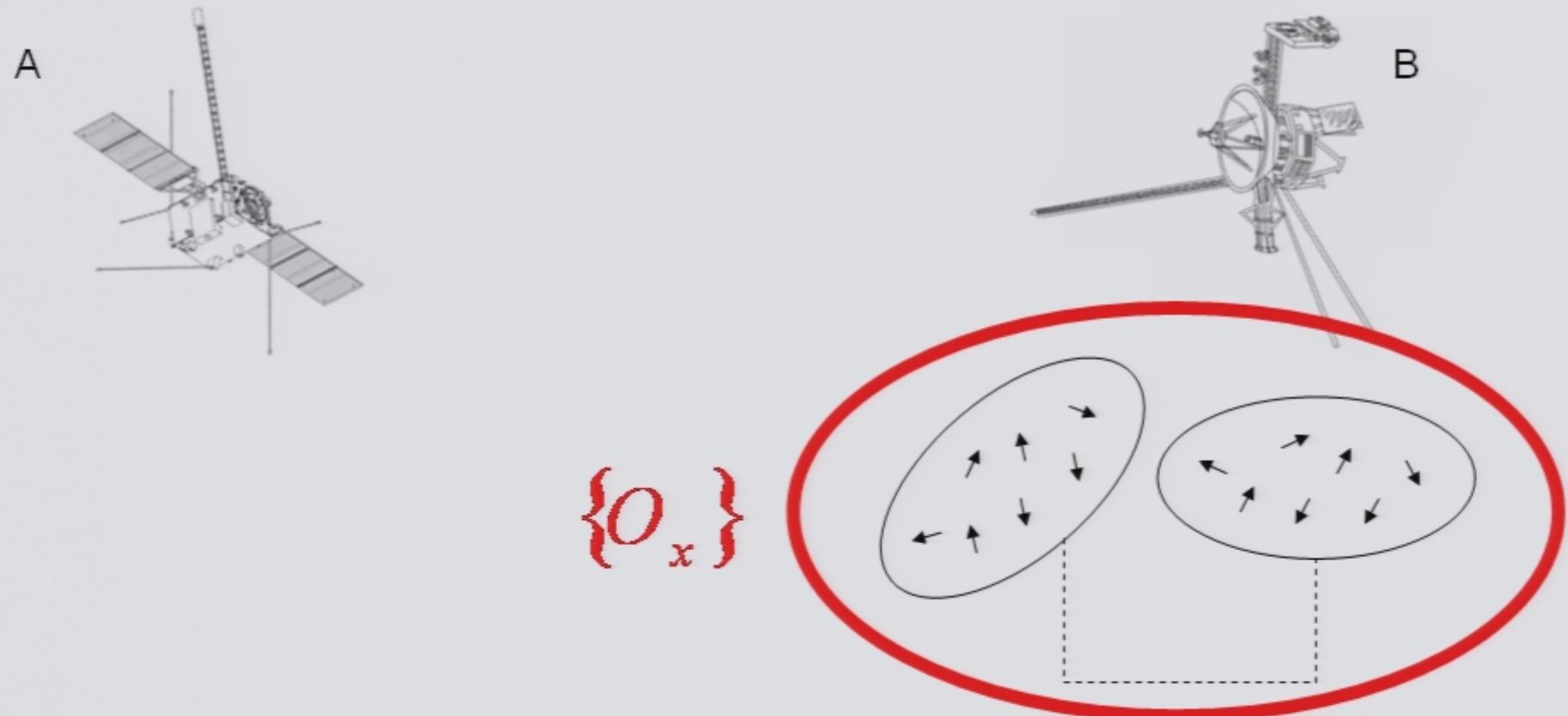
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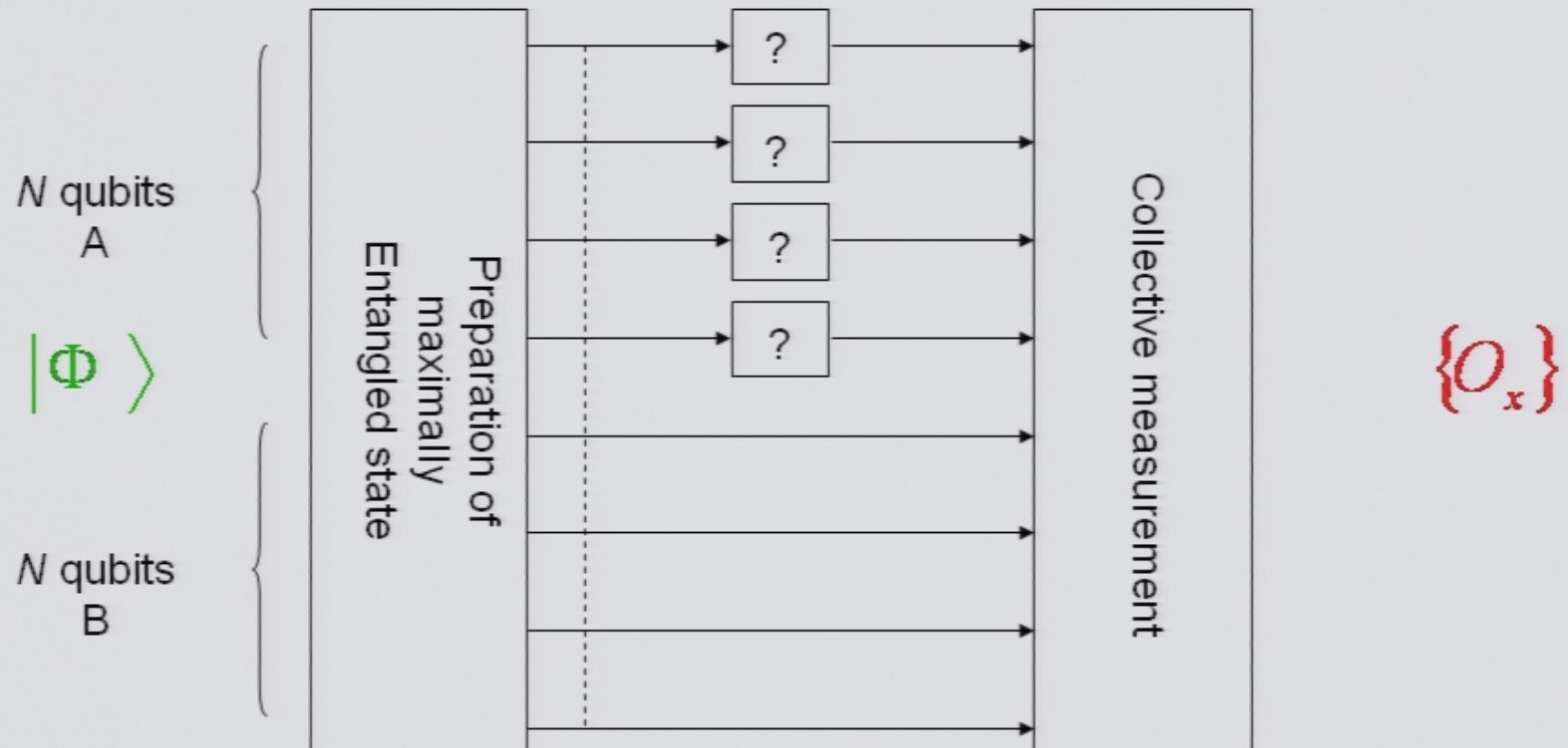
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Point of view 1: communication, measurement and alignment



Optimal scheme

Point of view 2:



Optimal scheme

Point of view 2:

N qubits
A

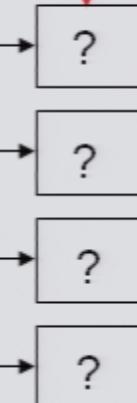
$|\Phi\rangle$

N qubits
B

{ }
N qubits A

Preparation of
maximally
Entangled state

Journey



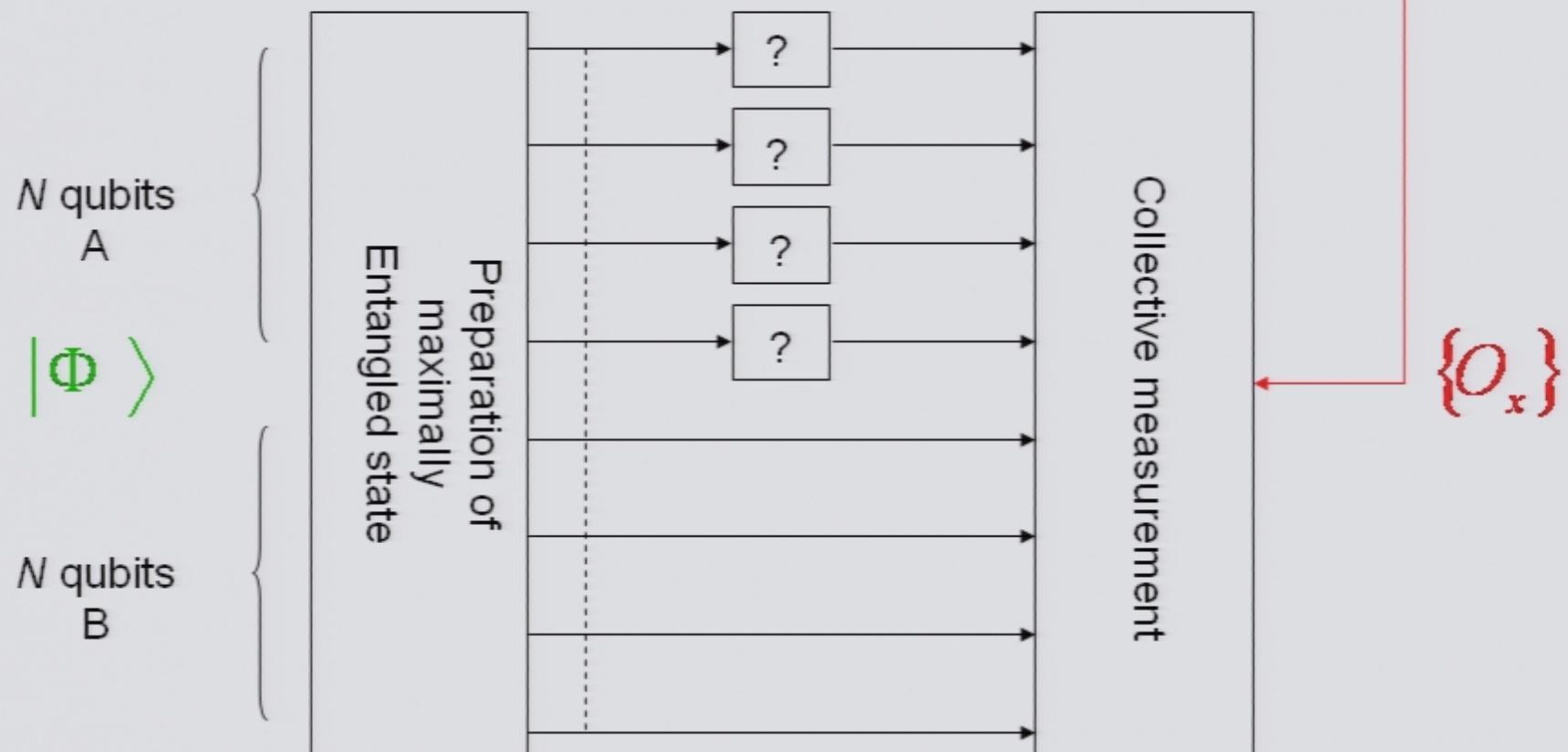
Collective measurement

$\{O_x\}$

Optimal scheme

Communication,
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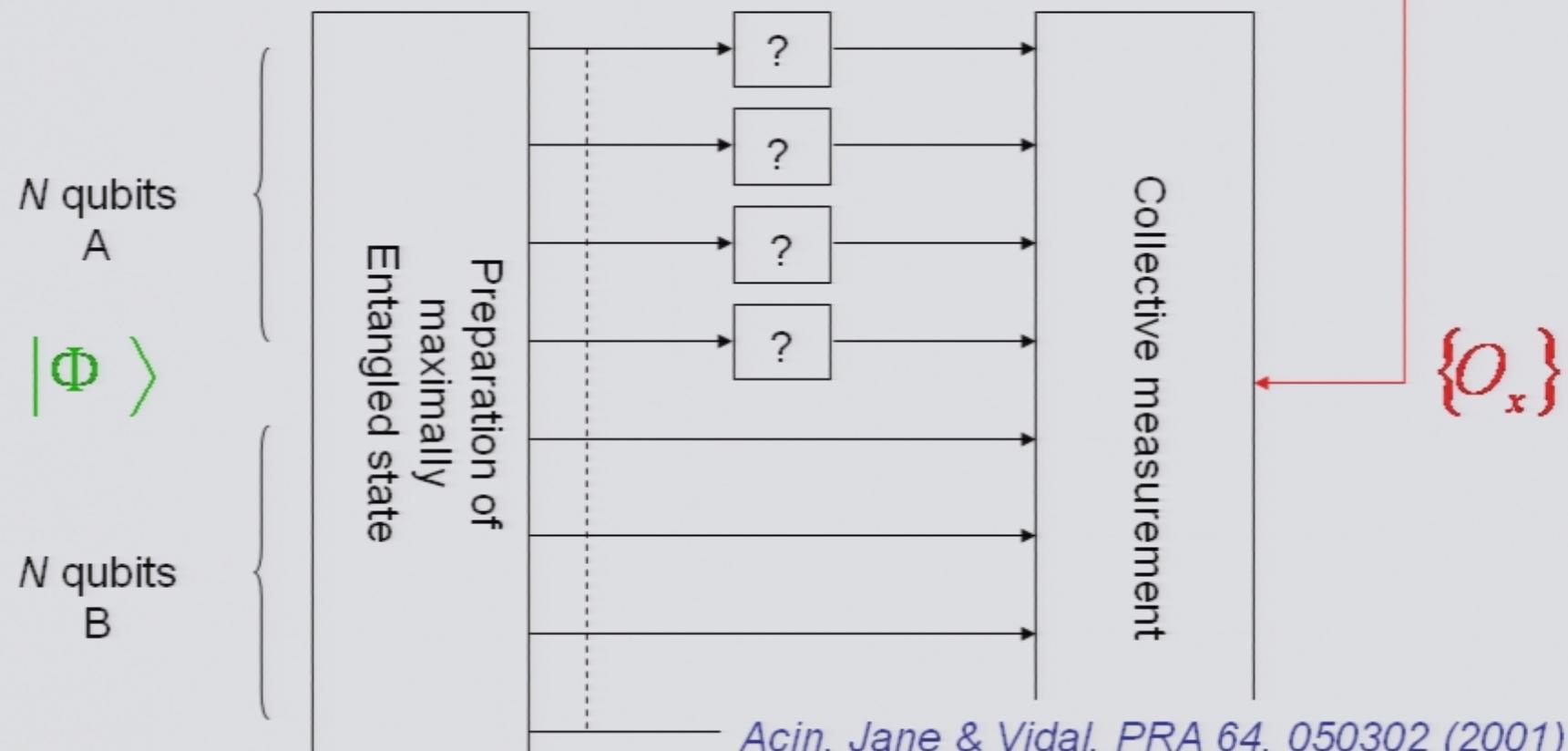
Point of view 2:



Optimal scheme

Communication,
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Point of view 2:



Acin, Jane & Vidal, PRA 64, 050302 (2001)



Figure of merit

Point of view 1: Holevo error

$$h(g, g_x) = \sum_{a=1}^3 |\vec{n}_a(g) - \vec{n}_a(g_x)|^2$$

x = outcomes of measurement $\{O_x\}$

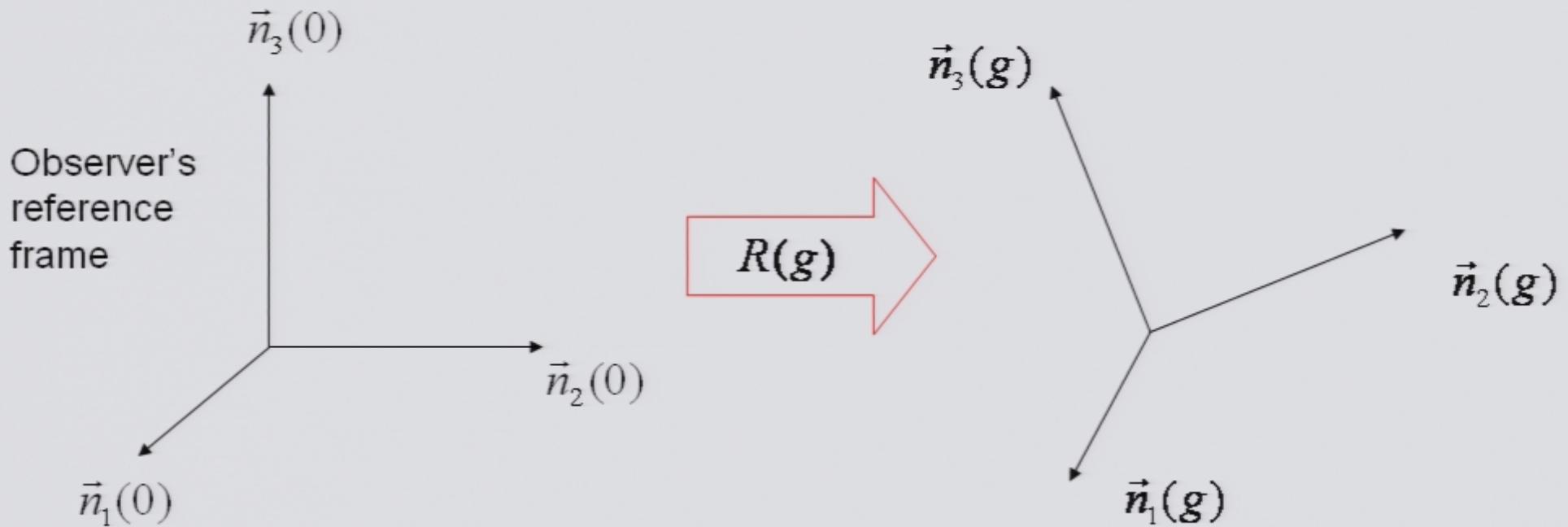




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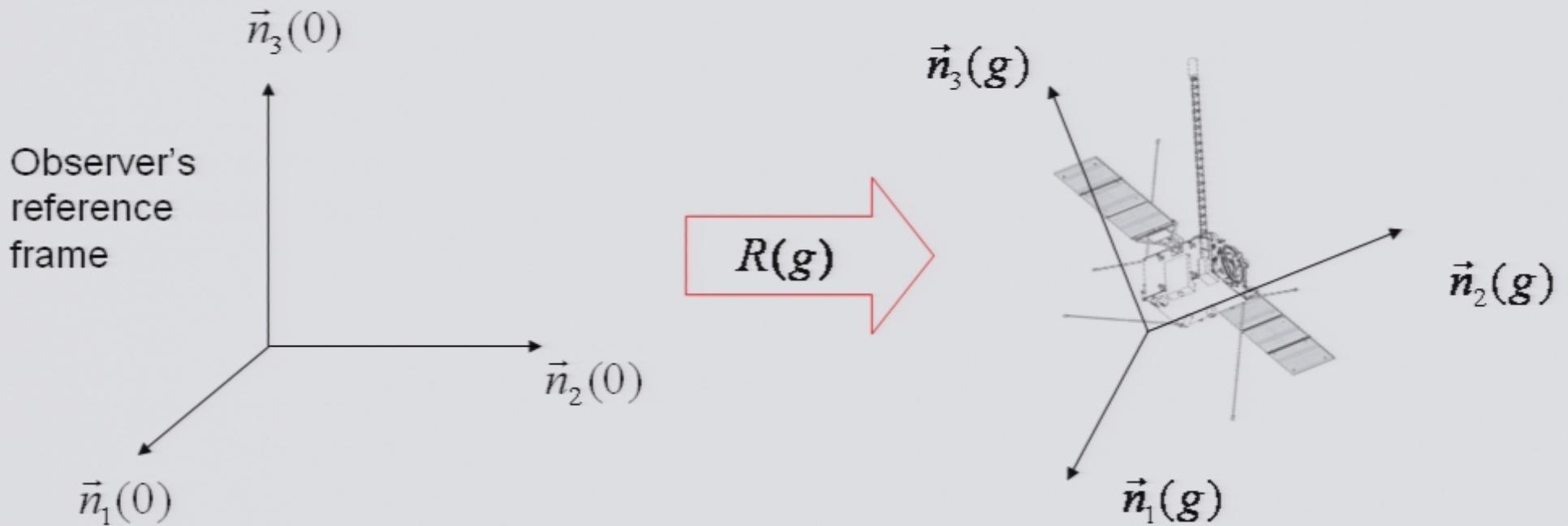




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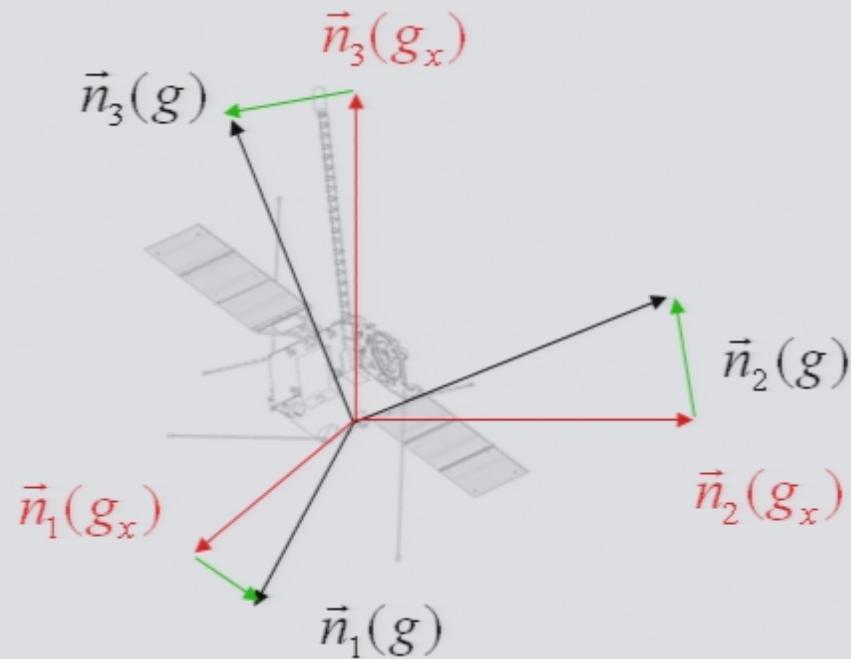




Figure of merit

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Averaged Holevo error

$$\langle h \rangle = \sum_x \int dg h(g, g_x) p(x | g)$$

x Outcomes of measurement; g_x Guess based on x

dg SU(2) Haar measure $d(gg') = d(g'g) = dg$

$p(x | g) = \text{tr}[O_x \rho(g)]; \quad \rho(g) = U(g) |\Phi\rangle\langle\Phi| U^+(g)$

Figure of merit

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Aim: minimize $\langle h \rangle$ over all POVMs
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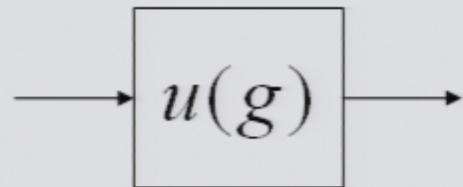
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Figure of merit

Point of view 2: *Fidelity*

$$F(g, g_x) = \left| \text{tr} [u^+(g_x) u(g)] \right|^2 / 4$$



x Outcomes of measurement; g_x Guess based on x

It is a measure of how well $|u(g)|\phi\rangle$ compares to $|u(g_x)|\phi\rangle$
in average (over all possible states $|\phi\rangle$)



Figure of merit

Point of view 2: *Fidelity*

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Figure of merit

- Point of view 1

$$h(g, g_x) = 6 - 2\chi_1(g_x^{-1}g)$$

- Point of view 2

$$F(g, g_x) = \frac{1 + \chi_1(g_x^{-1}g)}{4}$$



Figure of merit

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$$h(g, g_x) = 6 - 2\chi_1(g_x^{-1}g)$$

$$8 \geq h(g, g_x) \geq 0$$

- Point of view 2

$$F(g, g_x) = \frac{1 + \chi_1(g_x^{-1}g)}{4}$$

$$-1 \leq \chi_1(g) \leq 3$$

$$0 \leq F(g, g_x) \leq 1$$

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Aim: maximize
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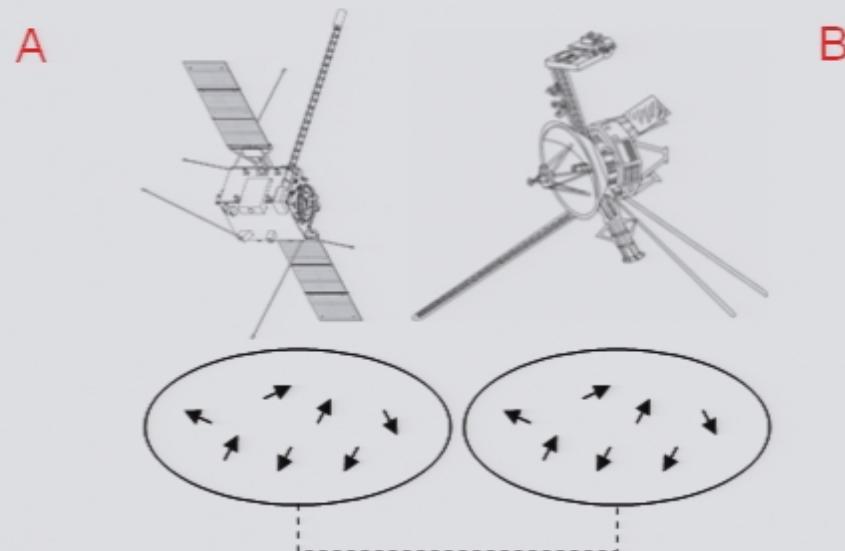
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Preparation

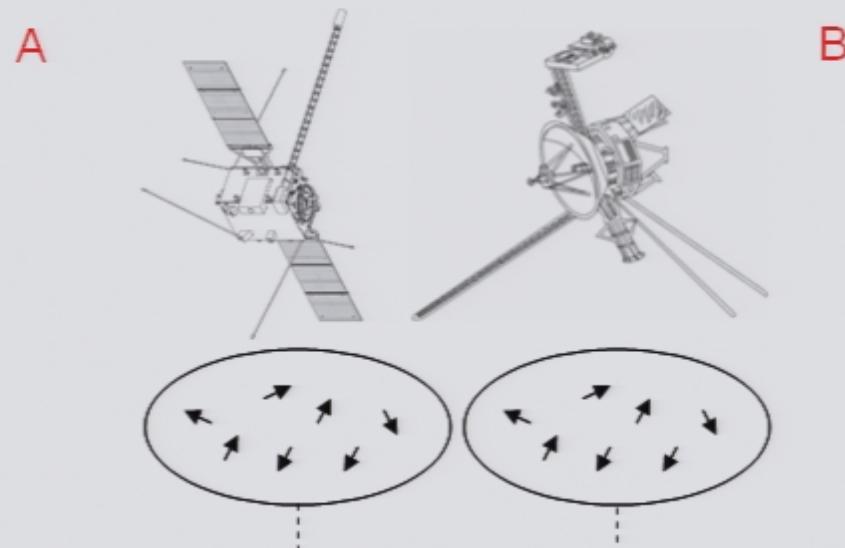
$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle |jm\rangle$$



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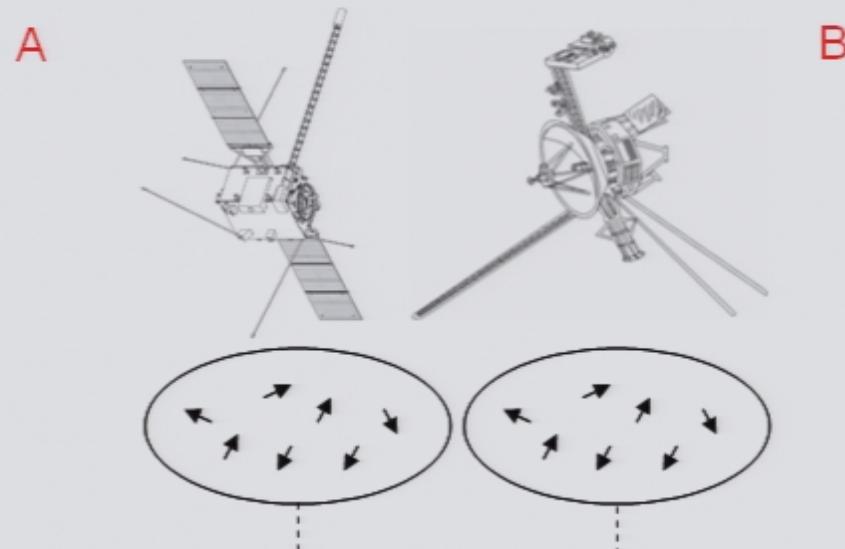
Referred to B reference frame



Preparation

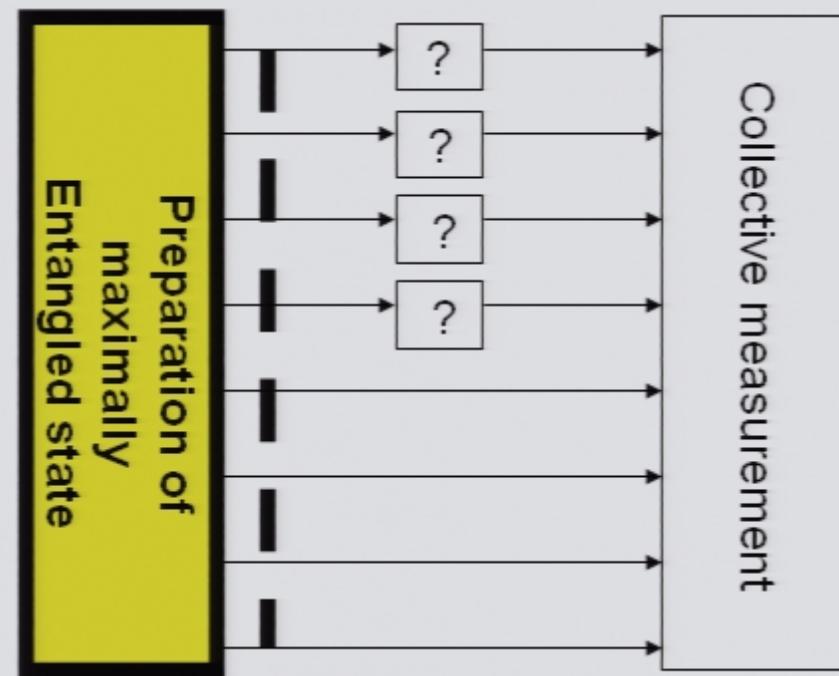
$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle_A |jm\rangle_B \quad \text{Referred to B reference frame}$$

Maximally entangled state in each irreducible representation of SU(2)



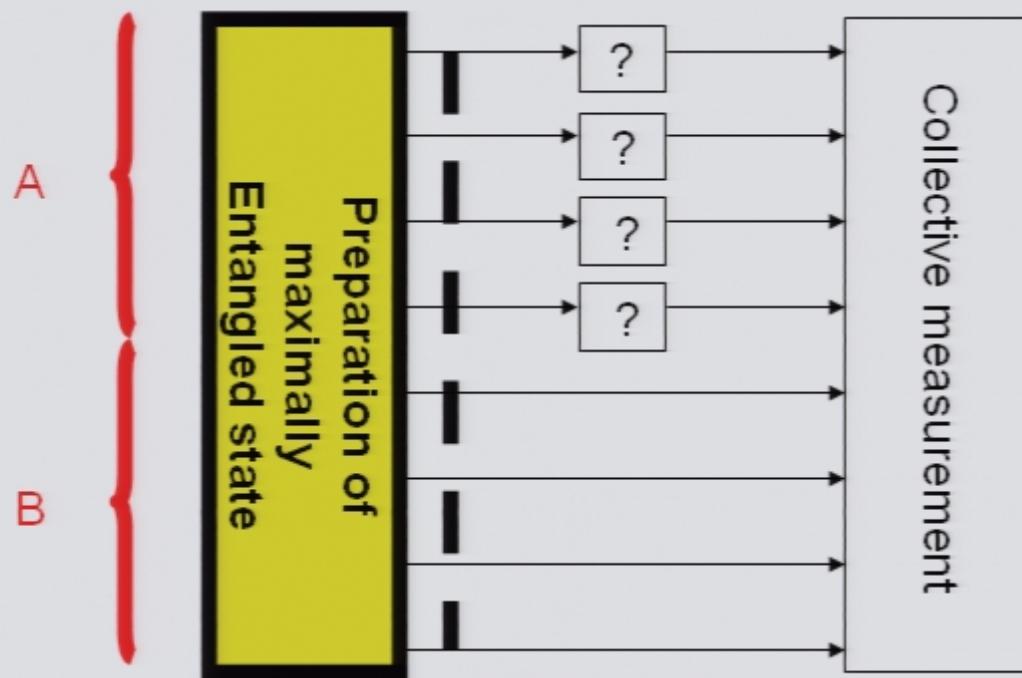
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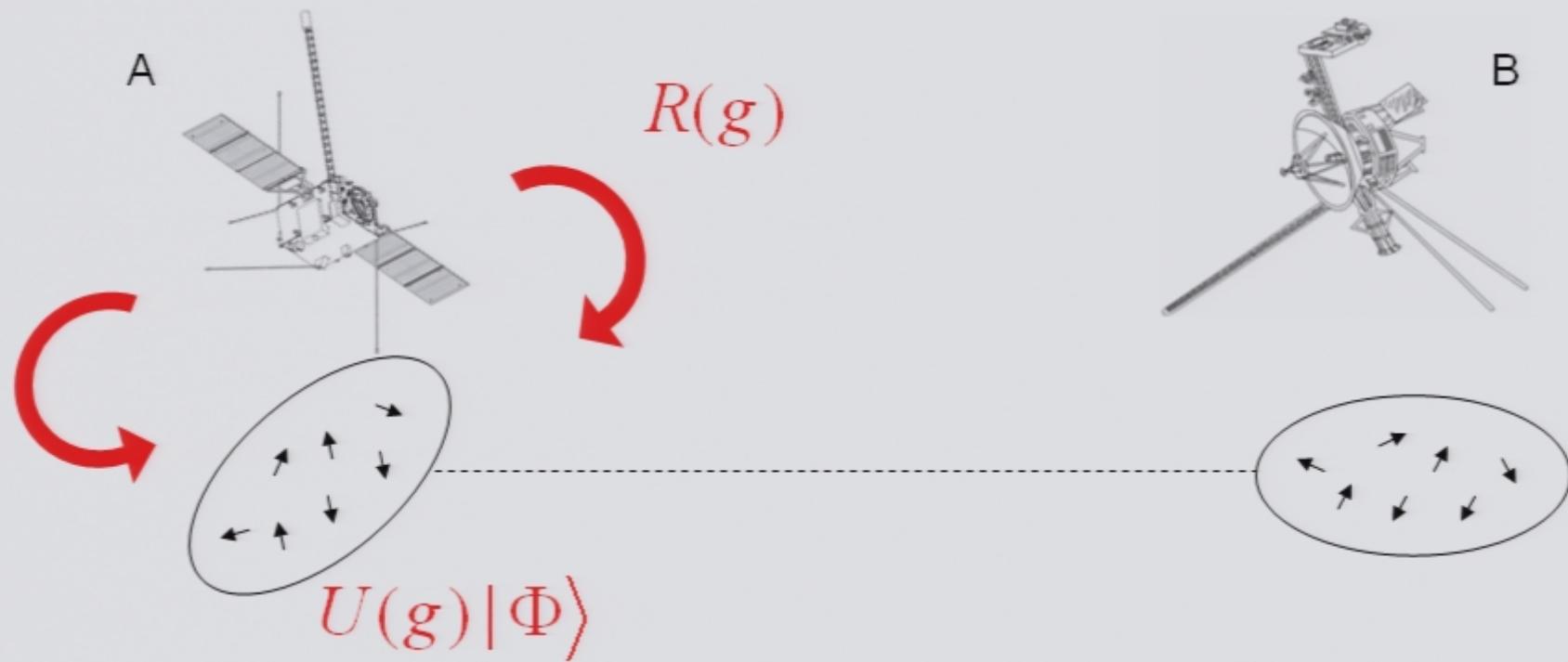
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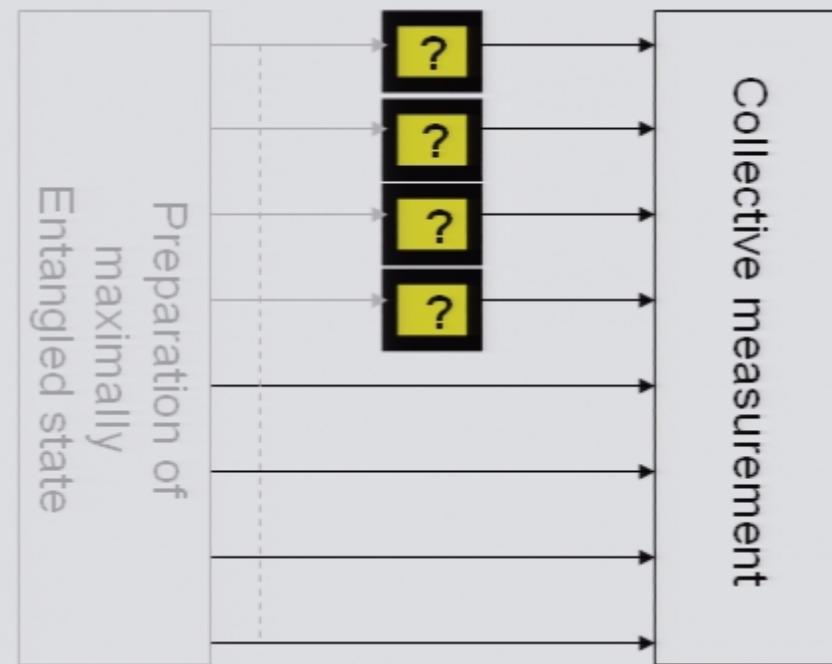
Journey

$$|\Phi\rangle \rightarrow U(g)|\Phi\rangle = u^{\otimes N}(g) \otimes I |\Phi\rangle$$



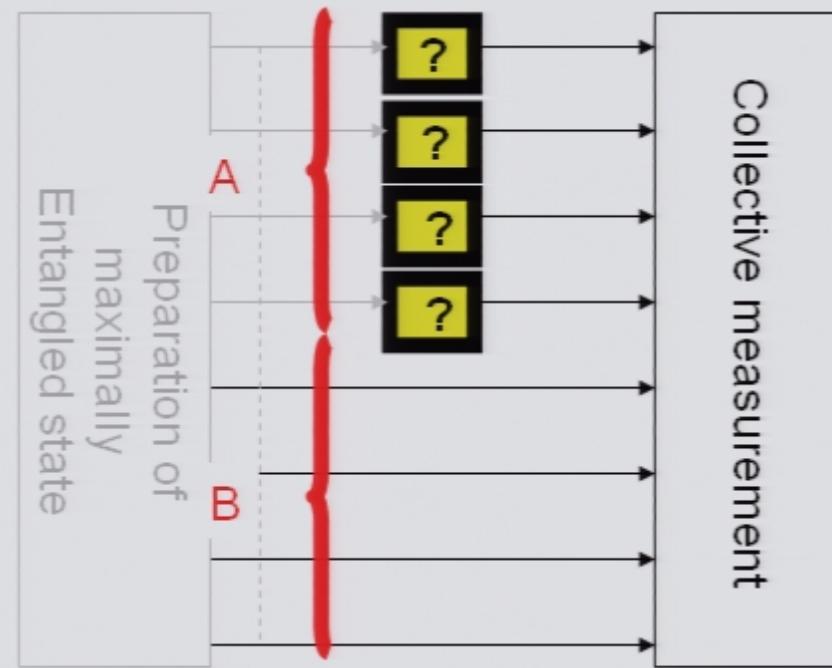
“Journey”

$$|\Phi\rangle \rightarrow U(g)|\Phi\rangle = u^{\otimes N}(g) \otimes I |\Phi\rangle$$



“Journey”

$$|\Phi\rangle \rightarrow U(g)|\Phi\rangle = u_A^{\otimes N}(g) \otimes I_B |\Phi\rangle$$

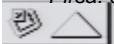
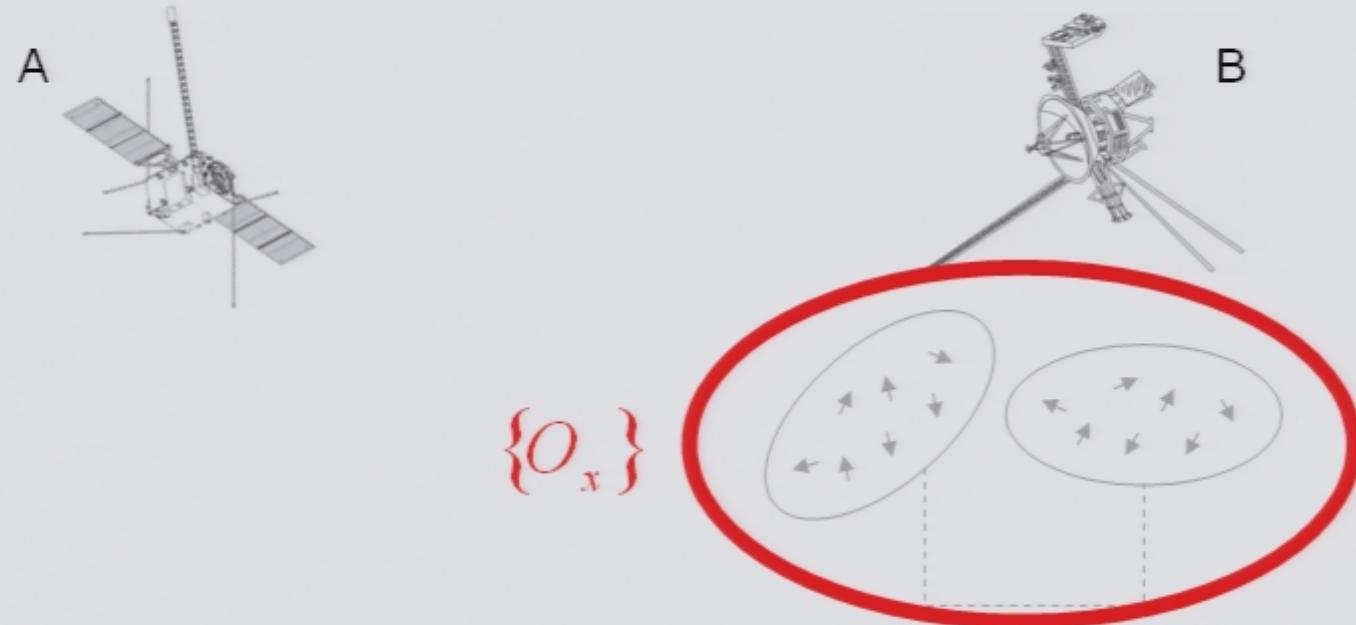




Measurement

$$O_x = c_x u^{\otimes N}(g_x) \otimes I - |\Psi\rangle\langle\Psi| - [u^{\otimes N}(g_x)]^+ \otimes I$$

$$|\Psi\rangle = \sum_j \sqrt{d_j} \sum_{m=-j}^j |jm\rangle |jm\rangle$$

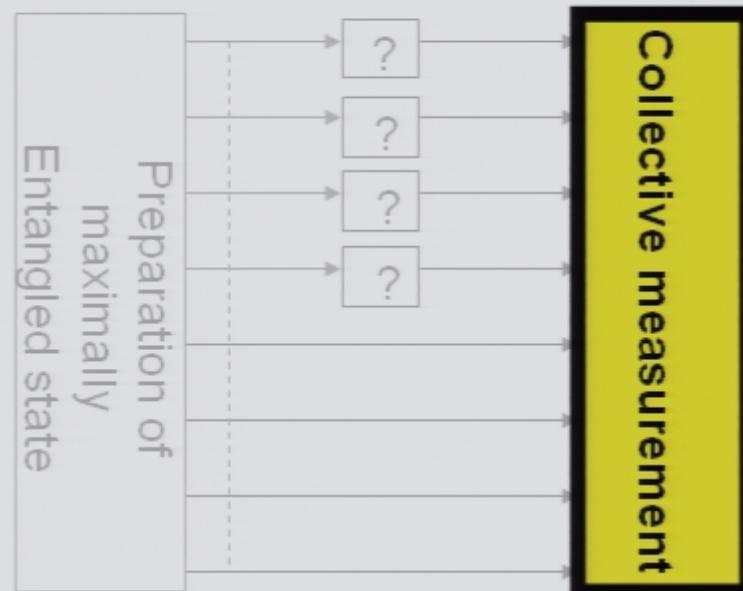




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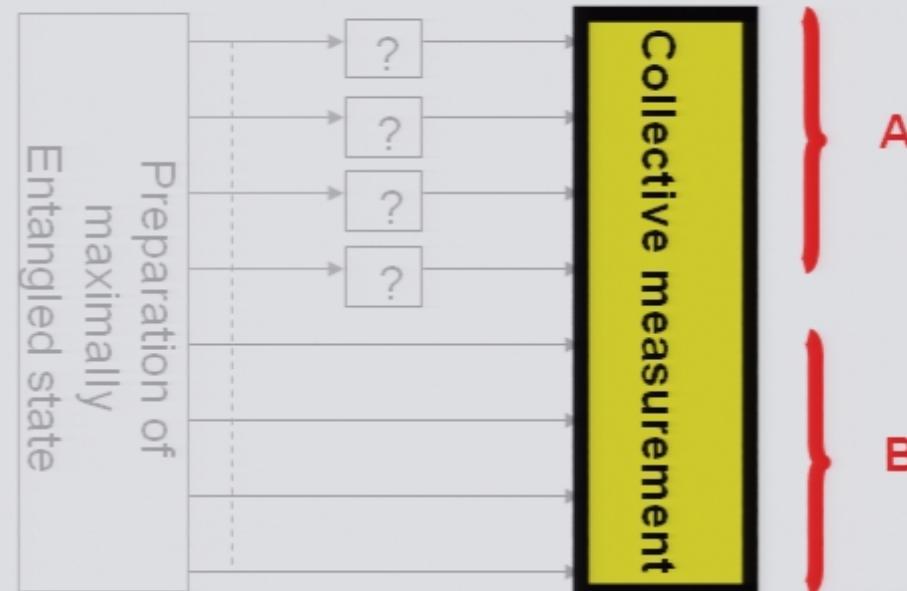




Measurement

$$O_x = c_x u_A^{\otimes N}(g_x) \otimes I_B |\Psi\rangle\langle\Psi| [u_A^{\otimes N}(g_x)]^+ \otimes I_B$$

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Measurement

A particular instance of

$$O_x = c_x u^{\otimes N}(g_x) \otimes I \quad |\Psi\rangle\langle\Psi| \quad [u^{\otimes N}(g_x)]^+ \otimes I$$

Is the continuous POVM

$$O(g) = u^{\otimes N}(g) \otimes I \quad |\Psi\rangle\langle\Psi| \quad [u^{\otimes N}(g)]^+ \otimes I$$

$$|\Psi\rangle = \sum_j \sqrt{d_j} \sum_{m=-j}^j |jm\rangle \langle jm|$$

Some technical details

N qubit states belong to $(\frac{1}{2})^{\otimes N}$

Clebsch-Gordan series for SU(2) $(\frac{1}{2})^{\otimes N} \rightarrow \bigoplus_{j=0, \frac{N}{2}}^{N/2} n_j \mathbf{j}$

where

$$n_j = \frac{2j+1}{N/2+j+1} \binom{N}{N/2+j}$$

$$n_j \geq d_j \quad \text{if} \quad j < \frac{N}{2} \quad \text{whereas} \quad n_j = 1 \quad \text{if} \quad j = \frac{N}{2}$$

Some technical details

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Acin, Jane & Vidal, PRA 64, 050302 (2001)

$$|\Phi\rangle = \sum_j a_j |\Phi^j\rangle = \sum_j \frac{a_j}{\sqrt{d_j n_j}} \sum_{m=-j}^j \sum_{\alpha=1}^{n_j} |jm;\alpha\rangle |jm;\alpha\rangle$$

But $\langle U^{(j)}(g) \otimes I | \Phi^j \rangle_{\forall g}$ has dimension d_j (rather than $n_j d_j$)

We can change bases so that $\langle U^{(j)}(g) \otimes I | \Phi^j \rangle_{\forall g}$ is in just 1 irreducible representation \mathbf{j} .

Holevo error/Fidelity

$$\langle \chi_1 \rangle_{\max}$$

$$\langle \chi_1 \rangle_{\max} = 1 + 2 \cos \frac{2\pi}{N+3}$$

$$a_j = \frac{2}{\sqrt{N+3}} \sin \frac{(2j+1)\pi}{N+3}$$

Holevo error/Fidelity

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Exact!!!

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Main ingredients of the proof:

- Schur's lemma
- Schwarz inequality
- $\langle \chi_1 \rangle \leq 1 + \alpha^T M \alpha$ M tridiagonal
- Characteristic polynomial is a combination of Chebyshev's polynomials
- Continuous POVM saturates bound

Holevo error/Fidelity

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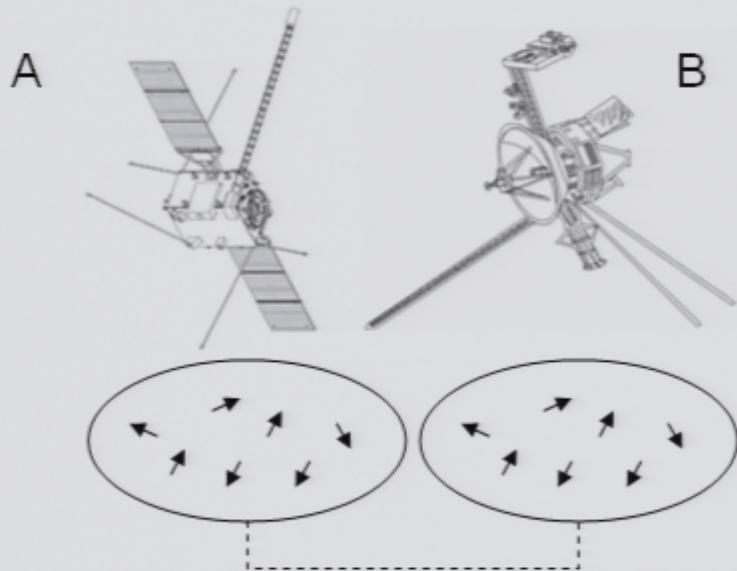
For large N we have

$$\langle \chi_1 \rangle_{\max} = 3 - \frac{4\pi^2}{N^2} + \dots$$

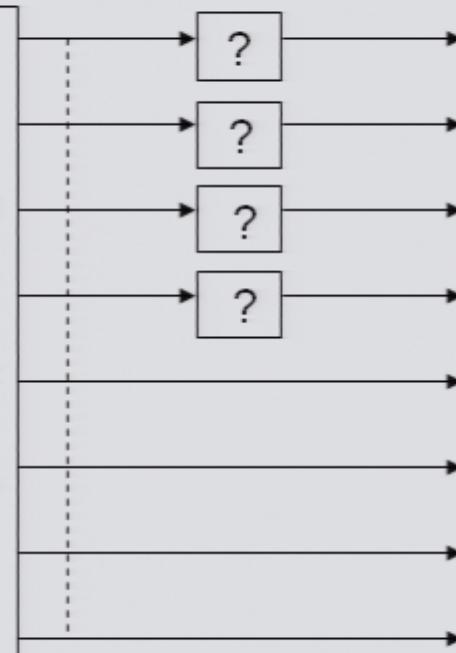
$$\langle h \rangle_{\min} = \frac{8\pi^2}{N^2} + \dots$$

$$\langle F \rangle_{\max} = 1 - \frac{\pi^2}{N^2} + \dots$$

Saving resources



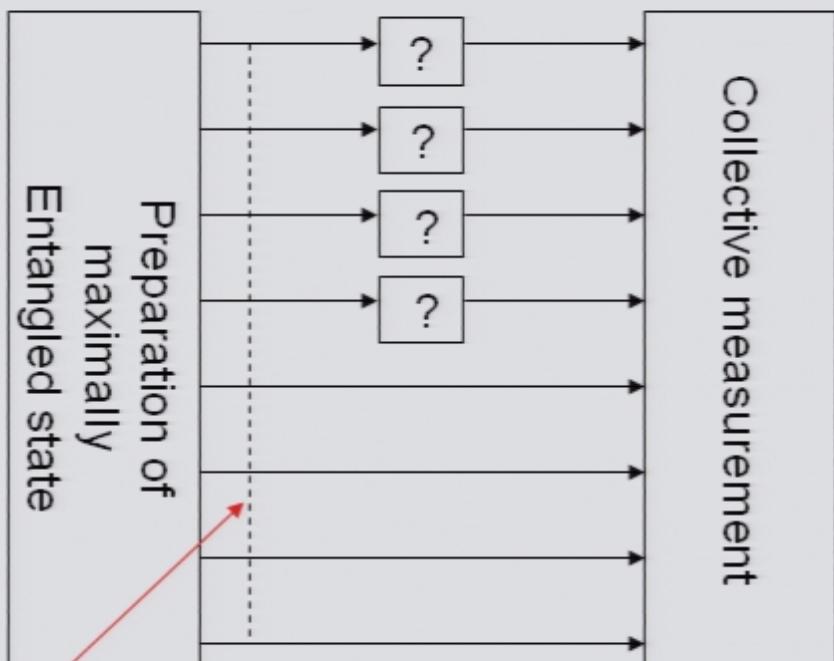
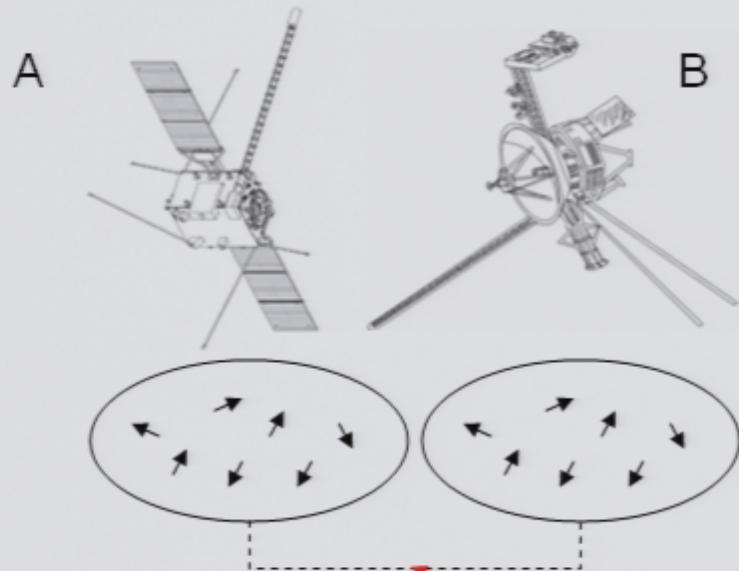
Preparation of
maximally
Entangled state



Collective measurement

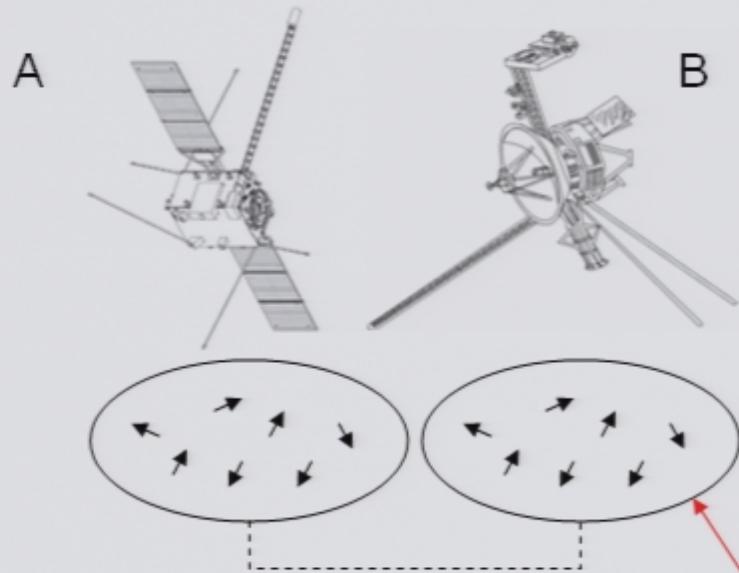
Shared entanglement + N extra spins

Saving resources

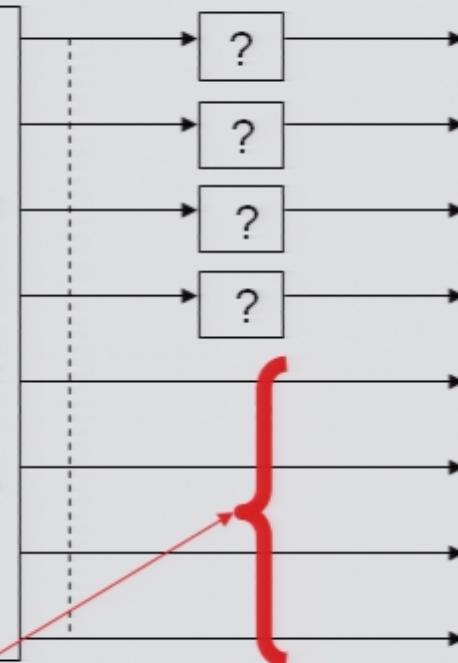


Shared entanglement + N extra spins

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Shared entanglement + N extra spins

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- How !!??

Saving resources

Recall

Clebsch-Gordan series for SU(2) $\left(\frac{1}{2}\right)^{\otimes N} \rightarrow \bigoplus_{j=0,\frac{1}{2}}^{\frac{N}{2}} n_j \mathbf{j}$

where

$$n_j \geq d_j \quad \text{if} \quad j < \frac{N}{2} \quad \text{whereas} \quad n_j = 1 \quad \text{if} \quad j = \frac{N}{2}$$

$$|\phi\rangle = \sum_j |\phi^j\rangle = \sum_j \sum_{m=-j}^j \sum_{\alpha=1}^{n_j} \phi_{m\alpha}^j |jm;\alpha\rangle$$

Saving resources

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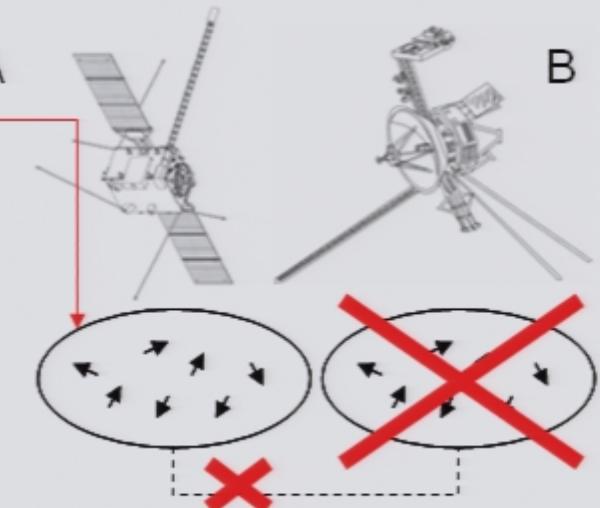
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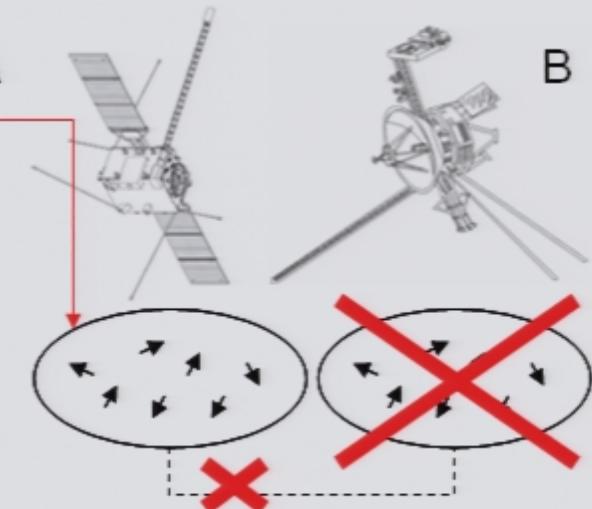
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But $\langle U^{(j)}(g) |\phi^j\rangle_{\forall g}$

belongs to d_j representations \mathbf{j} at most.



Saving resources *preparation*

$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle_A |jm\rangle_B \rightarrow |\phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle^{A \quad B}$$

α_m 1-1 map from

$$\{-j, -j+1, \dots, j\} \longrightarrow \{1, \dots, d_j\}$$

Rather than entangling two different parties, we entangle different degrees of freedom

Saving resources journey

$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle \langle jm| \rightarrow |\phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle$$

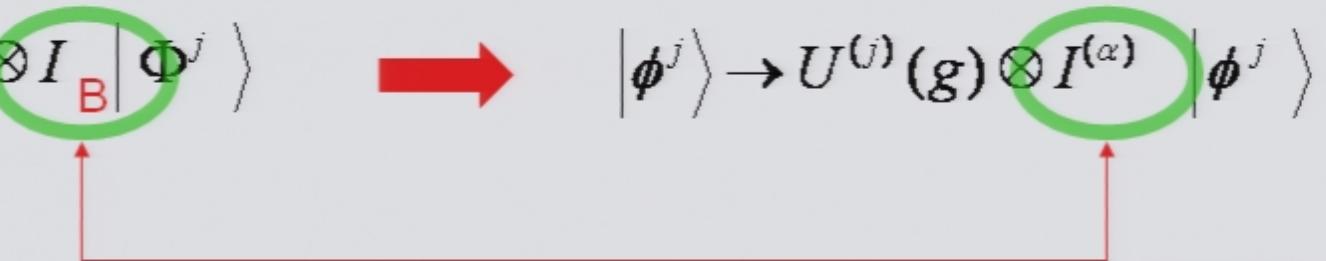
Just like j , α is a scalar under SU(2).

$$|\Phi^j\rangle \rightarrow U_A^{(j)}(g) \otimes I_B |\Phi^j\rangle \rightarrow |\phi^j\rangle \rightarrow U^{(j)}(g) \otimes I^{(\alpha)} |\phi^j\rangle$$

Saving resources *journey*

$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle \langle jm| \rightarrow |\phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle \langle jm; \alpha_m|$$

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Saving resources *measurement*

$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle \langle jm| \rightarrow |\phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle$$

$$|\Phi^j\rangle \xrightarrow{\textcolor{red}{U_A^{(j)}(g) \otimes I_B}} |\phi^j\rangle \rightarrow U^{(j)}(g) \otimes I^{(\alpha)} |\phi^j\rangle$$

$$|\Psi\rangle = \sum_j \sqrt{d_j} \sum_{m=-j}^j |jm\rangle \langle jm| \rightarrow |\psi\rangle = \sum_j \sqrt{d_j} \sum_{m=-j}^j |jm; \alpha_m\rangle$$

Saving resources *but.....*

$$n_j = 1 < d_j \quad \text{if} \quad j = N/2$$

$$|\Phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm\rangle \langle jm| \quad \rightarrow \quad |\phi\rangle = \sum_j \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle$$

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$$|\phi\rangle = \sum_{j < N/2} \frac{a_j}{\sqrt{d_j}} \sum_{m=-j}^j |jm; \alpha_m\rangle + a_{N/2} |N/2 \ N/2\rangle$$

Bagan, Baig & Muñoz-Tapia, PRL 87 257903 (2001)

Saving resources *as to the measurement.....*

$$|\psi\rangle = \sum_{j < N/2} \sqrt{d_j} \sum_{m=-j}^j |jm; \alpha_m\rangle + \sqrt{d_{N/2}} |N/2\ N/2\rangle$$

Saving resources as to the measurement.....

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This scheme is **not** optimal but saturates de optimal bound asymptotically

$$\langle \chi_1 \rangle_{\max} = 3 - \frac{4\pi^2}{N^2} + \dots$$

$$\langle h \rangle_{\min} = \frac{8\pi^2}{N^2} + \dots$$

$$\langle F \rangle_{\max} = 1 - \frac{\pi^2}{N^2} + \dots$$

Bagan, Baig & Muñoz-Tapia, quant-ph/0405082

Chiribella et al., quant-ph/0405095

Summary & conclusions

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- **Quantum systems can**
 - transmit spacial reference frames
 - be used for *reverse-engineering*



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Peres & Scudo, PRL 87, 167901 (2001); Bagan, Baig & Muñoz-Tapia, PRL 87, 257903 (2001)

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